

Einstein's Notation

Kartik Gokhale

April 2021

1 Introduction

The Einstein notation or Einstein summation convention is a notational convention that implies summation over a set of indexed terms in a formula, thus achieving notational brevity. When an index appears twice in a term and is not otherwise defined, it is taken as the sum over all possible values of said index. For example,

$$\sum_{i=1}^3 c_i x^i = c_1 x^1 + c_2 x^2 + c_3 x^3 \quad (1)$$

would just be written as $c_i x^i$. The summation is implicit. According to this convention:

1. The whole term is summed over any index which appears **twice** in that term
2. Any index appearing only once in the term is not summed over and must be present in both LHS and RHS

2 Two important symbols

2.1 Kronecker Delta

δ_{ij} is the **Kronecker Delta** symbol. It is defined as

$$\delta_{ij} = \begin{cases} 1 & \text{if } i = j \\ 0 & \text{if } i \neq j \end{cases}$$

2.2 Levi-Civita

ϵ_{ijk} is the **Levi-Civita** symbol (in 3 dimensions). It is defined as

$$\epsilon_{ijk} = \begin{cases} 1 & \text{if } ijk \text{ is a cyclic permutation of } xyz, \text{ eg } xyz, yzx, zxy \\ -1 & \text{if } ijk \text{ is an anticyclic permutation of } xyz, \text{ eg } xzy, zyx, yxz \\ 0 & \text{if in } ijk \text{ any index is repeating, eg } xxy, xyy \text{ etc} \end{cases}$$

3 Application to Vectors

3.1 Dot Product

Consider 2 vectors in 3-d space.

$$\begin{aligned}\vec{a} &= a_1\hat{i} + a_2\hat{j} + a_3\hat{k} \\ \vec{b} &= b_1\hat{i} + b_2\hat{j} + b_3\hat{k}\end{aligned}$$

Hence, their dot product is

$$\vec{a} \cdot \vec{b} = \sum_{i=1}^3 \sum_{j=1}^3 a_i b_j \delta_{ij}$$

This can be written quite succinctly using Einstein's notation as

$$\vec{a} \cdot \vec{b} = a_i b_j \delta_{ij} \quad (2)$$

As the summation over the indices is implicit

3.2 Cross Product

Consider 2 vectors again

$$\begin{aligned}\vec{a} &= a_1\hat{i} + a_2\hat{j} + a_3\hat{k} \\ \vec{b} &= b_1\hat{i} + b_2\hat{j} + b_3\hat{k}\end{aligned}$$

Hence, their cross product is

$$\vec{a} \times \vec{b} = \begin{vmatrix} \hat{i} & \hat{j} & \hat{k} \\ a_1 & a_2 & a_3 \\ b_1 & b_2 & b_3 \end{vmatrix}$$

This can be written quite succinctly using Einstein's notation as

$$\vec{a} \times \vec{b} = a_i b_j \epsilon_{ijk} \hat{e}_k \quad (3)$$

where \hat{e}_k is the unit vector in k direction.

As the summation over the indices is implicit

3.3 Product of Levi-Civitas

We can resolve the product of two Levi-Civitas(in 3-d) as

$$\epsilon_{ijk} \epsilon_{lpq} = \begin{vmatrix} \delta_{il} & \delta_{ip} & \delta_{iq} \\ \delta_{jl} & \delta_{jp} & \delta_{jq} \\ \delta_{kl} & \delta_{kp} & \delta_{kq} \end{vmatrix}$$

4 Problems

4.1 BAC-CAB Rule

Consider 3 vectors, $\vec{A}, \vec{B}, \vec{C}$. Our goal is to simplify

$$\vec{A} \times (\vec{B} \times \vec{C})$$

Solution

Let us define another vector $\vec{D} = \vec{B} \times \vec{C}$

Hence, our expression simplifies to

$$\begin{aligned}\vec{A} \times (\vec{B} \times \vec{C}) &= \vec{A} \times \vec{D} \\ &= a_i d_j \epsilon_{ijk} \hat{e}_k\end{aligned}$$

Since D is a cross product of 2 vectors

$$d_j = b_l c_m \epsilon_{lmj} \hat{e}_j$$

By back-substitution

$$a_i d_j \epsilon_{ijk} \hat{e}_k = a_i b_l c_m \epsilon_{lmj} \epsilon_{ijk} \hat{e}_k$$

