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# A simulation study of surface temperature influence on localized atmospheric flows and dispersion using shallow water equations

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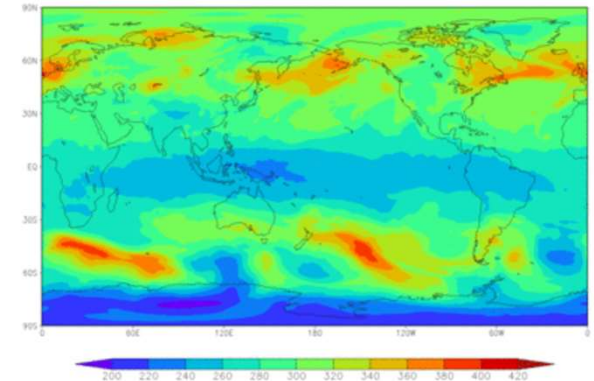
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RECAP..

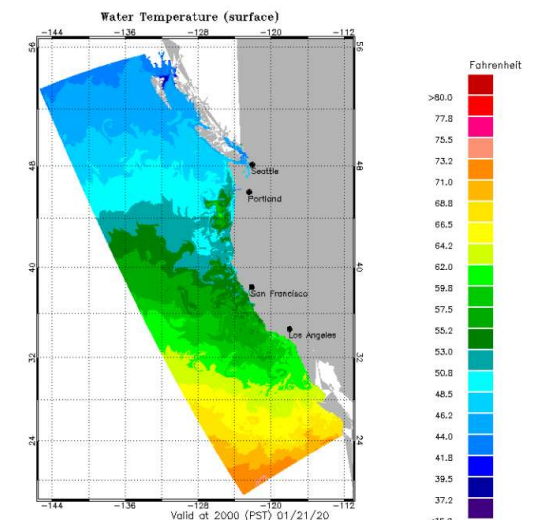
## RECAP

# Motivation

- Atmospheric models are used in
  - Simulation of earth's atmosphere
  - Weather forecast
  - Climate modelling (*Vallis, 2017*)
  - Global Models: GFS, ICON; Regional Models: ETA, COAMPS etc.
- Conventional Models:
  - Governing Equations: Navier-Stokes and Energy Balance Equations
  - Complex and Computationally Expensive.
- Shallow Water Model: Simple to solve and can be used in problems with smaller domains.
  - Can be applicable in Industrial Planning (Identifying Threat Zones and Accident Propagation)
  - Localized Atmospheric Modelling



*Fig: GFS Temperature contours (atmospheric)*



*Fig: COAMPS Sea Surface Temperature Contours*

## RECAP

# Introduction

- SWEs are hyperbolic/parabolic PDEs derived from the **Euler's Equation:**

$$\rho \left( \frac{\partial \vec{v}}{\partial t} + \vec{v} \cdot \nabla \vec{v} \right) = -\nabla P$$

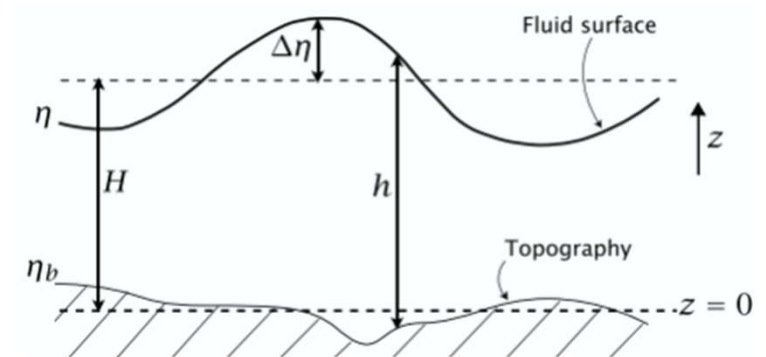
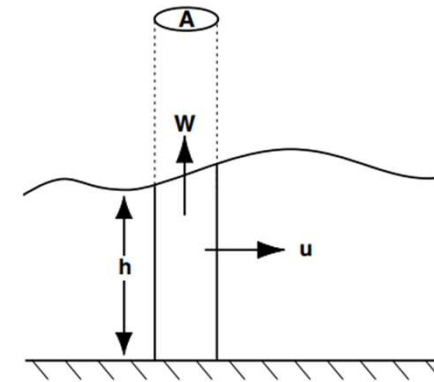
- For 1-D flow(s) the SWEs are written as (Vallis, 2017):

