# 1 Total Energy for GFN2-xTB

Most equations etc. are from the xTB review paper or the GFN2-xTB paper.

$$E_{GFN2-xTB} = E_{rep}^{(0)} + E_{disp}^{(0,1,2)} + E_{EHT}^{(1)} + E_{IES+IXC}^{(2)} + E_{AES+AXC}^{(2)} + E_{IES+IXC}^{(3)}$$
$$= E_{rep} + E_{disp}^{D4'} + E_{EHT} + E_{\gamma} + E_{AES} + E_{AXC} + E_{\Gamma}^{GFN2}$$
(1)

## 1.1 Repulsion Energy

$$E_{rep} = \frac{1}{2} \sum_{A,B} \frac{Z_A^{eff} Z_B^{eff}}{R_{AB}} e^{-\sqrt{a_A a_B} (R_{AB})^{(k_f)}}$$

$$k_f = \begin{cases} 1 & if A, B \in \{H, He\} \\ \frac{3}{2} & otherwise \end{cases}$$

$$(2)$$

 $Z^{eff}$  and a are variables fitted for each element. A,B are the labels of atoms. Since we only have C and H in our systems we can simplify this quite a bit in code.  $R_{AB}$  is the distance between the A and B atoms.

### 1.2 Extended Hückel Theory Energy

$$E_{EHT} = \sum_{\mu\nu} P_{\mu\nu} + H_{\mu\nu}^{EHT}$$
 (3)

where  $\mu$  and  $\nu$  are AO indecies, l and l' index shells. Both AO's are associated with an atom labled A and B.

$$P_{\mu\nu} = P_{\mu\nu}^{(0)} + \delta P_{\mu\nu}$$

$$P_{\mu\nu}^{(0)} = ??$$

$$\delta P_{\mu\nu} = ??$$

$$H_{\mu\nu}^{EHT} = \frac{1}{2} K_{AB}^{ll'} S_{\mu\nu} (H_{\mu\mu} + H_{\nu\nu})$$

$$\cdot X(EN_A, EN_B)$$

$$\cdot \Pi(R_{AB}, l, l')$$

$$\cdot Y(\zeta_l^A, \zeta_{l'}^B), \forall \mu \in l(A), \nu \in l'(B)$$
(4)

QUESTION: How do we calculate the density matrix or the zeroth order and delta terms?

### 1.3 Isotropic electrostatic and XC energy

### 1.3.1 Second order

$$E_{\gamma} = \frac{1}{2} \sum_{A,B}^{N_{atoms}} \sum_{l \in A} \sum_{l' \in B} q_l q_{l'} \gamma_{AB,ll'}$$

$$\gamma_{AB,ll'} = \frac{1}{\sqrt{R_{AB}^2 + \eta_{AB,ll'}^{-2}}}$$

$$\eta_{AB,ll'} = \frac{1}{2} \left[ \eta_A (1 + k_A^l) + \eta_B (1 + k_B^{l'}) \right]$$
(5)

 $q_l$  is a partial muliken charge.  $\eta_A$  and  $\eta_B$  are element-specific fit parameters, while  $k_A^l$  and  $k_B^{l'}$  are element-specific scaling factors for the individual shells  $(k_A^l=0 \text{ when } l=0)$ .

### 1.3.2 Third order

$$E_{\Gamma}^{GFN2} = \frac{1}{3} \sum_{A}^{N_{atoms}} \sum_{l \in A} (q_l)^3 K_l^{\Gamma} \Gamma_A \tag{6}$$

 $K_l^{\Gamma}$  is a shell specific constant common for all elements and  $\Gamma_A$  is an element specific constant.

QUESTION: We can do all of this, except get the partial charges. We assume we have to do Self Consistent Charge, but how we do that is not clear from the GFN2-xTB paper, the xTB review paper or Frank Jensen CH 7.6. Any hints?

