

DERIVATION OF POYNTING THEOREM

MAXWELL'S EQUATIONS

$$\nabla \times \mathbf{E} = -j\omega\mu\mathbf{H}$$

$$\nabla \times \mathbf{H} = j\omega\epsilon\mathbf{E} + \sigma\mathbf{E} + \mathbf{J} \quad \text{WHERE } \mathbf{J} \text{ IS THE IMPOSED CURRENT}$$

SCALAR PRODUCT BETWEEN \mathbf{H}^* AND THE FIRST EQUATION

$$\mathbf{H}^* \cdot \nabla \times \mathbf{E} = \mathbf{H}^* \cdot (-j\omega\mu\mathbf{H}) = -j\omega\mu |\mathbf{H}|^2$$

SCALAR PRODUCT BETWEEN \mathbf{E} AND THE COMPLEX CONJUGATE OF THE SECOND EQUATION

$$\mathbf{E} \cdot (\nabla \times \mathbf{H})^* = \mathbf{E} \cdot (j\omega\epsilon\mathbf{E} + \sigma\mathbf{E} + \mathbf{J})^* = \mathbf{E} \cdot (-j\omega\epsilon\mathbf{E}^* + \sigma\mathbf{E}^* + \mathbf{J}^*)$$

$$\mathbf{E} \cdot \nabla \times \mathbf{H}^* = -j\omega\epsilon |\mathbf{E}|^2 + \sigma |\mathbf{E}|^2 + \mathbf{E} \cdot \mathbf{J}^*$$

THE TWO EQUATIONS READ AS

$$\mathbf{H}^* \cdot \nabla \times \mathbf{E} = -j\omega\mu |\mathbf{H}|^2$$

$$\mathbf{E} \cdot \nabla \times \mathbf{H}^* = -j\omega\epsilon |\mathbf{E}|^2 + \sigma |\mathbf{E}|^2 + \mathbf{E} \cdot \mathbf{J}^*$$

WE SUBTRACT THE SECOND EQUATION FROM THE FIRST

$$\mathbf{H}^* \cdot \nabla \times \mathbf{E} - \mathbf{E} \cdot \nabla \times \mathbf{H}^* = -j\omega\mu |\mathbf{H}|^2 + j\omega\epsilon |\mathbf{E}|^2 - \sigma |\mathbf{E}|^2 - \mathbf{E} \cdot \mathbf{J}^*$$

$$\text{SINCE } \nabla \cdot (\mathbf{U} \times \mathbf{V}) = \mathbf{V} \cdot \nabla \times \mathbf{U} - \mathbf{U} \cdot \nabla \times \mathbf{V}$$

WE CAN REARRANGE THE LEFT-HAND SIDE TERM BY OBSERVING THAT:

$$\nabla \cdot (\mathbf{E} \times \mathbf{H}^*) = \mathbf{H}^* \cdot \nabla \times \mathbf{E} - \mathbf{E} \cdot \nabla \times \mathbf{H}^*$$

THE EXPRESSION BECOMES

$$\nabla \cdot (\mathbf{E} \times \mathbf{H}^*) = -j\omega\mu |\mathbf{H}|^2 + j\omega\epsilon |\mathbf{E}|^2 - \sigma |\mathbf{E}|^2 - \mathbf{E} \cdot \mathbf{J}^*$$

WE MOVE THE "SOURCE" TERM $-\mathbf{E} \cdot \mathbf{J}^*$ TO THE LEFT-HAND SIDE AND THE $\nabla \cdot (\mathbf{E} \times \mathbf{H}^*)$ TERM TO THE RIGHT HAND SIDE

$$\mathbf{E} \cdot \mathbf{J}^* = -\nabla \cdot (\mathbf{E} \times \mathbf{H}^*) - j\omega\mu |\mathbf{H}|^2 + j\omega\epsilon |\mathbf{E}|^2 - \sigma |\mathbf{E}|^2$$

WE CHANGE THE SIGN EVERYWHERE AND WE MULTIPLY BY $1/2$ TO OBTAIN TIME-AVERAGE VALUES (SINCE WE ARE WORKING WITH PHASORS)


$$-\frac{1}{2} \mathbf{E} \cdot \mathbf{J}^* = \frac{1}{2} \nabla \cdot (\mathbf{E} \times \mathbf{H}^*) + \frac{1}{2} \sigma |\mathbf{E}|^2 + \frac{1}{2} j\omega\epsilon |\mathbf{E}|^2 + \frac{1}{2} j\omega\mu |\mathbf{H}|^2$$

WE INTEGRATE THE PREVIOUS EXPRESSION OVER A VOLUME V BOUNDED BY A CLOSED SURFACE S

$$-\frac{1}{2} \int_V \mathbf{E} \cdot \mathbf{J}^* dV = \frac{1}{2} \int_V \nabla \cdot (\mathbf{E} \times \mathbf{H}^*) dV + \frac{1}{2} \int_V \sigma |\mathbf{E}|^2 dV + \frac{1}{2} j\omega \int_V (\mu |\mathbf{H}|^2 - \epsilon |\mathbf{E}|^2) dV$$

THANKS TO THE GAUSS THEOREM WE CAN WRITE THAT

$$\frac{1}{2} \int_V \nabla \cdot (\mathbf{E} \times \mathbf{H}^*) dV = \frac{1}{2} \int_S \mathbf{E} \times \mathbf{H}^* \cdot d\vec{S}$$

$d\vec{S} = \hat{n} dS$

 UNIT VECTOR
 NORMAL TO THE SURFACE
 AND DIRECTED OUT FROM THE VOLUME
 INCREMENT
 OF THE AREA

AND FINALLY WE CAN WRITE

$$-\frac{1}{2} \int_V \mathbf{E} \cdot \mathbf{J}^* dV = \frac{1}{2} \int_S \mathbf{E} \times \mathbf{H}^* \cdot d\vec{S} + \frac{1}{2} \int_V \sigma |\mathbf{E}|^2 dV + j2\omega \int_V \left(\frac{1}{4} \mu |\mathbf{H}|^2 - \frac{1}{4} \epsilon |\mathbf{E}|^2 \right) dV$$