

## Problem Set 11 (part 1)

### 1 Dust Temperatures and Opacity

The bulk of the interstellar radiation field in the Galaxy is from the light of O stars, which have mass  $\sim 100M_{\odot}$ . There are  $\sim 5 \times 10^4$  O stars in the galaxy, distributed over a cylindrical disk with radius 50 kpc and height 200 pc.

**1.1** Estimate the total energy density of starlight in the Galaxy in units of  $\text{eV cm}^{-3}$ . Use the fact that O stars radiate near the Eddington luminosity.

**1.2** This interstellar radiation field heats dust grains in the interstellar medium. Estimate the temperature  $t$  of the largest grains in the ISM, which have a radius of  $a \sim 0.1$  microns. Assume  $Q_{abs} \sim 1$  for wavelengths shorter than  $2\pi a$ , and for longer wavelengths, take  $m = 1 + 0.25i$ . Is the grain hotter, cooler, or equal to the temperature of an ideal blackbody placed in the interstellar radiation field.

**1.3** At what wavelength,  $\lambda_{peak}$ , does the energy density ( $\nu F_{\nu}$ ) of the grains peak?

**1.4** The largest grains carry most of the mass in the interstellar grain distribution. Given a dust-to-gas ratio of 0.01 and an average hydrogen number density  $n_H = 0.1 \text{ cm}^{-3}$ , calculate the specific intensity of the Milky Way at  $\lambda_{peak}$ .

**1.5** Typically, the number density of dust grains scales with size as  $dn/da \propto a^{-3.5}$  in the ISM, where  $dn$  is the number density of grains with radii between  $a$  and  $a + da$ . The law holds over  $a_{min} = 10^{-3}$  microns to  $a_{max} = 0.1$  microns. Plot the grain opacity  $\kappa(\lambda)$  contributed by all dust grains over the wavelength range  $\lambda = 0.1$  to 10 microns. Consider only absorption and neglect scattering. Indicate over each decade of wavelength which grain sizes dominate the opacity.

## Problem Set 10 (part 2)

### 2 Diatomic

Based on Rybicki & Lightman 11.1.

Consider an electrically neutral medium of diatomic molecules in thermal equilibrium at temperature  $T$ . Each molecule contains one nucleus of mass  $m_p$  and one of mass  $2m_p$  at an equilibrium separation  $r_0$ .

**2.1** Estimate  $r_0$  in terms of fundamental constants.

**2.2** Estimate the cross-section  $\sigma_c$  for collisions between molecules.

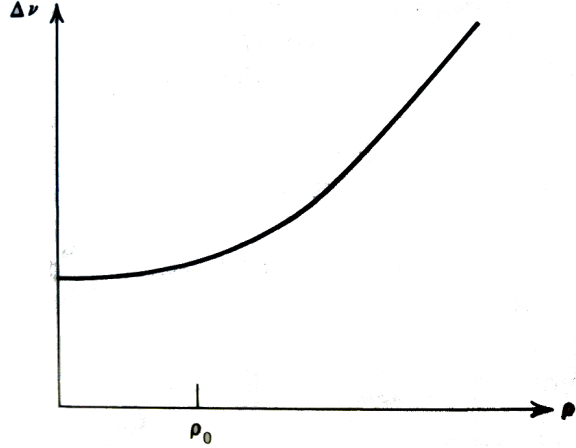


Figure 1: Line width as a function of density for emission from a medium of diatomic molecules.

**2.3** It is experimentally observed that, as a function of mass density  $\rho$  of the medium, the line width of the rotational lines has the form shown in Figure 1. If only Doppler and collisional broadening are present, estimate  $\rho_0$  and show that it may be written completely in terms of fundamental constants, independent of  $m_p$ .

We derived in class that for a synchrotron self-Compton (SSC) emitting source that is static in bulk, the ratio of the luminosity due to first-generation inverse Compton scattering of synchrotron photons to the luminosity due to synchrotron processes is

$$f \equiv \frac{L_{IC,1}}{L_{sync}} = C \nu_m T_{bm}^5 \quad (1)$$

where  $\nu_m$  is the frequency at which the synchrotron spectrum beaks, and  $T_{bm}$  is the brightness temperature of the synchrotron radiation at that frequency.

**2.4** Derive an estimate for the constant  $C$  in terms of fundamental constants. Assume whatever geometry you like for the source. I found I had to use the (usual) constants  $e$ ,  $k$ ,  $c$ , and  $m_e$ .

**2.5** If  $\nu_m \sim 1$  GHz (as it typically is for compact radio sources), what is the maximum  $T_{bm}$  above which  $f > 1$ ? Express in Kelvin.

**2.6** Given  $f$ , what is  $L_{IC,2}/L_{IC,1}$ ? That is, what is the ratio of second-generation to first-generation scattered power? Explain your answer.

**2.7** Consider the situation after  $N$  scatterings, where  $N \rightarrow \infty$ . What is the critical  $f$  for which the total luminosity of the source becomes unbounded? Careful estimates will be rewarded.