

Lecture 28

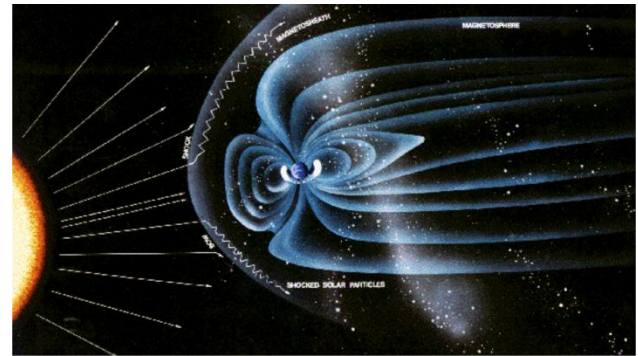
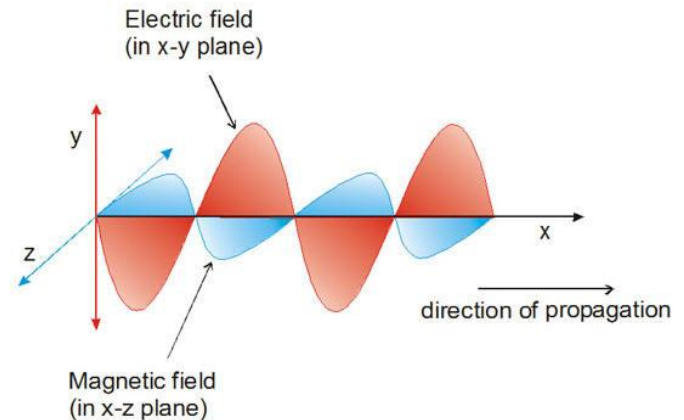
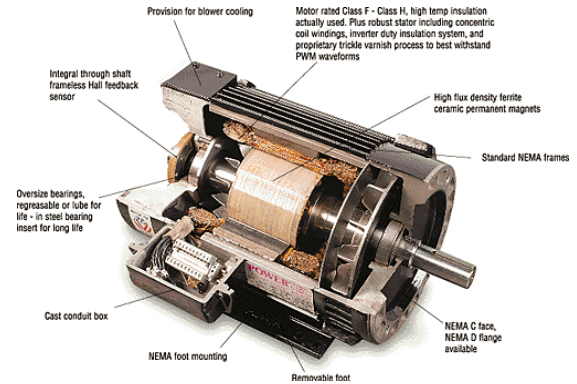
(Magnetic Force & Torque)

Physics 161-01 Spring 2012

Douglas Fields

Magnetic Forces

- Magnets attract and repel magnets.
- Magnets attract some metals.
- Magnets have some affect that can either turn electric current into work or turn work into electric current.
- What is the underlying, fundamental force that can describe all of these phenomena?



Magnetic Force

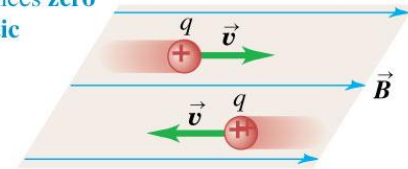
- Magnetic fields have an affect on *moving* charges.

$$\vec{F} = q\vec{v} \times \vec{B}$$

- Since it has been a while since we dealt with cross products, let's take a few minutes to review what this means.
- There is no force on a charge moving in the same direction as the magnetic field.
- The force on a moving charge in a magnetic field is perpendicular to both the field direction and the direction of motion.
- The direction of the force on a negative charge is opposite to that on a positive charge moving in the same direction.
- Since the force is perpendicular to the direction of motion, no work is done by the magnetic force.

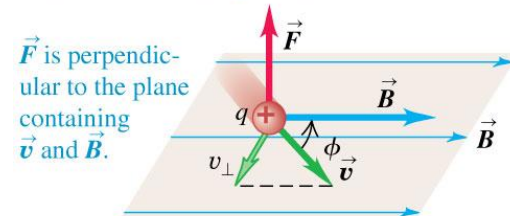
(a)

A charge moving **parallel** to a magnetic field experiences **zero** magnetic force.



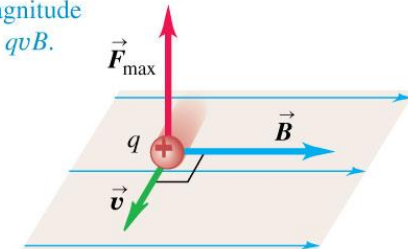
(b)

A charge moving at an angle ϕ to a magnetic field experiences a magnetic force with magnitude $F = |q|v_{\perp}B = |q|vB \sin \phi$.



(c)

A charge moving **perpendicular** to a magnetic field experiences a maximal magnetic force with magnitude $F_{\max} = qvB$.



Consequences and Applications

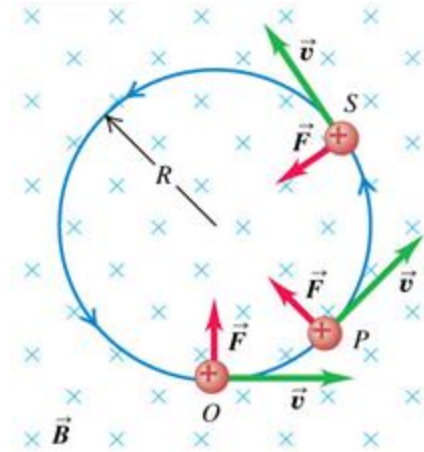
- A beam of charged particles will move in a circle at constant speed when they are sent into it perpendicular to a magnetic field.

