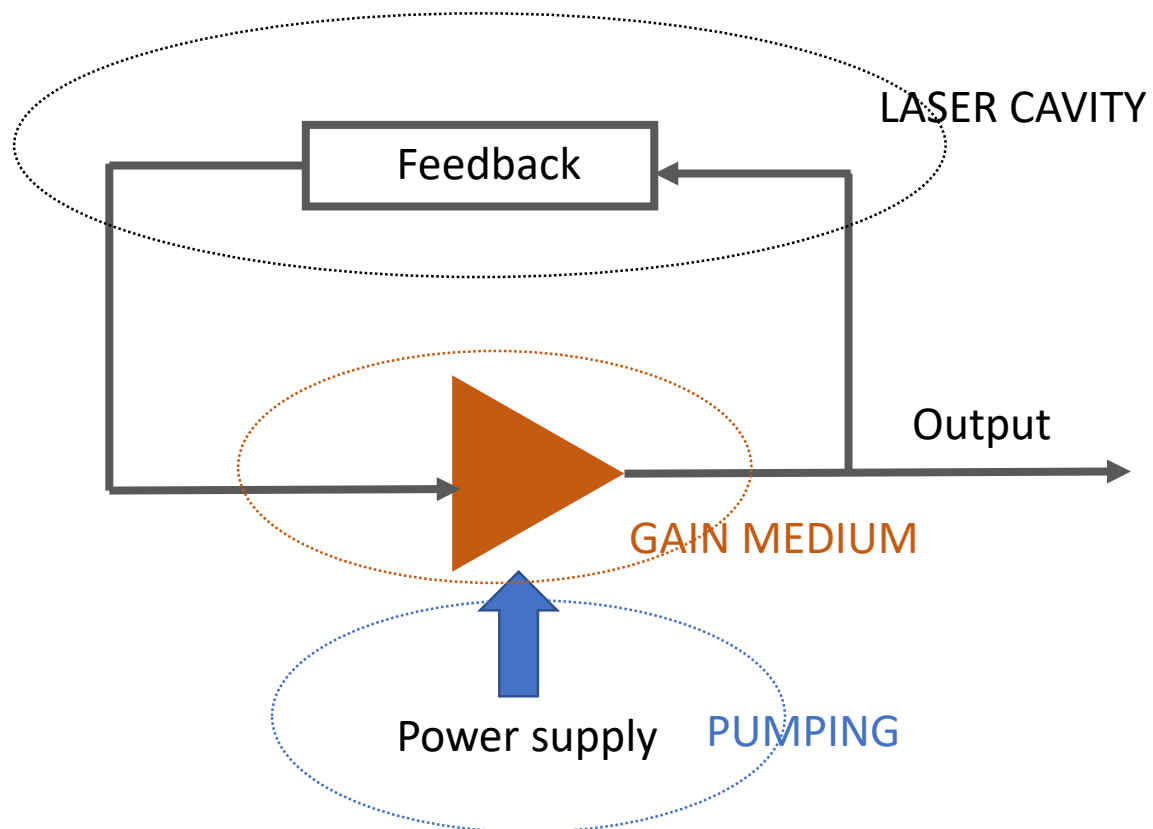


Introduction to lasers

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Chapter 3: Laser Oscillator

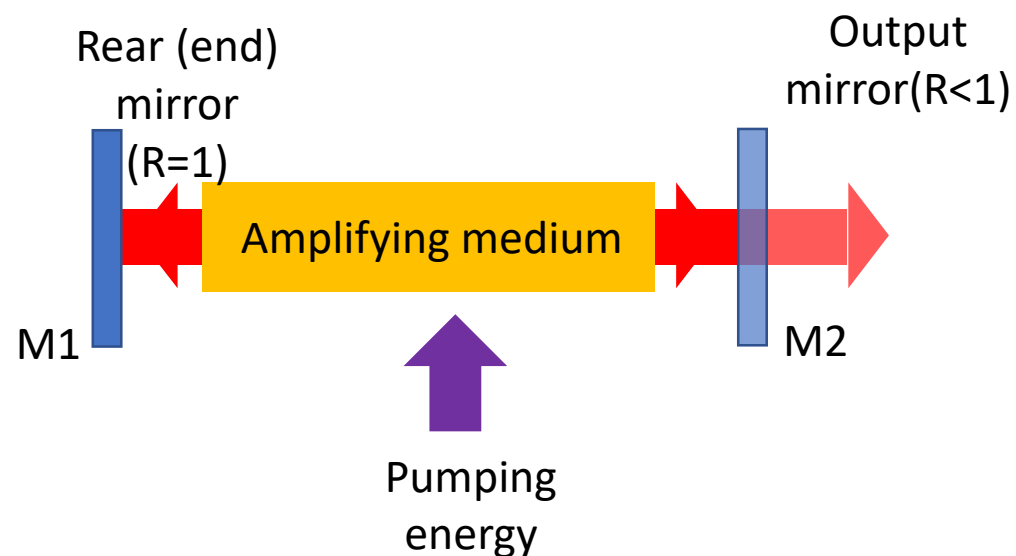




