CHAPTER 6
POWER COUPLING
AT FIBERS CONNECTIONS

Philippe DI BIN
University of Limoges
philippe.dibin@unilim.fr

Université de Limoges · Università degli Studi di Brescia · Aston University · Universidad del País Vasco / Euskal Herriko Unibertsitatea / The University of the Basque Country · Universitatea Tehnică din Cluj-Napoca

Introduction

Introduction

Connections between fibers will induce losses and power coupling between transmitted modes

→ Need to identify coupling conditions
→ Model of power coupling between fiber modes at connections

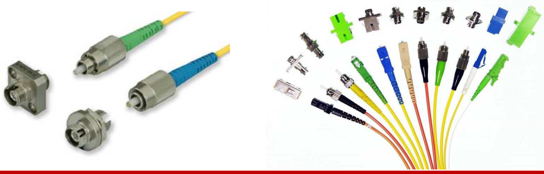
Aim of the chapter 6

→ Connections geometric defects definition
→ Power transfert between optical fibers
→ Connection losses in singlemode fibers (gaussian mode approximation)
→ Modal coupling at misaligned multimode fibers connections

Fiber connections

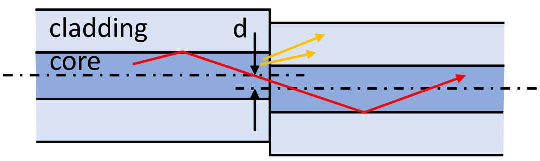
- Aims
  - Increasing the length of a fiber link
  - Allowing fiber network flexibility
  - Low loss
  - Reliability
- Technologies
  - Permanent link :
    Electrical arc fusion splices
  - Non-permanent link :
    Fiber connectors

http://www.next-up.org/France/Fibre_optique.php

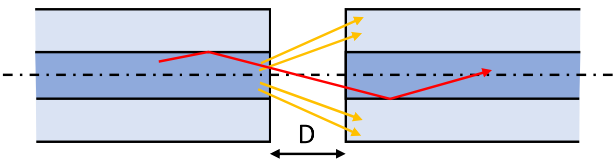


Fiber connections loss causes (1)

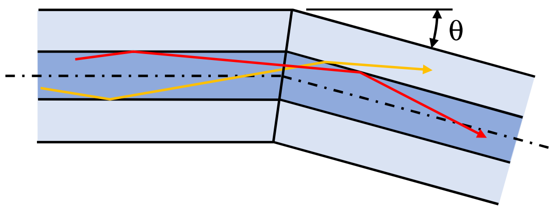
- Mechanical misalignments
  - Transversal offset d



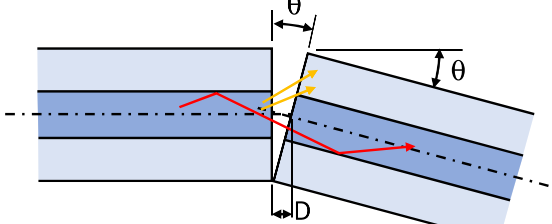
  - Longitudinal gap D



- Angular offset
  - ➢ Fusion splice (angle θ only)

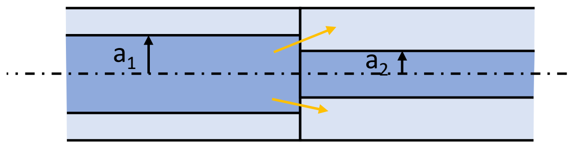


  - ➢ Connector (angle θ and gap D)

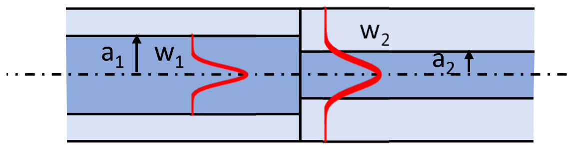


Fiber connections loss causes (2)

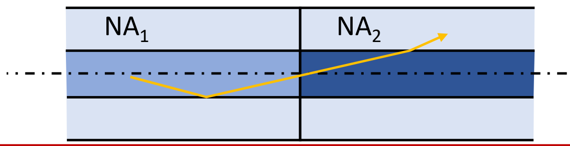
- Opto-geometrical mismatches
  - Core size mismatch (Multimode fiber)

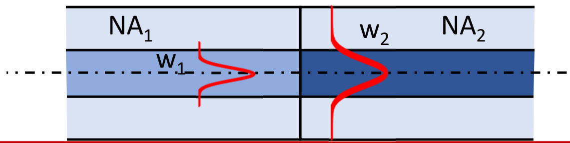


  - Mode field diameter mismatch (Singlemode fiber)



Numerical Aperture mismatch (Multimode fiber)





Power flow in optical fiber (1)