$$\vec{F} = q\vec{v} \times \vec{B} = m\vec{a} \Rightarrow$$

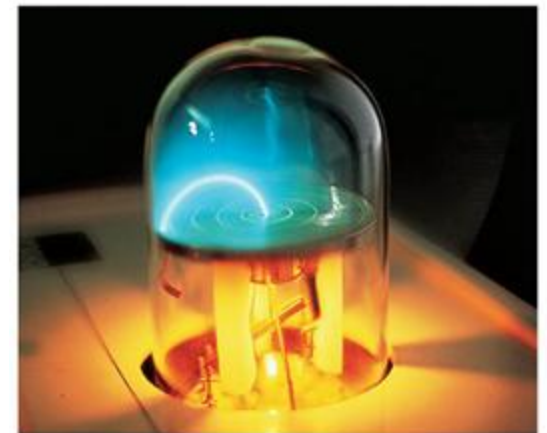
$$|\vec{a}| = \frac{|q\vec{v}\vec{B}|}{m} = \frac{v^2}{R} \Rightarrow$$

$$R = \frac{mv}{qB}$$

$$\left[\frac{\text{kg} \cdot \frac{\text{m}}{\text{s}}}{\text{C} \cdot \text{T}} \right] = \left[\frac{\text{kg} \cdot \frac{\text{m}}{\text{s}}}{\text{C} \cdot \frac{\text{N}}{\text{A} \cdot \text{m}}} \right] = \left[\frac{\text{kg} \cdot \frac{\text{m}}{\text{s}}}{\text{C} \cdot \frac{\text{kg} \cdot \text{m}}{\text{s}^2} \cdot \frac{\text{s}}{\text{C} \cdot \text{m}}} \right] = [\text{m}]$$



(b) An electron beam (seen as a white arc) curving in a magnetic field



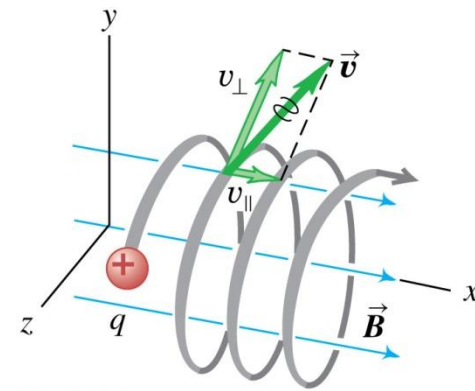
Consequences and Applications

- Magnetic Bottles

- If there is a component of the charged particle's velocity in the direction of the magnetic field, the path will be helical.
- By using a non-uniform field of the appropriate design, you can trap charged particles in a “magnetic bottle”.

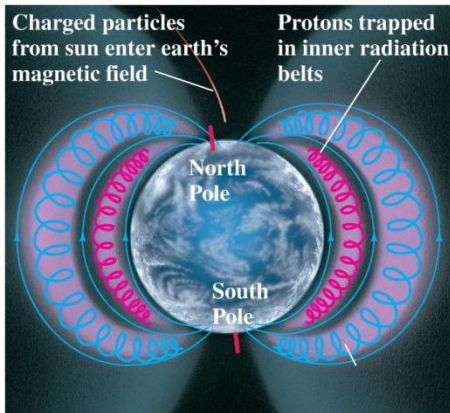
$$\vec{F} = q\vec{v} \times \vec{B}$$

This particle's motion has components both parallel (v_{\parallel}) and perpendicular (v_{\perp}) to the magnetic field, so it moves in a helical path.



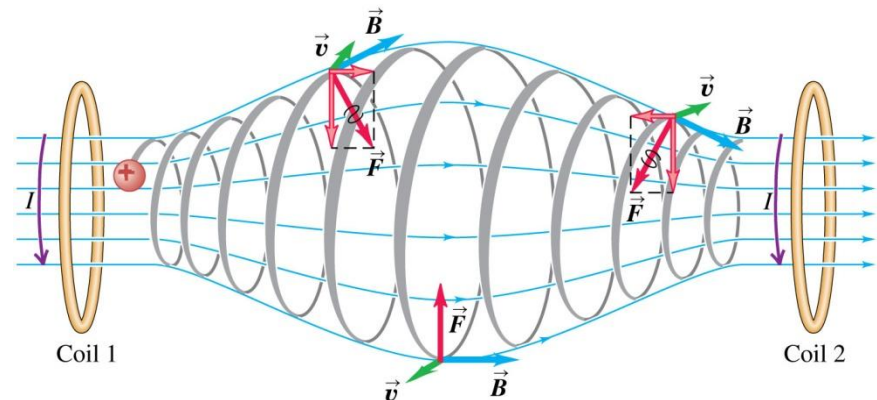
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(a)



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(b)



