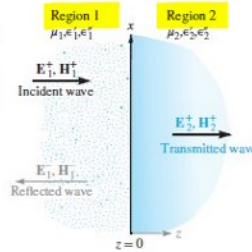


Chapter 12 Wave Reflection and Dispersion

Monday, February 22, 2021 12:18 PM

12.1 Reflection of Uniform Plane waves @ normal incidence

- uniform plane wave incident on boundary b/w 2 diff. regions of material
- **normal incidence**: wave propagation \perp boundary
 - \rightarrow reflected wave
 - \rightarrow transmitted wave through region 2



12.1.1 Reflected & Transmitted Waves @ boundary

- **incident wave** Region 1 (lin. pol.):

- **field wave \hat{z} :**

$$E_{x1}(z, t) = E_{x10} e^{-\alpha_1 z} \cos(\omega t - \beta_1 z)$$

phasor: $E_{x1}(z) = E_{x10} e^{-j\beta_1 z}$

- **magnetic field \hat{y} :**

$$H_{y1}(z) = \frac{1}{\eta_1} E_{x10} e^{-j\beta_1 z}$$

- **Transmitted wave** Region 2:

$$E_{x2}(z) = E_{x20} e^{-j\beta_2 z}$$

$$H_{y2}(z) = \frac{1}{\eta_2} E_{x20} e^{-j\beta_2 z}$$

- **@ Boundary $z=0$:**

- \vec{E} fields (\hat{z}) in regions: $E_{x10} = E_{x20}$

- \vec{H} fields (\hat{y}) in regions: $H_{y01} = H_{y02}$

- $\eta_1 = \eta_2 \rightarrow \frac{E_{x10}}{\eta_1} = \frac{E_{x20}}{\eta_2}$

- Satisfy Boundary conditions w/ reflected wave

- **Reflected wave $\hat{z}=0$:**

$$E_{x1}(z) = E_{x10} e^{j\beta_1 z}$$

$$H_{y1}(z) = -\frac{E_{x10}}{\eta_1} e^{j\beta_1 z}$$

\rightarrow field in $-\hat{z}$: $E_{x1} = -\eta_1 H_{y1}$

\rightarrow Poynt vector: $\langle \vec{S} \rangle = \vec{E}_1 \times \vec{H}_1 \rightarrow -\hat{z}$

12.1.2 Reflection & Transmission Coefficients

- **Boundary Conditions** for \vec{E} & \vec{H} tangential continuity:

$$E_{x10} = E_{x10} = E_{x20}$$

$$H_{y10} = H_{y10} = H_{y20}$$

$$\rightarrow \frac{E_{x10}}{\eta_1} - \frac{E_{x10}}{\eta_1} = \frac{E_{x20}}{\eta_2}$$

$$\therefore E_{x10} = E_{x20} \frac{\eta_2 - \eta_1}{\eta_2 + \eta_1}$$

- **Reflection Coefficient Γ :**

$$\Gamma = \frac{E_{x10}}{E_{x20}} = \frac{\eta_2 - \eta_1}{\eta_2 + \eta_1} = |\Gamma| e^{j\theta_\Gamma}$$

\rightarrow if η_1 or η_2 complex, Γ complex
 \therefore include reflective phase shift θ

- **Transmission coefficient τ :**

$$\tau = \frac{E_{x20}}{E_{x10}} = \frac{2\eta_2}{\eta_1 + \eta_2} = 1 + \Gamma = |\tau| e^{j\theta_\tau}$$

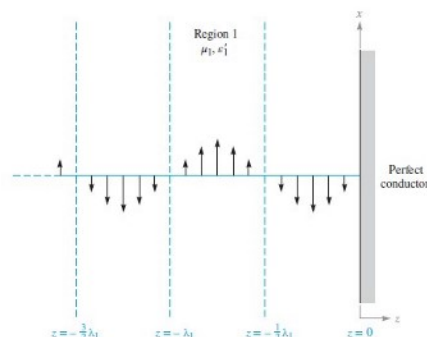


Figure 12.2 The instantaneous values of the total field E_1 are shown at $\omega t = \pi/2$. $E_1 = 0$ for all time at multiples of one half-wavelength from the conducting surface.