Continuity:

$$\frac{\partial h}{\partial t} + \frac{\partial (uh)}{\partial x} = 0$$

Momentum:

$$\frac{\partial (uh)}{\partial t} + \frac{\partial \left( u^2 h + \frac{1}{2} g h^2 \right)}{\partial x} = 0$$



$h = h(x, t)$ : thickness of liquid column

$H$ : mean height

$\eta$ : height of free surface

$h = \eta - \eta_b$ : height of floor of the container

$u$  = velocity of stream

Contd..

## RECAP

# Introduction

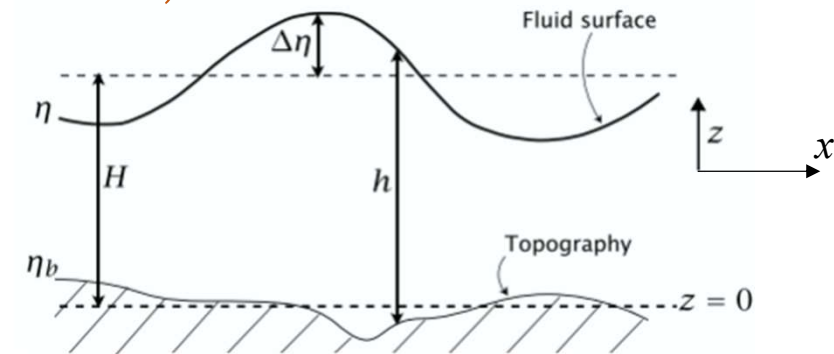
Euler's Equation:

$$\rho \left( \frac{\partial \vec{v}}{\partial t} + \vec{v} \cdot \nabla \vec{v} \right) = -\nabla P$$

Continuity: 
$$\frac{\partial h}{\partial t} + \frac{\partial(uh)}{\partial x} = 0$$

Momentum: 
$$\frac{\partial(uh)}{\partial t} + \frac{\partial\left(u^2h + \frac{1}{2}gh^2\right)}{\partial x} = 0$$

- Basic assumptions for the current set of equations (Vallis, 2017):
  - 1-D flow
  - Inviscid (Assumed due to high Re: Atmospheric flows)
  - Incompressible flow.
  - Negligible Coriolis (Body) forces.
  - Horizontal length scale is very large as compared to the height of the stream. ( $h \ll L$ )



$h = h(x, t)$ : thickness of liquid column

$H$ : mean height


$\eta$ : height of free surface

$h = \eta - \eta_b$ : height of floor of the container

$u$  = velocity of stream

Contd..

# Literature Review

- **Pritchard, D. and Hogg, A. (2002)**
  - **Cao et al. (2004)**
- 
- 1-D and 2-D Hydrodynamic Models built on SWEs to simulate sediment transport under dam break, using finite volume solvers.
  - Regular Terrains
- 
- **Shuangcai et al. (2012):**
    - Discussed about modelling of coupled system of PDEs (Dispersion and Shallow Water) and their outcomes over the un-coupled models.
    - A 2-D Shallow Water Model coupled with Dispersion equation was developed.
    - Simulated for Advection of Pollutant in Square Cavity and with an Asymmetric Dam Break using Finite Volume Solver.

Contd..

# Literature Review & Gap

- **Chistyakov et al.(2020):**
  - 3D Hydrodynamic Model for Salt dispersion and Heat Transfer.
  - 3D Navier-Stokes + Heat & Mass Transport.
- **Kissami Imad., et al. (2021):**
  - Pollutant dispersion using SWE and Dispersion
- **Staples et al. (2023):**
  - The model uses publicly available meteorological data and is applicable to shallow waterbodies less than 1 meter deep.
  - Empirical model that is highly reliant on the publicly available data.

