

## 18.335 Problem Set 4 Solutions

### Problem 1: (10 points)

This is a Galerkin method, very similar to the Rayleigh–Ritz method, and is known as the “FOM” algorithm (“full orthogonalization method”). There are some FOM references on the 18.335 lecture-summary page: in practice, it converges similarly to GMRES, but with some inconvenient quirks, so it is rarely used in practice.

We want to find  $x \in \mathcal{K}_n$  where  $b - Ax$  is  $\perp \mathcal{K}_n$ .  $x \in \mathcal{K}_n$  implies that  $x = Q_n z$  for some  $z \in \mathbb{C}^n$ , and  $b - Ax \perp \mathcal{K}_n$  means  $Q_n^*(b - Ax) = 0 = Q_n^*b - Q_n^*AQ_n z$ . From class (and the book),  $Q_n^*AQ_n = H_n$ . So, we have the  $n \times n$  system of equations

$$H_n z = Q_n^* b$$

whose solution gives  $x = Q_n z$ .

In the common case where our iterative method starts with the initial guess  $x_1 = 0$ , then the first vector in the Arnoldi iteration is  $q_1 = b/\|b\|_2$ , and the above equation simplifies further:  $Q_n^*b = e_1\|b\|_2$  where  $e_1$  is the coordinate vector  $(1, 0, 0, \dots)^T$  as usual, but this simplification isn’t that vital.

### Problem 2: (5+10 points)

Suppose  $A$  is a diagonalizable matrix with eigenvectors  $\mathbf{v}_k$  and eigenvalues  $\lambda_k$ , in decreasing order  $|\lambda_1| \geq |\lambda_2| \geq \dots$ . Recall that the power method starts with a random  $\mathbf{x}$  and repeatedly computes  $\mathbf{x} \leftarrow A\mathbf{x}/\|A\mathbf{x}\|_2$ .

- (a) After many iterations of the power method, the  $\lambda_1$  and  $\lambda_2$  terms will dominate:

$$\mathbf{x} \approx c_1 \mathbf{v}_1 + c_2 \mathbf{v}_2$$

for some  $c_1$  and  $c_2$ . However, this is not an eigenvector. Multiplying this by  $A$  gives  $\lambda_1 c_1 \mathbf{v}_1 + \lambda_2 c_2 \mathbf{v}_2 = \lambda_1 \left( c_1 \mathbf{v}_1 + \frac{\lambda_2}{\lambda_1} c_2 \mathbf{v}_2 \right)$ , which is not a multiple of  $\mathbf{x}$  and hence will be a different vector after normalizing, meaning that it does not converge to any fixed vector.

- (b) The key point is that if we look at the vectors  $\mathbf{x} \approx c_1 \mathbf{v}_1 + c_2 \mathbf{v}_2$  and  $\mathbf{y} \approx \lambda_1 c_1 \mathbf{v}_1 + \lambda_2 c_2 \mathbf{v}_2$  from **two subsequent** iterations, then after **many iterations** these are *linearly independent* vectors that *span the two desired eigenvectors*. We can then employ e.g. a Rayleigh–Ritz procedure to find  $\mathbf{v}_1$  and  $\mathbf{v}_2$ : use Gram–Schmidt to find an orthonormal basis  $\mathbf{q}_1 = \mathbf{x}/\|\mathbf{x}\|_2$  and  $\mathbf{q}_2 = (\mathbf{y} - \mathbf{q}_1 \mathbf{q}_1^* \mathbf{y})/\|\dots\|_2$ , form the matrix  $Q = (\mathbf{q}_1, \mathbf{q}_2)$  and find the  $2 \times 2$  matrix  $A_2 = Q^* A Q$ . The eigenvalues of  $A_2$  (the Ritz values) will then converge to the eigenvalues  $\lambda_1$  and  $\lambda_2$  and we obtain  $\mathbf{v}_1$  and  $\mathbf{v}_2$  (or some multiple thereof) from the corresponding Ritz vectors. The key point is that  $AQ$  is in the span of  $\mathbf{q}_1$  and  $\mathbf{q}_2$  (in the limit of many iterations so that other eigenvectors disappear), so the Ritz vectors are eigenvectors.