### 1.4 Anisotropic electrostatic energy

$$E_{AES} = E_{q\mu} + E_{q\Theta} + E_{\mu\mu}$$

$$= \frac{1}{2} \sum_{A,B} \{ f_3(R_{AB}) [q_A(\boldsymbol{\mu}_B^T \boldsymbol{R}_{BA}) + q_B(\boldsymbol{\mu}_A^T \boldsymbol{R}_{AB})]$$

$$+ f_5(R_{AB}) [q_A \boldsymbol{R}_{AB}^T \boldsymbol{\Theta}_B \boldsymbol{R}_{AB} + q_B \boldsymbol{R}_{AB}^T \boldsymbol{\Theta}_A \boldsymbol{R}_{AB}$$

$$- 3(\boldsymbol{\mu}_A^T \boldsymbol{R}_{AB})(\boldsymbol{\mu}_B^T \boldsymbol{R}_{AB}) + (\boldsymbol{\mu}_A^T \boldsymbol{\mu}_B) R_{AB}^2 ] \}$$

$$(7)$$

QUESTION: Is  $\mathbf{R}_{AB} = (x_i - x_j, y_i - y_j, z_i - z_j)^T$ , if atoms A and B are centered in  $(x_i, y_i, z_i)^T$  and  $(x_j, y_j, z_j)^T$ ?

 $\mu_A$  is the cumulative atomic dipole moment of atom A and  $\Theta_A$  is the corresponding traceless quadrupole moment. The curly braces and brackets are used in the same way as normal parenthesis for showing order of operations.  $q_A$  is

the atomic charge of atom A.

$$\Theta_A^{\alpha\beta} = \frac{3}{2}\theta_A^{\alpha\beta} - \frac{\delta_{\alpha\beta}}{2}\left(\theta_A^{xx} + \theta_A^{yy} + \theta_A^{zz}\right) \tag{8}$$

$$\theta_A^{\alpha\beta} = \sum_{\kappa \in A} \sum_{\lambda} P_{\lambda} \left( \alpha_A D_{\lambda\kappa}^{\beta} + \beta_A D_{\lambda\kappa}^{\alpha} - \alpha_A \beta_A S_{\lambda\kappa} - Q_{\lambda\kappa}^{\alpha\beta} \right) \tag{9}$$

$$q_A = Z_A - \sum_{\kappa \in A} \sum_{\lambda} P_{\kappa \lambda} S_{\lambda \kappa} \tag{10}$$

$$\mu_A^{\alpha} = \sum_{\kappa \in A} \sum_{\lambda} P_{\kappa\lambda} \left( \alpha_A S_{\kappa\lambda} - D_{\kappa\lambda}^{\alpha} \right) \tag{11}$$

$$D^{\alpha}_{\lambda\kappa} = \langle \phi_{\lambda} | \alpha_i | \phi_{\kappa} \rangle \tag{12}$$

$$D^{\alpha}_{\lambda\kappa} = \langle \phi_{\lambda} | \alpha_{i} | \phi_{\kappa} \rangle \tag{12}$$

$$Q^{\alpha\beta}_{\lambda\kappa} = \langle \phi_{\lambda} | \alpha_{i} \beta_{i} | \phi_{\kappa} \rangle \tag{13}$$

 $\alpha$  and  $\beta$  are Cartesian components.

QUESTION: Are we correct in the following, assuming that atom A is centered in  $(x_i, y_i, z_i)^T$ ?

$$\mathbf{\Theta}_{A} = \begin{pmatrix} \Theta_{A}^{xx} & \Theta_{A}^{xy} & \Theta_{A}^{xz} \\ \Theta_{A}^{yx} & \Theta_{A}^{yy} & \Theta_{A}^{yz} \\ \Theta_{A}^{zx} & \Theta_{A}^{zy} & \Theta_{A}^{zz} \\ \Theta_{A}^{zx} & \Theta_{A}^{zy} & \Theta_{A}^{zz} \end{pmatrix}$$

$$oldsymbol{\mu}_A = egin{pmatrix} \mu_A^x \\ \mu_A^y \\ \mu_A^z \end{pmatrix}$$

QUESTION: Is  $\delta$  1 if the labels match and 0 otherwise, i.e. a delta function?

#### Anisotropic XC energy 1.5

$$E_{AXC} = \sum_{A} (f_{XC}^{\mu_A} |\boldsymbol{\mu}_A|^2 + f_{XC}^{\Theta_A} ||\boldsymbol{\Theta}_A||^2)$$
 (14)

Where  $f_{XC}^{\mu_A}$  and  $f_{XC}^{\Theta_A}$  are fitted values. QUESTION: What norms are these, what are the formulas for calculating them?

