

# Dielectrics

Friday, December 4, 2020 8:33 PM

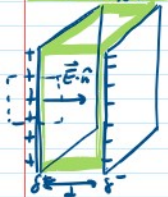
## Dielectrics

- Insulator  $\rightarrow$  large bandgap,  $E_g > 3.2 \text{ eV}$
- conductivity: Result of  $e^-$  & hole movement, depends on density of states  $f(E)$ , occupation of states in CB  $f(E)$  & VB  $1-f(E)$
- $\beta$  mobilities of carriers  $\mu_e, \mu_h$
- Increase charge storage of capacitors
- $\rightarrow$  Piezoelectrics

## Fermi Dirac Function

- $(1 + \exp[(E - E_F)/kT])^{-1}$
- $f(E)$  at bottom of CB is  $\approx 10^{-22} \rightarrow$  eg.  $E_g = 2.5 \text{ eV} \rightarrow f(E) \approx 10^{-22} = 1 - f(E)$

## Gauss law



$$\oint \vec{E} \cdot d\vec{A} = \frac{Q_{enc}}{\epsilon_0}$$

$\rightarrow$  dielectric const.  $\epsilon_r$

$$\epsilon = \epsilon_r \epsilon_0 \rightarrow \text{permittivity}$$

$$\oint \vec{E} \cdot d\vec{A} = EA = \frac{QA}{\epsilon_r \epsilon_0} \rightarrow E = \frac{Q}{\epsilon_r \epsilon_0 A}$$

$$V = -\int E dl = -\frac{Qd}{\epsilon_r \epsilon_0 A} = \frac{Q}{A \epsilon_r \epsilon_0} \rightarrow Q = \sigma A = V(\epsilon_r \epsilon_0 A/d)$$

- const. V,  $d \rightarrow \uparrow E$

## Capacitors:

$$C = \frac{Q}{V} = \frac{\epsilon_r \epsilon_0 A}{d}$$

$$E = \frac{1}{2} CV^2$$

$\rightarrow \uparrow \epsilon_r, \uparrow Q$  for Voltage,  $\uparrow E$  stored

- Dielectric const:  $\epsilon_r = (1 + \chi) \rightarrow \chi = \text{dielectric susceptibility} \rightarrow \text{deg. of polarization of a dielectric in response to electric field}$

## Atomic level

- (+) charged nucleus ( $q^+$ ) surrounded by  $e^-$  cloud
- When  $\vec{E}$  field app.  $\rightarrow$ , nucleus shifts  $\rightarrow$
- $e^-$  cloud shifts  $\leftarrow$
- @ equilibrium,  $\vec{E}$  balanced by attraction b/w nucleus &  $e^-$  cloud
- $\rightarrow \vec{E}$  felt by  $q^+$  towards center:  $\vec{E}_e = -\frac{q}{4\pi\epsilon_0 a^3} \hat{r}$ ,  $q^+ \leftarrow \uparrow \vec{E}$
- $\rightarrow$  @ eqm:  $\vec{E}_{tot} = -\vec{E}_e$

- Molecule has induced dipole moment:  $\vec{p} = qd$
- dipole moment  $\propto$  to external  $\vec{E}$  field:  $\vec{p} = 4\pi\epsilon_0 \alpha \vec{E}_{ext} = \alpha \vec{E}_{ext} \rightarrow \alpha = \text{polarizability}$
- $\vec{E}$  field experienced by dipole  $\vec{E}_{local} \neq \vec{E}_{ext}$ :  $\rightarrow$  When  $\vec{E}$  field applied, each atom has induced dipole  $\vec{p} = \alpha \vec{E}$

## Polarization

- For  $n$  atoms/V, induced dipole moment (Polarization):  $P = n\vec{p} = n\alpha \vec{E}$
- Induced dipole moment/V
- @ low fields:  $\vec{P} = \chi \epsilon_0 \vec{E} \rightarrow \chi$  in free space  $= 0$



$$\text{in Polarization: } P = \frac{\sum qd}{V} = \frac{12 qd}{12 (2a)^3} = \frac{qd}{2a^3}$$

## Surface Charge Density:

- Surface charge density:  $\sigma = Nm qd = N\alpha \vec{E} \rightarrow Nm = \text{# molecules/V}$
- Plane of polarized dielectric atoms @ surface:  $V = Ad$
- dipole moment  $\vec{p} = PV = PA d\epsilon = Q_{total} d\epsilon \hat{n}$
- $\therefore$  Surface charge density from polarization:  $\sigma_{bound} = \frac{Q_{total}}{A} = |\vec{P}| \rightarrow$  eq. to polarization!