Of course, since we don’t know  $\lambda_3$  then we don’t know how many iterations to run, but we can do the obvious convergence tests: every few iterations, find the Ritz values from the last two iterations, and stop when these Ritz values stop changing to our desired accuracy.

Alternatively, if we form the matrix  $X = (\mathbf{x}, \mathbf{y})$  from the vectors of two subsequent iterations, then we know that (after many iterations) the columns of  $AX$  are in  $C(X) = \text{span}\{\mathbf{x}, \mathbf{y}\}$ . Therefore, the problem  $AX = XS$ , where  $S$  is a  $2 \times 2$  matrix, has an exact solution  $S$ . If we then diagonalize  $S = Z\Lambda Z^{-1}$  and multiply both sides by  $Z$ , we obtain  $AXZ = XZ\Lambda$ , and hence the columns of  $XZ$  are eigenvectors of  $A$  and the eigenvalues  $\text{diag } \Lambda$  of  $S$  are the eigenvalues  $\lambda_1$  and  $\lambda_2$  of  $A$ . However, this is computationally equivalent to the Rayleigh–Ritz procedure above, since to solve  $AX = XS$  for  $S$  we would first do a QR factorization  $X = QR$ , and then solve the normal equations  $X^*XS = X^*AX$  via  $RS = Q^*AQ R = A_2 R$ . Thus,  $S = R^{-1}A_2 R$ : the  $S$  and  $A_2$  eigenproblems are similar; in exact arithmetic, the two approaches will give exactly the same eigenvalues and exactly the same Ritz vectors.

[As yet another equivalent alternative, we could write  $AXZ = XZ\Lambda$  as above, and then turn it into

$(X^*AX)Z = (X^*X)Z\Lambda$ , which is a  $2 \times 2$  *generalized* eigenvalue problem, or  $(X^*X)^{-1}(X^*AX)Z = Z\Lambda$ , which is an ordinary  $2 \times 2$  eigenproblem.]

### Problem 3 (15 points)

Trefethen, problem 27.5. The basic answer here is that there *is* a big roundoff error, but it is in the direction of the eigenvector we want, so we don't care (since overall scale factors are irrelevant to eigenvector computations).

In finite precision, instead of  $w = A^{-1}v$ , we will get  $\tilde{w} = w + \delta w$  where  $\delta w = -(A + \delta A)^{-1} \delta A w$  (from the formula on page 95), where  $\delta A = O(\epsilon_{\text{machine}}) \|A\|$  is the backwards error. [Note that we cannot use  $\delta w \approx -A^{-1} \delta A w$ , which neglects the  $\delta A \delta w$  terms, because in this case  $\delta w$  is not small.] The key point, however, is to show that  $\delta w$  is mostly parallel to  $q_1$ , the eigenvector corresponding to the smallest-magnitude eigenvalue  $\lambda_1$  (it is given that all other eigenvalues have magnitude  $\geq |\lambda_2| \gg |\lambda_1|$ ). Since  $w$  is also mostly parallel to  $q_1$ , this will mean that  $\tilde{w}/\|\tilde{w}\|_2 \approx q_1 \approx w/\|w\|_2$ .

First, exactly as in our analysis of the power method, note that  $w = A^{-1}v = \alpha_1 q_1 [1 + O(\lambda_1/\lambda_2)]$ , since  $A^{-1}$  amplifies the  $q_1$  component of  $v$  by  $1/|\lambda_1|$  which is much bigger than the inverse of all the other eigenvalues. Thus,  $w/\|w\|_2 = q_1 [1 + O(\lambda_1/\lambda_2)]$ .

