

Lecture 24

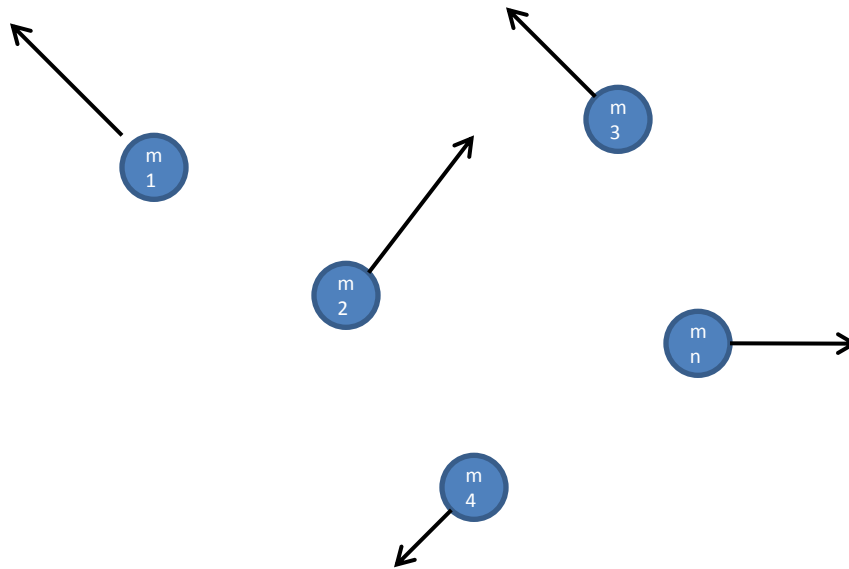
(Equations of Angular Motion)

Physics 160-01 Fall 2012

Douglas Fields

Rotational Kinetic Energy

- Let's examine a system of n particles:

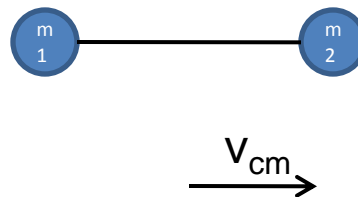


- The total kinetic energy of the system is:

$$\begin{aligned} KE &= \frac{1}{2}m_1v_1^2 + \frac{1}{2}m_2v_2^2 + \frac{1}{2}m_3v_3^2 + \cdots \frac{1}{2}m_nv_n^2 \\ &= \sum_{i=1}^n \frac{1}{2}m_iv_i^2 \end{aligned}$$

Rotational Kinetic Energy

- But, for a rigid body, there are constraints on what the velocities can be.
- Let's examine a system of 2 constrained particles (not rotating):



- The total kinetic energy of the system is:

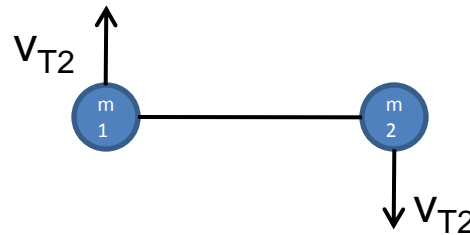
$$KE_{cm} = \frac{1}{2} m_1 v_1^2 + \frac{1}{2} m_2 v_2^2$$

$$\text{but, } v_1 = v_2 = v_{cm}$$

$$KE_{cm} = \frac{1}{2} (m_1 + m_2) v_{cm}^2 = \frac{1}{2} M v_{cm}^2$$

Rotational Kinetic Energy

- Now let's examine a system of 2 constrained particles (no cm motion, but rotating):

