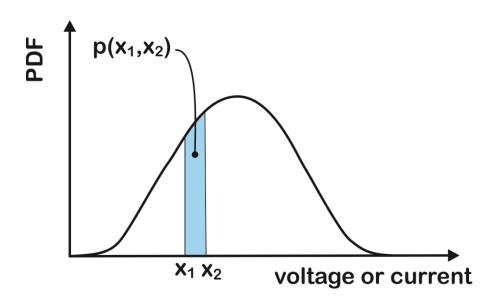
Probability density function



- □ a random signal (the noise) can be described just statistically: the PDF is a good descriptor
- \Box by definition PDF $p(x_1, x_2)$ is the probability that the actual value of x is included in the interval $x_1 \div x_2$

☐ therefore the expected value, the average, is: $\bar{x} = E(x) = \int_{-\infty}^{+\infty} x \cdot PDF(x) \cdot dx$ in case of electronic noise the average can be assumed equal to zero

 \Box the variance is $\sigma^2 = E(x - \bar{x})^2 = \int_{-\infty}^{+\infty} (x - \bar{x})^2 \cdot PDF(x) \cdot dx = \overline{x^2} - (\bar{x})^2$



Noise power

- \square a voltage V over a resistive load of resistance R dissipates o power equal to V^2/R , in the same way a current I flowing in a resistor of resistance R dissipates a power equal to RI^2
- \Box if $R = 1 \Omega$ the dissipated powers becomes V^2 and I^2
- \Box by convention, the noise power is implicitly defined as the power dissipated by the noise when $R=1~\Omega$, therefore we assume V^2 or I^2 as the power of a voltage noise or current noise respectively

$$\Box E(x - \bar{x})^2 = \int_{-\infty}^{+\infty} (x - \bar{x})^2 \cdot PDF(x) \cdot dx = \overline{x^2} - (\bar{x})^2$$
 (1)

(x represent either a voltage or a current)

 \square assuming $\bar{x} = 0$ (offset = 0) we can say that $\overline{x^2}$ represents the average power of the noise statistically described by the variable x of eq. (1)



Power spectral density

- ☐ a random signal (the noise) can be described in the frequency-domain, however, since the time-domain waveform is not defined, also the Fourier transform of the random signal can not be defined
- \square we can define the function *power spectral density* $S_x(f)$ (one sided) of a signal x such that the signal power P_{12} associated to the frequency band $f_1 \div f_2$ is given as:

$$P_{12} = \int_{f_1}^{f_2} \mathbf{S}_{x}(f) \cdot df$$

- \Box the total power is: $P_{tot} = \int_0^{+\infty} S_x(f) \cdot df$
- ☐ the total power can also be evaluated in the time-domain as

$$P_{tot} = \lim_{T \to \infty} \frac{1}{2T} \int_{-T}^{T} x^2(t) dt = \overline{x^2(t)}$$
 (offset = 0)

$$\Box$$
 therefore $\int_0^{+\infty} S_x(f) \cdot df = \overline{x^2(t)}$



Power spectral density

- □ the Wiener–Khinchin theorem (wide-sense stationary process)
 - given the autocorrelation function

$$R_{xx}(\tau) = \lim_{T \to \infty} \frac{1}{2T} \int_{-T}^{T} x(t)x(t+\tau)dt$$

 the power spectral density (two-sided) coincides with the Fourier transform of that function:

$$S_{x}(f) = \int_{-\infty}^{\infty} R_{xx}(\tau) e^{-j2\pi f \tau} d\tau$$

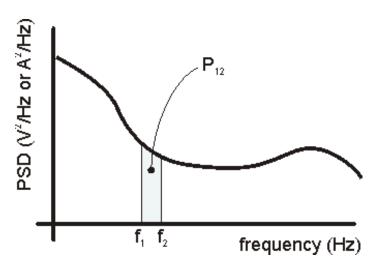
□total noise power:

$$\overline{x^2(t)} = \int_{-\infty}^{+\infty} S_x(f) df = 2 \int_0^{+\infty} S_x(f) df = \int_0^{\infty} \mathbf{S}_x(f) df$$

□ noise power in a given bandwidth: $BW: \overline{x_{BW}^2(t)} = \int_{f_1}^{f_2} S_{\chi}(f) df_{\chi}(f)$

Frequency distribution of the noise

□ the power spectral density can be seen as the power dissipated on a resistive load with a resistance of 1Ω per unit band, that is in a band of 1Hz



