

NUMERICAL TECHNIQUES

In particular, we consider the split-step Fourier method. To understand the philosophy behind the split-step Fourier method, it is useful to write the NLSE in the form

$$\frac{\partial F}{\partial z} = \left(\hat{D} + \hat{N} \right) F$$

where \hat{D} is a differential operator that accounts for dispersion (but it can take also into account for absorption and/or higher order dispersions) in a linear medium;

\hat{N} is a nonlinear operator that governs the effect of fiber nonlinearity on pulse propagation, i.e. Kerr nonlinearity (but it can take also into account other nonlinear phenomena, like Raman effects, etc.).

Thus

$$\hat{D} = -i \frac{\beta''}{2} \frac{\partial^2}{\partial t^2}$$

$$\hat{N} = i \gamma |F|^2$$

In general, dispersion and nonlinearity
act together along the fiber

The split-step Fourier Method obtains an
approximate solution by assuming that

In propagating the envelope over a small
distance h , the dispersive and nonlinear
effects can be considered to act independently.

More specifically, propagation from (z) to
 $(z+h)$ is carried out in two steps.

In the first step, the nonlinearity acts alone,
and $\hat{D} = C$.

In the second step, dispersion acts alone,
and $\hat{N} = 0$.

Mathematically:

$$\bar{F}(z+h, t) \approx \exp(h\hat{D}) \exp(h\hat{N}) F(z, t)$$

It's interesting to underline that the exponential operator $\exp(h\hat{D})$ can be evaluated in the Fourier domain using the prescription

$$\exp(h\hat{D}) F(z, t) = F_T^{-1} \exp[h\hat{D}(i\omega)] F_T F(z, t)$$

where F_T denotes the Fourier-Transform operation, $\hat{D}(i\omega)$ is obtained by replacing the differential operator $\frac{\partial}{\partial t} \rightarrow i\omega$

$$\frac{\partial}{\partial t} \xrightarrow{\text{by}} i\omega \quad \text{and} \quad \frac{\partial^2}{\partial t^2} \xrightarrow{\text{by}} (i\omega)^2 = -\omega^2$$

where ω is the frequency in the Fourier domain.

As $\hat{D}(i\omega)$ is just a number in the Fourier space, the evaluation of the dispersive term is simple.

The use of the FFT algorithm makes numerical evaluation fast.

To estimate the accuracy of the split-step method, we note that formally exact solution of the NLSE is given by the expression

$$F(z+h, t) = \exp [h(\hat{D} + \hat{N})] F(z, t)$$

assuming \hat{N} z -independent.

AT This point, it's useful To recall Thc
Baker-Hausdorff Formula For operators \hat{a} and \hat{b} :

$$\exp(\hat{a}) \exp(\hat{b}) = \exp\left(\hat{a} + \hat{b} + \frac{1}{2} [\hat{a}, \hat{b}] + \frac{1}{12} [\hat{a} - \hat{b}, [\hat{a}, \hat{b}]] + \dots\right)$$

where $[\hat{a}, \hat{b}] = \hat{a}\hat{b} - \hat{b}\hat{a}$.

IF we have two commuting operators $[\hat{a}, \hat{b}] = 0$

IF NOT, $[\hat{a}, \hat{b}] \neq 0$

A comparison of the expression

$$F(z+h, t) \approx \exp(h\hat{D}) \exp(h\hat{N}) F(z, t) *$$

and the expression

$$\bar{F}(z+h, t) = \exp[h(\hat{D} + \hat{N})] F(z, t)$$

shows that the split-step method (a) ignores the noncommuting nature of operators \hat{D} and \hat{N} .

By using the equation

$$\exp(\hat{a}) \exp(\hat{b}) = \exp\left(\hat{a} + \hat{b} + \frac{1}{2} [\hat{a}, \hat{b}] \dots\right)$$

with $\hat{a} = h \hat{D}$ and $\hat{b} = h \hat{N}$, The dominant error is found to result from the term

$$\frac{1}{2} [h^2, [\hat{D}, \hat{N}]]$$

Thus, the split-step Fourier method is accurate to second order in the step size $\frac{h}{\pi}$.