Recall product of 2 Levi-Civitas

$$\epsilon_{ijk} \epsilon_{lpq} = \begin{vmatrix} \delta_{il} & \delta_{ip} & \delta_{iq} \\ \delta_{jl} & \delta_{jp} & \delta_{jq} \\ \delta_{kl} & \delta_{kp} & \delta_{kq} \end{vmatrix}$$

Substituting Values and simplifying

$$\begin{aligned}a_i b_l c_m \epsilon_{lmj} \epsilon_{ijk} \hat{e}_k &= a_i b_j c_m (\delta_{lk} \delta_{mi} - \delta_{li} \delta_{mk}) \hat{e}_k \\ &= a_i b_j c_m \delta_{lk} \delta_{mi} \hat{e}_k - a_i b_j c_m \delta_{li} \delta_{mk} \hat{e}_k\end{aligned}$$

Recall these are just dot products

$$\begin{aligned}&= b_k (\vec{A} \cdot \vec{C}) \hat{e}_k - c_k (\vec{A} \cdot \vec{B}) \hat{e}_k \\ &= \vec{B} (\vec{A} \cdot \vec{C}) - \vec{C} (\vec{A} \cdot \vec{B})\end{aligned}$$

Hence, we obtain the very popular result

$$\vec{A} \times (\vec{B} \times \vec{C}) = \vec{B} (\vec{A} \cdot \vec{C}) - \vec{C} (\vec{A} \cdot \vec{B}) \quad (4)$$

4.2 More Problems

Question The gradient operator ∇ behaves like a vector in “some sense”. For example, divergence of a curl ($\nabla \cdot \nabla \times \vec{A} = 0$) for any \vec{A} , may suggest that it is just like $\vec{A} \cdot \vec{B} \times \vec{C}$ being zero if any two vectors are equal. Prove that $\nabla \times \nabla \times \vec{F} = \nabla(\nabla \cdot \vec{F}) - \nabla^2 \vec{F}$. To what extent does this look like the well known expansion of $\vec{A} \times \vec{B} \times \vec{C}$?

Solution

This question is very similar to the previous one.

Now we will represent each component of ∇ , $\frac{\partial}{\partial i}$ by ∂_i ,

$$\begin{aligned}
 \nabla \times (\nabla \times \vec{F}) &= \epsilon_{ijk} \partial_j (\nabla \times \vec{F})_k \hat{e}_i = \epsilon_{ijk} \epsilon_{klm} \partial_j \partial_l F_m \hat{e}_i = \epsilon_{kij} \epsilon_{klm} \partial_j \partial_l F_m \hat{e}_i \\
 &= (\delta_{il} \delta_{jm} - \delta_{im} \delta_{jl}) \partial_j \partial_l F_m \hat{e}_i & (\because \epsilon_{ijk} \epsilon_{ilm} = \delta_{jl} \delta_{km} - \delta_{jm} \delta_{kl}) \\
 &= (\partial_m \partial_i F_m - \partial_l \partial_l F_i) \hat{e}_i & (\because \delta_{ij} A_j = A_i) \\
 &= \partial_i (\partial_m F_m) \hat{e}_i - \partial_l \partial_l (F_i \hat{e}_i) & (\because \partial_m \partial_i = \partial_i \partial_m) \\
 &= \nabla(\nabla \cdot \vec{F}) - \nabla^2 \vec{F} \quad \blacksquare
 \end{aligned}$$

Notice that this proof works because $\partial_m \partial_i = \partial_i \partial_m$ which is just fancy notation for $\frac{\partial^2}{\partial x \partial y} = \frac{\partial^2}{\partial y \partial x}$

We know that $\vec{A} \times \vec{B} \times \vec{C} = \vec{B}(\vec{A} \cdot \vec{C}) - (\vec{A} \cdot \vec{B})\vec{C}$. If we write $\vec{A} = \nabla$, $\vec{B} = \nabla$, $\vec{C} = \vec{F}$, we get $\nabla \times \nabla \times \vec{F} = \nabla(\nabla \cdot \vec{F}) - (\nabla \cdot \nabla)\vec{F} = \nabla(\nabla \cdot \vec{F}) - \nabla^2 \vec{F}$. So this serves as a fake proof of the identity

Question As a more involved example, show that the operator $\mathbf{L} = -i\vec{r} \times \nabla$ where ($i = \sqrt{-1}$) satisfies $\mathbf{L} \times \mathbf{L}f = i\mathbf{L}f$ where f is an arbitrary test function. (Notice that the cross product of an operator with itself does not necessarily vanish, can you see why?)

