MATLAB version number 5 of the APBS

(Dated: December 4, 2009)

Abstract

This solver uses the biconjugate gradient stabilized method and the inexact LU decomposition to numerically solve the linearized PB equation on a cartesian 3D-grid. This version requires the shifted dielectric and the ion accessibility coefficient (kappa function) maps as generated by the APBS code as well as the corresponding pqr file generated by the pdb2pqr code. It uses standard three-linear splines (spl0) to spread the charge density along the nearest grid points if needed. It is able to solve the linear PB eq with either Dirichlet or focus boundary conditions. In the later case, this code solves the PB equation in a large (low resolution) domain and the resulting electrostatic potential solution is subsequently used to evaluate the Dirichlet boundary condition to solve the PB equation in a (higher resolution) sub-domain region. The resulting electrostatic potential and charge maps for both the coarse and target grids are saved in dx format in different folders. This version includes the option to calculate the energy for point-like charge systems. For visualization purpose, this code also generates two files (fig and fifth) corresponding to the graphical representation of the electrostatic potential surface. This version was used to solve the pka and point-pmf examples provided by the apbs package. If the Dirichlet Boundary Condition is required, then this version basically provides the same results than the previous one. The main difference is that the user doesn't have to edit the source files in this version but only have to provide the target input-file name and the corresponding full path as the only argument of the (new) matlab function MAPBS (x). See the pka example and the MATLAB_PB_SOLVER_6 package for more details.

Description

This code is based on Michel Holst's thesis and Nathan Baker's APBS approach. The box-method is used to discretize the following (linearized) PB equation

$$-\nabla \cdot (\epsilon(\mathbf{r}) \nabla u(\mathbf{r})) + \bar{\kappa}(\mathbf{r})^{2} u(\mathbf{r}) = magic \sum_{i=1}^{N} z_{i} \delta(\mathbf{r} - \mathbf{r}_{i})$$
(1)

where $u(\mathbf{r}) = e_c \Phi(\mathbf{r}) / K_B T$ and $magic = 4\pi e_c^2 / K_B T$. For a diagonal dielectric tensor, the resulting discretized linear PB equations at the nodes $u_{ijk} = u(x_i, y_j, z_k)$ for $1 \le i \le N_x$, $1 \le j \le N_y$ and $1 \le k \le N_z$ reads

$$\left[\epsilon_{i-1/2,j,k}^{x} \frac{\left(h_{j-1}+h_{j}\right) \left(h_{k-1}+h_{k}\right)}{4h_{i-1}}+\epsilon_{i+1/2,j,k}^{x} \frac{\left(h_{j-1}+h_{j}\right) \left(h_{k-1}+h_{k}\right)}{4h_{i}}+\right.\right.$$

$$\epsilon_{i,j-1/2,k}^{y}\frac{\left(h_{i-1}+h_{i}\right)\left(h_{k-1}+h_{k}\right)}{4h_{j-1}}+\epsilon_{i,j+1/2,k}^{y}\frac{\left(h_{i-1}+h_{i}\right)\left(h_{k-1}+h_{k}\right)}{4h_{j}}+$$

$$\epsilon_{i,j,k-1/2}^{k} \frac{\left(h_{i-1} + h_{i}\right)\left(h_{j-1} + h_{j}\right)}{4h_{k-1}} + \epsilon_{i,j,k+1/2}^{k} \frac{\left(h_{i-1} + h_{i}\right)\left(h_{j-1} + h_{j}\right)}{4h_{k}} +$$

$$\kappa_{ijk} \frac{(h_{i-1} + h_i)(h_{j-1} + h_j)(h_{k-1} + h_k)}{8} u_{ijk} +$$