12.1.3 Total Reflection: Standing Wave Ratio

- **Region 1** \rightarrow perfect dielectric

Region 2 \rightarrow perfect conductor

$$\rightarrow \epsilon_2 = \frac{\sigma_2}{\omega} \rightarrow \eta_2 = \sqrt{\frac{j\omega\mu_2}{\sigma_2 + j\omega\epsilon_2}} = 0 \quad \therefore \sigma_2 \rightarrow \infty$$

$$\therefore E_{x20} = 0, \Gamma = -1, E_{x10} = -E_{x10}$$

\rightarrow skin depth $\delta = 0$

\rightarrow All incident energy is reflected by perfect conductor

\rightarrow Reflected field shifted in phase by 180° relative to incident field

$$\therefore \vec{E} \text{ field in R1: } E_{x1}(z) = E_{x10} e^{-j\beta_1 z} + E_{x10} e^{j\beta_1 z}$$

- All incident energy is reflected by perfect conductor
- Reflected field shifted in phase by 180° relative to incident field
- ∴ **E field in R1**: $E_{x1} = E_{x10} + E_{x1-} = E_{x10} e^{-i\beta_1 z} - E_{x10} e^{i\beta_1 z}$
 $j\mathbf{k}_1 = 0 + j\beta_1 \rightarrow -j2 \sin(\beta_1 z) E_{x10}$ } $\text{Re} \{ E_{x10} e^{j\omega t} \}$
- ∴ **Real Inst. Form**: $E_{x1}(z,t) = 2 E_{x10} \sin(\beta_1 z) \sin(\omega t)$
- **Incident wave**: $E_{x1}(z,t) = E_{x10} \cos(\omega t - \beta_1 z)$
- Null locations occur @ $z = n\lambda/2$
- ∴ $E_{x1} = 0$ @ boundary $z=0$ & every $\lambda/2$ from bound to $z < 0$
- ∴ **Magnetic Field**: $H_{y1} = \frac{E_{x10}}{\eta_1} (e^{-i\beta_1 z} + e^{i\beta_1 z})$
- ∴ **Real Inst Form**: $H_{y1}(z,t) = 2 \frac{E_{x10}}{\eta_1} \cos(\beta_1 z) \cos(\omega t)$

12.1.4 Partial Reflection & Power Reflectivity

- **Region 1 & 2** → Perfect dielectrics
- η_1 & $\eta_2 \rightarrow (+)$, $\sigma_1 = \sigma_2 = 0$
- **Reflection coeff**: $\Gamma = \frac{\eta_2 - \eta_1}{\eta_2 + \eta_1} \rightarrow E_{x10} = \Gamma E_{x10}$
- **Magnetic Field Intensities**: $H_{y10} = \frac{E_{x10}}{\eta_1}$
 $H_{y1-} = -\frac{E_{x10}}{\eta_1}$
- **Power density**:
 → **Region 1**: Incident → $\langle S_{i1} \rangle = \frac{1}{2} \text{Re} \{ \vec{E}_1 \times \vec{H}_1^* \} = \frac{1}{2} E_{x10} H_{y10}$
 Reflected → $\langle S_{r1} \rangle = -\frac{1}{2} E_{x10} H_{y10}$
 → **Region 2**: $E_{x20} = \eta_2 E_{x10}$
 $H_{y20} = \frac{E_{x20}}{\eta_2}$
- Avg power density transmitted to R2:
 $\langle S_{t2} \rangle = \frac{1}{2} E_{x20} H_{y20}$
 $\langle S_{it} \rangle = \langle S_{ir} \rangle + \langle S_{t2} \rangle$
- **General Relation b/w reflected & incident power**
 $\langle S_{ir} \rangle = |\Gamma|^2 \langle S_{i1} \rangle$
 $\langle S_{t2} \rangle = (1 - |\Gamma|^2) \langle S_{i1} \rangle$