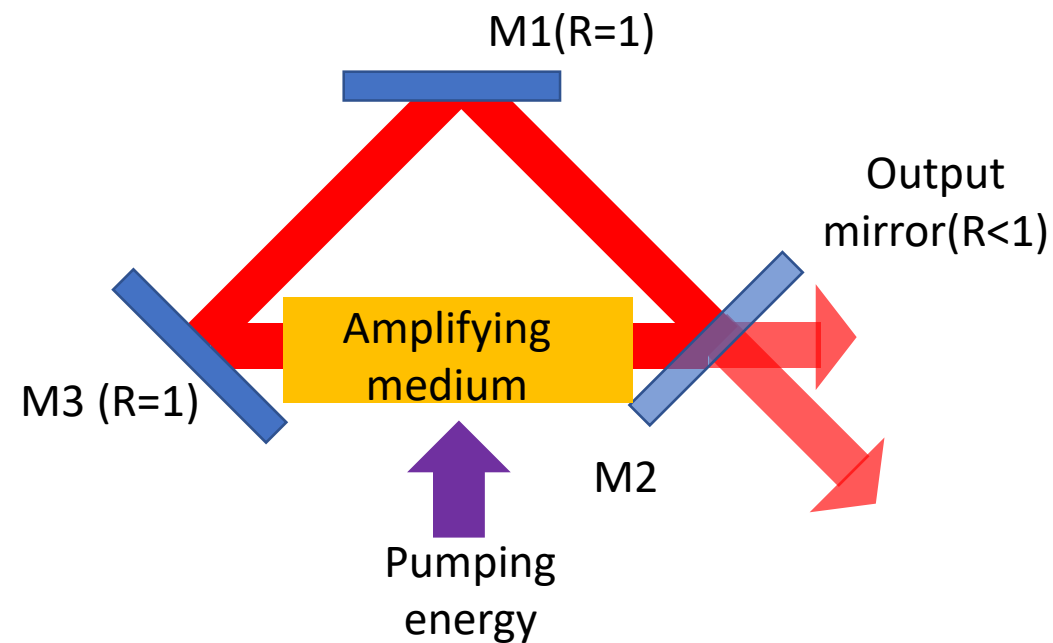
Oscillator: amplifier with a positive feedback

I. Types of cavities (see chapter VI for actual architectures)

Fabry-Perot cavity (schematic diagram)