$$\left[-\epsilon_{i-1/2,j,k}^{x} \frac{\left(h_{j-1} + h_{j}\right)\left(h_{k-1} + h_{k}\right)}{4h_{i-1}} \right] u_{i-1jk} + \left[-\epsilon_{i+1/2,j,k}^{x} \frac{\left(h_{j-1} + h_{j}\right)\left(h_{k-1} + h_{k}\right)}{4h_{i}} \right] u_{i+1jk} + \left[-\epsilon_{i+1/2,j,k}^{x} \frac{\left(h_{j-1} + h_{j}\right)\left(h_{k-1} + h_{k}\right)}{4h_{i}} \right] u_{i+1jk} + \left[-\epsilon_{i+1/2,j,k}^{x} \frac{\left(h_{j-1} + h_{j}\right)\left(h_{k-1} + h_{k}\right)}{4h_{i}} \right] u_{i+1jk} + \left[-\epsilon_{i+1/2,j,k}^{x} \frac{\left(h_{j-1} + h_{j}\right)\left(h_{k-1} + h_{k}\right)}{4h_{i}} \right] u_{i+1jk} + \left[-\epsilon_{i+1/2,j,k}^{x} \frac{\left(h_{j-1} + h_{j}\right)\left(h_{k-1} + h_{k}\right)}{4h_{i}} \right] u_{i+1jk} + \left[-\epsilon_{i+1/2,j,k}^{x} \frac{\left(h_{j-1} + h_{j}\right)\left(h_{k-1} + h_{k}\right)}{4h_{i}} \right] u_{i+1jk} + \left[-\epsilon_{i+1/2,j,k}^{x} \frac{\left(h_{j-1} + h_{j}\right)\left(h_{k-1} + h_{k}\right)}{4h_{i}} \right] u_{i+1jk} + \left[-\epsilon_{i+1/2,j,k}^{x} \frac{\left(h_{j-1} + h_{j}\right)\left(h_{k-1} + h_{k}\right)}{4h_{i}} \right] u_{i+1jk} + \left[-\epsilon_{i+1/2,j,k}^{x} \frac{\left(h_{j-1} + h_{j}\right)\left(h_{k-1} + h_{k}\right)}{4h_{i}} \right] u_{i+1jk} + \left[-\epsilon_{i+1/2,j,k}^{x} \frac{\left(h_{j-1} + h_{j}\right)\left(h_{k-1} + h_{k}\right)}{4h_{i}} \right] u_{i+1jk} + \left[-\epsilon_{i+1/2,j,k}^{x} \frac{\left(h_{j-1} + h_{j}\right)\left(h_{k-1} + h_{k}\right)}{4h_{i}} \right] u_{i+1jk} + \left[-\epsilon_{i+1/2,j,k}^{x} \frac{\left(h_{j-1} + h_{j}\right)\left(h_{k-1} + h_{k}\right)}{4h_{i}} \right] u_{i+1jk} + \left[-\epsilon_{i+1/2,j,k}^{x} \frac{\left(h_{j-1} + h_{j}\right)\left(h_{k-1} + h_{k}\right)}{4h_{i}} \right] u_{i+1jk} + \left[-\epsilon_{i+1/2,j,k}^{x} \frac{\left(h_{j-1} + h_{j}\right)\left(h_{k-1} + h_{k}\right)}{4h_{i}} \right] u_{i+1jk} + \left[-\epsilon_{i+1/2,j,k}^{x} \frac{\left(h_{j-1} + h_{j}\right)\left(h_{k-1} + h_{k}\right)}{4h_{i}} \right] u_{i+1jk} + \left[-\epsilon_{i+1/2,j,k}^{x} \frac{\left(h_{j-1} + h_{j}\right)\left(h_{k-1} + h_{k}\right)}{4h_{i}} \right] u_{i+1jk} + \left[-\epsilon_{i+1/2,j,k}^{x} \frac{\left(h_{j-1} + h_{j}\right)\left(h_{k-1} + h_{k}\right)}{4h_{i}} \right] u_{i+1jk} + \left[-\epsilon_{i+1/2,j,k}^{x} \frac{\left(h_{j-1} + h_{j}\right)\left(h_{k-1} + h_{k}\right)}{4h_{i}} \right] u_{i+1jk} + \left[-\epsilon_{i+1/2,j,k}^{x} \frac{\left(h_{j-1} + h_{j}\right)\left(h_{k-1} + h_{k}\right)}{4h_{i}} \right] u_{i+1jk} + \left[-\epsilon_{i+1/2,j,k}^{x} \frac{\left(h_{j-1} + h_{j}\right)\left(h_{k-1} + h_{k}\right)}{4h_{i}} \right] u_{i+1jk} + \left[-\epsilon_{i+1/2,j,k}^{x} \frac{\left(h_{j-1} + h_{j}\right)\left(h_{k-1} + h_{k}\right)}{4h_{i}} \right] u_{i+1jk} + \left[-\epsilon_{i+1/2,j,k}^{x} \frac{\left(h_{j-1} + h_{j}\right)\left(h_{k-1} + h_{k}\right)}{4h_{i}} \right] u_{i+1jk} + \left[-\epsilon_{i+1/2,j,k}^{x} \frac{\left(h_{j-1} + h_{j}$$