12.2 Standing Wave Ratio

- $|\Gamma| < 1 \rightarrow$ Energy transmitted into R2 & reflected
- **Region 1** → Perfect dielectric ($\sigma_1 = 0$)
- **Region 2** → Any
- **Region 1**:
 E-field phasor: $E_{x1T} = E_{x1+} + E_{x1-} = E_{x10} e^{-i\beta_1 z} + \Gamma E_{x10} e^{i\beta_1 z}$
 → $\Gamma = \frac{\eta_2 - \eta_1}{\eta_2 + \eta_1} = |\Gamma| e^{i\phi} \rightarrow$ possible reflection coeff.
 ∴ include phase ϕ
 → lossless: $\eta_1 \rightarrow \text{Real } (+)$, $\eta_2 \rightarrow \text{complex}$
 → perfect conductor: $\eta_2 = 0$, $\phi = \pi$
 → If $\eta_2 \rightarrow \text{Real } < \eta_1$, $\phi = \pi$
 → If $\eta_1 \rightarrow \text{Real } > \eta_2$, $\phi = 0$
- ∴ **Total Field**: $E_{x1T} = (e^{-i\beta_1 z} + |\Gamma| e^{i(\beta_1 z + \phi)}) E_{x10}$
- **Maximum**: $|E_{x1T}|_{\max} = (1 + |\Gamma|) E_{x10}$ where $z_{\max} = -\frac{1}{2\beta_1} (\phi + 2m\pi)$
- **Minimum**: $|E_{x1T}|_{\min} = (1 - |\Gamma|) E_{x10}$ where $z_{\min} = -\frac{1}{2\beta_1} (\phi + (2m+1)\pi)$
- **Total field in Region 1**: $E_{x1T}(z,t) = \frac{(1 - |\Gamma|) E_{x10} \cos(\omega t - \beta_1 z)}{\text{traveling wave}} + \frac{2|\Gamma| E_{x10} \cos(\beta_1 z + \phi/2) \cos(\omega t + \phi/2)}{\text{standing wave}}$
- **Standing Wave Ratio**
 → Ratio of max: min Amplitudes
 $S = \frac{|E_{x1T}|_{\max}}{|E_{x1T}|_{\min}} = \frac{1 + |\Gamma|}{1 - |\Gamma|}$

12.3 Wave Reflection from multiple interfaces

12.3.1 Two Interface Problem

- Uniform plane wave \hat{z} incident
- 2nd region of impedance → 2nd interface @ $z=0$
 → 1st interface @ $z=-l$
- @ 1st interface: Reflect & transmit
- @ 2nd interface: transmit to R3 & reflected to 1st





- 1st interface $z=0$
- @ 1st interface: Reflect & transmit
 - @ 2nd interface: transmit into R3 & reflected to 1st
 - Again partially reflected then combine w/ transmitted Energy from R1, etc.
 - Steady state eventually reached

12.3.2 Wave Impedance

- All Regions → lossless media
- In Region 2 (x-polarized):

$$E_{x2} = E_{x20} e^{-j\beta_2 z} + E_{x20} e^{j\beta_2 z} \rightarrow \beta_2 = \frac{\omega \sqrt{\epsilon_2}}{c}$$

$$H_{y2} = H_{y20} e^{-j\beta_2 z} + H_{y20} e^{j\beta_2 z}$$
- Reflection Coefficient

$$\Gamma_{12} = \frac{\eta_3 - \eta_2}{\eta_3 + \eta_2} \rightarrow E_{x20} = \Gamma_{12} E_{x10} + E_{x20}$$

$$H_{y20} = \frac{E_{x20}}{\eta_2}$$

$$H_{y20} = -\frac{E_{x20}}{\eta_2} = -\Gamma_{12} \frac{E_{x10}}{\eta_2}$$
- Wave impedance: Ratio of total E-Field

$$\eta_0(z) = \frac{E_{x2}}{H_{y2}} = \frac{E_{x20} e^{-j\beta_2 z} + E_{x20} e^{j\beta_2 z}}{H_{y20} e^{-j\beta_2 z} + H_{y20} e^{j\beta_2 z}}$$

$$= \eta_2 \frac{e^{-j\beta_2 z} + \Gamma_{12} e^{j\beta_2 z}}{e^{-j\beta_2 z} - \Gamma_{12} e^{j\beta_2 z}} = \eta_2 \frac{\eta_3 \cos(\beta_2 z) - j\eta_2 \sin(\beta_2 z)}{\eta_2 \cos(\beta_2 z) - j\eta_3 \sin(\beta_2 z)}$$

$$= \eta_2 \frac{\eta_3 \cos(\beta_2 z) - j\eta_2 \sin(\beta_2 z)}{\eta_2 \cos(\beta_2 z) - j\eta_3 \sin(\beta_2 z)}$$
- ∴ $\frac{E_{x10}}{E_{x10}^+} = \Gamma = \frac{\eta_{in} - \eta_1}{\eta_{in} + \eta_1}$
- Input impedance: $\eta_{in} = \eta_2 \frac{\eta_3 \cos(\beta_2 l) + j\eta_2 \sin(\beta_2 l)}{\eta_2 \cos(\beta_2 l) + j\eta_3 \sin(\beta_2 l)}$