- Field inside  $\downarrow$ : (+) (-) cancels
- @ const V:  $\sigma_{free} = \frac{Q_{free}}{A} = \epsilon_0 \frac{V}{d} = \epsilon_0 E$   $\sigma_{total} = \sigma_{free} + \sigma_{bound} \rightarrow$  due to polarization
- $\rightarrow$  due to charge on electrodes

$$\oint \vec{E} \cdot d\vec{A} = \frac{Q_{enc}}{\epsilon_0}, \text{ charge on electrode: } Q = CV$$

$$EA = \frac{QA}{\epsilon_0} \quad E = \frac{-V}{\epsilon_0 d} \therefore \text{Field inside: } E = \frac{-\sigma_{free} + \sigma_{bound}}{\epsilon_0} \rightarrow \sigma_{bound} = P = P/V_{volume} \rightarrow \sigma_{free} = |\vec{D}| = |\epsilon_r \vec{E}| = (\epsilon_r + 1)\epsilon_0 \vec{E}$$

- If field applied, pull (+)  $\rightarrow$  (-) plate & pull (-)  $\rightarrow$  (+) plate  $\therefore$  changing @ material edges
- $\rightarrow$  Reduces field inside
- @ const V,  $\uparrow \sigma_{free}$  to compensate for  $\sigma_{bound}$

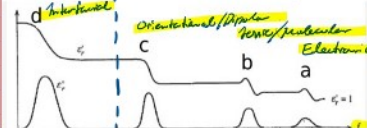
- Electric displacement: free charge density on capacitor plates

$$|\vec{D}| = |\chi \epsilon_0 \vec{E}| + |\epsilon_0 \vec{E}| = |\vec{P}| + |\epsilon_0 \vec{E}| = \sigma_{bound} + \sigma_{free} = \sigma_{total}$$

- Dipole moments can result from polar molecules, ionic polarization, or electronic polarization

$\rightarrow \epsilon_r \downarrow$  as  $f \uparrow$   $\therefore$  Rotations & Nuclei can't keep up w/ freq,  $e^-$  motion too large

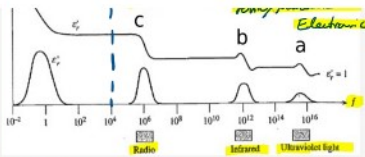
## Polarization @ diff frequencies



$$\text{Dielectric const } \epsilon_r: \epsilon_r = 1 + \chi_{total} = 1 + \chi_e + \chi_i + \chi_o$$

- @ low freq:  $\uparrow \epsilon_r$ , all 3 types can keep up w/ changing  $\vec{E}$  field
- @ mid freq:  $\downarrow \epsilon_r$ ,  $\chi_o = 0$   $\therefore$  can't respond fast enough  $\therefore \epsilon_r = 1 + \chi_e + \chi_i$
- @ high freq:  $\downarrow \epsilon_r$ ,  $\chi_i = \chi_o = 0$   $\therefore$  can't respond fast enough  $\therefore \epsilon_r = 1 + \chi_e$





- @ low freq:  $\uparrow \epsilon_r$ , all 3 types can keep up w/ changing  $\vec{E}$  field
- @ mid freq:  $\downarrow \epsilon_r$ ,  $X_0 = 0$   $\therefore$  can't respond fast enough  $\therefore \epsilon_r = 1 + X_0 + X_0$
- @ high freq:  $\downarrow \epsilon_r$ ,  $X_0 = X_0 = 0$   $\therefore$  can't respond fast enough  $\therefore \epsilon_r = 1 + X_0$
- @ optical freq: use refractive index  $n$  instead of  $\epsilon_r \rightarrow n^2 \approx \epsilon_r$

- **Electronic:** All materials consist of ions surrounded by  $e^-$  clouds  $\rightarrow$  rapid response to field changes eg. ( $H_2O$ )

- **Electronic polarization in solid:**

$\rightarrow$  Polarization of 1 atom:  $\vec{p} = q \cdot d = 4\pi\epsilon_0 a^3 \epsilon_{ext} = \alpha \epsilon_{ext}$

$\rightarrow$  Polarization/Volume:  $P = 4\pi\epsilon_0 a^3 \epsilon_{ext} \cdot N \rightarrow N = \# \text{ atoms/V}$  eg. cubic lattice spacing  $a$ , 1 atom =  $(2a)^3$