**Gap**  
No heat transfer +  
SWE coupled  
Models

**Objective**

To develop an SWE+HT Model  
to be used in problems with  
smaller domains.  
▪Localized Atmospheric  
Modelling



# Numerical Analysis

## System of Governing Equations

Continuity:  $\frac{\partial h}{\partial t} + \frac{\partial(uh)}{\partial x} = 0$

Momentum:  $\frac{\partial(uh)}{\partial t} + \frac{\partial(u^2h + \frac{1}{2}gh^2)}{\partial x} = 0$

$h = h(x, t); u = u(x, t)$

Finite difference method

“pdepe” solver subroutine:

`sol = pdepe(m, pdefun, icfun, bcfun, xmesh, tspan, options)`

Code equation:  $c\left(x, t, u, \frac{\partial u}{\partial x}\right) \left(\frac{\partial u}{\partial t}\right) = x^{-m} \frac{\partial}{\partial x} \left(x^m f\left(x, t, u, \frac{\partial u}{\partial x}\right)\right) + s\left(x, t, u, \frac{\partial u}{\partial x}\right)$

$m = 1$

$c = \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}$  (diagonal values only)

$f = \begin{bmatrix} -uh \\ -(u^2h + \frac{1}{2}gh^2) \end{bmatrix}$

$s = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$

## RECAP

# 1D Dispersion Model

$$\frac{\partial h}{\partial t} + \frac{\partial(uh)}{\partial x} = 0 \quad - (1)$$

$$\frac{\partial(uh)}{\partial t} + \frac{\partial\left(u^2h + \frac{1}{2}gh^2\right)}{\partial x} = 0 \quad - (2)$$

$$\frac{\partial(\psi h)}{\partial t} + \frac{\partial(\psi uh)}{\partial x} = \frac{\partial\left(\mu \frac{\partial c}{\partial x}\right)}{\partial x} \quad - (3)$$

— (1)

SWE

— (2)

Dispersion Equation

Heat Dispersion Equation (After  
Non dimensionalisation)

$$\frac{\partial(\psi_h h)}{\partial t} + \frac{\partial(\psi_h uh)}{\partial x} = \frac{\partial\left(h_t \frac{\partial T}{\partial x}\right)}{\partial x}$$

T: Temperature

$h_t$ : Overall Heat Transfer Coefficient

- $\psi$  = depth averaged volumetric concentration
- $c$  = local concentration

**Ref:** Shuangcai et al. (2012)

CURRENT WORK..

# Solution Methods

- Some common numerical schemes:
  - Finite Difference Approach
    - First Order Upwind Scheme
    - Lax-Wendroff Scheme
  - Finite Element Approach
  - Finite Volume Approach
- FDM: Approximates the derivatives using finite difference approximations, depending on the scheme, it may be forward difference, backward difference or central difference.
- FEM: Approximations using Piecewise continuous functions, (linear or polynomial) on a spatially discretized domain.

**Ref: Jain & Iyengar**

Contd..

# Solution Methods

- FOU Scheme:
  - First Order Scheme
  - Ensures solution stability (Further discussed in CFL Criteria)
  - It approximates the derivatives using one-sided differences
- Lax- Wendroff Scheme:
  - Second Order Finite Difference Scheme
  - Used for Hyperbolic PDEs, especially for conservation laws.
  - Uses a second order central difference at current time level → For time derivative
  - Uses a second order central difference at current node → For spatial derivative

**Ref: Jain & Iyengar**

# Lax-Wendroff Scheme

- Approximates the solution twice: Once at the half time step, and again at the full time step, as a predictor-corrector approach.

