

# EPHY105L (Fall Semester 2018-2019)

## Solutions to Problem Sheet 4

$$1. (a) \quad \oint \vec{E} d\vec{a} = \frac{Q_{\text{enclosed}}}{\epsilon_0} = \frac{q_2}{\epsilon_0},$$

$$\oint \vec{D} d\vec{a} = Q_{\text{free (enclosed)}} = q_2$$

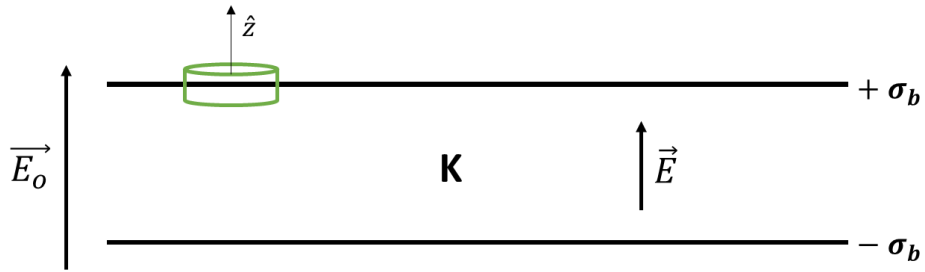
$$(b) \quad \oint \vec{E} d\vec{l} = 0,$$

$$(c) \quad \text{At point A, } \vec{\nabla} \cdot \vec{D} = 0$$

$$(d) \quad \text{At point B, } \vec{\nabla} \cdot \vec{E} = 0$$

2. The electric field inside the dielectric can be calculated using Gauss's law involving  $\vec{D}$ :

$$\oint \vec{D} \cdot d\vec{a} = q_{f, \text{encl}} \quad (1)$$



Let us consider a cylindric pillbox of area \$A\$ half way submerged into the surface of a dielectric. The electric field inside (\$\vec{E}\$) and outside the dielectric (\$\vec{E}\_0\$) are shown in the Figure; the \$\vec{D}\$ will be parallel to the \$\vec{E}\$ with \$D\_0 = \epsilon\_0 E\_0\$ and \$D = \epsilon\_0 K E\$. Since there are no free charges enclosed by the Gaussian surface we must have

$$\oint \vec{D} \cdot d\vec{a} = 0 \quad (2)$$

Since \$\vec{D}\$ is normal to the flat surfaces of the pillbox, we obtain

$$D = D_0$$

Thus we have

$$E = \frac{E_0}{K}$$

Now, the polarization vector is obtained from

$$\vec{D} = \epsilon_0 \vec{E} + \vec{P}. \quad (4)$$

Since \$\vec{D} = \epsilon \vec{E}\$, therefore

$$\vec{P} = (\epsilon - \epsilon_0) \vec{E}. \quad (5)$$

Then the bound surface charge density is

$$\sigma_b = \vec{P} \cdot \hat{n} = (\epsilon - \epsilon_0)E = \frac{(K - 1)}{K} E_0$$

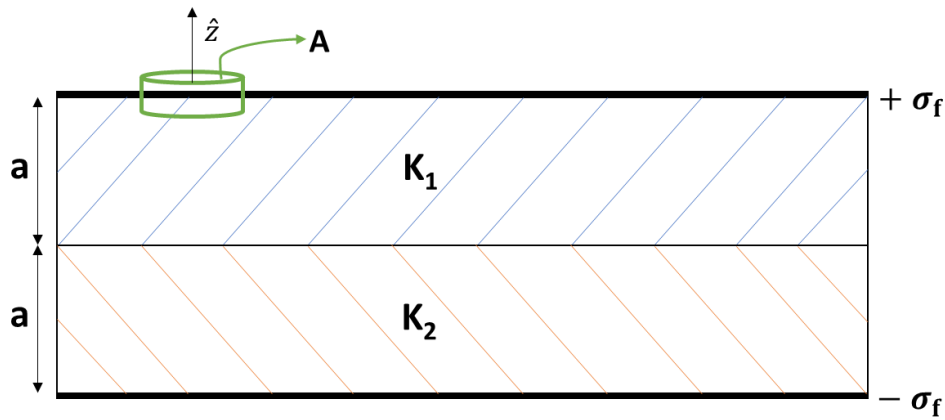
The surface charge density on the lower surface would be  $-\sigma_b$ .