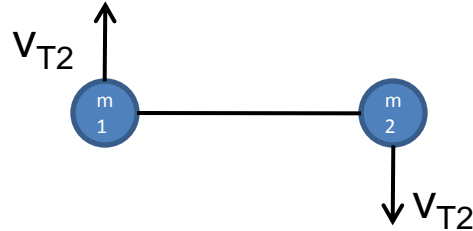


- Then the total kinetic energy of the system is:

$$\begin{aligned} KE_{rot} &= \frac{1}{2} m_1 v_{T1}^2 + \frac{1}{2} m_2 v_{T2}^2 \\ KE_{rot} &= \frac{1}{2} m_1 (\omega r_1)^2 + \frac{1}{2} m_2 (\omega r_2)^2 \\ &= \frac{1}{2} \omega^2 (m_1 r_1^2 + m_2 r_2^2) \end{aligned}$$

Rotational Kinetic Energy

- Notice that the term in parenthesis is only a function of the geometry of the system!



$$KE_{rot} = \frac{1}{2} \omega^2 (m_1 r_1^2 + m_2 r_2^2) = \frac{1}{2} I \omega^2,$$

where,

$$I = (m_1 r_1^2 + m_2 r_2^2)$$

- Or, in general,

$$I = \sum_i m_i r_i^2$$

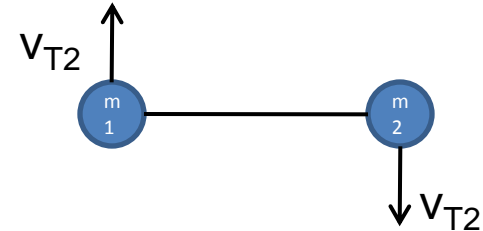
- This is called the moment of inertia.

Rotational Kinetic Energy

- But what is r_i ? Remember where it came from:

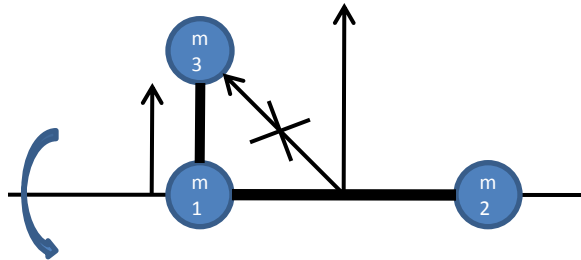
$$KE_{rot} = \frac{1}{2}m_1v_{T1}^2 + \frac{1}{2}m_2v_{T2}^2$$

$$KE_{rot} = \frac{1}{2}m_1(\omega r_1)^2 + \frac{1}{2}m_2(\omega r_2)^2$$
$$= \frac{1}{2}\omega^2(m_1r_1^2 + m_2r_2^2)$$

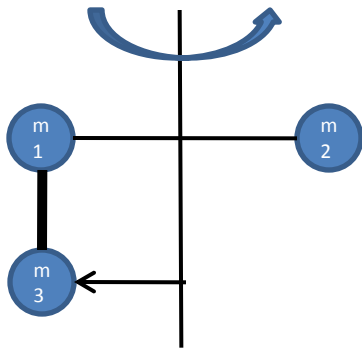


Rotational Kinetic Energy

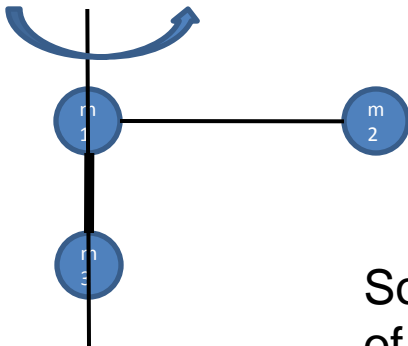
- So the r_i are the distance of the masses to the axis of rotation



$$I = \sum_i m_i r_i^2 = m_1 (0)^2 + m_2 (0)^2 + m_3 (d)^2 \\ = m_3 d^2$$



$$I = \sum_i m_i r_i^2 = m_1 (d)^2 + m_2 (d)^2 + m_3 (d)^2 \\ = (m_1 + m_2 + m_3) d^2$$



$$I = \sum_i m_i r_i^2 = m_1 (0)^2 + m_2 (2d)^2 + m_3 (0)^2 \\ = (m_2) 4d^2$$

So, the same object can have a different moment of inertia depending upon the axis of rotation!