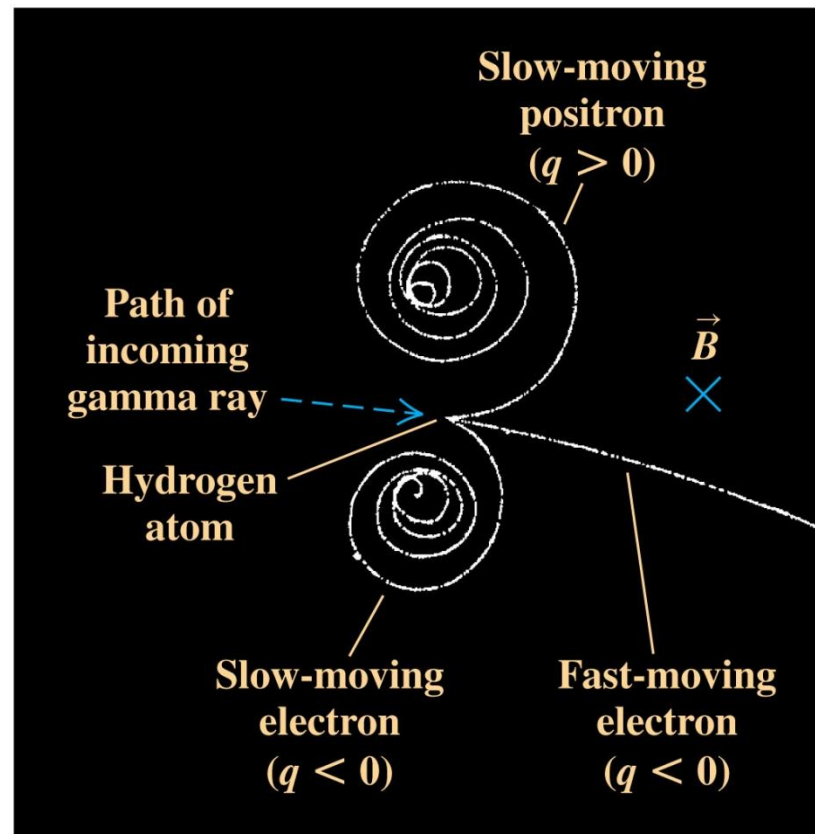
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Consequences and Applications

- Momentum Determination

- In my line of work, magnetic fields are used to give information on the mass, charge and velocity of particles that are created when other particles collide together.

$$\vec{F} = q\vec{v} \times \vec{B}$$



Consequences and Applications

- Velocity Filter

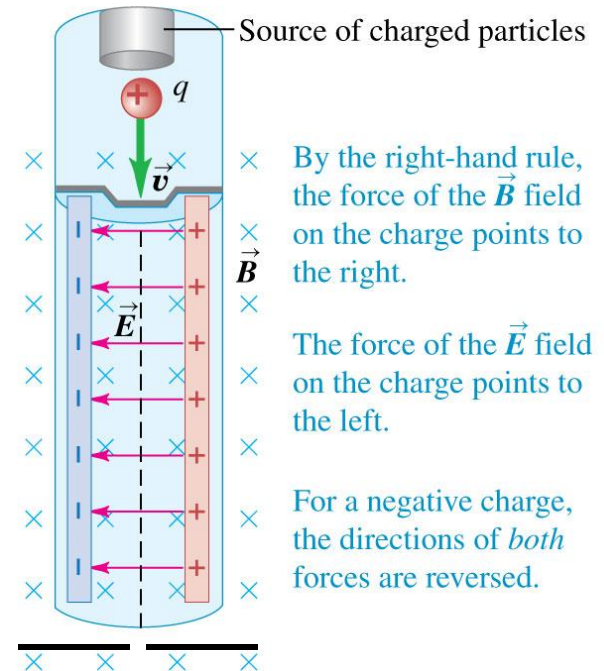
- We can use the magnetic force in conjunction with the electric force to filter out particles of a certain velocity (or just determine velocity).

$$\vec{F}_B = q\vec{v} \times \vec{B}$$

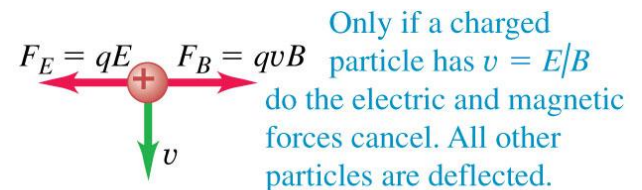
$$\vec{F}_E = q\vec{E}$$

- When the forces are equal, there is no deflection and $v = E/B$.
- By moving a slit that blocks particles except that go through the hole, you can pick out different velocities.

(a) Schematic diagram of velocity selector



(b) Free-body diagram for a positive particle



Consequences and Applications

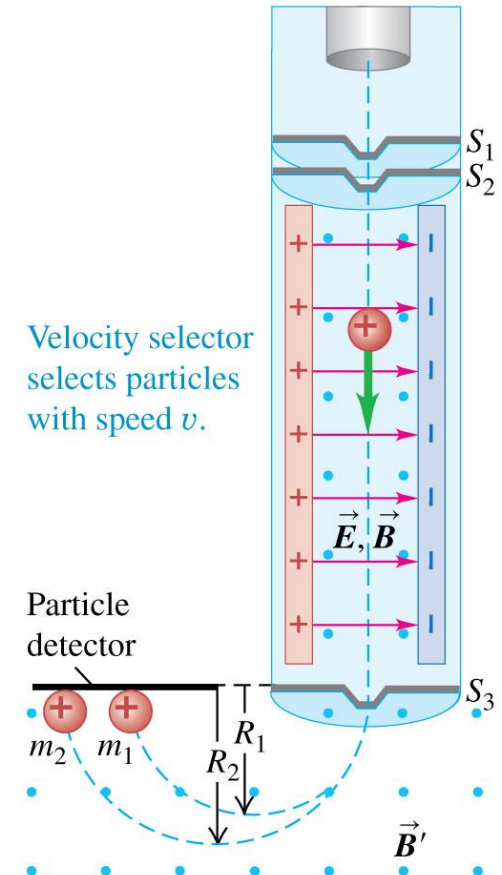
- Mass Spectrometer

- By putting another magnetic field outside of the velocity filter, and then detecting the radius of curvature, one can separate particles out by mass.
- This is how a mass spectrometer works.

$$\vec{F}_B = q\vec{v} \times \vec{B}$$

$$\vec{F}_E = q\vec{E}$$

$$R = \frac{mv}{qB'}$$



Consequences and Applications

- Thomson's e/m Apparatus

- In 1897, J.J. Thompson used a velocity filter to determine the ratio of the charge to mass of particles emitted from a cathode.
- He found that, regardless of cathode material, the ratio was always constant, and thus discovered the electron as a universal particle.