12.3.3 Half & Quarter Waves (Total Transmission/No reflection)

- Half wave Matching
 - $\Gamma = 0$ or $\eta_{in} = \eta_1 \rightarrow$ input impedance matched to incident medium
 - $\eta_2 = \eta_3 \rightarrow l = n \frac{\lambda_2}{2} \rightarrow$ int multiple of $\frac{1}{2}$ wavelengths
 - Matched input impedance: $\eta_3 = \eta_1 \rightarrow \eta_{in} = \eta_3$
 - choosing RL thickness → half-wave matching
- Refractive index: $n = \sqrt{\epsilon_r}$
- Lossless Media:

$$\beta = k = \omega \sqrt{\mu_0 \epsilon_0} \sqrt{\epsilon_r} = \frac{\omega n}{c}$$

$$\eta = \frac{1}{\sqrt{\epsilon_r}} \sqrt{\frac{\mu_0}{\epsilon_0}} = \frac{\eta_0}{n}$$

$$v_p = \frac{c}{n}$$

$$\lambda = \frac{v_p}{f} = \frac{\lambda_0}{n}$$
- Quarter wave Matching

$$l = (2m-1) \frac{\lambda_2}{4} \rightarrow \eta_{in} = \frac{\eta_1^2}{\eta_3}$$
 - choose η_2 for matching b/w η_1 & η_3
 - total transmission: $\eta_{in} = \eta_1$
 - ∴ $\eta_2 = \sqrt{\eta_1 \eta_3}$
 - Anti-reflective coating for optical devices



12.3.4 Multilayer Problems: Impedance Transformation

- Transform R4 impedance → input impedance @ boundary R2 & R3

$$\eta_{in,b} = \eta_3 \frac{\eta_{in,b} \cos \beta_3 l_b + j\eta_2 \sin \beta_3 l_b}{\eta_2 \cos \beta_3 l_b + j\eta_{in,b} \sin \beta_3 l_b}$$
- Reflected power fraction $|\Gamma|^2$

$$\Gamma = \frac{\eta_{in,a} - \eta_1}{\eta_{in,a} + \eta_1}$$
- Fraction of power transmitted into R4: $1 - |\Gamma|^2$

12.4 Plane Wave Propagation in General Directions

- Lossless medium, prop. const $\beta = k = \omega \sqrt{\mu \epsilon}$

- phase $\phi(x, z) = \vec{k} \cdot \vec{r}$

↳ k = prop. const. vector

- Wave (phasor): $\vec{E}_s = \vec{E}_0 e^{-j\vec{k} \cdot \vec{r}}$

↳ position vector: $\vec{r} = x\hat{a}_x + z\hat{a}_z$

$$\lambda = \frac{2\pi}{k} = \frac{2\pi}{\sqrt{k_x^2 + k_z^2}}$$

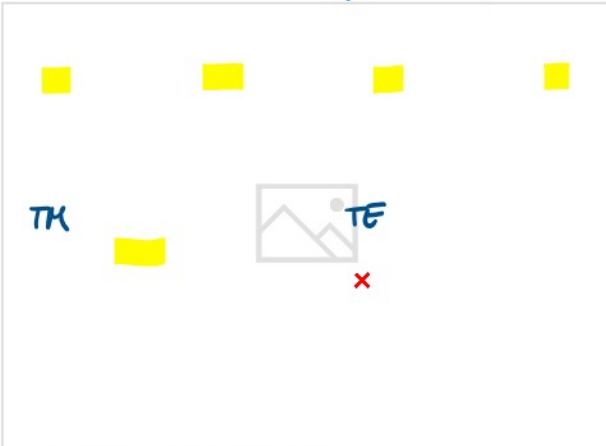
$$v_p = \frac{\omega}{k} = \frac{\omega}{\sqrt{k_x^2 + k_z^2}}$$

$$f = \frac{v_p}{\lambda}$$



12.5 Plane Wave Reflection @ Oblique Incidence Angles

- Relation b/w incident, reflected, transmitted angles



- Incident wave d.r. & phase $\rightarrow \vec{k}_1^+$

- angle of incidence θ_1

- Reflected wave \vec{k}_1^- propagate away @ θ_1'

