

Problem Set 2

1 Blackbody Flux

Derive the blackbody flux formula

$$F = \sigma T^4 \quad (1)$$

from the Planck function for specific intensity,

$$B_\nu = \frac{2h\nu^3}{c^2} \frac{1}{e^{\frac{h\nu}{kT}} - 1} \quad (2)$$

You may use the integral $\int_0^\infty \frac{dx}{x^5(e^{1/x}-1)} = \pi^4/15$. Those who know what they have to do can stop reading here. Those who need a bit of help should read on.

Imagine a perfectly flat, blackbody patch at the center of a sphere of radius R . The patch has tiny surface area dA and is lying flat at $r = 0$ in the $\theta = \pi/2$ plane (in spherical coordinates, where r , θ , ϕ are the radius, polar angle, and azimuth). Only one side of the patch is warm; the other side is at absolute zero. Assume that the warm side radiates as a perfect thermal blackbody.

A thermal blackbody emits radiation isotropically — i.e., with no preference for direction. In other words, imagine each point on the patch beaming out a tiny, uniform, hemispherical dome of rays (like an umbrella). It is a hemispherical dome and not a full sphere because only one side of the patch is warm. Note that the Planck blackbody function has no variable inside it that specifies direction of emission; the direction of emission doesn't matter for purely thermal emission. The only variables that need to be specified for the Planck function are T and ν .

Despite the isotropic nature of the emission, because the patch is geometrically flat, has definite area, and emits only on one side of itself, a detector glued to the inside of the sphere at $r = R$ detects different amounts of radiation depending on where it is placed. For example, if the detector is glued on the half of the sphere that can't see the bright side of the patch, the detector detects nothing. If the detector is glued to the sphere directly above the patch, the detector picks up the most number of photons, because it sees the full face-on area of the patch (namely, dA). If the detector is glued at an angle to the pole, then it sees less than area dA . If it is glued in the plane of the patch, it sees nothing but an infinitesimally thin line segment.

This problem asks you to place detectors everywhere on the inside of the sphere and sum up the radiation collected; this operation is equivalent to “integrating the Planck function over all solid angles into which the radiation is beamed.”

Calculate the total luminosity (in units of energy/time) emitted by the patch of area dA . You must integrate the Planck function over the entire emitting area (which is maximally dA but in general will be less, depending on the viewing geometry), over all solid angles into which the radiation is beamed, and over all frequencies. Think about placing tiny detectors all over the inside of the sphere and asking how much energy/time each of the detectors receives.

Then divide the luminosity by dA to calculate the total flux (in units of energy/time/area) emitted by the patch. You should recover the usual blackbody flux formula, σT^4 . By definition, σT^4 is the total amount of energy radiated per time per unit area of a blackbody surface, radiated into all solid angles and over all frequencies.

2 Flat Disks

This problem forms the foundation for understanding the spectral energy distributions of circumstellar disks, e.g., those surrounding pre-main-sequence stars and AGN.

Consider a perfectly flat, blackbody disk encircling a blackbody star. The star has radius R_* and effective temperature T_* . The disk begins at a stellocentric radius far from the star, i.e. the inner disk radius $r_i \gg R_*$. The disk extends to infinity.

Calculate the temperature of the disk, $T(r)$, as a function of radius, r . Make whatever approximations you deem necessary in light of the fact that $r_i \gg R_*$. Be sure you get the scaling of T with r correctly. It is less important that you get the numerical coefficient correctly.

This disk is PERFECTLY flat. Do not consider individual particles in the disk.

Also, neglect radial transport of energy. Each annulus is independent of neighboring annuli. This is a fine approximation for thin disks since they transport more radiation vertically than radially (the temperature gradient is steeper vertically than radially).

3 Hyperfine $^3\text{He}^+$

Observations of the hyperfine transition in $^3\text{He}^+$ are used to probe the $^3\text{He}^+/H$ abundance in the galaxy. This abundance reflects the primordial yield from big bang nucleosynthesis and galactic chemical evolution.

3.1

Estimate, using the scaling relations presented in class and whatever facts you remember, the wavelength of the ground-state electronic (i.e. Ly- α equivalent) transition of $^3\text{He}^+$.

3.2

Estimate, using the scaling relations presented in class and whatever facts you remember, the wavelength of the ground-state hyperfine transition of $^3\text{He}^+$. Compare to the true answer of 3.46 cm.