Solution

$$\begin{aligned}\mathbf{L} \times \mathbf{L}f &= \epsilon_{ijk} L_j L_k f \hat{e}_i \\ &= -i\epsilon_{ijk}\epsilon_{jlm} r_l \partial_m L_k f \hat{e}_i \\ &= -i\epsilon_{jki}\epsilon_{jlm} r_l \partial_m L_k f \hat{e}_i \\ &= -i(\delta_{kl}\delta_{im} - \delta_{km}\delta_{il}) r_l \partial_m L_k f \hat{e}_i \\ &= -i(r_l \partial_i L_k f - r_i \partial_k L_l f) \hat{e}_i \\ &= -i(A_i - B_i) \hat{e}_i\end{aligned}$$

$$A_i = r_k \partial_i L_k f = -i\epsilon_{kpq} r_k \partial_i r_p \partial_q f$$

$$B_i = r_i \partial_k L_k f = -i\epsilon_{kpq} r_i \partial_k r_p \partial_q f$$

Now we can't directly switch the order of ∂_i & ∂_q ,

we have to use the product rule: $(\partial_a(r_b \partial_c))f = (r_b \partial_a \partial_c)f + ((\partial_a r_b) \partial_c)f$

$$\begin{aligned}\Rightarrow A_i &= \cancel{-i\epsilon_{kpq} r_k r_p \partial_i \partial_q f}^0 - i\epsilon_{kpq} r_k (\partial_i r_p) \partial_q f = -i\epsilon_{kpq} r_k (\delta_{ip}) \partial_q f = -i\epsilon_{k iq} r_k \partial_q f \\ (\because \epsilon_{kpq} r_k r_p \partial_i \partial_q f &= -\epsilon_{pkq} r_k r_p \partial_i \partial_q f = -\epsilon_{kpq} r_p r_k \partial_i \partial_q f = -\epsilon_{kpq} r_k r_p \partial_i \partial_q f \\ &\Rightarrow \epsilon_{kpq} r_k r_p \partial_i \partial_q f = 0)\end{aligned}$$

The first term cancels because in the summation 1 term will have $k, q = x, y$ and another term $k, q = y, z$. Both differ only in their sign, which is due to the ϵ_{kpq} .

So they cancel out

$$\begin{aligned}\Rightarrow B_i &= \cancel{-i\epsilon_{kpq} r_i r_p \partial_k \partial_q f}^0 - i\epsilon_{kpq} r_i (\partial_k r_p) \partial_q f = -i\epsilon_{kpq} r_i (\delta_{kp}) \partial_q f = -i\epsilon_{ppq} r_i \partial_q f = 0 \\ (\because \epsilon_{kpq} r_i r_p \partial_k \partial_q f &= -\epsilon_{qpk} r_i r_p \partial_k \partial_q f = -\epsilon_{kpq} r_i r_k \partial_q \partial_k f = -\epsilon_{kpq} r_i r_p \partial_k \partial_q f \\ &\Rightarrow \epsilon_{kpq} r_i r_p \partial_k \partial_q f = 0)\end{aligned}$$

Similar reasoning as above is used here

$$\Rightarrow \mathbf{L} \times \mathbf{L}f = -\epsilon_{k iq} r_k \partial_q f \hat{e}_i = \epsilon_{ikq} r_k \partial_q f \hat{e}_i = i\mathbf{L}f \quad \blacksquare$$

Notice that we have used $\partial_a r_b = \delta_{ab}$. This is because r_a is simply x if $a = x$, y if $a = y$ and so on. And if we differentiate that w.r.t $b = x$ or y or z , then it is 1 if $a = b$, and 0 otherwise, because $\frac{\partial}{\partial x} x = 1$, $\frac{\partial}{\partial x} y = 0$

The cross product of the operator with itself doesn't vanish here because the operators act on different objects. One \mathbf{L} acts on $\mathbf{L}f$, and the other on f .