- Transmitted wave \vec{k}_2 @ θ_2

- Incident & reflected angles are equal ($\theta_1 = \theta_1'$)

- p-polarized (TM)

- \vec{E} polarized in page plane

- $\vec{H} \perp$ page pointing outward, parallel/transverse to interface \rightarrow transverse magnetic

- \vec{E} lying in the plane of incidence \rightarrow parallel polarization

- s-polarized (TE)

- Fields rotated by $90^\circ \uparrow$

- \vec{H} in incidence plane

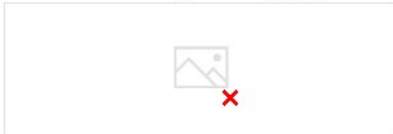
- $\vec{E} \perp$ plane \rightarrow perpendicular polarization

↳ parallel to interface \rightarrow transverse electric

- wave vector magnitudes:

$$k_1 = \frac{\omega \sqrt{\epsilon_1}}{c} = \frac{n_1 \omega}{c}$$

$$k_2 = \frac{\omega \sqrt{\epsilon_2}}{c} = \frac{n_2 \omega}{c}$$



- @ $x=0$: $E_{z\bar{s}1} = E_{z\bar{s}2} = E_{z\bar{s}2}$

- $\theta_1 = \theta_1' \rightarrow k_1 \sin \theta_1 = k_2 \sin \theta_2$

- Snell's Law of Refraction

$$n_1 \sin \theta_1 = n_2 \sin \theta_2$$

- Relation b/w Amplitudes

- @ $x=0$: $H_{1o} = H_{1r} = H_{2o}$

- effective impedances:

$$\eta_{1p} = \eta_1 \cos \theta_1$$

$$\eta_{2p} = \eta_2 \cos \theta_2$$

$$\Gamma_p = \frac{E_{1r}}{E_{1o}} = \frac{\eta_{2p} - \eta_{1p}}{\eta_{2p} + \eta_{1p}}$$

$$\tau_p = \frac{E_{2o}}{E_{1o}} = \frac{2\eta_{2p}}{\eta_{1p} + \eta_{2p}} \left(\frac{\cos \theta_1}{\cos \theta_2} \right)$$

p-polarization

$$\begin{aligned}
 r_p &= \frac{E_{1p}}{E_{i0}} = \frac{n_{2p} - n_{1p}}{n_{2p} + n_{1p}} \\
 r_p &= \frac{E_{2p}}{E_{i0}} = \frac{2n_{1p}}{n_{1p} + n_{2p}} \left(\frac{\cos \theta_1}{\cos \theta_2} \right) \\
 n_{1s} &= n_1 \sec \theta_1 \\
 n_{2s} &= n_2 \sec \theta_2 \\
 r_s &= \frac{E_{1s}}{E_{i0}} = \frac{n_{2s} - n_{1s}}{n_{2s} + n_{1s}} \\
 r_s &= \frac{E_{2s}}{E_{i0}} = \frac{2n_{2s}}{n_{2s} + n_{1s}}
 \end{aligned}$$

p-polarization

s-polarization

12.6 Total Reflection & Total transmission of obliquely incident waves

- **Total Reflection** → total Power Reflection

- $|r|^2 = r r^* = 1 \rightarrow r = r_p \text{ or } r_s$

- $\sin \theta_1 \geq \frac{n_2}{n_1}$

- **Critical Angle**: $\sin \theta_c = \frac{n_2}{n_1}$

↳ For total Reflection: $\theta_1 > \theta_c$

- **Optical waveguides**: 3 layers of glass, mid layer $n >$ outer two

- **Total transmission**

- $r = 0$

- **s-polarization**:

↳ $r_s = 0$: $n_{2s} = n_{1s}$ or $n_2 \sec \theta_2 = n_1 \sec \theta_1$

↳ $r_p = 0 \rightarrow$ Snell's law solve

↳ $\sin \theta_1 = \sin \theta_2 = \frac{n_2}{\sqrt{n_1^2 + n_2^2}} \rightarrow n_1 = \frac{n_0}{n_1}, n_2 = \frac{n_0}{n_2}$

- **Brewster Angle** (polarization Angle)

↳ if incident @ $\theta_1 = \theta_B$, p-component transmitted & s-polarized reflected

↳ Polaroid Sunglasses → block transmission of horizontally polarized light passing vertically polarized light