- General scheme of estimation of solution:

- At half time-step:

$$y_m = y_0 + \frac{1}{2} \Delta T \cdot f(y_0)$$

- At full time step

$$y_{m+1} = y_0 + \Delta T \cdot f(y_m)$$

Governing Equations

$$U = \begin{bmatrix} h \\ uh \\ \psi h \end{bmatrix}; E = \begin{bmatrix} uh \\ u^2 h + \frac{1}{2} g h^2 \\ \psi u h \end{bmatrix}; S = \begin{bmatrix} 0 \\ 0 \\ \mu \left( \frac{\partial c}{\partial x} \right) \end{bmatrix}$$

$$\frac{\partial U}{\partial t} + \frac{\partial E}{\partial x} = S$$

**Ref.: C. B. Moler.**

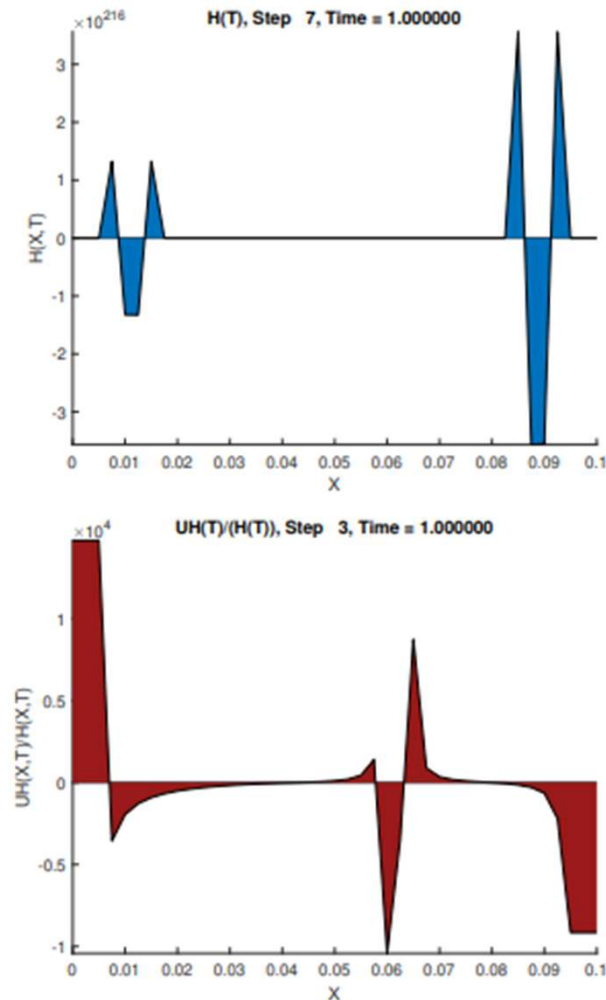
# CFL criteria for numerical analysis

- CFL : Courant-Friedrichs-Lewy condition
- To ensure the number of time-steps and the number of nodes are complying with the following inequality to prevent the loss of convergence at the boundaries.

$$dt \leq \frac{dx}{|u| + \sqrt{gh}}$$

- If we reduce  $dx$ , we may have to reduce  $dt$  as well, if we are near the CFL limit.
- Here  $\sqrt{gh}$  measures motion due to gravity waves.

**Ref.: S. C. Chapra and R. P. Canale.**



*Figure: Output plots for a condition where  $dx$  and  $dt$  do not satisfy the CFL Criteria*

# Model Validation

## Governing Equations

$$U = \begin{bmatrix} h \\ uh \\ \psi h \end{bmatrix}; E = \begin{bmatrix} uh \\ u^2h + \frac{1}{2}gh^2 \\ \psi uh \end{bmatrix}; S = \begin{bmatrix} 0 \\ 0 \\ hK \left( \frac{\partial c}{\partial x} \right) \end{bmatrix}$$

$$\frac{\partial U}{\partial t} + \frac{\partial E}{\partial x} = S$$

Scarcity of relevant data for the SWE+HT modeling.  
Consequently compared the Lax-Wendroff and ANSYS-Fluent CFD solutions.

Estimation of each variable using discretized equations, for each node at each time step.