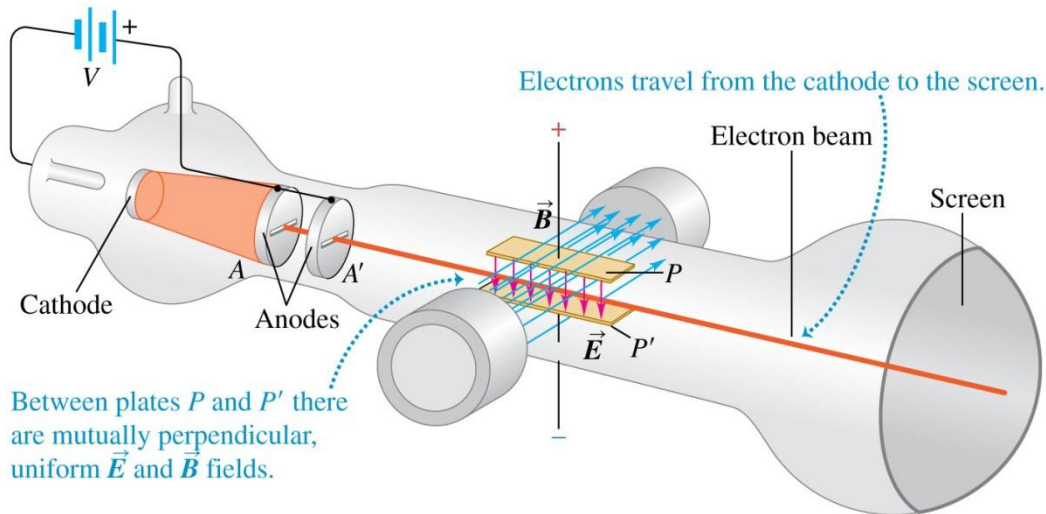
$$\frac{1}{2}mv^2 = eV \Rightarrow v = \sqrt{\frac{2eV}{m}}$$

$$\vec{F}_B = q\vec{v} \times \vec{B}$$

$$\vec{F}_E = q\vec{E}$$

$$v = \frac{E}{B} \Rightarrow$$

$$\frac{E}{B} = \sqrt{\frac{2eV}{m}} \Rightarrow \underline{\underline{\frac{e}{m} = \frac{E^2}{2B^2V}}}$$



CPS 28-1

A charged particle moves through a region of space that has both a uniform electric field and a uniform magnetic field. In order for the particle to move through this region at a constant velocity,

- A. the electric and magnetic fields must point in the same direction.
- B. the electric and magnetic fields must point in opposite directions.
- C. the electric and magnetic fields must point in perpendicular directions.
- D. The answer depends on the sign of the particle's electric charge.

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Magnetic Force on a Current Element

- In electronics, we rarely deal with “beams” of charged particles, but rather deal with current in a wire.

- But current is just moving charged particles.

$$\vec{F}_B = q\vec{v}_d \times \vec{B}$$

- The total force on the wire segment of length dl is just the sum of the forces on all the moving charges:

$$d\vec{F} = n(V)q\vec{v}_d \times \vec{B}$$

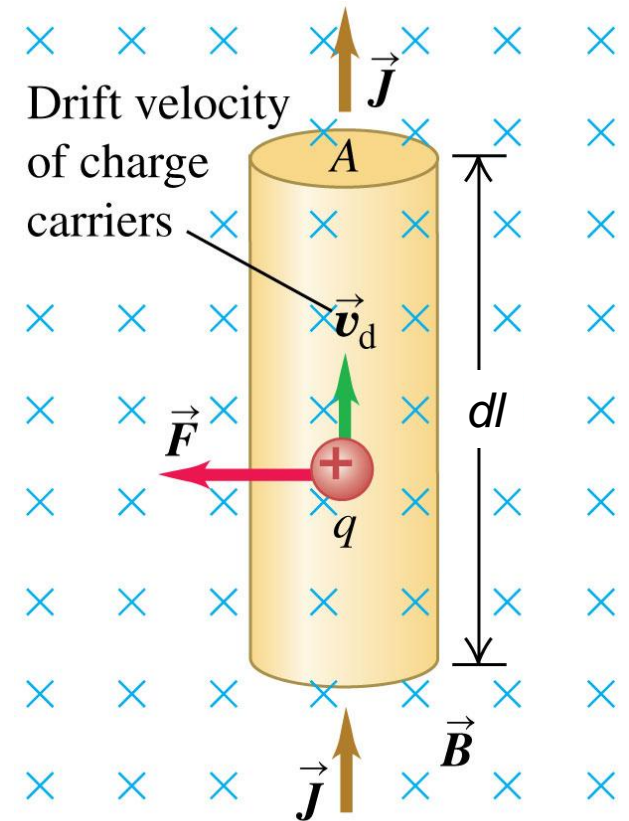
$$= n(Adl)q\vec{v}_d \times \vec{B}$$

- But $nq\vec{v}_d$ is just the current density, and the current density times the area is just the current:

$$d\vec{F} = n(Adl)q\vec{v}_d \times \vec{B}$$

$$d\vec{F} = I d\vec{l} \times \vec{B}$$

- Where we have designated the direction of $d\vec{l}$ to be in the same direction as the current.