Second, if we Taylor-expand  $(A + \delta A)^{-1}$  in powers of  $\delta A$ , i.e. in powers of  $\epsilon_{\text{machine}}$ , we obtain:<sup>1</sup>  $(A + \delta A)^{-1} = A^{-1} - A^{-1} \delta A A^{-1} + O(\epsilon_{\text{machine}}^2)$ . Since all of the terms in this expansion are multiplied on the *left* by  $A^{-1}$ , when multiplied by *any* vector they will again amplify the  $q_1$  component much more than any other component. In particular, the vector  $\delta A w$  is a vector in a random direction (since  $\delta A$  comes from roundoff and is essentially random) and hence will have some nonzero  $q_1$  component. Thus,  $\delta w = -(A + \delta A)^{-1} \delta A w = \beta_1 q_1 [1 + O(\lambda_1/\lambda_2)]$  for some constant  $\beta_1$ .

Putting these things together, we see that  $\tilde{w} = (\alpha_1 + \beta_1) q_1 [1 + O(\lambda_1/\lambda_2)]$ , and hence  $\tilde{w}/\|\tilde{w}\|_2 = q_1 [1 + O(\lambda_1/\lambda_2)] = \frac{w}{\|w\|_2} [1 + O(\lambda_1/\lambda_2)]$ . Q.E.D.

### Problem 4 (5+5+5+5+5 pts):

Trefethen, problem 33.2:

- (a) In this case, the  $q_{n+1}$  vector is multiplied by a zero row in  $\tilde{H}_n$ , and we can simplify 33.13 to  $AQ_n = Q_n H_n$ . If we consider the full Hessenberg reduction,  $H = Q^* A Q$ , it must have a “block Schur” form:

$$H = \begin{pmatrix} H_n & B \\ 0 & H' \end{pmatrix},$$

where  $H'$  is an  $(m-n) \times (m-n)$  upper-Hessenberg matrix and  $B \in \mathbb{C}^{n \times (m-n)}$ . (It is *not* necessarily the case that  $B = 0$ ; this is only true if  $A$  is Hermitian.)

- (b)  $Q_n$  is a basis for  $\mathcal{K}_n$ , so any vector  $x \in \mathcal{K}_n$  can be written as  $x = Q_n y$  for some  $y \in \mathbb{C}^n$ . Hence, from above,  $Ax = AQ_n y = Q_n H_n y = Q_n (H_n y) \in \mathcal{K}_n$ . Q.E.D.
- (c) The  $(n+1)$  basis vector,  $A^n b$ , is equal to  $A(A^{n-1} b)$  where  $A^{n-1} b \in \mathcal{K}_n$ . Hence, from above,  $A^n b \in \mathcal{K}_n$  and thus  $\mathcal{K}_{n+1} = \mathcal{K}_n$ . By induction,  $\mathcal{K}_\ell = \mathcal{K}_n$  for  $\ell \geq n$ .
- (d) If  $H_n y = \lambda y$ , then  $AQ_n y = Q_n H_n y = \lambda Q_n y$ , and hence  $\lambda$  is an eigenvalue of  $A$  with eigenvector  $Q_n y$ .
- (e) If  $A$  is nonsingular, then  $H_n$  is nonsingular (if it had a zero eigenvalue,  $A$  would too from above). Hence, noting that  $b$  is proportional to the first column of  $Q_n$ , we have:  $x = A^{-1} b = A^{-1} Q_n e_1 \|b\| = A^{-1} Q_n H_n^{-1} e_1 \|b\| = A^{-1} A Q_n H_n^{-1} e_1 \|b\| = Q_n H_n^{-1} e_1 \|b\| \in \mathcal{K}_n$ . Q.E.D.

<sup>1</sup>Write  $(A + \delta A)^{-1} = [A(I + A^{-1} \delta A)]^{-1} = (I + A^{-1} \delta A)^{-1} A^{-1} \approx (I - A^{-1} \delta A) A^{-1} = A^{-1} - A^{-1} \delta A A^{-1}$ . Another approach is to let  $B = (A + \delta A)^{-1} = B_0 + B_1 + \dots$  where  $B_k$  is the  $k$ -th order term in  $\delta A$ , collect terms order-by-order in  $I = (B_0 + B_1 + \dots)(A + \delta A) = B_0 A + (B_0 \delta A + B_1 A) + \dots$ , and you immediately find that  $B_0 = A^{-1}$ ,  $B_1 = -B_0 \delta A A^{-1} = -A^{-1} \delta A A^{-1}$ , and so on.