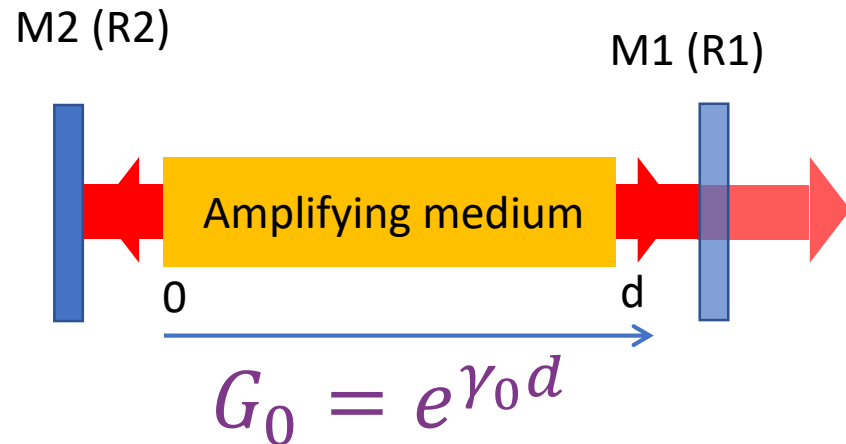


Ring cavity (schematic diagram)



II. First condition for laser oscillation: Gain condition

1. Laser threshold for Fabry Perot cavity



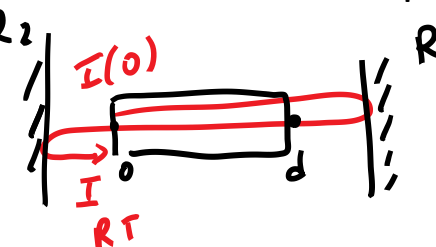
After one round trip in the cavity, i.e. two passages through

the amplifying medium, the small signal gain is : $e^{2\gamma_0 d}$

And the cavity losses (except the ones due to mirrors, i.e diffusion, spontaneous emission, finite size of cavity

components) : $e^{-2\alpha d}$

After one round-trip inside the cavity:



$$I_{AT} = I(0) \cdot R_1 \cdot \underbrace{e^{-\gamma_0 d} \cdot e^{-\gamma_0 d}}_{e^{-2\gamma_0 d}} \cdot R_2 \cdot e^{-2\alpha d}$$

To complete

$$I_{RT} = R_1 R_2 e^{2(\gamma_0 - \alpha)d} \cdot I(0)$$

The initiation of the oscillations requires $I_{RT} \geq I(0)$

$$R_1 R_2 e^{2(\gamma_0 - \alpha)d} \geq 1 \Rightarrow e^{2(\gamma_0 - \alpha)d} \geq \frac{1}{R_1 R_2} \Rightarrow 2(\gamma_0 - \alpha)d \geq -\ln R_1 R_2$$