$$\left[-\epsilon_{i,j-1/2,k}^{y} \frac{\left(h_{i-1}+h_{i}\right)\left(h_{k-1}+h_{k}\right)}{4h_{j-1}}\right] u_{ij-1k} + \left[-\epsilon_{i,j+1/2,k}^{y} \frac{\left(h_{i-1}+h_{i}\right)\left(h_{k-1}+h_{k}\right)}{4h_{j}}\right] u_{ij+1k} + \left[-\epsilon_{i,j+1/2,k}^$$

$$\left[-\epsilon_{i,j,k-1/2}^{k}\frac{\left(h_{i-1}+h_{i}\right)\left(h_{j-1}+h_{j}\right)}{4h_{k-1}}\right]u_{ijk-1}+\left[-\epsilon_{i,j,k+1/2}^{k}\frac{\left(h_{i-1}+h_{i}\right)\left(h_{j-1}+h_{j}\right)}{4h_{k}}\right]u_{ijk+1}=0$$

$$magic \frac{(h_{i-1} + h_i)(h_{j-1} + h_j)(h_{k-1} + h_k)}{8} f_{ijk}$$
 (2)

in which

$$h_i = x_{i+1} - x_i, h_i = y_{i+1} - y_i h_k = z_{k+1} - z_k$$

The delta functions appearing in the right hand side of the starting equations are approximated with linear B-splines (spl0) which spread the point like charge along the nearest neighborhood. The resulting f_{ijk} represent the smearing of the point charges along the grid points.

For more details, including used unit system, please refer to the Michel Holst's thesis and the APBS user guide online. To visualize more clearly the problem, let's explicitly write the first equations for a cubic grid of 5x5x5 containing general coefficients

$$a_{222}u_{222} + a_{122}u_{122} + a_{322}u_{322} + a_{212}u_{212} + a_{232}u_{232} + a_{221}u_{221} + a_{223}u_{223} = f_{222}u_{222} + a_{222}u_{222} + a_{222}u_{222} + a_{312}u_{312} + a_{332}u_{332} + a_{321}u_{321} + a_{323}u_{323} = f_{322}u_{322} + a_{122}u_{122} + a_{322}u_{322} + a_{212}u_{212} + a_{232}u_{232} + a_{221}u_{221} + a_{223}u_{223} = f_{422}u_{232} + a_{132}u_{132} + a_{332}u_{332} + a_{222}u_{222} + a_{242}u_{242} + a_{231}u_{231} + a_{233}u_{233} = f_{232}u_{232} + a_{132}u_{132} + a_{332}u_{332} + a_{222}u_{222} + a_{242}u_{242} + a_{231}u_{231} + a_{233}u_{233} = f_{232}u_{232} + a_{132}u_{132} + a_{332}u_{332} + a_{222}u_{222} + a_{242}u_{242} + a_{231}u_{231} + a_{233}u_{233} = f_{232}u_{232} + a_{232}u_{232} + a_{2$$

. .

in which the nodes are arranged using the natural ordering

$$U = [u_{111}, u_{211}, ..., u_{N_x 11}, u_{121}, ..., u_{221}, u_{321}, ..., u_{N_x 21}..., u_{N_x N_y N_z}]^T$$