CPS Question 23-1

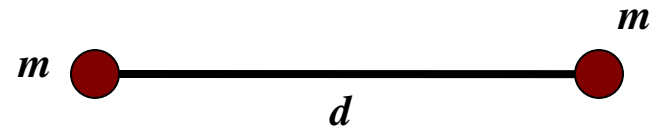
- Two point masses of mass m are attached to a long massless rod of length d . What is the moment of inertia that you would use in the calculation of the system's kinetic energy?

A) $2md^2$.

B) md^2 .

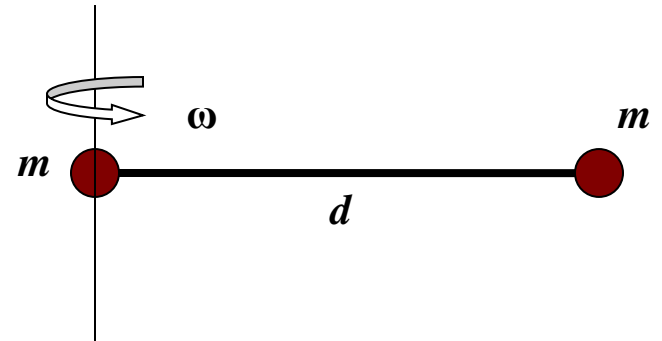
C) $\frac{1}{2} \omega^2 md^2$.

D) Not enough information to solve.



CPS Question 23-2

- Two point masses of mass m are attached to a long massless rod of length d . The system is rotating about one end. What is the moment of inertia that you would use in the calculation of the system's kinetic energy?



A) $2md^2$.

B) md^2 .

C) $\frac{1}{2} \omega^2 md^2$.

D) Not enough information to solve.

Continuous Distributions

- Generalization of the sum

$$I = \sum_i m_i r_i^2 \rightarrow \int r^2 dm$$

- But what is dm ???

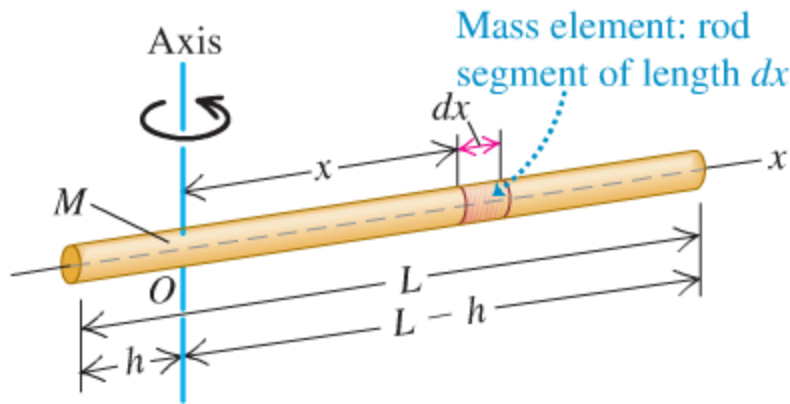
$$dm = \rho dV$$

- In Cartesian coordinates then we can write

$$I = \int r^2 \rho dV = \int r^2 \rho dx dy dz$$

Example

- Thin rod of uniform density (mass/unit length)