// I want to emphasize that taking \underline{h} small,
we have small errors

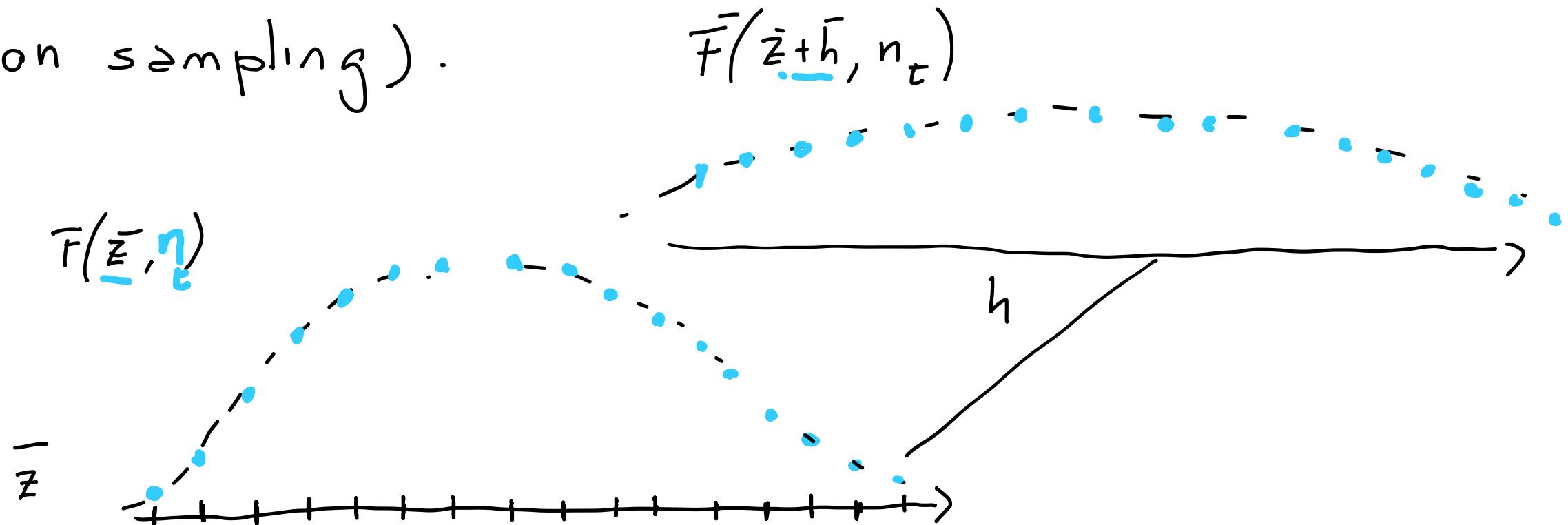
In general, the implementation of the method is

$$\bar{F}(z+h, t) = \exp[h(\hat{D} + \hat{N})] F(z, t) \underset{\approx}{\sim} \exp(h\hat{D}) \exp(h\hat{N}) F(z, t)$$

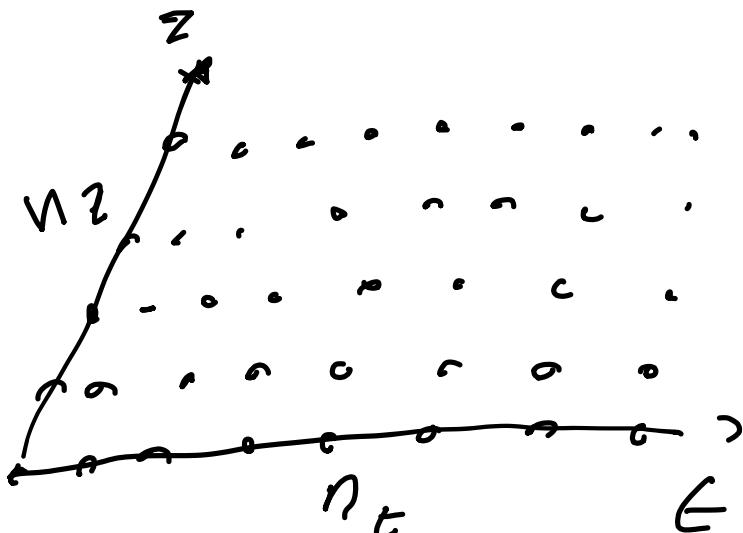
To propagate from (z) to $(z+h)$ we consider