Note that the prescribed values of nodes u_{1jk} , $u_{N_x,j,k}$, $u_{i,1,k}$, $u_{i,N_y,k}$, u_{ij1} and u_{ijN_z} along the faces of the box coming from the Dirichlet boundary conditions will have their corresponding elements removed in such a way that only equations for the interior nodes remain. In other words, we will only consider the following set of unknown nodes

$$U = [u_{222}, u_{322}, ..., u_{N_x - 1, 22}, u_{232}, ..., u_{332}, u_{432}, ..., u_{N_x - 2, 32}, ..., u_{N_x - 1, N_y - 1, N_z - 1}]^T$$

in such a way that the previous equations become

$$a_{222}u_{222} + a_{322}u_{322} + a_{232}u_{232} + a_{223}u_{223} = f_{222} - a_{122}u_{122} - a_{212}u_{212} - a_{221}u_{221} \equiv b_{222}$$

$$a_{322}u_{322} + a_{222}u_{222} + a_{422}u_{422} + a_{332}u_{332} + a_{323}u_{323} = f_{322} - a_{312}u_{312} - a_{321}u_{321} \equiv b_{322}$$

$$a_{422}u_{422} + a_{322}u_{322} + a_{232}u_{232} + a_{223}u_{223} = f_{422} - a_{122}u_{122} - a_{212}u_{212} - a_{221}u_{221} \equiv b_{422}u_{122} + a_{122}u_{122} + a_{122}u_{122} = b_{122}u_{122} + a_{122}u_{122} + a_{122}u_{122} = b_{122}u_{122} + a_{122}u_{122} + a_{122}u_{122} + a_{122}u_{122} + a_{122}u_{122} = b_{122}u_{122} + a_{122}u_{122} + a_{122}u_{1$$

$$a_{232}u_{232} + a_{332}u_{332} + a_{222}u_{222} + a_{242}u_{242} + a_{233}u_{233} = f_{232} - a_{132}u_{132} - a_{231}u_{231} \equiv b_{232}u_{232} + a_{232}u_{232} + a_{232}u_{2$$

in which the boundary u's are conveniently brought to the right-hand-side of the equations. The resulting left-hand side equations can be written in compact form in term of matrix vector product as follows

$$Au = b$$

in which

$$u(p) = u_{ijk},$$
 $b(p) = b_{ijk},$ $p = (k-2)(N_x - 2)(N_y - 2) + (j-2)(N_x - 2) + i - 1$
 $i = 2, ..., N_x - 2,$ $j = 2, ..., N_y - 2,$ $k = 2, ..., N_z - 2$

and A is a (seven banded block tri-diagonal form) $(N_x - 2)(N_y - 2)(N_z - 2)$ by $(N_x - 2)(N_y - 2)(N_z - 2)$ squared symmetric positive definite matrix containing the following nonzero elements (see figure 1):

• The main diagonal elements

$$d_{0}(p) = \left[\epsilon_{i-1/2,j,k}^{x} \frac{(h_{j-1} + h_{j}) (h_{k-1} + h_{k})}{4h_{i-1}} + \epsilon_{i+1/2,j,k}^{x} \frac{(h_{j-1} + h_{j}) (h_{k-1} + h_{k})}{4h_{i}} + \epsilon_{i,j-1/2,k}^{y} \frac{(h_{i-1} + h_{i}) (h_{k-1} + h_{k})}{4h_{j-1}} + \epsilon_{i,j+1/2,k}^{y} \frac{(h_{i-1} + h_{i}) (h_{k-1} + h_{k})}{4h_{j}} + \epsilon_{i,j,k-1/2}^{k} \frac{(h_{i-1} + h_{i}) (h_{j-1} + h_{j})}{4h_{k-1}} + \epsilon_{i,j,k+1/2}^{k} \frac{(h_{i-1} + h_{i}) (h_{j-1} + h_{j})}{4h_{k}} + \epsilon_{i,j,k+1/2}^{k} \frac{(h_{i-1} + h_{i}) (h_{i-1} + h_{i})}{4h_{k}} + \epsilon_{i,j,k+1/2}^{k} \frac{(h_{i-1} + h_{i})}{4h_{k}} + \epsilon_{i,j,k+1/2}^{k} \frac{(h_{i-1} + h_{i}) (h_{i-1} + h_{i})}{4h_{k}} + \epsilon_{i,j,k+1/2}^{k} \frac{(h_{i-1} + h_{i}) (h_{i-1} + h_{i})}{4h_{k}} + \epsilon_{i,j,k+1/2}^{k} \frac{(h_{i-1} + h$$