- □ the unit of measurement of the power spectral density is $\frac{\lfloor V \rfloor^2}{\lfloor Hz \rfloor}$ for a voltage noise and $\frac{\lfloor A \rfloor^2}{\lfloor Hz \rfloor}$ for a current noise
- □ in case we want to represent the noise with amplitude spectra the unit of measurement will be $\frac{[V]}{[\sqrt{Hz}]}$ for voltage noise and $\frac{[A]}{[\sqrt{Hz}]}$ for current noise



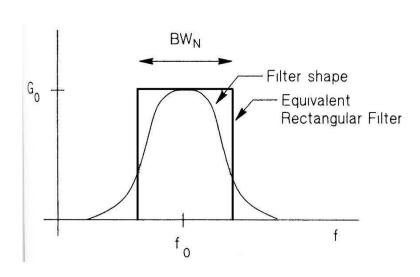
Noise equivalent bandwidth

☐ How much noise will be present at the output of a filter?

$$P_n = \int_0^\infty G(f) \cdot N_0 \cdot df = N_0 \cdot \int_0^\infty G(f) \cdot df$$

 N_0 = input noise PSD

G(f) =power gain of the filter



 \Box at the output of an ideal (rectangular response) filter having a power gain G_0 and a bandwidth BW_N the noise power should be:

$$P_n = N_0 \cdot G_0 \cdot BW_N$$

□ the noise equivalent bandwidth is: $BW_N = \frac{1}{G_0} \int_0^\infty G(f) \cdot df$

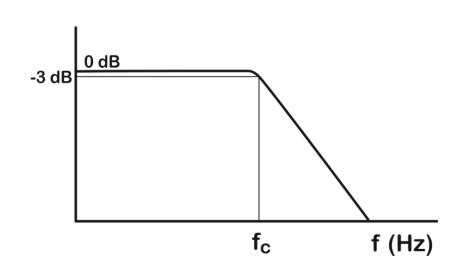


Example

□one pole low-pass filter:

$$T(s) = \frac{1}{1+s\tau}$$

$$T(f) = \frac{f_c}{f_c + jf}$$



□power gain:

$$G(f) = |T(f)|^2 = \frac{f_c^2}{f_c^2 + f^2}$$

□noise equivalent bandwidth:

$$BW_n = \frac{1}{G_0} \int_0^\infty G(f) df = \int_0^\infty \frac{f_c^2}{f_c^2 + f^2} df = \frac{\pi}{2} \cdot f_c$$



Noise and decibel

 \Box for a noise bandwidth of 1 Hz we can write:

$$P_n(dBm, 1 Hz) = 10 \log \left(\frac{\overbrace{N_0 \cdot 1}^{noisepower(W)}}{\underbrace{0.001}_{reference(1mW)}} \right)$$

 \Box for a generic noise bandwidth BW_N we can write:

$$P_n(dBm, BW_N) = 10 \log\left(\frac{N_0 \cdot BW_N}{0.001}\right) = 10 \log(BW_N) + P_n(dBm, 1Hz)$$

 \square for a given constant $PSD\ N_0$ we can switch from a bandwidth BW_1 to another bandwidth BW_2 :

$$P_{\mathbf{n}}(dBm, BW_2) = 10 \log \left(\frac{BW_2}{BW_1}\right) + P_{\mathbf{n}}(dBm, BW_1)$$



Noise measurement with a spectrum analyzer

- □ by construction, a spectrum analyzer provides a trace representing the power content of the input signal included in the resolution bandwidth as a function of frequency
- The output of the spectrum analyzer measuring chain $P_{out}(f)$ is always transformed into a PSD reading by normalizing the measurement chain output power $P_{mc}(f)$ by using the measuring chain signal bandwidth BW_S :

$$P_{readout}(f) = \frac{P_{out}(f)}{BW_S}$$
 (1)

- □ for noise measurements the normalization is based on the knowledge of the equivalent noise bandwidth of the measuring chain, so in (1) we use BW_N instead of BW_S
- □typically, the equivalent noise bandwidth of the spectrum analyzer filter is wider than the signal bandwidth (-3 dB) of about 15-20%

Noise measurement with a spectrum analyzer

- ☐ for a given frequency, the equivalent noise power density at the input of the analyzer must be significantly lower than the noise power density we want to measure
- most spectrum analyzer has dedicated programs for noise measurements
- □the noise level of the analyzer measuring chain depends on the frequency resolution of the measurement
- □ the rms meter of the analyzer is optimized to work with sinusoids, if the input is noise (random signal):
 - the noise level measurement must be corrected
 - the best analyzers are equipped with a calibration and auto-compensation mechanism



cont.