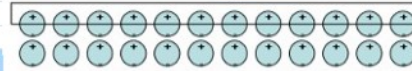


$Q_{bound} = PA = \frac{1}{2} \epsilon_0 \epsilon_{ext} A = \frac{1}{2} \epsilon_0 \epsilon_{ext} A \therefore N = \frac{1}{(2a)^3} \therefore P = \frac{1}{2} \epsilon_0 \epsilon_{ext}$

$EA = Q_{bound} + Q_{free} = -PA + \epsilon_0 E A = -X \epsilon_0 A E + \epsilon_0 E A$

$E(\epsilon_0 + X \epsilon_0) = \epsilon_0 E \therefore \epsilon_r = 1 + X = 1 + \frac{1}{2} = 2.5$

- Index of Refraction:  $v = c/n$ , For materials w/  $\downarrow X$ :  $n = \sqrt{\epsilon_r} = \sqrt{2.5}$



- **Ionic/Molecular:** Lattice ions charged, material has permanent dipoles (polar)

- w/  $\vec{E}$  field, lattice ions are stretched  $\rightarrow$  increase of dipole moment of lattice

$\therefore$  NaCl, transparent, slower than electronic eg. ( $NaCl$  &  $KCl$ )

- @  $\downarrow f$ , force due to  $\vec{E}_{app}$  balanced by electrostatic restoring force

- @  $\uparrow f$ , resonant & attached (like mass/spring)

- Beyond resonant  $f$ , force consumed w/ acc. mass,  $\downarrow \epsilon_r$  &  $\uparrow f$ :  $\downarrow P$



- **Orientation/Dipole:**  $\vec{E}$  field exerts torque on molecule, aligning them to  $\vec{E}$  field dir.

- Involves stretching & rotating  $\therefore$  responds slower than ions

Torque:  $T = E_0 (\frac{1}{2} \sin \Theta + \frac{1}{2} \sin \Theta) = E_0 \sin \Theta \rightarrow$  Apply  $\vec{E}$  field, Torque  $\vec{p}$  align w/ Field

avg deg. of alignment  $\bar{p}$  avg  $\vec{p}$ :  $\bar{p} = \frac{\vec{p}^2 E}{kT}$



- 3D dipole aligns counter to field, add. transl. PE:  $U = - \int T d\Theta = - \int \vec{p} \sin \Theta d\Theta = - \vec{E} \cdot \vec{p}$

- **Debye Equation:** describes freq. response, justify drop in  $\epsilon_r$  due to orientation polarization

$\epsilon' = \frac{(\epsilon_s - \epsilon_\infty) \epsilon_\infty}{1 + \omega^2 \tau^2} \quad \epsilon'' = \frac{(\epsilon_s - \epsilon_\infty) \omega \tau}{1 + \omega^2 \tau^2} \rightarrow \tau = \text{time const. at which molecule responds to } \vec{E} \text{ field}$

$\epsilon_s = \text{dielectric const. below transition freq. } (\omega \ll 1/\tau)$

$\epsilon_\infty = \text{rel. dielectric const. @ freq. above } 1/\tau \text{ but } \ll \text{ next transition}$

- friction due to collisions  $\rightarrow$  can't keep up with  $\vec{E}$  field oscillations

-  $\uparrow f$ , dipole rotates faster, drag force from collisions prevent dipoles from aligning

- freq. dependence of polarization:  $m = F - Kx - b\dot{x} \rightarrow m = i\omega m$  or  $e^-$  cloud mass,  $\omega = \text{accl.}$ ,  $K = \text{restoring } F$

Transfer from  $\vec{p}$  &  $\vec{E}$  (susceptibility):  $\chi = \frac{1}{m^2 \omega^2 + b^2 + K^2} \quad x = \text{displacement}$ ,  $b = \text{damping coeff}$ ,  $v = \text{velocity}$

@  $\downarrow f$ : reaches resonance

@  $\uparrow f$ : drops @ rate of  $1/\omega^2$

$F = q\vec{E}$

## Dielectric Loss

- Energy loss from collisions (heating material), Joule heating (small conduction currents)

- Given  $\epsilon_r = \epsilon' + j\epsilon''$ , Im. of  $\epsilon_r$ :  $\epsilon'' = \frac{\sigma}{\omega}$