- Complex fields

$$\vec{\mathcal{E}}(x,y,z,t) = \vec{E}(x,y)e^{j(wt-\beta z)} = \vec{E}(x,y)e^{j\phi(z,t)}$$
$$\vec{\mathcal{H}}(x,y,z,t) = \vec{H}(x,y)e^{j(wt-\beta z)} = \vec{H}(x,y)e^{j\phi(z,t)}$$

- Poynting vector

$$\vec{S} = \mathcal{R}e\left(\vec{\mathcal{E}}\right) \wedge \mathcal{R}e\left(\vec{\mathcal{H}}\right) = \frac{1}{2}\left(\vec{\mathcal{E}} + \vec{\mathcal{E}}^*\right) \wedge \frac{1}{2}\left(\vec{\mathcal{H}} + \vec{\mathcal{H}}^*\right)$$
$$\vec{S} = \frac{1}{4}\left(\vec{E} \wedge \vec{H}^* + \vec{E}^* \wedge \vec{H} + \vec{E} \wedge \vec{H} e^{j2\phi(z,t)} + \vec{E}^* \wedge \vec{H}^* \, e^{-j2\phi(z,t)}\right)$$

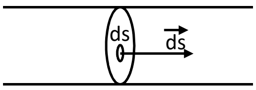
- Time averaged Poynting vector

$$< \vec{S} >= \frac{1}{T}\int_0^T \vec{S}\,dt = \frac{1}{4}\left(\vec{E} \wedge \vec{H}^* + \vec{E}^* \wedge \vec{H}\right) = \boxed{\frac{1}{2}\mathcal{R}e\big(\underbrace{\vec{E}(x,y)}_{V\,m^{-1}} \wedge \underbrace{\vec{\mathcal{H}}^*(x,y)}_{A\,m^{-1}}\big)}$$

$W\,m^{-2}$ *Surfacic power density*

Power flow in optical fiber (2)

- Power flow across the fiber section



$$< \vec{S} >= \begin{vmatrix} S_r \\ S_\theta \\ S_z \end{vmatrix}$$

The mean guided power

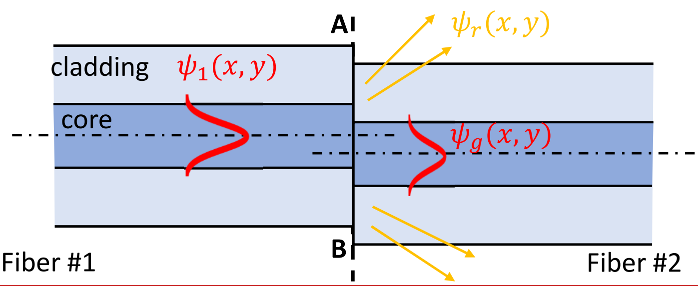
$$\bar{P} = \iint_{-\infty}^{+\infty} < \vec{S} > \vec{ds} = \iint_{-\infty}^{+\infty} S_z \, dx \, dy = \iint_{-\infty}^{+\infty} \frac{n}{2}\sqrt{\frac{\varepsilon_0}{\mu_0}} |E(x,y)|^2 \, dx \, dy$$

We define $\quad |\psi(x,y)|^2 = S_z = \frac{n}{2}\sqrt{\frac{\varepsilon_0}{\mu_0}} |E(x,y)|^2$

$$\boxed{\bar{P} = \iint_{-\infty}^{+\infty} |\psi(x,y)|^2 dx \, dy}$$

Fields overlap integral (1)

- Fields at the connection (plane AB)
  - Incident optical wave field : $\boldsymbol{\psi_1(x,y)}$ carrying a mean power $P_1 = \iint_{-\infty}^{+\infty} |\psi_1(x,y)|^2 dx \, dy$
  - Output mode field : $\boldsymbol{\psi_2(x,y)}$ with unknown amplitude and power
  - Guided output field : $\boldsymbol{\psi_g(x,y) = \varepsilon \, \psi_2(x,y)}$ carrying a power $P_g = |\varepsilon|^2 \iint_{-\infty}^{+\infty} |\psi_2(x,y)|^2 dx \, dy$
  - Radiative (unguided) field : $\boldsymbol{\psi_r(x,y)}$ carrying a mean power $P_r = \iint_{-\infty}^{+\infty} |\psi_r(x,y)|^2 dx \, dy$

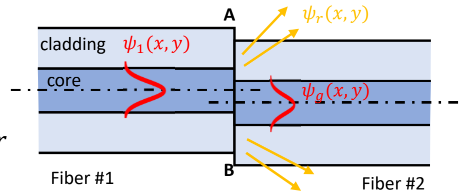


In AB plane

$$\boldsymbol{\psi_1(x,y) = \psi_g(x,y) + \psi_r(x,y)}$$

$$\boldsymbol{= \varepsilon \, \psi_2(x,y) + \psi_r(x,y)} \quad \text{(Eqn 1)}$$

Fields overlap integral (2)



