

MICROWAVE ENGINEERING

Lecture 9:

Smith Chart,
Quarter Wave
Transformer
Lossy Transmission
Lines

The Smith's Chart (1939 by P. Smith @ Bells lab)

↳ POLAR PLOT of the reflection coefficient Γ

$$\Gamma = |\Gamma| e^{j\theta}$$

$|\Gamma| \leq 1$
is plotted
as a radius

$-180^\circ \leq \theta \leq 180^\circ$

Measured from the right hand
side of the horizontal diameter
on the graph.

Reflection coeff. can be converted

→ Impedance
→ Admittance

On the chart we always plot NORMALIZED IMPEDANCES

$$z = \frac{Z}{Z_0}$$

If we have a lossless line with char. imp. to end terminated with a load Γ :

$$\Gamma = \frac{z_L - 1}{z_L + 1} = |\Gamma| e^{j\theta}$$

$$z_L = \frac{z_0}{\Gamma}$$



Solving for $z_L \Rightarrow \Gamma(z_L + 1) = (z_L - 1) \Rightarrow \Gamma z_L + \Gamma = z_L - 1$

$$z_L = \frac{1 + |\Gamma| e^{j\theta}}{1 - |\Gamma| e^{j\theta}}$$

$$\Gamma = \Gamma_r + j \Gamma_i$$

$$z_L = r_L + j x_L$$

$$\begin{aligned}
 r_e + jx_{UL} &= \frac{1 + \Gamma_r + j\Gamma_i}{1 - \Gamma_r - j\Gamma_i} = \frac{[(1 + \Gamma_r) + j\Gamma_i][1 - \Gamma_r + j\Gamma_i]}{[1 - \Gamma_r - j\Gamma_i][1 - \Gamma_r + j\Gamma_i]} = \\
 &= \frac{1 - \cancel{\Gamma_r} + j\Gamma_i + \cancel{\Gamma_r} - \Gamma_r^2 + j\Gamma_r\Gamma_i + j\Gamma_i - j\Gamma_i\Gamma_r - \Gamma_i^2}{(1 - \Gamma_r)^2 + \Gamma_i^2} =
 \end{aligned}$$

$$\begin{aligned}
 &= \frac{1 - \Gamma_r^2 - \Gamma_i^2 + 2j\Gamma_i}{(1 - \Gamma_r)^2 + \Gamma_i^2}
 \end{aligned}$$

$$\boxed{r_L = \frac{1 - \Gamma_r^2 - \Gamma_i^2}{(1 - \Gamma_r)^2 + \Gamma_i^2}}$$

LOAD
RESISTANCE

LOAD
REACTANCE

$$X_L = \frac{2r_i}{(1-r_f)^2 + r_i^2}$$

We can rearrange in eq. of circles :

Resistance
circles

$$\left(r_f - \frac{r_i}{1+r_L}\right)^2 + r_i^2 = \left(\frac{1}{1+r_L}\right)^2$$

Reactance
circles

$$(r_f - 1)^2 + \left(r_i - \frac{1}{X_L}\right)^2 = \left(\frac{1}{X_L}\right)^2$$

For ex: if $r_i = 1$ \rightarrow Circle is centered in $r_f = 0.5$
 $r_i = 0$

$$\left(\Gamma_r - \frac{1}{2}\right)^2 + \Gamma_i^2 = \left(\frac{1}{2}\right)^2$$

↑
 offset
 with respect
 to O

↑ radius

Smith chart can be used to graphically solve transmission line eq.

$$Z_{IN} = Z_0 \frac{1 + \Gamma e^{-2j\beta l}}{1 - \Gamma e^{-2j\beta l}}$$

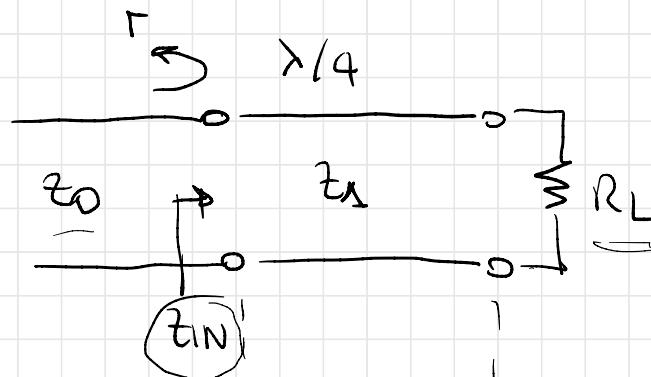
By plotting Γ on the chart the normalized Z_{IN} can be found by rotating clockwise an amount $2\beta l$ around the center.

(subtract $2\beta l$ from θ)

Max rotation on chart $0.5\lambda \rightarrow 2\pi$



THE QUARTER WAVE TRANSFORMER



loss less
 T_L
with z_L charact. impedance.