Since the polarization vector has a constant magnitude, therefore bound volume charge density  $\rho_b = -\nabla \cdot \vec{P} = 0$  and bound surface charge density on the upper surface is  $\sigma_b = \vec{P} \cdot \hat{n} = (K - 1) \epsilon_0 \frac{E_0}{K}$ .

The total bound charge will be zero.

3. (a) Gauss theorem in terms of electric displacement vector  $\vec{D}$  is

$$\oint \vec{D} \cdot d\vec{a} = q_{free}$$



where  $q_{free}$  is enclosed free charge. To obtain electric displacement vector within the dielectric slabs 1 and 2, we again consider the cylindrical pillbox Gaussian surface with surface area  $A$ . The free charge enclosed within the Gaussian surface is  $A \cdot \sigma_f$ . Electric field and displacement vector would be pointing down (i.e., towards positive to negatively charged plate). Therefore

$$\oint \vec{D} \cdot d\vec{a} = -D \cdot A = A \cdot \sigma_f$$

or

$$\vec{D} = \sigma_f (-\hat{z}).$$

The displacement vector depends only on free charge on the metallic plates and is same for both dielectrics.

(b) The electric polarization is obtained using

$$\vec{D} = \epsilon_0 \vec{E} + \vec{P}$$

or

$$\vec{P} = \epsilon \vec{E} - \epsilon_0 \vec{E} = (\epsilon - \epsilon_0) \vec{E}.$$

The electric field in each of the dielectric can be obtained using relation

$$\vec{D} = \epsilon \vec{E}.$$

Using the displacement vector from the solution of part (a) as

$$\vec{D} = \sigma_f (-\hat{z}) = \epsilon \vec{E},$$

the electric field inside the dielectric 1 and 2 is

$$\vec{E}_1 = \frac{\sigma_f}{\epsilon_1} (-\hat{z}) \text{ and } \vec{E}_2 = \frac{\sigma_f}{\epsilon_2} (-\hat{z}).$$

Using  $K_1 = \epsilon_1/\epsilon_0$  and  $K_2 = \epsilon_2/\epsilon_0$ , we obtain

$$\vec{P}_1 = \frac{(K_1-1)}{K_1} \sigma_f (-\hat{z}) \text{ and } \vec{P}_2 = \frac{(K_2-1)}{K_2} \sigma_f (-\hat{z}).$$

(c) Since the polarization vector has a constant magnitude, therefore volume bound charge density