- Energy conservation
  - The total output power is equal to the input power : $P_{output} = P_1 = P_g + P_r$

$$\iint_{-\infty}^{+\infty} |\psi_1(x,y)|^2 dx\, dy = |\varepsilon|^2 \iint_{-\infty}^{+\infty} |\psi_2(x,y)|^2 dx\, dy + \iint_{-\infty}^{+\infty} |\psi_r(x,y)|^2 dx\, dy$$

  - From equation 1 $\quad \psi_1(x,y) = \psi_g(x,y) + \psi_r(x,y) = \varepsilon\, \psi_2(x,y) + \psi_r(x,y)$

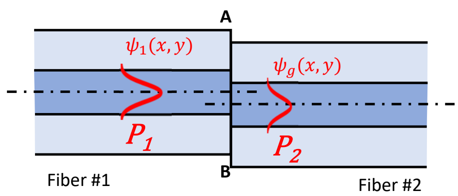
$$\iint_{-\infty}^{+\infty} |\psi_1(x,y)|^2 dx\, dy = \iint_{-\infty}^{+\infty} |\varepsilon\, \psi_2(x,y) + \psi_r(x,y)|^2 dx\, dy$$

$$= |\varepsilon|^2 \iint_{-\infty}^{+\infty} |\psi_2(x,y)|^2 dx\, dy + \iint_{-\infty}^{+\infty} |\psi_r(x,y)|^2 dx\, dy + \underbrace{\varepsilon \iint_{-\infty}^{+\infty} \psi_2(x,y)\, \psi_r^*(x,y)\, dx\, dy + \varepsilon^* \iint_{-\infty}^{+\infty} \psi_2^*(x,y)\, \psi_r(x,y)\, dx\, dy}_{=0}$$

$$\Leftrightarrow \iint_{-\infty}^{+\infty} \boldsymbol{\psi_2(x,y)\, \psi_r^*(x,y)\, dx\, dy = 0}$$

$\Leftrightarrow$ **$\psi_2(x,y)$ and $\psi_r(x,y)$ are orthogonal fields**
Guided and radiative modes are orthogonal

Fields overlap integral (3)



- Power transfer
  - Coupling power coefficient definition

$$\alpha^2 = \frac{P_g}{P_1} = \frac{P_2}{P_1} = \frac{|\varepsilon|^2 \iint_{-\infty}^{+\infty} |\psi_2(x,y)|^2 dx\, dy}{\iint_{-\infty}^{+\infty} |\psi_1(x,y)|^2 dx\, dy} \qquad \text{Eqn(2)}$$

  - Value of ε

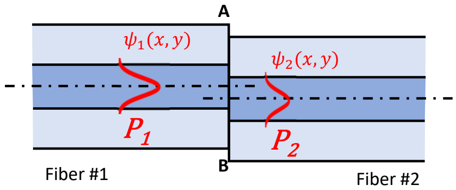
Multiplying by $\psi_2^*(x,y)$ and integrating eqn. 1 $\psi_1(x,y) = \varepsilon\, \psi_2(x,y) + \psi_r(x,y)$

$$\Leftrightarrow \iint_{-\infty}^{+\infty} \psi_1(x,y)\, \psi_2^*(x,y)\, dx\, dy = \varepsilon \iint_{-\infty}^{+\infty} |\psi_2(x,y)|^2 dx\, dy + \underbrace{\iint_{-\infty}^{+\infty} \psi_r(x,y)\, \psi_2^*(x,y) dx\, dy}_{=0}$$

$$\Leftrightarrow \varepsilon = \frac{\iint_{-\infty}^{+\infty} \psi_1(x,y)\, \psi_2^*(x,y)\, dx\, dy}{\iint_{-\infty}^{+\infty} |\psi_2(x,y)|^2 dx\, dy}$$

By introducing this expression in Eqn(2) we get the final expression of α^2

Fields overlap integral (4)



- Power transfer
  - Coupling power coefficient – Fields overlap integral

$$\alpha^2 = \frac{P_2}{P_1} = \frac{\left|\iint_{-\infty}^{+\infty} \psi_1(x,y)\,\psi_2^*(x,y)\,dx\,dy\right|^2}{\iint_{-\infty}^{+\infty}|\psi_1(x,y)|^2 dx\,dy \;\; \iint_{-\infty}^{+\infty}|\psi_2(x,y)|^2 dx\,dy}$$

Fields overlap integral

Power normalization terms

Fields are expressed in the plane AB and in the same coordinates axes

For normalized fields

$$\iint_{-\infty}^{+\infty} |\psi_1(x,y)|^2 dx\,dy = 1 \;\text{ and }\; \iint_{-\infty}^{+\infty} |\psi_2(x,y)|^2 dx\,dy = 1$$