• The Next upper band diagonal, which is shifted in one column to the left from the first column, contains the following elements

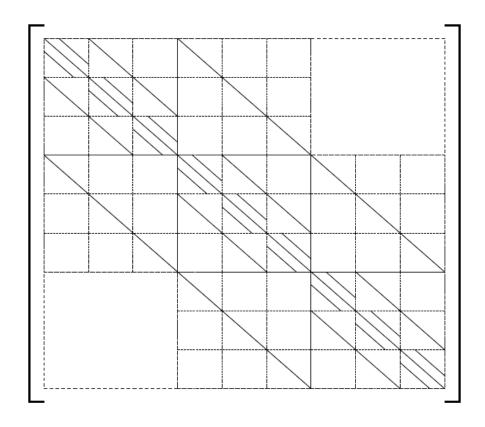


Figure 1: A Matrix representation

$$d_1(p) = -\epsilon_{i+1/2,j,k}^x \frac{(h_{j-1} + h_j)(h_{k-1} + h_k)}{4h_i}$$
(4)

ullet The second upper band diagonal which is shifted N_x-2 columns from the first column

$$d_2(p) = \left[-\epsilon_{i,j+1/2,k}^y \frac{(h_{i-1} + h_i)(h_{k-1} + h_k)}{4h_j} \right]$$
 (5)

• The third upper band diagonal which is shifted $(N_x - 2)(N_y - 2)$ columns from the first column

$$d_3(p) = \left[-\epsilon_{i,j,k+1/2}^k \frac{(h_{i-1} + h_i)(h_{j-1} + h_j)}{4h_k} \right]$$
 (6)

The remaining elements of the upper triangular squared matrix A are set equal to zero. By symmetry we obtain the lower triagonal elements of the matrix A. Because the matrix A is sparse and large, we can implement efficient methods that optimally solve the linear system for U. Specifically, we use the biconjugate gradient stabilized method combined with the inexact LU decomposition of the matrix A. Having the numerical values for the nodes in the interior of the box, we finally add the previously removed prescribed values along the six faces to get the solution over the complete set of grid points.

Dirichlet Boundary Condition

The values of nodes u_{1jk} , $u_{N_x,j,k}$, $u_{i,1,k}$, $u_{i,N_y,k}$, u_{ij1} and u_{ijN_z} along the six faces of the box is set to the values prescribed by a Debye-Hückel model for a multiple, non-interacting spheres with a point charges. The sphere radii are set to the atomic radii of the biomolecule and the sphere charges are set to the total charge of the protein.

Focus Boundary Condition

Our code uses linear interpolation to obtain the value of the potential at the six faces of the target box from the value of the potential obtained at larger domain.

Computational algorithm for Dirichlet boundary conditions

- 1. The code reads the (target) input file .inm to get the APBS input files (shifted dielectric coefficients, kappa function and pqr data file) as well as the number of (target) grid points, the box Lengths, temperature, and bulk properties (ionic strength and solvent dielectric coefficient) among other parameters.
- 2. The center of the grid is evaluated from the corresponding pqr file.
- 3. By using linear B-splines, the charge density is discretized to get f_{ijk} for $i = 1, ..., N_x$, $j = 1, ..., N_y$, $k = 1, ..., N_z$. The Dirichlet boundary condition along the six faces of the box u_{1jk} , $u_{N_x,j,k}$, $u_{i,1,k}$, $u_{i,N_y,k}$, u_{ij1} and u_{ijN_z} are calculated by using the temperature, the value of the bulk dielectric coefficient (usually water) and ionic strength.