$$dm = \frac{M}{L} dx$$

$$I = \int r^2 dm = \int x^2 \frac{M}{L} dx$$

$$= \frac{M}{L} \int x^2 dx$$

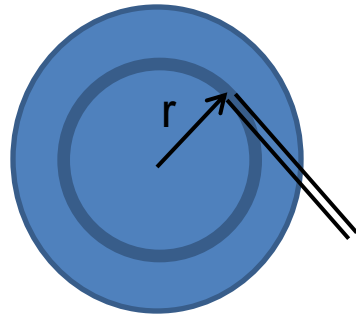
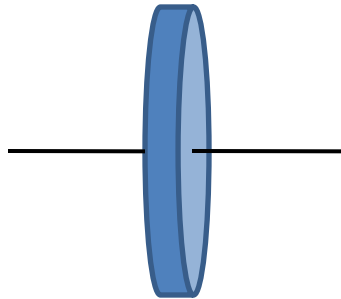
$$= \frac{M}{L} \int_{-h}^{L-h} x^2 dx = \frac{M}{L} \left[\frac{x^3}{3} \right]_{-h}^{L-h} = \frac{1}{3} M (L^2 - 3Lh + 3h^2)$$

$$\text{when } h = 0, \quad I = \frac{1}{3} ML^2$$

$$\text{when } h = \frac{L}{2}, \quad I = \frac{1}{12} ML^2$$

Example

- Thin disk of uniform density (mass/unit area)



$$dm = \frac{M}{A} dxdy = \frac{M}{A} 2\pi r dr$$

$$I = \int r^2 dm = \int r^2 \frac{M}{A} 2\pi r dr = 2\pi \frac{M}{A} \int r^3 dr$$

$$= 2\pi \frac{M}{A} \left[\frac{r^4}{4} \right]_0^R = 2\pi \frac{M}{A} \frac{R^4}{4}$$

but,

$$A = \pi R^2$$

$$I = 2\pi \frac{M}{\pi R^2} \frac{R^4}{4} = \frac{1}{2} MR^2$$

Example

- Sphere of uniform density (mass/unit volume)

$$dm = \rho dV = \rho \pi r^2 dx = \rho \pi (R^2 - x^2) dx$$

$$dI = \frac{1}{2} r^2 dm = \frac{1}{2} (R^2 - x^2) \rho \pi (R^2 - x^2) dx$$

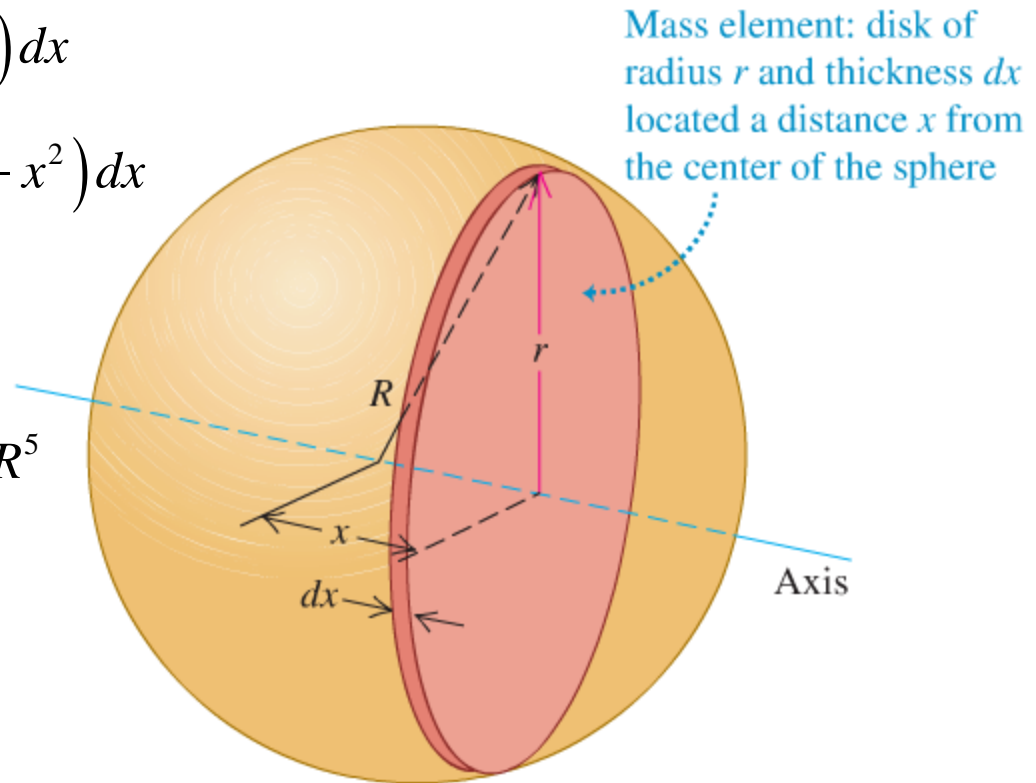
$$= \frac{\rho \pi}{2} (R^2 - x^2)^2 dx \Rightarrow$$