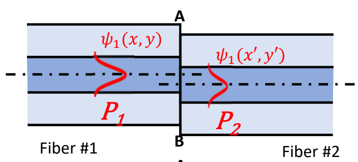
$$\alpha^2 = \frac{P_2}{P_1} = \left|\iint_{-\infty}^{+\infty} \psi_1(x,y)\,\psi_2^*(x,y)\,dx\,dy\right|^2$$

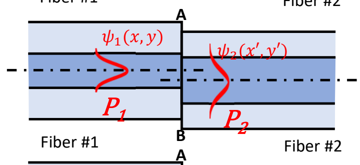
Normalized fields overlap integral

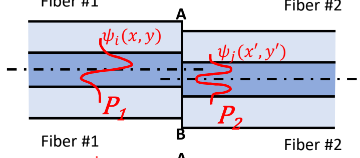
Fields overlap integral (5)

$$P_2 = \alpha^2\, P_1$$

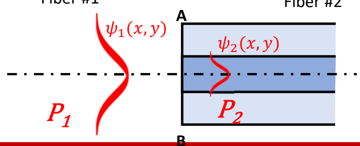
- Power transfer
  - Between fibers
    - Identical singlemode fibers :
      - mode of fiber #1 : $\boldsymbol{\psi_1(x,y)}$ with (x,y) coordinates of the fiber #1
      - mode of fiber #2 : $\boldsymbol{\psi_2(x',y')}$=$\boldsymbol{\psi_1(x',y')}$ with (x',y') coordinates of the fiber #2
    - Non – identical singlemode fibers
      - mode of fiber #1 : $\boldsymbol{\psi_1(x,y)}$
      - mode of fiber #2 : $\boldsymbol{\psi_2(x',y')}$
    - Multimodes fibers
      - mode #i of fiber #1 : $\boldsymbol{\psi_1(x,y) = \psi_i(x,y)}$
      - mode #j of fiber #2 : $\boldsymbol{\psi_2(x',y') = \psi_j(x',y')}$







- At fiber input face
  - Power injection in fiber
    - input field : $\boldsymbol{\psi_1(x,y)}$ with (x,y) coordinates of the fiber
    - mode of fiber : $\boldsymbol{\psi_2(x,y) = \psi_j(x,y)}$



Fields overlap integral (6)

- Power transfer
  - If $\psi_1(x,y) = \psi_2(x,y)$ then
    $\alpha^2 = 1 \Leftrightarrow P_2 = P_1$: no coupling loss
  - If $\psi_2(x,y)$ is orthogonal to $\psi_1(x,y)$ (guided or radiative modes)
    $\alpha^2 = 0 \Leftrightarrow P_2 = 0$: no coupled power to the second fiber mode ψ_2
  - Schwarz inequality

$$0 \leq \left| \iint_{-\infty}^{+\infty} \psi_1(x,y)\, \psi_2^*(x,y)\, dx\, dy \right|^2 \leq \iint_{-\infty}^{+\infty} |\psi_1(x,y)|^2 dx\, dy \iint_{-\infty}^{+\infty} |\psi_2(x,y)|^2 dx\, dy$$

$$\Leftrightarrow \boxed{0 \leq \alpha^2 \leq 1} \quad \Leftrightarrow \boxed{P_2 \leq P_1}$$

  - Independence of propagation direction (from fiber #1 to #2 or #2 to #1)

Connexion losses between singlemode fibers (1)

- Gaussian guided mode fields :

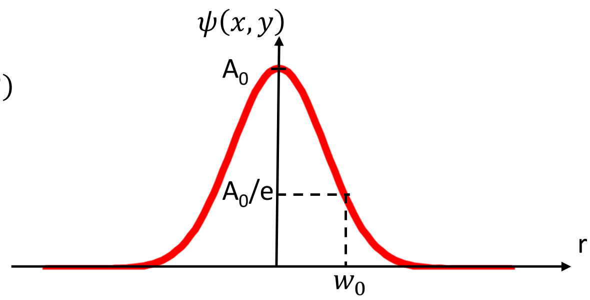
$$\psi(x,y) = A_0\, e^{-\frac{x^2+y^2}{w_0^2}} = A_0 e^{-\frac{r^2}{w_0^2}}$$

w_0 is the mode field radius @ 1/e of the maximal amplitude (1/e² of the maximal intensity)