For Instance, at Half time-step, estimation of  $h$  becomes:

$$hm(i) = \frac{h(i) + h(i+1)}{2} - \frac{\Delta t}{2} \frac{uh(i+1) - uh(i)}{\Delta x}$$

At full time-step:

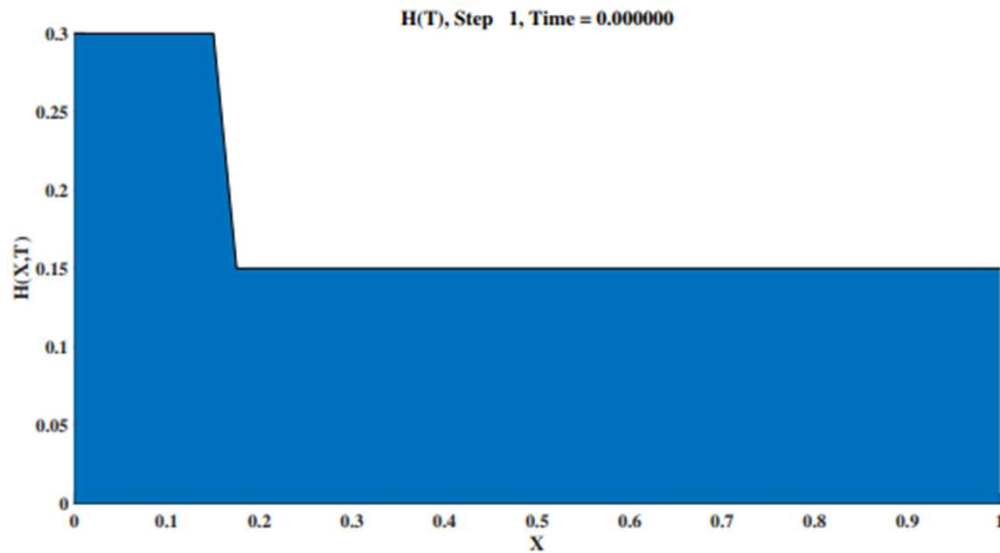
$$\begin{aligned} hm(i) \\ = h(i) + (h(i+1)) - \Delta t \frac{(uh(i+1) - uh(i))}{\Delta x} \end{aligned}$$

*The estimation of the other variables in the system can be estimated similarly*

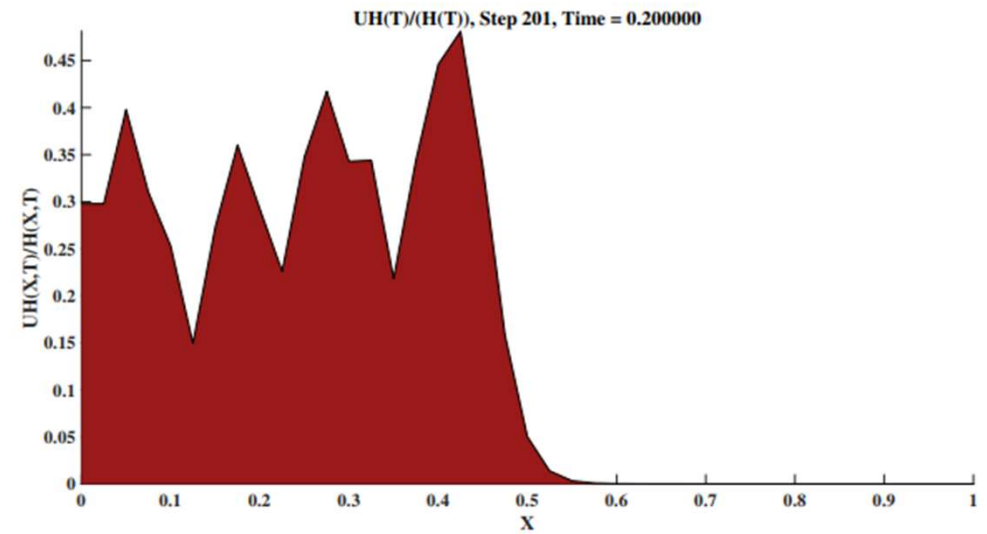
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# Model Validation