$$\Rightarrow \boxed{\gamma_0 \geq \alpha - \frac{\ln R_1 R_2}{2d}}$$

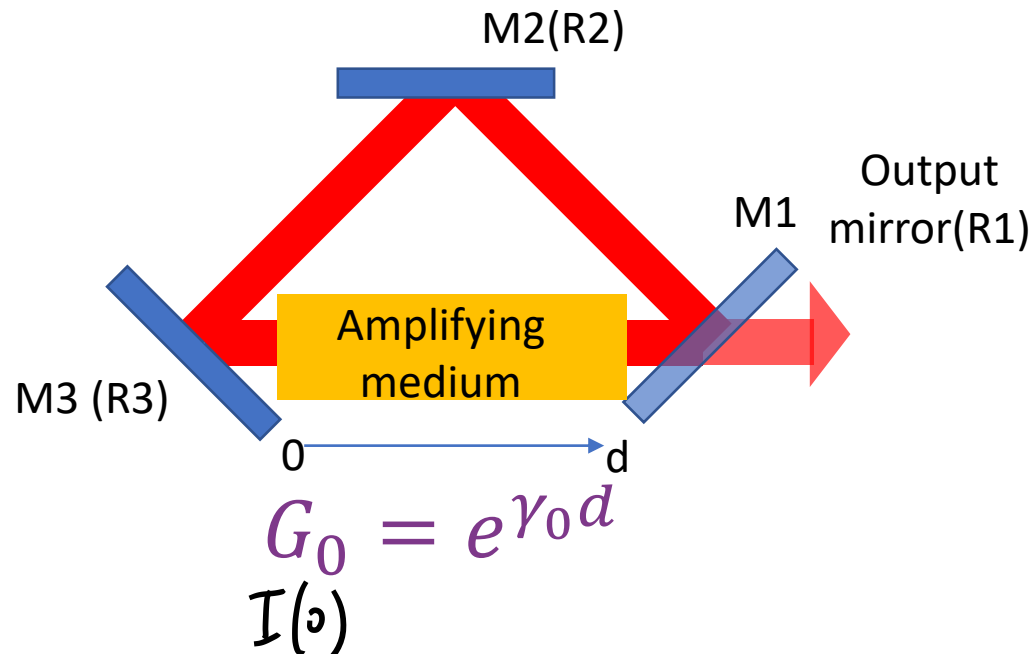
GAIN CONDITION - LASER THRESHOLD

$$\alpha_t = \text{loss coefficient} = \alpha - \frac{\ln R_1 R_2}{2d}$$

$$\boxed{\gamma_0 \geq \alpha_t}$$

The small-signal gain must be greater than the loss-coefficient

2. Laser threshold for ring cavity



To complete

$$I_{RT} = I(0) \cdot e^{\gamma_0 d} R_1 R_2 R_3 \cdot e^{-\alpha d}$$

$$I_{RT} \geq I(0) \quad \text{For oscillations}$$

$$e^{(\gamma_0 - \alpha) d} R_1 R_2 R_3 \geq 1$$

$$(\gamma_0 - \alpha) d \geq -\ln R_1 R_2 R_3$$

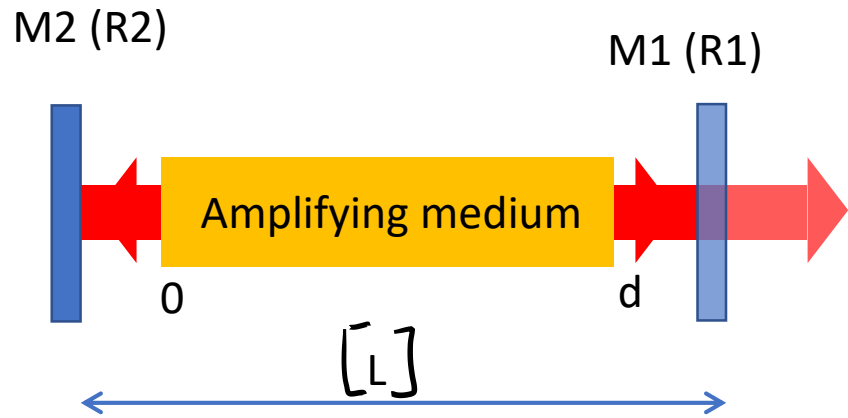
$$\gamma_0 \geq \alpha - \frac{\ln \sum R}{d}$$

GAIN
CONDITION

III - Second condition for laser oscillation: Phase condition

To complete

→ Laser frequencies



$[L]$ = optical path
 $= n \cdot L$
 \hookrightarrow refractive index = $\frac{c}{v}$

→ It requires that the phase shift imparted to the laser wave completing a round-trip within the cavity must be a multiple of 2π

Phase shift after one round-trip:

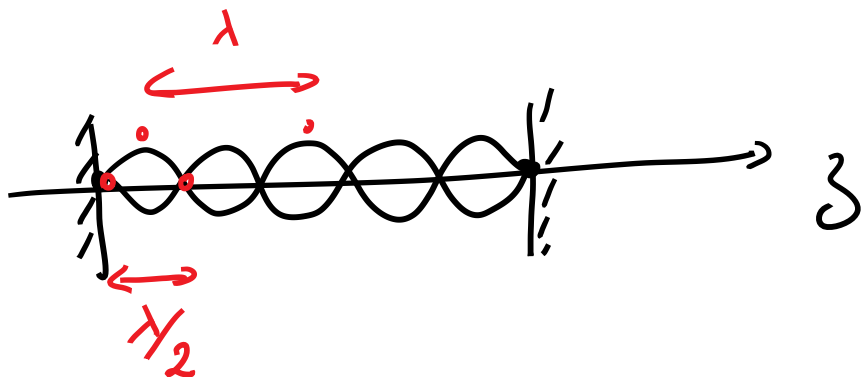
$$\Delta\varphi = \frac{2\pi}{\lambda_k} \cdot 2[L] = 2k\pi \quad k \text{ integer}$$

$$\frac{2\pi v_k}{c} \cdot 2[L] = 2k\pi$$

$$\Rightarrow \boxed{v_k = \frac{1}{2} \left(\frac{c}{2[L]} \right)}$$

Laser frequencies

Free spectral range of the cavity



$$\Rightarrow k \cdot \frac{\lambda}{2} = [L] \Rightarrow k \cdot \frac{c}{2\nu_k} = [L]$$

$$\Rightarrow \boxed{\nu_k = k \cdot \frac{c}{2[L]}}$$

$$L = 1 \text{ m} \quad \frac{c}{2[L]} = \frac{3 \cdot 10^8}{2} = 1.5 \cdot 10^8 \text{ Hz}$$

$$\lambda = 1 \mu\text{m} \Rightarrow \nu = \frac{3 \cdot 10^8}{10^{-6}} = 3 \cdot 10^{14} \text{ Hz}$$

To complete

IV - Steady-state operation

→ Operating point

Laser pumped above threshold ($\gamma_0 > \alpha_t$)

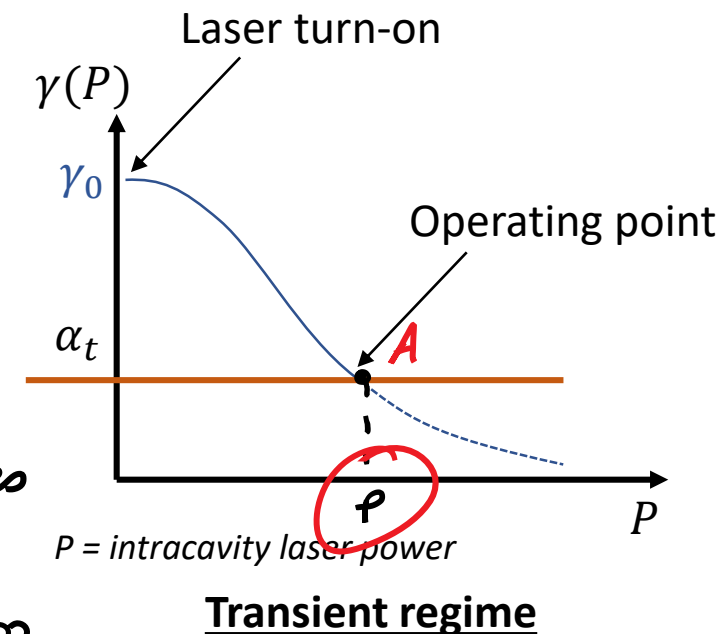
$$\gamma(P) = \frac{\gamma_0}{1 + P/P_{sat}} \quad (1)$$

As the power inside the cavity increases, the gain coefficient decreases due to relation (1)

As long as the gain coefficient remains greater than α_t (loss level), the photon flux continues to increase.