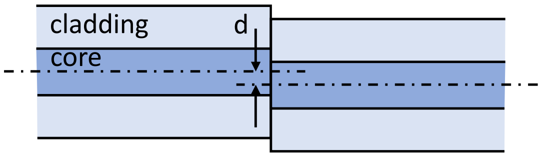
$$w_0 \approx a\,(0,65 + 1.619\, V^{-\frac{3}{2}} + 2.879\, V^{-6})$$

with $V = k\, a \sqrt{n_1^2 - n_2^2}$

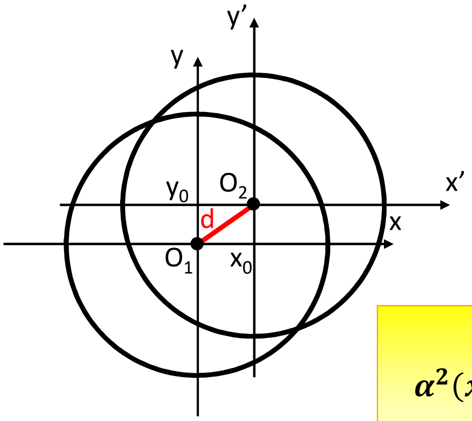
For $1.2 \leq V \leq 4$



Connexion losses between singlemode fibers (2)



- Tranversal offset (1)



$$\psi_1(x,y) = \psi(x,y)$$
$$\psi_2(x,y) = \psi(x - x_0, y - y_0)$$

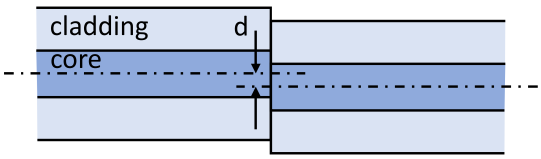
are real fields

$$d = \sqrt{x_0^2 + y_0^2}$$

$$\alpha^2 = \frac{P_2}{P_1} = \frac{\left|\iint_{-\infty}^{+\infty} \psi(x,y)\psi(x-x_0,y-y_0)\,dx\,dy\right|^2}{\iint_{-\infty}^{+\infty} |\psi(x,y)|^2 dx\,dy \quad \iint_{-\infty}^{+\infty} |\psi(x-x_0,y-y_0)|^2 dx\,dy}$$

$$\boldsymbol{\alpha^2(x_0,y_0) = \frac{\left(\psi(x,y) * \psi(x,y)\right)^2_{x_0,y_0}}{\left(\psi(x,y) * \psi(x,y)\right)^2_{0,0}}}$$

Connexion losses between singlemode fibers (3)

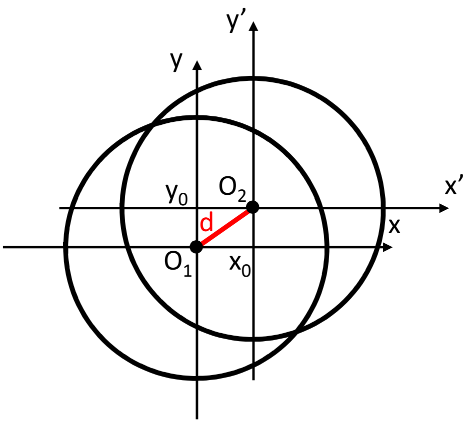


- Tranversal offset (2)

with

$$e^{-\pi x^2} \rightleftharpoons e^{-\pi u^2}$$

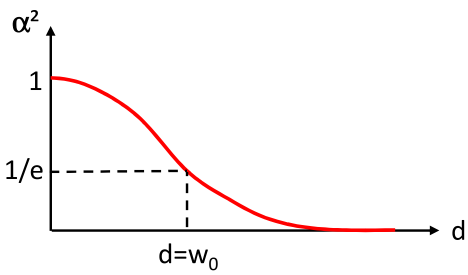
$$f(ax) \rightleftharpoons \frac{1}{|a|} F\left(\frac{u}{a}\right)$$



$$\left(\psi(x,y) * \psi(x,y)\right)^2_{x_0,y_0} = \frac{\pi w_0^2}{2} e^{-\frac{x_0^2+y_0^2}{w_0^2}}$$

$$\left(\psi(x,y) * \psi(x,y)\right)^2_{0,0} = \frac{\pi w_0^2}{2}$$

$$\boldsymbol{\alpha^2(x_0,y_0) = \frac{P_2}{P_1} = e^{-\frac{x_0^2+y_0^2}{w_0^2}} = e^{-\frac{d^2}{w_0^2}}}$$



Offset $\frac{d}{w_0}$	Case w_0=4.75µm	Losses (dB)
0.1	0.95 µm	0.04
0.15	1.4 µm	0.1
0.2	1.9 µm	0.17
0.3	2.85 µm	0.4
0.5	4.75 µm	1