IMPORTANT:

R_L and z_0 are given parameters and assumed to be REAL

To match a load to a line $\Rightarrow \Gamma=0$ at the input of the transformer

$$Z_{IN} = Z_1 \frac{R_L + jZ_1 \tan \phi}{Z_1 + jR_L \tan \phi}$$

$$fL = \frac{2\pi}{2} \cdot \frac{X}{4} = \frac{\pi}{2}$$

$$Z_{IN} = Z_1 \frac{\frac{R_L}{\tan \phi} + jZ_1}{\frac{Z_1}{\tan \phi} + jR_L} = \frac{Z_1^2}{R_L}$$

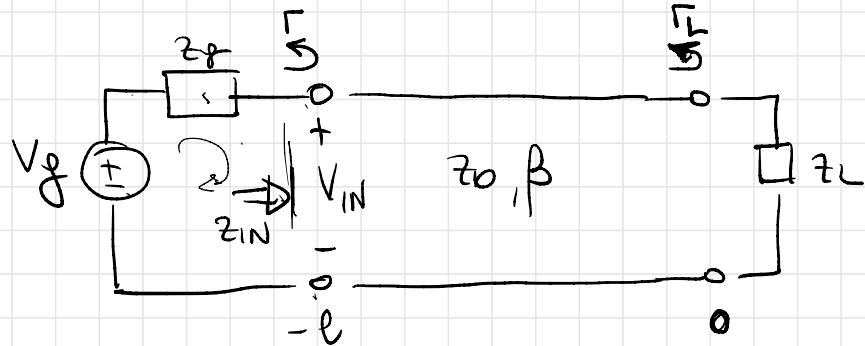
$$\text{to get } \Gamma=0 \Rightarrow Z_{IN} = Z_0 = \frac{Z_1^2}{R_L}$$

$$Z_1 = \sqrt{R_L Z_0}$$

\uparrow
SWR=1 for this condition

IMPORTANT: Matching is valid at 1 frequency

GENERATOR AND LOAD MISMATCHES



- Γ_L , Z_L can be complex

- Z_L is lossless

$$Z_{IN} = Z_0 \frac{1 + \Gamma_L e^{-2j\beta L}}{1 - \Gamma_L e^{-2j\beta L}} = Z_0 \frac{Z_L + j Z_0 \tan \beta L}{Z_0 + j Z_L \tan \beta L}$$
$$\Gamma_L = \frac{Z_L - Z_0}{Z_L + Z_0}$$

The voltage on the line is:

$$V(z) = V_0^+ \left(e^{-j\beta z} + r_L e^{j\beta z} \right)$$

In $z = -l$

$$V(-l) = V_0^+ \frac{\cancel{z_{IN}}}{\cancel{z_{IN}} + z_g} = V_0^+ \left[e^{j\beta l} + r_L e^{-j\beta l} \right]$$

$$V_0^+ = V_g \frac{z_{IN}}{z_{IN} + z_g} \frac{l}{e^{j\beta l} + r_L e^{-j\beta l}}$$

$$z_{IN} = z_0 \frac{1 + r_L e^{-2j\beta l}}{1 - r_L e^{-2j\beta l}}$$

Replacing z_{IN} in the expression of V_0^+

$$V_o^+ = V_g z_0 \frac{1 + \Gamma_L e^{-2j\beta e}}{1 - \Gamma_L e^{-2j\beta e}} \cdot \frac{1}{z_0 + z_g \frac{1 + \Gamma_L e^{-2j\beta e}}{1 - \Gamma_L e^{-2j\beta e}}} \cdot \frac{1}{e^{j\beta e} + \Gamma_L e^{j\beta e}}$$

$$= V_g z_0 \frac{1 + \Gamma_L e^{-2j\beta e}}{1 - \Gamma_L e^{-2j\beta e}} \cdot \frac{1 - \Gamma_L e^{-2j\beta e}}{z_0 (1 + \Gamma_L e^{-j\beta e}) + z_g (1 - \Gamma_L e^{-2j\beta e})} \cdot \frac{1}{e^{j\beta e} + \Gamma_L e^{j\beta e}}$$

multiply by $e^{-j\beta e}$

$$= V_g z_0 \frac{(1 + \Gamma_L e^{-2j\beta e})}{z_0 (1 + \Gamma_L e^{-2j\beta e}) + z_g (1 - \Gamma_L e^{-2j\beta e})} \cdot \frac{e^{-j\beta e}}{1 + \Gamma_L e^{-2j\beta e}} =$$

$$= V_g z_0 \frac{e^{-j\beta e}}{z_0 + z_g + (z_0 - z_g) \Gamma_L e^{-2j\beta e}}$$