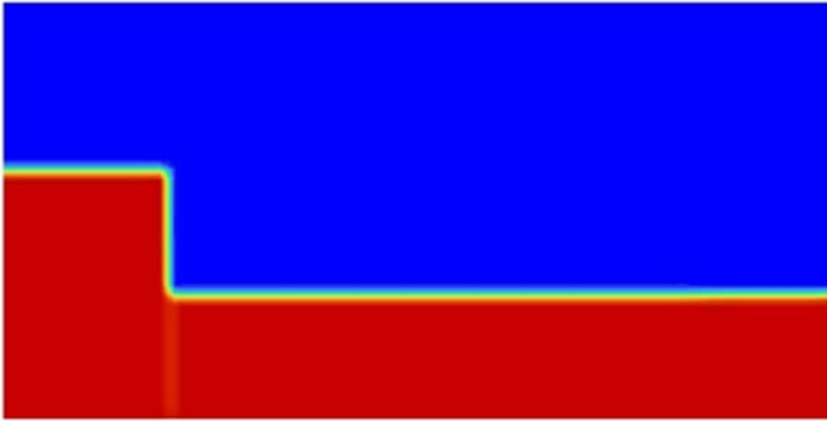
*Figure: Initial time-steps after dam-break (Lax-Wendroff)*



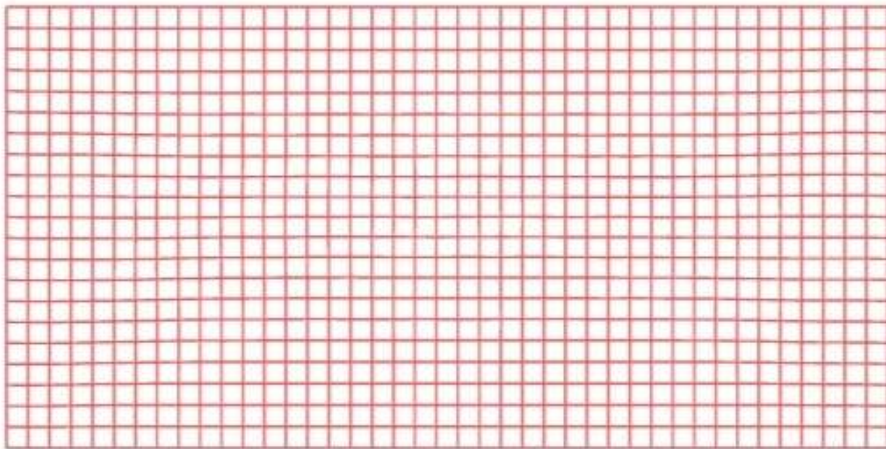
*Figure: Velocity plot at the final time-step (Lax-Wendroff)*

Contd..