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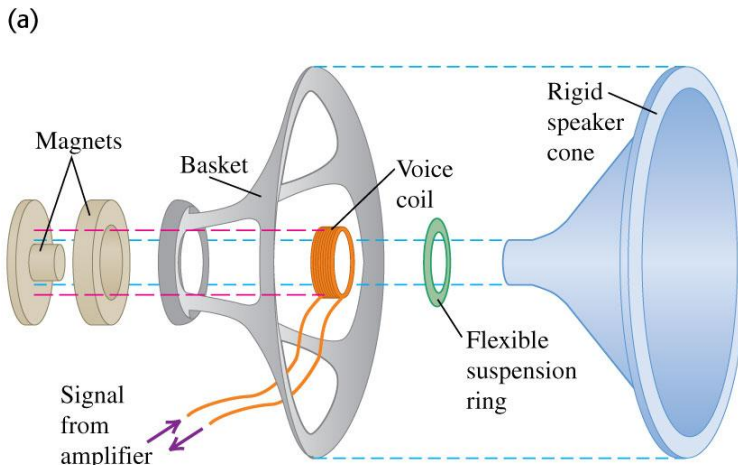
Magnetic Force on a Current

- One common use of this force is in speakers.

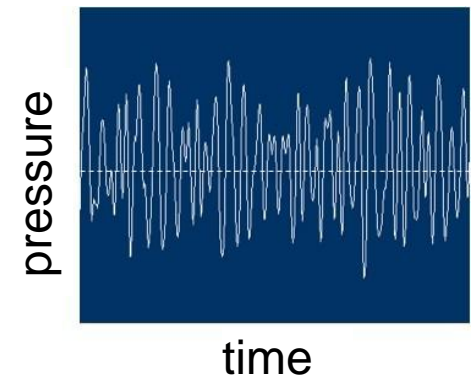
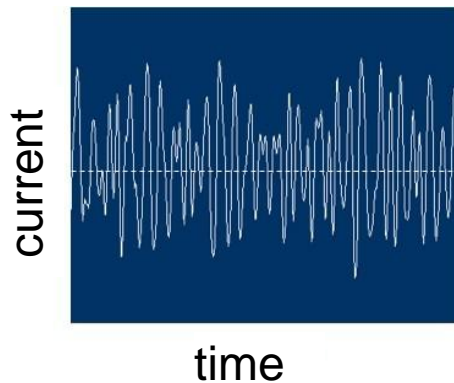
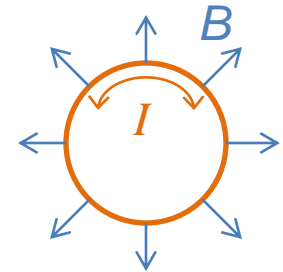
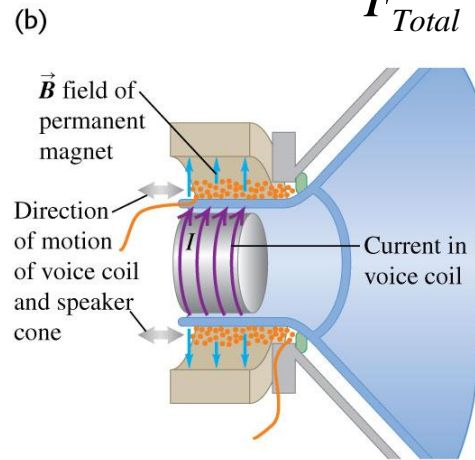
$$d\vec{F} = n(Adl)q\vec{v}_d \times \vec{B}$$

$$d\vec{F} = Id\vec{l} \times \vec{B} \Rightarrow$$

$$F_{Total} = IL_{Total}B$$



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Magnetic Force on a Current-Carrying Wire

- If the force isn't constant (either direction or magnitude), we must use our calculus tools:

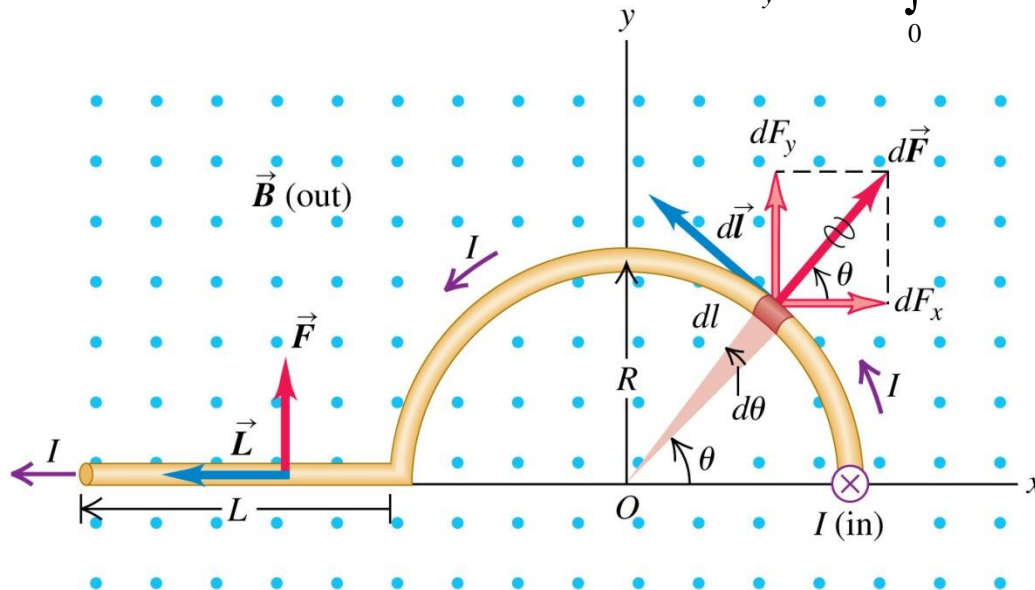
$$d\vec{F} = n(Adl)q\vec{v}_d \times \vec{B}$$