Finally, when $\gamma(P) = \alpha_t$, the power ceases to increase \Rightarrow **STEADY STATE REGIME.**

\rightarrow Gain clamping at the value of losses



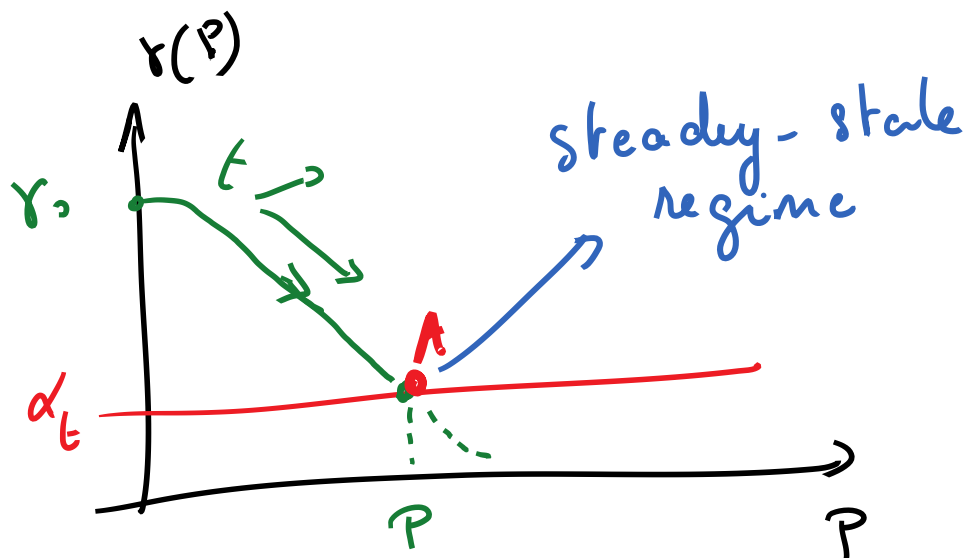
GAIN SATURATION \equiv STABILIZER

To complete

if $P \uparrow \Rightarrow \gamma(P) \downarrow \Rightarrow \gamma(P) < \alpha_t \Rightarrow P \downarrow$
 $P \downarrow \Rightarrow \gamma(P) \uparrow \Rightarrow \gamma(P) > \alpha_t \Rightarrow P \uparrow$

Stabilization on point **A** for which

$$\gamma(P) = \alpha_t$$



$$\gamma(P) = \frac{\gamma_0}{1 + P/P_{sat}} = \alpha_t$$

$$\frac{P}{P_{sat}} + 1 = \frac{\gamma_0}{\alpha_t} \Rightarrow$$

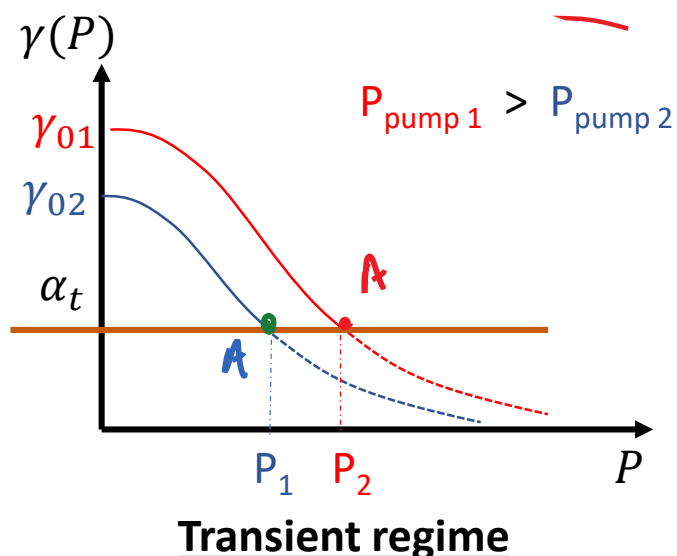
$$P = P_{sat} \left(\frac{\gamma_0}{\alpha_t} - 1 \right)$$

Power inside the cavity

Assumption: P does not depend on $z \Leftrightarrow R_1$ and R_2 close to 1.

To complete

Comment



At steady-state, the intracavity power increases as the pump power increases due to the increase of γ_0

P = intracavity laser power