# Model Validation



*Figure: Dam Break Initialization using VOF in Ansys*

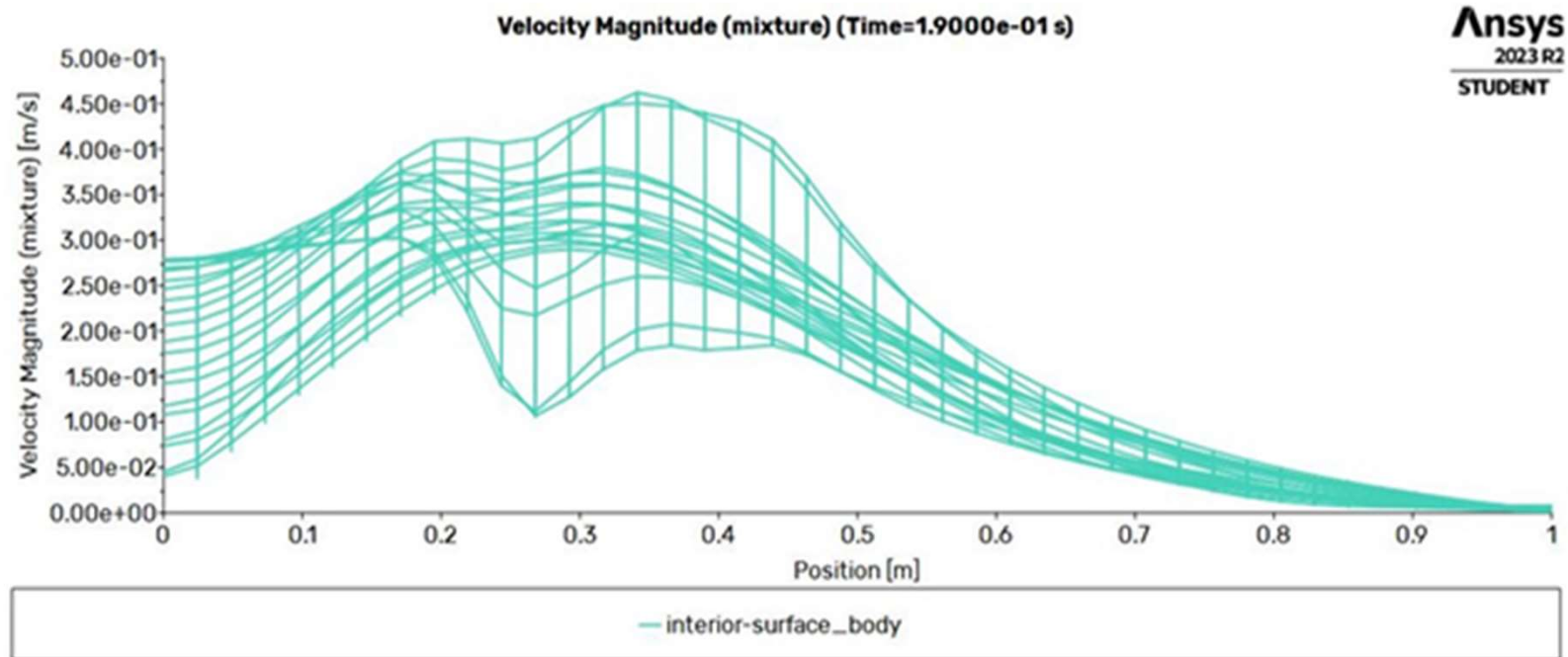


*Figure: Fluent Meshing*

- The mesh is generated in such a way that there are 41 nodes in the X-Direction, as discretized in the Lax-Wendroff Solver discussed in the previous section.
- 200 timesteps each of a size 0.1 seconds
- “SIMPLEC Pressure Velocity Coupling scheme” along with a First order Implicit method for the Transient formulation.
- Inviscid Volume of Fluids approach.

Contd..

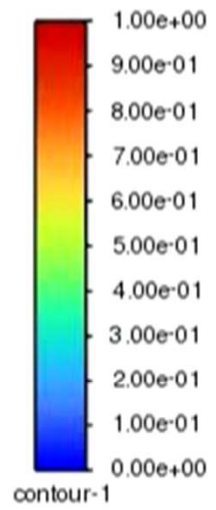
# Model Validation



*Figure: Velocity Plot for timestep corresponding to the timestep of the Lax-Wendroff Scheme*

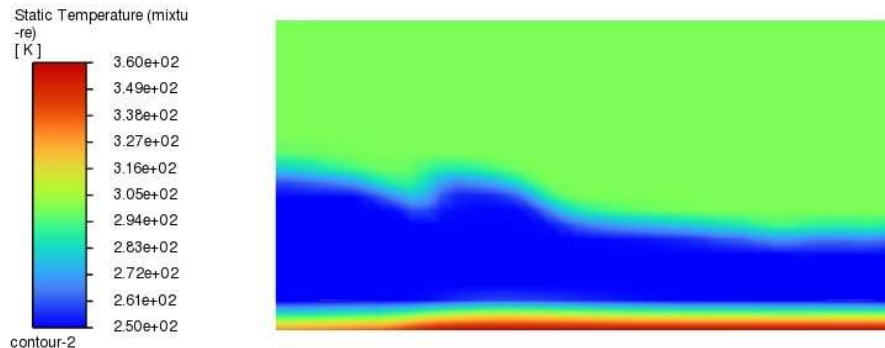
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Volume fraction (water)