$C = Q/V \quad C = \epsilon A/d \rightarrow$  Find Capacitive:  $\frac{1}{C} = \frac{1}{\epsilon A/d} + \dots \rightarrow \frac{1}{C_{eff}} = \frac{1}{\epsilon A/d} + \frac{1}{\sigma A/d} \rightarrow \frac{1}{C_{eff}} = \frac{1}{\epsilon A/d} + \frac{1}{\sigma A/d}$

$C_{eff} = \epsilon_0 \epsilon' \frac{A}{d} \rightarrow$  Find  $\epsilon_{eff}$

Laplace:  $V = I/R = \frac{I \epsilon}{\sigma C} \rightarrow I = I_1 + I_2 \therefore V = \frac{I}{\sigma C} \rightarrow V(\epsilon C + \frac{1}{R}) = I, V(C + \frac{1}{R\sigma}) = Q$

$\therefore C + \frac{1}{R\sigma} = C_{eff} \rightarrow R = \frac{1}{\sigma C} \therefore C_{eff} = \frac{1}{\sigma} [\epsilon + \frac{1}{j\omega}]$

$\epsilon_{eff} = \epsilon' - j\epsilon'' \epsilon_0 \rightarrow \epsilon'' \epsilon_0 = \sigma/\omega$

-  $C$  of to  $Re(\epsilon_r)$ , stored  $E$ :  $\frac{1}{2} C V^2 = \frac{1}{2} \epsilon A/d \cdot V^2$

$E$  density stored in  $C$ :  $\epsilon_r \epsilon_0 \vec{E}^2 = \frac{1}{2} \epsilon (V/d)^2 = \frac{1}{2} \epsilon E^2$

- To maximize stored  $E/V$ ,  $\uparrow \epsilon_r$  &  $\vec{E}$  (limited by dielectric breakdown) &  $\downarrow \epsilon''$

## Dielectric Breakdown

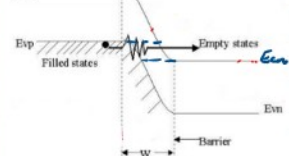
- Sudden increase in current when  $V$  beyond critical level

### Zener Breakdown

- Extreme reverse bias in diode, tunneling of  $e^-$  carriers from VB  $\rightarrow$  CB across depletion region

-  $V_A \uparrow \rightarrow E$  of VB on one side overlap w/ CB on other side

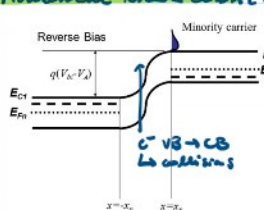
P-side n-side



- when biased  $\uparrow$ , VB top on p-side is level w/ CB bottom on n-side  $\rightarrow$  tunneling

to occur in heavily doped diodes (thin depletion region)

### Avalanche Breakdown (same as intrinsic)



- PE gained by minority  $e^- \rightarrow$  to n-side of highly reverse biased diode

gets more energy than  $E_g$  band gap

$\therefore$  Transfer Energy to VB  $e^-$  through collisions

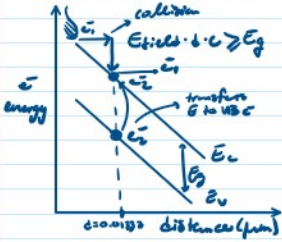
-  $\uparrow$  carriers through depletion region

-  $\uparrow I \rightarrow$  breakdown

- Avalanche occurs if:  $q(V_{bi} - V_A) \gg E_g$

### Intrinsic Breakdown

- Same as avalanche breakdown, but across dielectric
- $e^-$  in CB of dielectric highly accelerated  $\rightarrow$  collision w/ VB  $\therefore$  promoted to CB



- $e^-$  accelerated by  $E_{field}$ , lose energy in collision, then promotes  $e^-$  into CB
- KE gained b/w collisions:  $KE = \frac{1}{2}mv^2 = E_c \lambda > E_g$
- #  $e^-$  promoted  $\propto$  #  $e^-$  in CB
- exponential rise:  $10^8$
- $\uparrow$  depletion width  $\rightarrow$   $\downarrow$  probability

### Thermal Breakdown:

- Heating  $\rightarrow$  easier promotion of  $e^-$  into CB, pos. feedback loop
- $e^-$  in CB  $\rightarrow$  Joule heating  
 $\rightarrow$  May be due to dielectric loss
- $\uparrow$  Heat,  $\uparrow$  vibration  $\rightarrow$  dielectric failure

### Discharge Breakdown

- Gas in porous materials is ionized, damages material; accelerates breakdown esp. (ceramics & mica)
- $\rightarrow$  From applied field