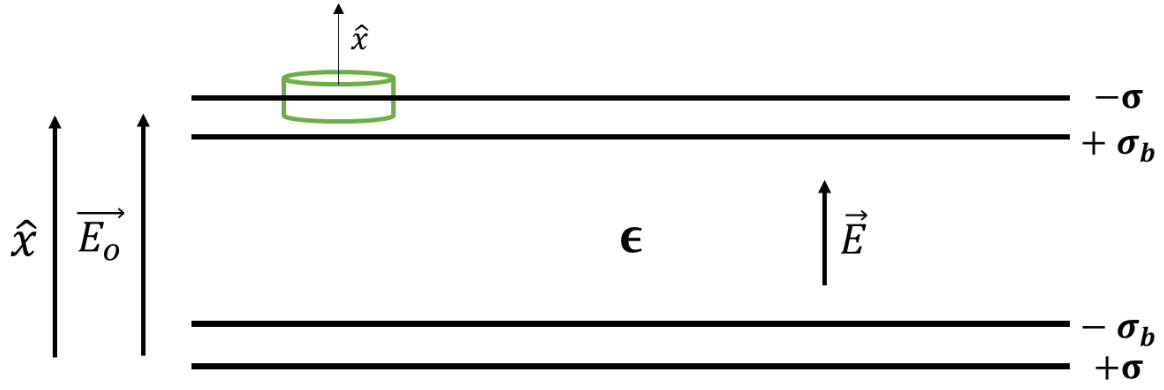
$$\rho_b = -\nabla \cdot \vec{P} = 0.$$

The surface bound charge density can be obtained using  $\sigma_b = \vec{P} \cdot \hat{n}$ , where  $\hat{n}$  is the outward vector normal to the surface of dielectric.



The vector  $\hat{n}$  for the top and bottom layers of the both dielectrics are in  $\hat{z}$  and  $-\hat{z}$  direction respectively. The polarization vector is always pointing in  $-\hat{z}$  direction. Therefore, the surface charge density at the top and bottom layer of the upper dielectric (with dielectric constant  $K_1$ ) is  $\sigma_b = -P_1$  and  $\sigma_b = +P_1$ , whereas at the top and bottom layer of the lower dielectric (with dielectric constant  $K_2$ ) is  $\sigma_b = -P_2$  and  $\sigma_b = +P_2$ .

4. The electric displacement vector  $\vec{D}$  can be calculated using  $\oint \vec{D} \cdot d\vec{a} = q_{\text{free}}$ , which means that displacement vector  $\vec{D} = \sigma (\hat{x})$



Electric field in the slab is  $\vec{E} = \frac{\vec{D}}{\epsilon} = \frac{\sigma}{\epsilon} (\hat{x})$

To find bound surface charge and volume charge densities, we need to calculate polarization vector  $\vec{P}$  as follows

$$\vec{D} = \epsilon_0 \vec{E} + \vec{P} \Rightarrow \epsilon \vec{E} = \epsilon_0 \vec{E} + \vec{P} \Rightarrow \vec{P} = \frac{(\epsilon - \epsilon_0)\sigma}{\epsilon} (\hat{x})$$

a) Volume bound charge density:  $\rho_b = -\nabla \cdot \vec{P}$

$$\begin{aligned} &= -\frac{\partial}{\partial x} (\vec{P}) = -\frac{\partial}{\partial x} \left( 1 - \frac{\epsilon_0}{\epsilon} \right) \sigma = \frac{\partial}{\partial x} \left( \frac{\epsilon_0}{\epsilon} \right) \sigma \\ &= \sigma \epsilon_0 \frac{\partial}{\partial x} \left( \frac{1}{\epsilon} \right) = -\frac{\sigma \epsilon_0}{\epsilon^2} \frac{\partial}{\partial x} (\epsilon) \\ &= -\frac{\sigma \epsilon_0}{\epsilon^2} \left( \frac{\epsilon_2 - \epsilon_1}{d} \right) \end{aligned}$$

b) Bound surface charge density:  $\sigma_b = \vec{P} \cdot \hat{n}$

$$= \frac{(\epsilon - \epsilon_0)\sigma}{\epsilon}$$

At bottom plate,  $x = 0$ , therefore  $\epsilon = \epsilon_1$

At top plate,  $x = d$ , therefore  $\epsilon = \epsilon_2$

Surface charge densities at bottom and top plates will be  $\frac{(\epsilon_1 - \epsilon_0)\sigma}{\epsilon_1}$  and  $\frac{(\epsilon_2 - \epsilon_0)\sigma}{\epsilon_2}$ , respectively.