$$d\vec{F} = Id\vec{l} \times \vec{B} = I(Rd\theta)B\hat{r}$$

$$d\vec{F} = I(Rd\theta)B[\cos\theta\hat{i} + \sin\theta\hat{j}]$$

$$dF_y = IRB \sin\theta d\theta, \quad dF_x = IRB \cos\theta d\theta$$

$$F_y = IRB \int_0^\pi \sin\theta d\theta$$



Magnetic Force and Torque on a Current Loop

- Let's look at the Net force and net torque on a current loop:

$$d\vec{F} = I d\vec{l} \times \vec{B} \Rightarrow$$

$$F = IaB \quad (\text{top and bottom})$$

$$F = IbB \quad (\text{sides})$$

- But, the forces on opposite sides are opposing, so:

$$F_{\text{Net}} = 0$$

- Take the axis of rotation to be the y-axis, then the torque is:

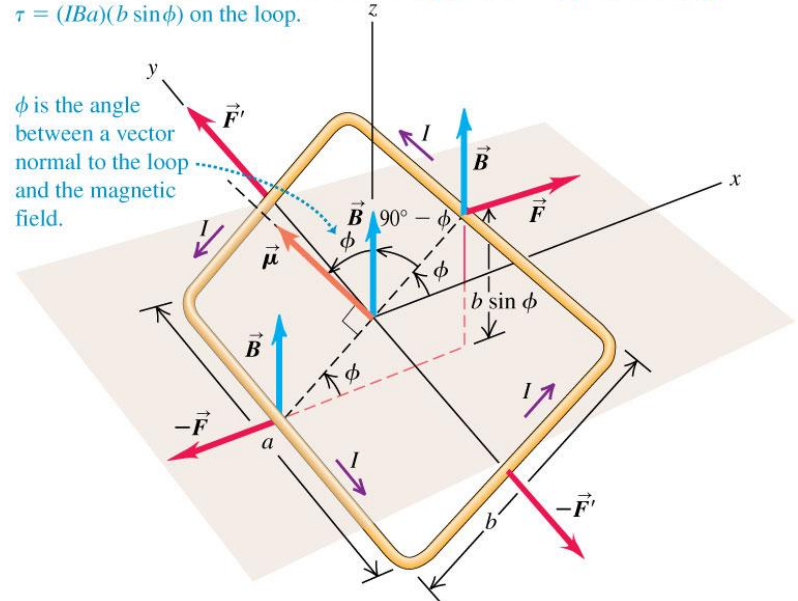
$$|\vec{\tau}| = IaB \left(\frac{b}{2} \sin \phi \right) + IaB \left(\frac{b}{2} \sin \phi \right)$$

$$|\vec{\tau}| = IabB \sin \phi = IAB \sin \phi$$

(a)

The two pairs of forces acting on the loop cancel, so no net force acts on the loop.

However, the forces on the a sides of the loop (\vec{F} and $-\vec{F}$) produce a torque $\tau = (IaB)(b \sin \phi)$ on the loop.



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Magnetic Torque on a Current Loop

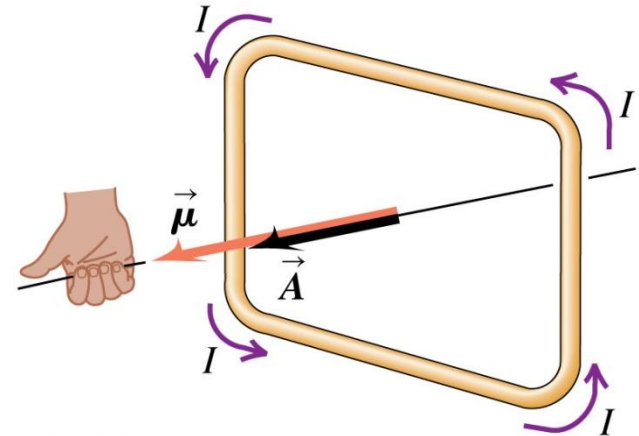
- We can rewrite this as:

$$|\vec{\tau}| = IAB \sin \phi = \mu B \sin \phi$$

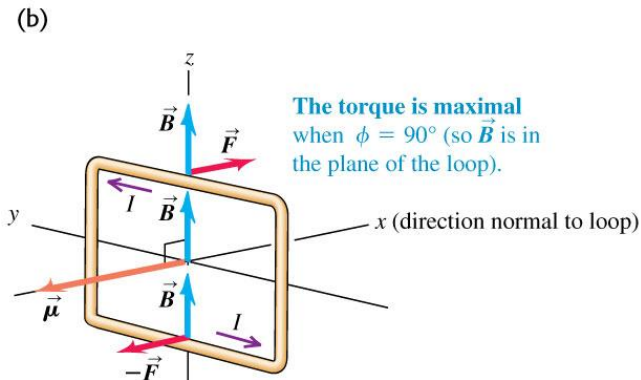
or

$$\vec{\tau} = \vec{\mu} \times \vec{B}$$

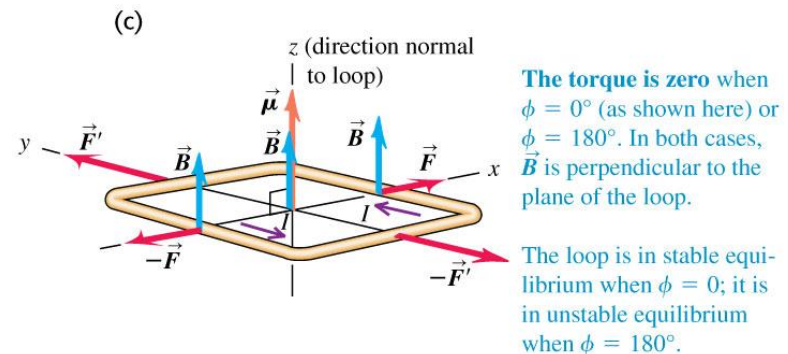
- Where the direction of the magnetic moment, μ , is given by the right-hand rule.



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CPS 28-2

A circular loop of wire carries a constant current. If the loop is placed in a region of uniform magnetic field, the *net magnetic force* on the loop is

- A. perpendicular to the plane of the loop, in a direction given by a right-hand rule.
- B. perpendicular to the plane of the loop, in a direction given by a left-hand rule.
- C. in the same plane as the loop.
- D. zero.
- E. The answer depends on the magnitude and direction of the current and on the magnitude and direction of the magnetic field.

CPS 28-2

A circular loop of wire carries a constant current. If the loop is placed in a region of uniform magnetic field, the *net magnetic force* on the loop is

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B. perpendicular to the plane of the loop, in a direction given by a left-hand rule.

C. in the same plane as the loop.

✓ D. zero.

E. The answer depends on the magnitude and direction of the current and on the magnitude and direction of the magnetic field.

CPS 28-3

A circular loop of wire carries a constant current. If the loop is placed in a region of uniform magnetic field, the *net magnetic torque* on the loop

- A. tends to orient the loop so that its plane is perpendicular to the direction of the magnetic field.
- B. tends to orient the loop so that its plane is edge-on to the direction of the magnetic field.
- C. tends to make the loop rotate around its axis.
- D. is zero.
- E. The answer depends on the magnitude and direction of the current and on the magnitude and direction of the magnetic field.

CPS 28-3

A circular loop of wire carries a constant current. If the loop is placed in a region of uniform magnetic field, the *net magnetic torque* on the loop



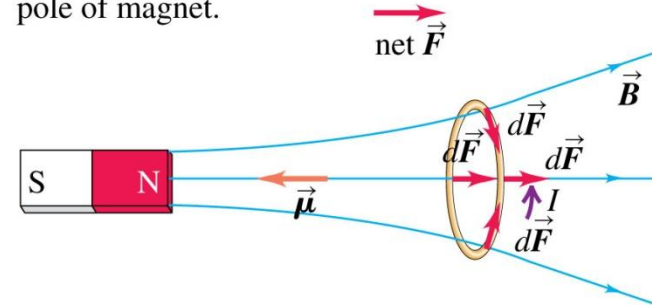
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- C. tends to make the loop rotate around its axis.
- D. is zero.
- E. The answer depends on the magnitude and direction of the current and on the magnitude and direction of the magnetic field.

Non-Uniform Magnetic Field Force on a Current Loop

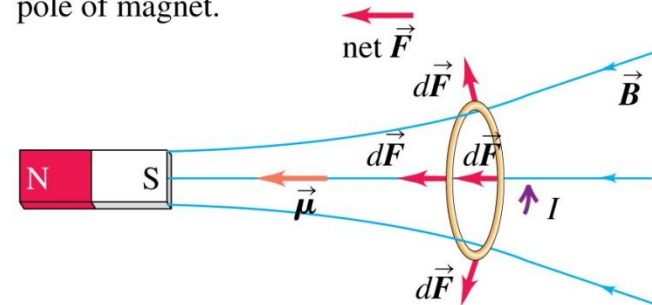
- Let's return to the question of how magnets attract each other and also non-magnetized metals.
- If the magnetic field is not uniform, then there CAN be a net force on a current loop (read: magnetic moment).
- The net force attracts if the magnetic moment aligns with the field direction, and repels if it is in the opposite direction.

$$d\vec{F} = I d\vec{l} \times \vec{B}$$

(a) Net force on this coil is away from north pole of magnet.



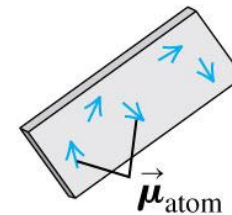
(b) Net force on same coil is toward south pole of magnet.



Magnetic Force on a Magnetic Moment

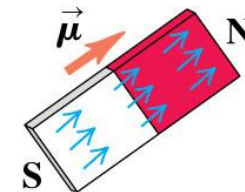
- Now, the atoms of iron (and other metals) have a magnetic moment because of the current loops of the electrons.
- In non-magnetized objects, they are randomly oriented.
- But in a magnet, they are (mostly) aligned with each other, giving a net total magnetic moment.

(a) Unmagnetized iron: magnetic moments are oriented randomly.

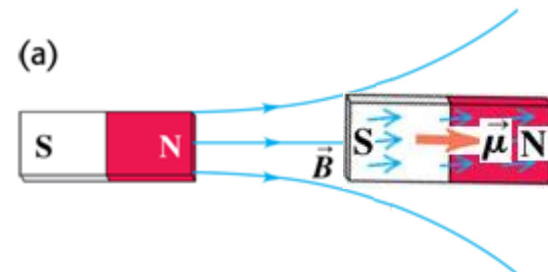


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(b) In a bar magnet, the magnetic moments are aligned.