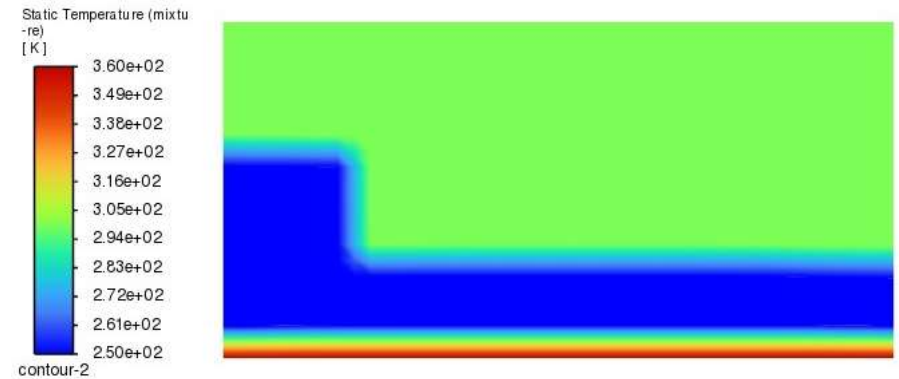


# Heat Transfer Incorporation

- Flat Plate steady state heat transfer: Analytical Solution.
- Numerical unsteady state problem studied using the current scheme.



*Figure: Final timestep for Dam Break along with Surface Heat Transfer*



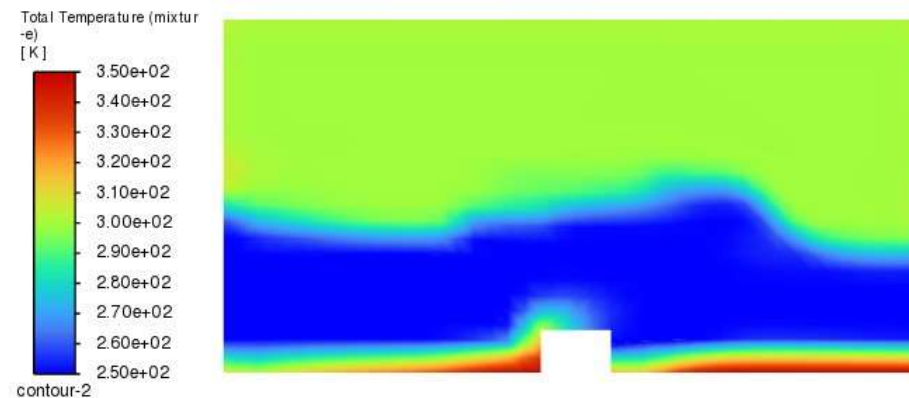
*Figure: Initial Condition before Dam Break for Heat Transfer*

Poor heat transfer from the surface to the fluid as time proceeds, because of the absence of vertical advection in the model.

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# Heat transfer performance by introducing obstacle

- Introducing an obstacle has led to variation in the local velocity magnitude more than for the smooth bottom case. Hence slightly better heat transfer performance is observed.



*Figure: Evaluation of Heat transfer performance using obstacle*

# Future Work & Recommendations

- The project should shift from a first-order approach to a second-order one to properly evaluate this model's effectiveness.
- The project focused solely on smooth bottom scenarios. However, results improved with obstacles, indicating a need for broader study incorporating realistic scenarios.
- To enhance heat transfer predictions, a new coupled system derived from thermodynamic Equations of States (EOS) can incorporate vertical advection, crucial for accurate local temperature and pressure estimation.
- Because of insufficient prior research, the model's validation and verification require expansion through experimental evaluations in specific cases.
- While the Dam-Break Problem is well-researched for hydrodynamics, studying it for atmospheric air flow is impractical. Alternative relevant scenarios are necessary.

# References

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