- □since the noise is random, it would be necessary to take an infinite number of measurements to obtain a perfect measurement
- □ important note: during the noise measurements it is assumed that the noise is white inside the noise equivalent bandwidth of the measuring chain
- \Box the displayed measurement result is always normalized to a noise equivalent bandwidth of 1~Hz



Phase noise

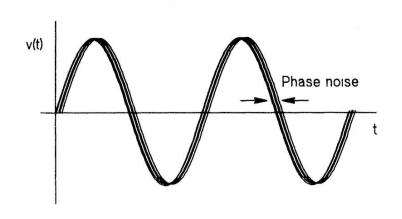
□ a sinusoid which is noisy in phase is represented as:

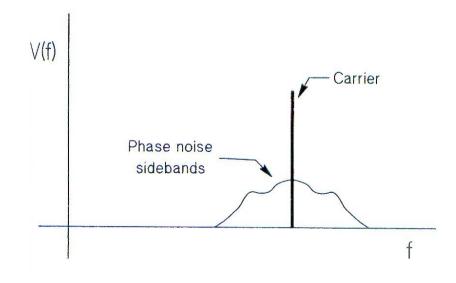
$$v(t) = V_0 \sin[2\pi f_0 t + \Phi_N(t)]$$

where $\Phi_N(t)$ is the phase noise



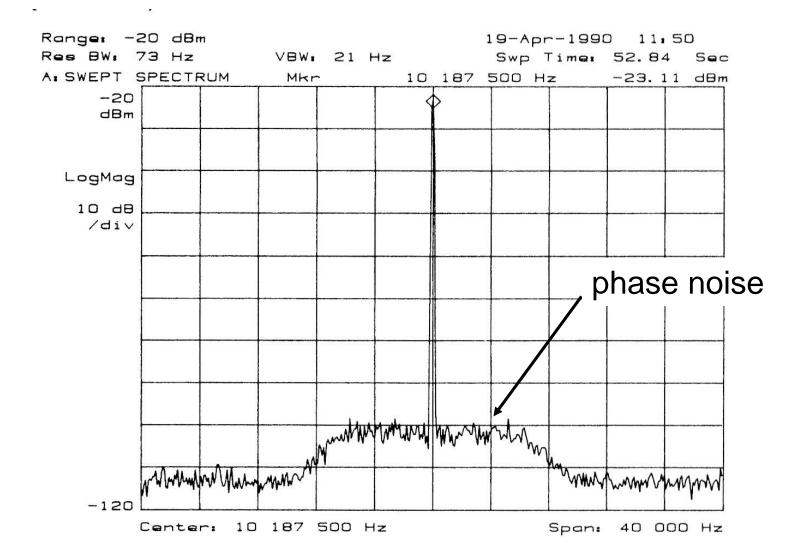
□ the result is the appearance of the symmetrical sidebands around the carrier







Phase noise: spectrum





Phase noise

□ the phase noise in the frequency domain can be expressed as single-sideband (SSB) phase-noise:

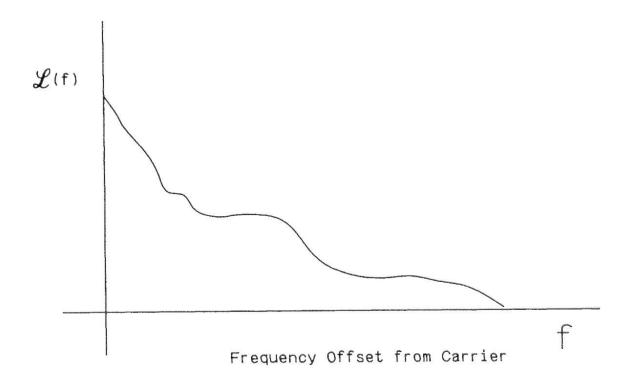
$$L(f) = \frac{V_N(1 Hz BW)}{V_0}$$

where $V_N(1HzBW)$ is the rms noise level in a bandwidth of $1\,Hz$ at $f\,Hz$ away from the carrier, V_0 is the rms amplitude of the carrier

$$\Box \text{in decibel } L(f)_{dB} = 20 \log \frac{V_N(1HzBW)}{V_0}$$



Phase noise: example



$$L(f)_{dB} = 20 \log \frac{V_N(1HzBW)}{V_0}$$