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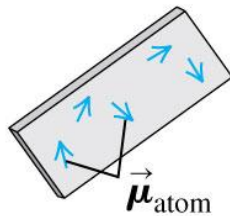


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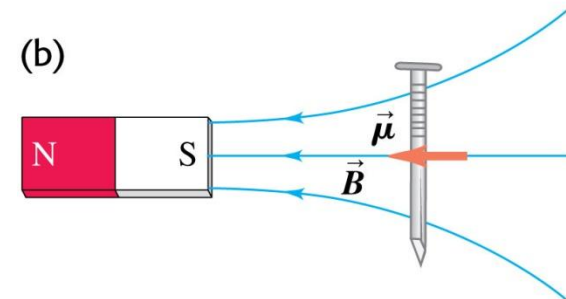
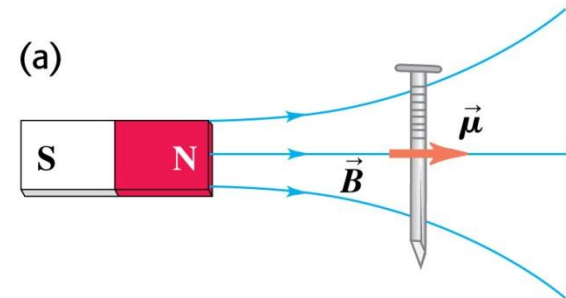
Magnetic Force on a Magnetic Moment

- For a non-magnetized object, like a nail, the field polarizes (rotates) the moments in the direction of the field.
- Then the object has a net magnetic moment and is attracted to the magnet.
- When the field is removed, the atoms randomize their magnetic moments again, and the object returns to be non-magnetized.
- UNLESS, you heat and then cool it, or shock it by hitting it with a hammer...

(a) Unmagnetized iron: magnetic moments are oriented randomly.



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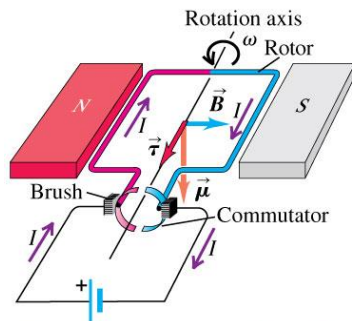


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Application – DC Motor

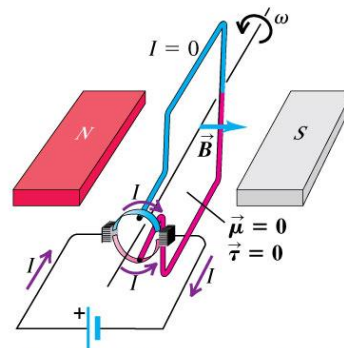
- We can now understand how a DC motor works.
- We put a current loop between the poles of a magnet.
- The torque on the loop causes rotation.
- We have to switch the direction of the current every time it rotates to keep it rotating in one direction.

(a) Brushes are aligned with commutator segments.



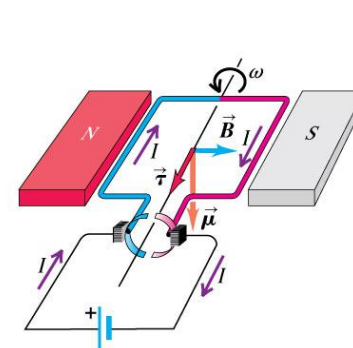
- Current flows into the red side of the rotor and out of the blue side.
- Therefore the magnetic torque causes the rotor to spin counterclockwise.

(b) Rotor has turned 90°.



- Each brush is in contact with both commutator segments, so the current bypasses the rotor altogether.
- No magnetic torque acts on the rotor.

(c) Rotor has turned 180°.



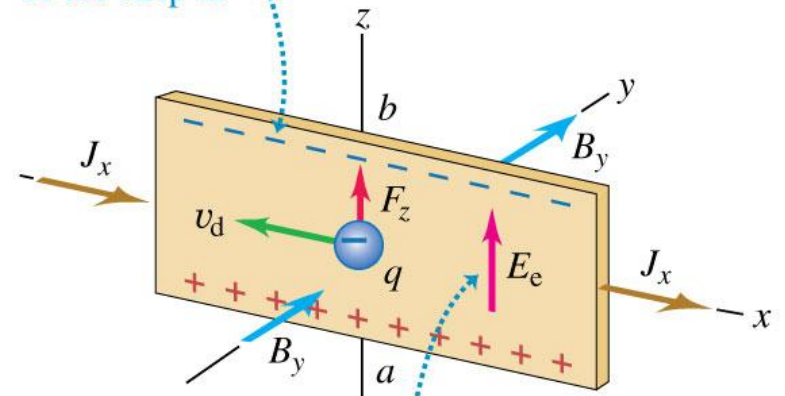
- The brushes are again aligned with commutator segments. This time the current flows into the blue side of the rotor and out of the red side.
- Therefore the magnetic torque again causes the rotor to spin counterclockwise.

Application – Hall Probe

- If we want to measure the magnetic field in a certain location, we can put a current through some (well understood) piece of conductor.
- The magnetic forces of the charges will tend to build up charges on opposite ends of the conductor.
- This charging will stop when the electric field force on the moving charges counters the magnetic force.
- By measuring the potential difference across the conductor in the z -direction, you can measure that charge, and thus, the field strength.

(a) Negative charge carriers (electrons)

The charge carriers are pushed toward the top of the strip ...



... so point a is at a higher potential than point b .

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