$$c) \vec{E} = \frac{\vec{D}}{\epsilon} = \frac{\sigma}{\epsilon} (\hat{x}) = \frac{\sigma}{\epsilon_1 + \left( \frac{\epsilon_2 - \epsilon_1}{d} \right) x} (\hat{x})$$

$$d) V = -\int_0^d \vec{E} \cdot d\mathbf{x} = -\int_0^d \frac{\sigma}{\epsilon_1 + \left( \frac{\epsilon_2 - \epsilon_1}{d} \right) x} \cdot dx$$

$$= -\frac{\sigma d}{\epsilon_2 - \epsilon_1} \ln \left[ \epsilon_1 + \left( \frac{\epsilon_2 - \epsilon_1}{d} \right) x \right]_0^d$$

$$= -\frac{\sigma d}{\epsilon_2 - \epsilon_1} [\ln(\epsilon_2) - \ln(\epsilon_1)]$$

$$= \frac{\sigma d}{\epsilon_2 - \epsilon_1} [\ln(\epsilon_1/\epsilon_2)]$$

5. Follow the answer of question 3 with following substitutions:

$$\sigma = 30 \frac{\text{mC}}{\text{m}^2}; K_1 = 2 \text{ \& } K_2 = 3$$

6.

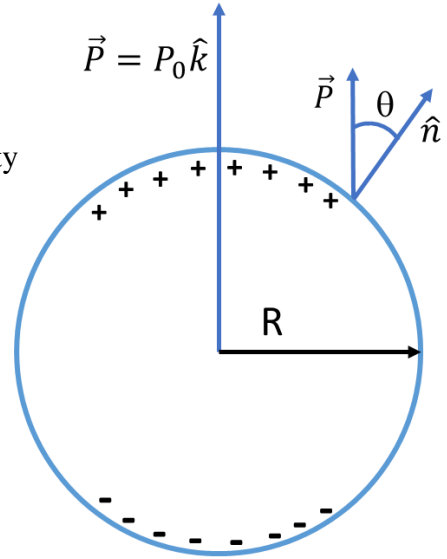
$\vec{P}$  is constant in magnitude. So, bound volume charge density

$$\sigma_b = -\vec{\nabla} \cdot \vec{P} = 0$$

Bound surface charge density,

$$\sigma_b = \vec{P} \cdot \hat{n}$$

$$\sigma_b = P_0 \cos \theta \text{ (in spherical polar coordinate)}$$



Total bound surface charge,

$$q_b = \int \sigma_b da \text{ (area element in spherical polar coordinate, } da = r^2 \sin \theta d\theta d\phi \text{)}$$

$$q_b = \int_{\phi=0}^{2\pi} \int_{\theta=0}^{\pi} P \cos \theta r^2 \sin \theta d\theta d\phi$$

$$q_b = r^2 P 2\pi \int_{\theta=0}^{\pi} \cos \theta \sin \theta d\theta d\phi = 0$$

7. For surface charge density, the polarization is given by:

$$\vec{P} = \vec{D} - \epsilon_0 \vec{E} \quad (1)$$

To obtain electric displacement inside sphere,  $\int \vec{D} \cdot d\vec{a} = q_{free}$

On the surface of radius r,

$$\vec{D} \cdot 4\pi r^2 = Q \hat{r}$$

Electric field inside the dielectric at radius r,

$$\vec{E} = \frac{Q}{4\pi \epsilon r^2} \hat{r}$$

Substituting the values of E and D in equation 1,

$$\vec{P} = \frac{Q}{4\pi r^2} \left[ 1 - \frac{\epsilon_o}{\epsilon} \right] \hat{r} = \frac{Q}{4\pi r^2} \left[ \frac{K-1}{K} \right] \hat{r}$$

At the surface of the sphere  $r = R$ ,

$$\vec{P} = \frac{Q}{4\pi R^2} \left[ \frac{K-1}{K} \right] \hat{r}$$

Therefore, bound surface charge density,

$$\sigma_b = \vec{P} \cdot \hat{n} = \frac{Q}{4\pi R^2} \left[ \frac{K-1}{K} \right]$$