obviate
z₀ + z_g

$$= V_g \frac{z_0}{z_0 + z_g} \frac{e^{-j\beta e}}{1 - \Gamma_L \Gamma_g e^{-j\beta e}}$$

$$\Gamma_g = \frac{z_g - z_0}{z_g + z_0}$$

On the Line

$$\text{SWR} = \frac{1 + |Z_L|}{1 - |Z_L|}$$

The power delivered at the load:

$$P = \frac{1}{2} \operatorname{Re} \left\{ V_{IN} I_N^* \right\} = \frac{1}{2} |V_{IN}|^2 \operatorname{Re} \left\{ \frac{1}{Z_N} \right\} = \frac{1}{2} |V_g|^2 \left| \frac{Z_{IN}}{Z_{IN} + Z_g} \right|^2 \operatorname{Re} \left\{ \frac{1}{Z_N} \right\}$$

$$I_N = \frac{V_{IN}}{Z_{IN}}$$

$$\frac{R_{IN}^2 - jX_{IN}}{R_{IN}^2 + X_{IN}^2}$$

$$\left. \begin{array}{l} \text{If } Z_{IN} = R_{IN} + jX_{IN} \\ Z_g = R_g + jX_g \end{array} \right\} \Rightarrow$$

$$P = \frac{1}{2} |V_g|^2 \frac{R_{IN}}{(R_{IN} + R_g)^2 + (X_{IN} + X_g)^2}$$

Three special cases are:

① LOAD IS MATCHED TO THE LINE $Z_L = Z_0$

In this case $\Gamma_L = 0$, SWR = 1 $\Rightarrow Z_{IN} = Z_0$

Therefore

$$P = \frac{1}{2} |V_g|^2 \frac{Z_0}{(Z_0 + R_g)^2 + X_g^2}$$

② GENERATOR IS MATCHED TO THE LOADED LINE: $Z_{IN} = Z_g$

$$\Gamma = 0 \Rightarrow \frac{Z_{IN} - Z_g}{Z_{IN} + Z_g} = 0$$

We know that $\Gamma_L \neq 0 \Rightarrow$ there will be a standing wave on the line

$$P = \frac{1}{2} |V_g|^2 \frac{R_g}{R_g^2 + X_g^2}$$

③ CONJUGATE MATCHING

If we want to max power at the load we can assume Z_g is fixed and we vary Z_{IN} until max power is delivered:

$$\boxed{\frac{\partial P}{\partial R_{IN}} = 0} \Rightarrow \frac{1}{(R_{IN} + R_g)^2 + (X_{IN} + X_g)^2} + \frac{-2 R_{IN} (R_{IN} + R_g)}{(R_{IN} + R_g)^2 + (X_{IN} + X_g)^2} = 0$$

$$\cancel{R_{IN}^2 + R_g^2 + 2 R_{IN} R_g - 2 R_{IN}^2 - 2 R_{IN} R_g + (X_{IN} + X_g)^2} = 0$$

$$\boxed{R_g^2 - R_{IN}^2 + (X_{IN} + X_g)^2 = 0}$$

or

$$\frac{\partial P}{\partial X_{IN}} = 0 \rightarrow -2 R_{IN} (X_{IN} + X_g) \left[(R_{IN} + R_g)^2 + (X_{IN} + X_g)^2 \right]^2 = 0$$

$$R_{IN} (X_{IN} + X_g) = 0$$

If we solve simultaneously for R_{IN} and X_{IN} :

$$R_{IN} = R_g$$

$$X_{IN} = -X_g$$

$$Z_{IN} = Z_g^*$$

CONJUGATE MATCHING
CONDITION

In this case

$$P = \frac{1}{2} |V_g|^2 \frac{1}{4R_g}$$

MAX AVAILABLE
POWER FOR A SPECIFIC
GENERATOR

NOTE : Γ_1 , Γ_g and Γ_e can be non-zero

NOTE 2 : If X_g is Real \Rightarrow Max power transfer $\Rightarrow R_N = R_g$

NOTE 3 : Even if we achieve max power transfer we still have max 50 % efficiency $\rightarrow z_g$ needs

to be
as small
as possible!