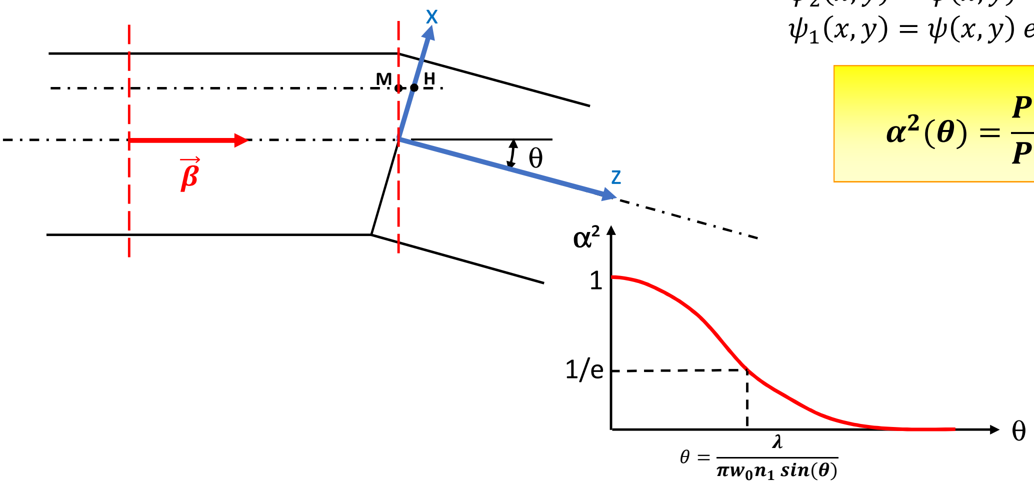
Connexion losses between singlemode fibers (4)

- Angular offset

In the coordinates axes of fiber #2 :

$\psi_2(x,y) = \psi(x,y) = \psi_2^*(x,y)$ because is real

$\psi_1(x,y) = \psi(x,y)\, e^{j\beta\, z(x,y))}$ with $z(x,y) = HM$

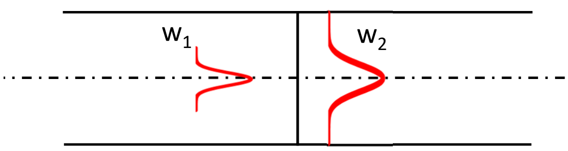


$$\alpha^2(\theta) = \frac{P_2}{P_1} = e^{-\left(\frac{\pi w_0 n_1 \sin(\theta)}{\lambda}\right)^2}$$

θ (deg) w_0=4.75µm L=1.55µm, n_1=1,45	Losses (dB)
0.1	0.0025
0.2	0.01
0.5	0.06
1	0,27
2	0,92

Connexion losses between singlemode fibers (5)

- Mode field diameter mismatch



In fiber #1

$$\psi_1(x,y) = A_0\, e^{-\frac{x^2+y^2}{w_1^2}}$$

In fiber #2

$$\psi_2(x,y) = A_0\, e^{-\frac{x^2+y^2}{w_2^2}}$$

$$\alpha^2 = \frac{P_2}{P_1} = \left(\frac{2w_1w_2}{w_1^2+w_2^2}\right)^2$$

$\frac{w_2}{w_1}$	Losses (dB)
1.1	0.04
1.3	0.3
1.63	1
2	2

Connexion losses between singlemode fibers (6)

- Triple defects connection (transversal+angular+mode mismatch)

$$\boldsymbol{\alpha^2(d,\theta,w_1,w_2) = \frac{P_2}{P_1} = \left(\frac{2w_1w_2}{w_1^2+w_2^2}\right)^2 e^{-\left(\frac{\pi w_0 n_1 \sin(\theta)}{\lambda}\right)^2} e^{-\frac{d^2}{w_m^2}}}$$

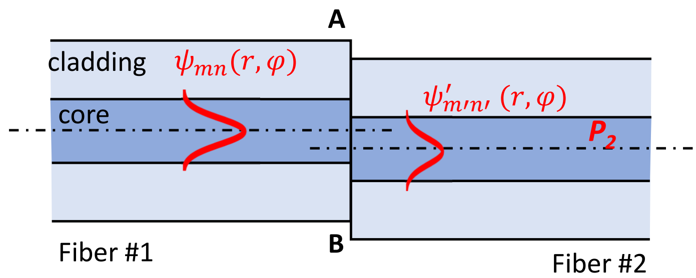
With

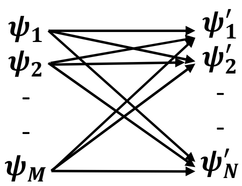
$$w_m^2 = \frac{w_1^2+w_2^2}{2}$$

$$w_0^2 = \frac{2w_1^2w_2^2}{w_1^2+w_2^2}$$

Connexion losses between multimode fibers (1)

- Power transfer at connection
  - Multimode fibers connection
    - M $LP_{m,n}$ modes guided in fiber #1
      $\psi_{mn}(r,\varphi)$
    - N $LP_{m',n'}$ modes guided in fiber #2
      $\psi'_{m'n'}(r,\varphi)$
    - MxN power coupling coefficients