$$I = (2) \frac{\rho \pi}{2} \int_0^R (R^2 - x^2)^2 dx = \frac{8\pi\rho}{15} R^5$$

but,

$$M = \rho V = \rho \frac{4}{3} \pi R^3 \Rightarrow$$

$$I = \frac{2}{5} MR^2$$



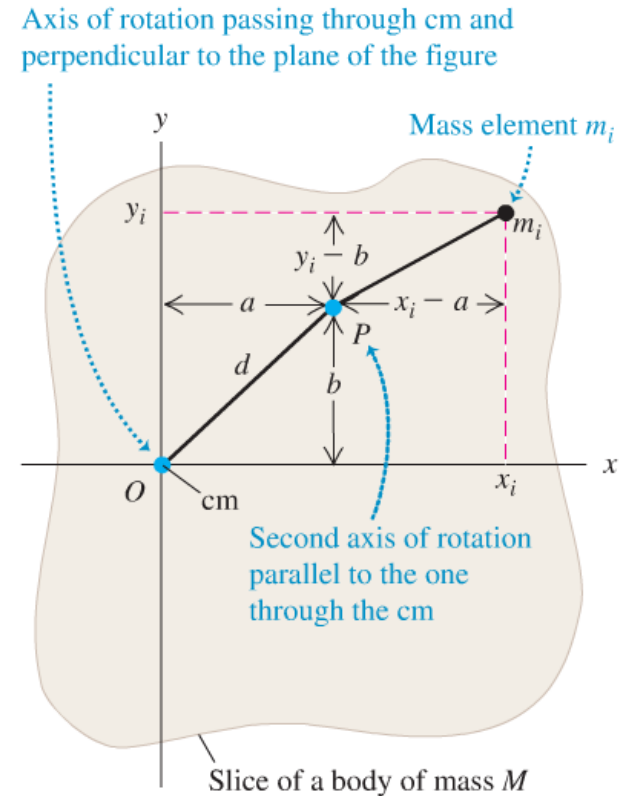
Parallel Axis Theorem

- If we know the moment of inertia about an axis that passes through the center of mass of an object, then the moment of inertia about any axis parallel to that a distance d away is given by:

$$I_P = I_{cm} + Md^2$$

$$I_{cm} = \sum_i m_i (x_i^2 + y_i^2)$$

$$\begin{aligned} I_P &= \sum_i m_i \left[(x_i - a)^2 + (y_i - b)^2 \right] \\ &= \sum_i m_i (x_i^2 + y_i^2) - 2a \sum_i m_i x_i - 2b \sum_i m_i y_i + (a^2 + b^2) \sum_i m_i \end{aligned}$$



Review

- Linear

- For const a :

$$x_f = x_i + v_i t + \frac{1}{2} a t^2$$

$$v_f = v_i + a t$$

- Kinetic energy:

$$KE = \frac{1}{2} M v^2$$

- Comes from:

$$W = \int \vec{F} \cdot d\vec{s}$$

- Rotational

- For const α :

$$\theta_f = \theta_i + \omega_i t + \frac{1}{2} \alpha t^2$$

$$\omega_f = \omega_i + \alpha t$$

- Kinetic energy:

$$KE = \frac{1}{2} I \omega^2$$

- Comes from:

???