LOSSY TRANSMISSION LINES

LOW-LOSS LINE

$$\gamma = \alpha + j\beta = \sqrt{(R+j\omega L)(G+j\omega C)} = \dots$$

$$= \sqrt{1 - j \left(\frac{R}{\omega L} + \frac{G}{\omega C} \right) - \frac{RG}{\omega^2 LC}}$$

If losses are small :

$$\left. \begin{array}{l} R \ll \omega L \\ G \ll \omega C \end{array} \right\} \Rightarrow RG \ll \omega^2 LC$$

$$\gamma \approx j\omega \sqrt{\omega C} \underbrace{\sqrt{1 - j \left(\frac{R}{\omega L} + \frac{G}{\omega C} \right)}}_{} = \text{expanding in Taylor series}$$

$$= j \omega \sqrt{LC} \left[1 - \frac{j}{2} \left(\frac{R}{WL} + \frac{G}{WC} \right) \right]$$

↙

$$d \approx \frac{1}{2} \left(R \sqrt{\frac{C}{L}} + G \sqrt{\frac{L}{C}} \right) = \frac{1}{2} \left(\frac{R}{Z_0} + G Z_0 \right)$$

$$Z_0 = \sqrt{\frac{L}{C}}$$

because

$$Z_0 = \sqrt{\frac{R+j\omega L}{G+j\omega C}} \approx \sqrt{\frac{L}{C}}$$

for small
losses

THE DISTORTIONLESS LINE

unless line is lossless β is NOT
a linear function of ω

$$\gamma = j\omega \sqrt{LC} \sqrt{1 - j\left(\frac{R}{\omega L} + \frac{G}{\omega C}\right) - \frac{RG}{\omega^2 LC}}$$

↓

the PHASE VELOCITY $V_p = \frac{\omega}{\beta}$

will be different for different frequencies

↓

Different signal wavelength traveling on the line will arrive at different times!

DISPERSION

A distortionless line has :

$$\boxed{\frac{R}{L} = \frac{G}{C}}$$

so the propagation constant is :

$$\gamma = j\omega\sqrt{LC} \sqrt{1 - j\left(\frac{R}{WL} + \frac{R}{WL}\right) - \frac{R^2}{\omega^2 L^2}} =$$

$$= j\omega\sqrt{LC} \sqrt{1 - 2j\frac{R}{WL} - \frac{R^2}{\omega^2 L^2}} = j\omega\sqrt{LC} \left(1 - j\frac{R}{WL}\right) =$$

$\underbrace{\left(1 - j\frac{R}{\omega L}\right)^2}$

$$= j\omega\sqrt{LC} + R\sqrt{\frac{C}{L}} = \alpha + j\beta$$

$\underbrace{j\omega\sqrt{LC}}_{\beta}$
 $\underbrace{R\sqrt{\frac{C}{L}}}_{\alpha}$

$$\boxed{\beta = \omega\sqrt{LC}}$$

LINEAR FUNCTION

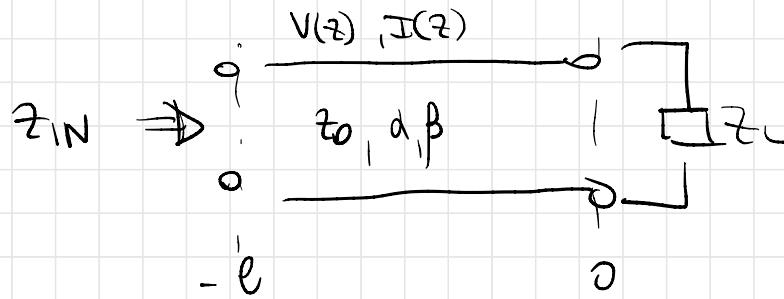
OF ω

$$\boxed{\alpha = R\sqrt{\frac{C}{L}}}$$

IT DOES
NOT DEPEND
OF FREQUENCY

 NO DISTORTION

THE TERMINATED LOSSY LINE



Let's assume low losses $\rightarrow z_0 \approx \sqrt{\frac{L}{G}}$

$$V(z) = V_0^+ [e^{-\gamma z} + \Gamma e^{\gamma z}]$$

$$I(z) = \frac{V_0^+}{z_0} [e^{-\gamma z} - \Gamma e^{\gamma z}]$$

At a distance ℓ from the load:

$$\Gamma(\ell) = \Gamma e^{-2j\beta\ell} e^{-2\alpha\ell} = \Gamma e^{-2\gamma\ell}$$

The input impedance

$$Z_{IN} = \frac{V(-l)}{I(-l)} = Z_0 \frac{Z_L + Z_0 \tanh \gamma l}{Z_0 + Z_L \tanh \gamma l}$$

The power delivered at $z = -l$ is

$$P_{IN} = \frac{1}{2} \operatorname{Re} \{ V(-l) I(-l)^* \} - \frac{N_0^+ R}{2 Z_0} [e^{2 \alpha l} - |\Gamma(l)|^2 e^{-2 \alpha l}]$$

The power at the load is

$$P_L = \frac{1}{2} \operatorname{Re} \{ V(0) + I_0^* \} = \frac{|V_0^+|^2}{2 Z_0} (1 - |\Gamma|^2)$$

The power loss :

$$\boxed{P_{LOSS} = P_{IN} - P_L}$$