- 4. The nonzero components of the matrix A, e.g., the diagonal elements $d_0(p), d_1(p), d_2(p)$, and $d_3(p)$, for $p = (k-2)(N_x-2)(N_y-2) + (j-2)(N_x-2) + i-1$ and $i = 2, ..., N_x 1$, $j = 2, ..., N_y 1$, $k = 2, ..., N_z 1$ are evaluated by using the expressions (3), (4),(5), and (6). The values for the shifted dielectric coefficients and kappa function elements are obtained from the APBS input files. The values of the mesh size h_i, h_j and h_k are obtained from the number of grid points and the Length of the box. Next, the sparse upper triangular matrix A is constructed by filling with zeros the remaining elements of the matrix A. Next, the lower triangular elements of the matrix A are obtained by using the following symmetry property $A_{pq} = A_{qp}$ for $q = 1, ..., (N_x 2)(N_y 2)(N_z 2)$ and $p = q, ..., (N_x 2)(N_y 2)(N_z 2)$.
- 5. The elements of b_{ijk} are evaluated by using the values obtained for the discretized charge density f_{ijk} and the values of the Dirichlet boundary elements multiplied by the appropriate shifted dielectric coefficient values. The natural ordering $p = (k-2)(N_x 2)(N_y 2) + (j-2)(N_x 2) + i 1$ and $i = 2, ..., N_x 1$, $j = 2, ..., N_y 1$, $k = 2, ..., N_z 1$ is used to construct the corresponding vector b(p) (one index) from the data array structure (three indices) b_{ijk} .
- 6. The inexact LU decomposition of the matrix A is performed. The default tolerance value is set equal to 0.25 which provides a fast evaluation of the matrices L and U.
- 7. The resulting L and U matrices, the matrix A and the vector b are used to approximately solve Au = b for the vector u using the biconjugate gradient stabilized method. The default accuracy is set equal to 10^-9 and the maximum number of iteration equal to 800.
- 8. The natural ordering relationship is used to convert the resulting vector u(p) to data array structure to get the numerical solution for u_{ijk} for $i=2,...,N_x-1, \quad j=2,...,N_y-1, \quad k=2,...,N_z-1.$
- 9. Finally the previously removed values of the nodes at the faces of the box are used to obtain the solution for the nodes u_{ijk} over the complete set of grid points, namely for $i = 1, ..., N_x$, $j = 1, ..., N_y$, $k = 1, ..., N_z$.
- 10. The electrostatic potential u_{ijk} and the charge f_{ijk} maps are saved in dx format files.

- 11. If required, it calculates the energy using linear interpolation to evaluate the solution obtained in the grid at the exact location of the point-like charges.
- 12. The electrostatic potential surface $u_{ij(N_z+1)/2}$ is saved in tiff and fig format files for visualization purpose.

Computational algorithm for Focus boundary conditions

The algorithm reads the target input file finding that the boundary condition line says "focusname.inm" instead of "sdh". Then the matlab code automatically first reads that input file "focusname.inm" to solve the PB equation in the specified coarse grid using Dirichlet boundary condition as explained previously. It saves the resulting electrostatic potential solution in a temporary dx formatted file and then perform the following steps;

- 1. The code reads the (target) input file .inm to get the APBS input files (shifted dielectric coefficients, kappa function and pqr data file) as well as the number of (target) grid points, the box Lengths, temperature, and bulk properties (ionic strength and solvent dielectric coefficient) among other parameters.
- 2. The center of the grid is evaluated from the corresponding pqr file.
- 3. By using linear B-splines, the charge density is discretized to get f_{ijk} for $i = 1, ..., N_x$, $j = 1, ..., N_y$, $k = 1, ..., N_z$. The Dirichlet boundary condition along the six faces of the target (smaller) box u_{1jk} , $u_{N_x,j,k}$, $u_{i,1,k}$, $u_{i,N_y,k}$, u_{ij1} and u_{ijN_z} are calculated by using a three-linear interpolation for the electrostatic potential solution obtained previously at larger boxsides (Focus boundary Condition).
- 4. The nonzero components of the matrix A, e.g., the diagonal elements $d_0(p), d_1(p), d_2(p)$, and $d_3(p)$, for $p = (k-2)(N_x-2)(N_y-2) + (j-2)(N_x-2) + i-1$ and $i = 2, ..., N_x 1$, $j = 2, ..., N_y 1$, $k = 2, ..., N_z 1$ are evaluated by using the expressions (3), (4),(5), and (6). The values for the shifted dielectric coefficients and kappa function elements are obtained from the APBS input files. The values of the mesh size h_i, h_j and h_k are obtained from the number of grid points and the Length of the box. Next, the sparse upper triangular matrix A is constructed by filling with zeros the remaining elements of the matrix A. Next, the lower triangular elements of

the matrix A are obtained by using the following symmetry property $A_{pq} = A_{qp}$ for $q = 1, ..., (N_x - 2)(N_y - 2)(N_z - 2)$ and $p = q, ..., (N_x - 2)(N_y - 2)(N_z - 2)$.