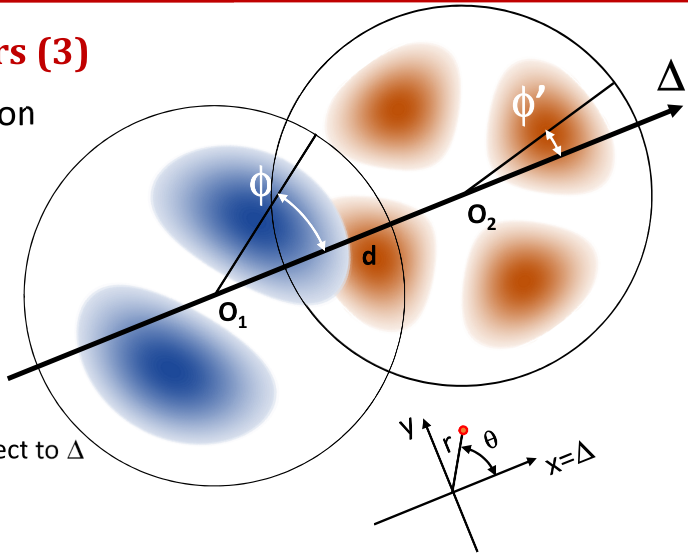


$$\left(\alpha_{mn}^{m'n'}\right)^2 = \frac{P_{m'n'}}{P_{mn}} = \left|\iint_{-\infty}^{+\infty} \psi_{mn}(r,\varphi)\,\psi'^{*}_{m'n'}(r,\varphi)\, r\, dr\, d\varphi\right|^2$$

$$P_2 = \sum_{j=1}^{N} P_{m'n'} = \sum_{i=1}^{M}\sum_{j=1}^{N} (\alpha_{mn}^{m'n'})^2 P_{mn}$$

Connexion losses between multimode fibers (3)

- Power coupling at transversally misaligned connection



  - Connection description
    - $d=O_1O_2$
    - Δ : axis of the transversal misalignment
    - ϕ : azimuthal orientation of mode LP_{mn} (fiber #1) with respect to Δ

$$\Psi_{mn} = R(r).\cos(m(\theta - \phi))$$

    - ϕ' : azimuthal orientation of mode $LP_{m'n'}$ (fiber #2) with respect to Δ

$$\Psi'_{m'n'} = R'(r).\cos(m'(\theta - \phi'))$$

r, θ } Cylindrical coordinates

Connexion losses between multimode fibers (3)

- Fields decomposition

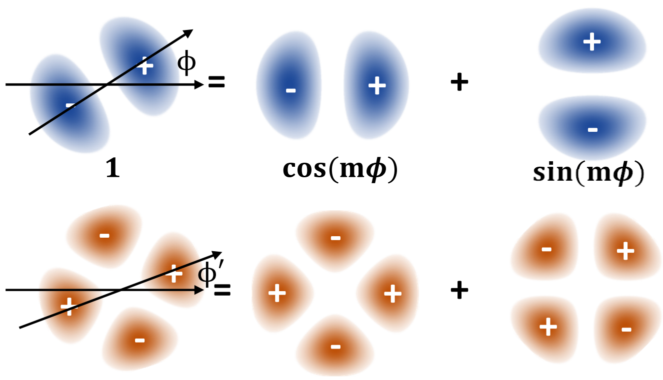
$$\Psi_{mn} = R(r)\,.\,(\cos(m\theta)\,\mathbf{\cos(m\phi)} + \sin(m\theta)\,\mathbf{\sin(m\phi)})$$

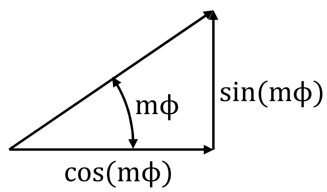
$$\Psi'_{m'n'} = R'(r).(\cos(m'\theta)\,\mathbf{\cos(m'\phi')} + \sin(m'\theta)\,\mathbf{\sin(m'\phi')})$$