- 1) First step: dispersion (or viceversa)
- 2) second step: nonlinear effect

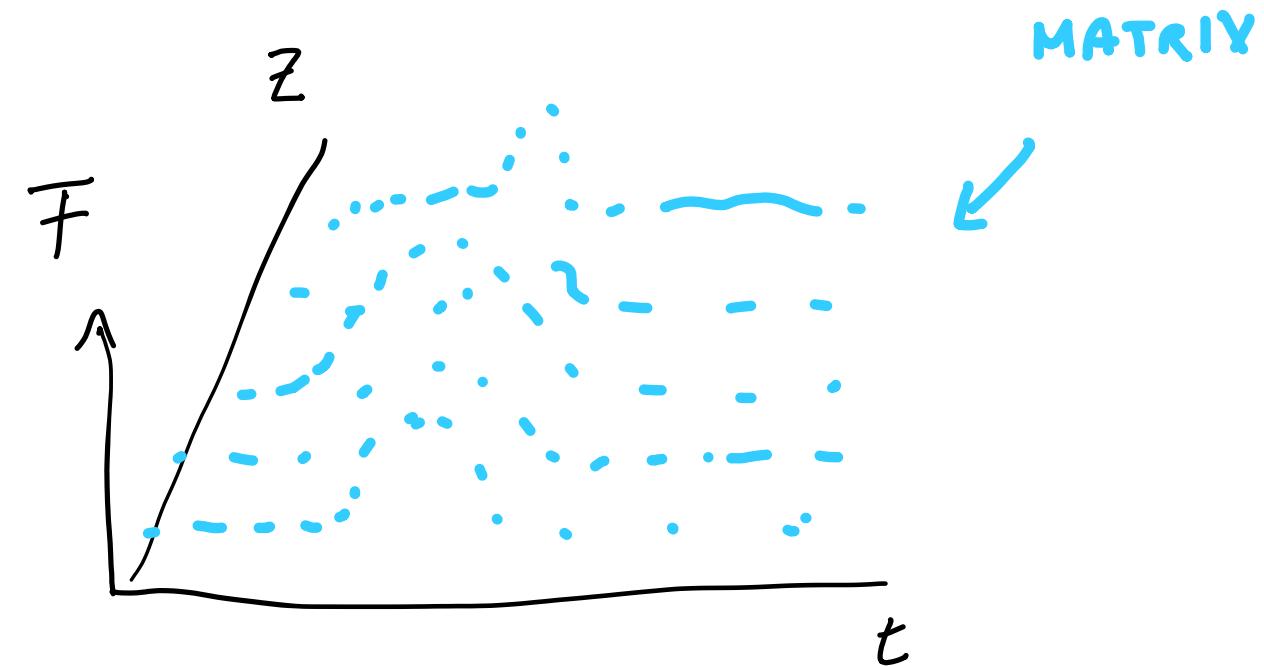
As to the envelope $F(z, t)$, take into account
that the profile is usually recorded in a
vector or matrix, by sampling F in a
good manner (i.e., ref. to Shannon Theorem
on sampling).



In general I want to represent and elaborate $F(z, t)$ in the space (z, t) , thus we have to take into account matrices



$F(z, t)$ REPRESENTED BY



$F(n_z, n_t)$

In details, as to the dispersive step, how we solve the equation

$$\frac{\partial F(z,t)}{\partial z} = -i \frac{B''}{2} \frac{\partial^2 F(z,t)}{\partial t^2}$$

We consider the Fourier Transform of F

$$\hat{F}(z,\omega) = \int_{-\infty}^{+\infty} F(z,t) e^{-i\omega t} dt$$

Thus

$$\frac{\partial \hat{F}(z, \omega)}{\partial z} = \left(i \frac{\omega^2 \beta''}{2} \right) \hat{F}(z, \omega)$$

The equation can be expressed only by $F(z, t)$

$$\hat{F}(z+h, \omega) = \hat{F}(z, \omega) \cdot \exp \left[+i \frac{\omega^2}{2} \beta'' \cdot h \right]$$

$$F(z+h, t) = \mathcal{J}^{-1} \left[\hat{F}(z+h, \omega) \right]$$

As to nonlinear term, how to solve

$$\frac{\partial F(z, t)}{\partial z} = i \gamma |F|^2 F$$

It's possible to demonstrate that $|F|^2$ is
z-independent

$$F(z+h, t) = F(z, t) \cdot \exp \left[i \gamma |F(z, t)|^2 \cdot h \right]$$

↑
step in z .

SCHEME FOR NUMERICAL ALGORITHM

- 1) INITIALIZATION
- 2) DEFINITION OF TEMPORAL AND SPATIAL COORDINATES
DEFINITION OF FREQUENCY DOMAIN
DEFINITION OF QUANTITIES (ENVELOPES ETC.)
- 3) FGR : $m:1:N_z$
 - dispersive step
 - nonlinear step
 - save data $F(n_t)$, every n_z
- 4) DISPLAY THE RESULTS

In next lectures our goal will be the implementation, the numerical implementation of this Technique.