- 5. The elements of b_{ijk} are evaluated by using the values obtained for the discretized charge density f_{ijk} and the values of the Dirichlet boundary elements multiplied by the appropriate shifted dielectric coefficient values. The natural ordering $p = (k-2)(N_x 2)(N_y 2) + (j-2)(N_x 2) + i 1$ and $i = 2, ..., N_x 1$, $j = 2, ..., N_y 1$, $k = 2, ..., N_z 1$ is used to construct the corresponding vector b(p) (one index) from the data array structure (three indices) b_{ijk} .
- 6. The inexact LU decomposition of the matrix A is performed. The default tolerance value is set equal to 0.25 which provides a fast evaluation of the matrices L and U.
- 7. The resulting L and U matrices, the matrix A and the vector b are used to approximately solve Au = b for the vector u using the biconjugate gradient stabilized method. The default accuracy is set equal to 10^-9 and the maximum number of iteration equal to 800.
- 8. The natural ordering relationship is used to convert the resulting vector u(p) to data array structure to get the numerical solution for u_{ijk} for $i=2,...,N_x-1, \quad j=2,...,N_y-1, \quad k=2,...,N_z-1.$
- 9. Finally the previously removed values of the nodes at the faces of the box are used to obtain the solution for the nodes u_{ijk} over the complete set of grid points, namely for $i = 1, ..., N_x$, $j = 1, ..., N_y$, $k = 1, ..., N_z$.
- 10. The electrostatic potential u_{ijk} and the charge f_{ijk} maps are saved in dx format files.
- 11. If required, it calculates the energy using linear interpolation to evaluate the solution obtained in the grid at the exact location of the point-like charges.
- 12. The electrostatic potential surface $u_{ij(N_z+1)/2}$ is saved in tiff and fig format files for visualization purpose.

Comment1: Note that in this case the user have to provide two inm.files and the corresponding dx and pqr files for both the coarse and target grids.

Comments2: In this version the user have to provide two pqr files, one representing the molecule by which the PB eq is solved and, the second one to define the center of the grid. It may be the same than the first one, but in general, for complex systems they are not.

Comments3: Fianlly, the user have to provide both directories for the input and output files respectively. In this way the user doesn't have to edit the source files at all. Just need to provide the name of the input file and the full path as the only argument in the Matlab function MPABS (x).

Input File structure

The .inm file parsing is strict. Input must contain the value of these parameters in exactly this order in a column:

```
dime
glen
T
bulk
bc
digpres
dielx_str
diely_str
dielz_str
kappa_str
pqr_str
pqr_cent_str
energy
in_name_str
name_str
```

This line is left to specify the number of grid points. The values for this keyword are:

nx ny nz

dime

the (integer) number of grid points in the x-, y-, and z-directions, respectively.

glen

This line is left to specify the mesh domain lengths; this may be different in each direction. The values for this keyword are:

xlen ylen zlen

the (floating point) grid lengths in the x-, y-, and z-directions (respectively) in Å.

T

This line is left to specify the temperature of the system in Kelvin. The value for this keyword is:

 \mathbf{T}

bulk

This line is left to specify the bulk properties. The values for this keyword are:

I Solv-epsilon

where **I** is the ionic strength defined by $I = 0.5 \sum_{i} c_i z_i^2$ where the dummy sum is over all different ionic species, z_i and c_i are the valence and ionic concentration in moles respectively. The other parameter "**Solv-epsilon**" represents the value of the solvent dielectric coefficient (usually equal to 78.54 for water).

bc

This line is left to specify the boundary condition to be used to solve the linear PB equation. The words for this keyword can be either

sdh

if the user requires Dirichlet boundary condition, or

focusname.inm

if the user requires focus boundary condition.

digpres

This line is left to specify the number of digits of presicion in the residual error obtained in the solution of the linear eq generated by the biconjugated gradient method. The value for this keyword is:

N

where usually this number is set equal to 6.