Even modes **Odd modes**

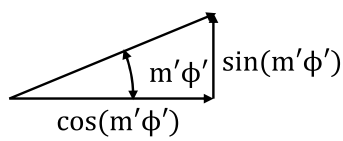
Input fiber

Output fiber





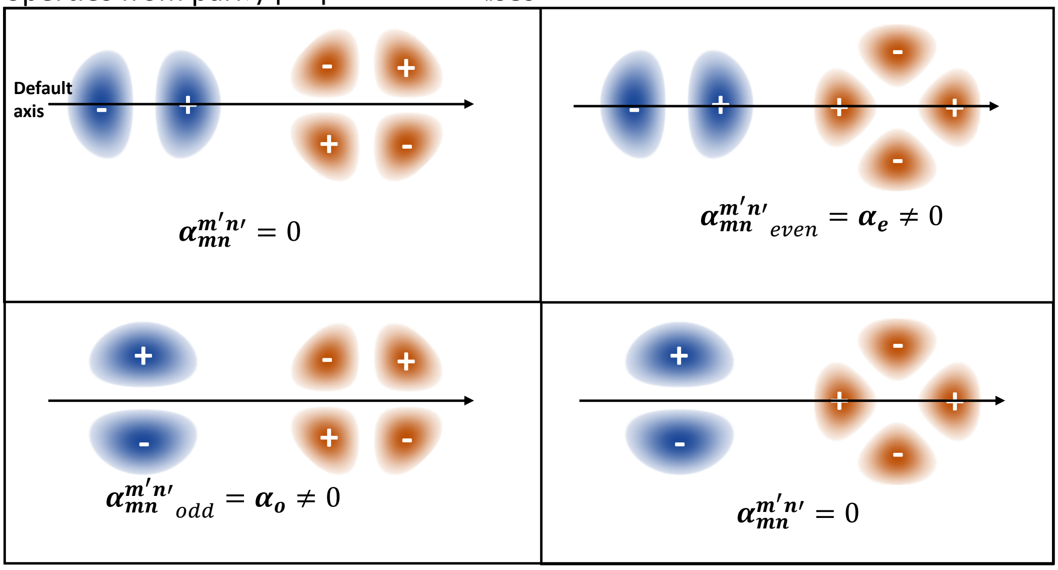
Fresnel representation



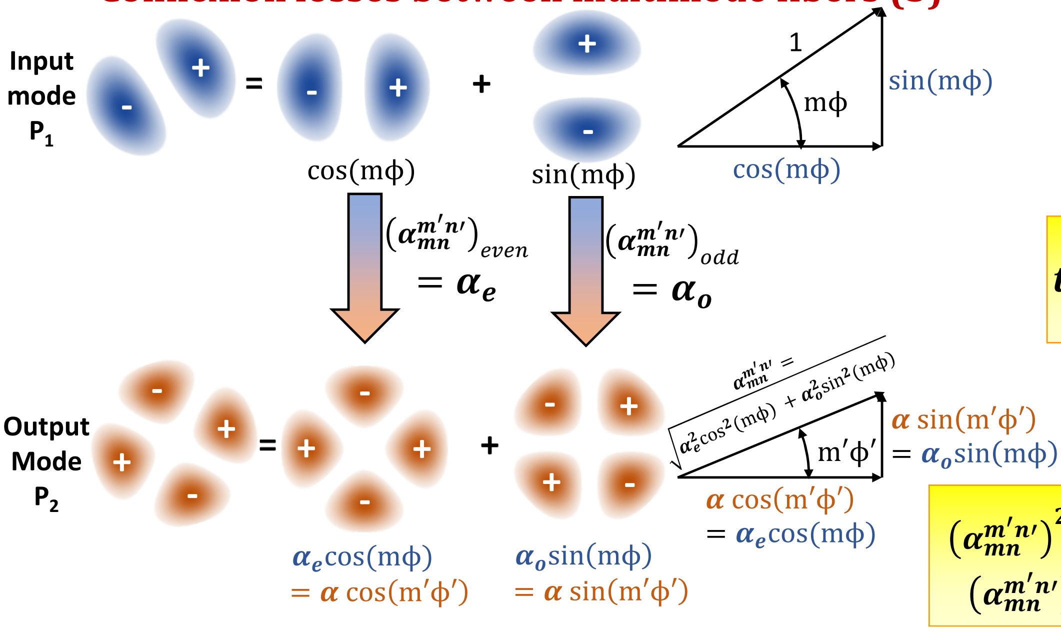
Connexion losses between multimode fibers (4)

$$\alpha^2 = \frac{P_2}{P_1} = \left| \iint_{-\infty}^{+\infty} \psi_1(x,y)\, \psi_2^*(x,y)\, dx\, dy \right|^2$$

- Coupling properties from parity properties – 4 cases



Connexion losses between multimode fibers (5)



$$P_2 = \left(\alpha_{mn}^{m'n'}\right)^2 P_1$$

$$\tan(m'\phi') = \frac{\alpha_o}{\alpha_e} \tan(m\phi)$$

The « tangents law »

$$\left(\alpha_{mn}^{m'n'}\right)^2 = \alpha_e^2 \cos^2(m\phi) + \alpha_o^2 \sin^2(m\phi)$$
$$\left(\alpha_{mn}^{m'n'}\right)^2 = \alpha_o^2 + \left(\alpha_e^2 - \alpha_o^2\right) \cos^2(m\phi)$$