 $dielx_str$

This line is left to specify the name of the shifted x-component of the dielectric coefficients in dx format as generated by apbs. For instances

xfilename.dx

 $diely_\,str$

This line is left to specify the name of the shifted y-component of the dielectric coefficients in dx format as generated by apbs. For instances

yfilename.dx

dielz str

This line is left to specify the name of the shifted z-component of the dielectric coefficients in dx format as generated by apbs. For instances

zfilename.dx

kappa str

This line is left to specify the name of the ionic accessibility coefficients kappa in dx format (.dx extension) as generated by apbs. For instances

kappafilename.dx

 pqr_str

This line is left to specify the name of the pqr file (.pqr extension) generated by pdbtopqr determining the target molecule by which the PB eq will be solved. For instances

moleculetarget.pqr

 pqr_cent_str

This line is left to specify the name of the pqr file (.pqr extension) generated by pdbtopqr determining the molecule defining the center of grid. It may be the same than pqr_str if there is only one molecule. In complex systems having more than one molecular species it uses to define one of them as the molecule reference. In such cases the center of grid would be defined by the coordinates of such molecule for all the molecular species contained in the complex system. For instances

moleculereference.pqr

In some cases the user prefers to define the three coordinates (in A units) of the center of grid explicitly, namely xcent ycent zcent, instead of calculating these coordinates from the location of the atoms defined in an specific molecule pqr file. Because this line only reads pqr file, the easy way to do this is by writing a pqr file for only one atom having as coordinates those defined by the center of grid. For instances, if the user want to define the center of grid at the origin, namely at 0 0 0, then the user should write a pqr file like this one

ATOM 1 I ION 1 0.000 0.000 0.000 1.00 1.00 or in general like this

ATOM 1 I ION 1 xcent ycent zcent 1.00 1.00

energy

This line is left to specify the calculation of the energy. The words for this line can be either

calceneryes

if the user requires the calculation of the enrgy, or any other word, for instances calcenerno

if the user doesn't want to.

```
in\_name\_str
```

This line is left to specify the full path to the input file directory containing all the dx and pqr files as well as focusname.inm if focus boundary condition is required. For instances

```
C:\User\myname\matlabworkspace\Input_Files for Windows users or 
/Users/myname/matlabworkspace/Input_Files for linux users.
```

```
name\_\,str
```

This line is left to specify the full path to the output file directory that will contain the resulting dx, fig and jpg files. For instances

```
C:\User\myname\matlabworkspace\systemname for Windows users or 
/Users/myname/matlabworkspace/systemname for linux users.
```

Input file examples

Just one inm. file is required if the user is not using focus boundary condition as you can see in the following example:

```
%|-example solvated-born.inm file using Dirichlet Boundary Condition-|
65 65 65
12 12 12
298.15
0.0 78.54
sdh
6
solvated-born-dielx.dx
```

```
solvated-born-diely.dx
  solvated-born-dielz.dx
  solvated-born-kappa.dx
  born-ion.pqr
  born-ion.pqr
  calceneryes
  C:\Users\Marce\Matlab work space\Input Files
  C:\Users\Marce\Matlab work space\born model
  Otherwise, two inm. files are required if the user use the focus boundary condition as
you can see in the following example:
  % - example solvated-born.inm file using focus boundary condition-
  65\ 65\ 65
  12 12 12
  298.15
  0.0 78.54
  focusname.inm
  6
  solvated-born-dielx.dx
  target-solvated-born-diely.dx
  target-solvated-born-dielz.dx
  target-solvated-born-kappa.dx
  born-ion.pqr
  born-ion.pqr
  calcenerno
  C:\Users\Marce\Matlab\_work\_space\Input\_Files
  % - example focus name in m file for the calculation of the elect pot in the coarse grained
calculation
  65 65 65
  50 50 50
```

```
298.15
0.0 78.54
sdh
6
coarse-born-dielx.dx
coarse-born-diely.dx
coarse-born-dielz.dx
coarse-born-kappa.dx
born-ion.pqr
complex.pqr
calceneryes
C:\Users\Marce\Matlab_work_space\Input_Files
C:\Users\Marce\Matlab_work_space\coarse_born_model
```