## 1.6 Dispersion Energy

$$E_{disp}^{D4'} = -\sum_{A>B} \sum_{n=6,8} s_n \frac{C_n^{AB}(q_A, CN_{cov}^A, q_B, CN_{cov}^B)}{R_{AB}^n} f_{damp,BJ}^{(n)}(R_{AB})$$

$$-s_9 \sum_{A>B>C} \frac{(3cos(\theta_{ABC})cos(\theta_{BCA})cos(\theta_{CAB}) + 1)C_9^{ABC}(CN_{cov}^A, CN_{cov}^B, CN_{cov}^C)}{(R_{AB}R_{AC}R_{BC})^3} (15)$$

$$\times f_{damp,zero}^{(9)}(R_{AB}, R_{AC}, R_{BC}).$$

The term in the second line is the three-body Axilrod– Teller–Muto (ATM) (What is this??????) term and the last line is the corresponding zero-damping function for this term.

 $C_6^{AB}$  is the pairwise dipole-dipole dispersion coefficients calculated by numerical integration via the Casimir-Polder relation.

$$C_6^{AB} = \frac{3}{\pi} \sum_j w_j \overline{\alpha}_A(i\omega_j, q_A, CN_{cov}^A) \overline{\alpha}_B(i\omega_j, q_B, CN_{cov}^B)$$
 (16)

 $w_j$  are the integration weights, which are derived from a trapeziodal partitioning between the grid points  $j(j \in [1, 23])$ .

The isotropically averaged, dynamic dipole-dipole polarizabilites  $\overline{\alpha}$  at the jth imaginary frequency  $i\omega_j$  are obtained from the self-consistent D4 model; i.e., they are depending on the covalent coordination number and are also charge dependent.

$$\overline{\alpha}_A(i\omega_j, q_A, CN_{cov}^A) = \sum_r^{N_{A,ref}} \xi_A^r(q_A, q_{A,r}) \overline{\alpha}_{A,r}(i\omega_j, q_{A,r}, CN_{cov}^{A,r}) W_A^r(CN_{cov}^A, CN_{cov}^{A,r})$$

$$(17)$$

The Gaussian weighting for each reference system is given by:

$$W_A^r(CN_{cov}^A, CN_{cov}^{A,r}) = \sum_{j=1}^{N_{gauss}} \frac{1}{N} \exp\left[-6j \cdot (CN_{cov}^A - CN_{cov}^{A,r})^2\right]$$
(18)

with

$$\sum_{N_{A,ref}} W_{A}^{r}(CN_{cov}^{A}, CN_{cov}^{A,r}) = 1$$
 (19)

 $\mathcal{N}$  is a normalization constant.

The number of Gaussian function per reference system  $N_{gauss}$  is mostly one, but equal to three for  $CN_{cov}^{A,r} = 0$  and reference systems with similar coordination number.

The carge-dependency is included via the empirical scaling function  $\xi_A^r$ .

$$\xi_A^r(q_A, q_{A,r}) = \exp\left[3\left\{1 - \exp\left[4\eta_A\left(1 - \frac{Z_A^{eff} + q_{A,r}}{Z_A^{eff} + q_A}\right)\right]\right\}\right]$$
 (20)

where  $\eta_A$  is the chemical hardness taken from ref 98.

 $Z_A^{eff}$  is the effective nuclear charge of atom A, which has been determined by subtracting the number of core electrons represented by the def2-ECPs in the time-dependent DFT reference calculations.

 $C_8^{AB}$  is calculated recursively from the lowest order  $C_6^{AB}$  coefficients.

$$C_8^{AB} = 3C_6^{AB}\sqrt{\mathcal{Q}^A \mathcal{Q}^B} \tag{21}$$

$$Q^{A} = s_{42} \sqrt{Z^{A}} \frac{\langle r^{4} \rangle^{A}}{\langle r^{2} \rangle^{A}} \tag{22}$$

 $\sqrt{Z^A}$  is the ad hoc nuclear charge dependent factor. (How do we get Z??????)  $\langle r4 \rangle$  and  $\langle r2 \rangle$  are simple multipole-type expectation values derived from atomic densities which are averaged geometrically to get the pair coefficients. (What is 'r', how we get???????) (what is  $s_{42}$ ???)

 $CN_{cov}^A$  is the covalent coordination number for atom A.

q is the atomic charge, so  $q_A$  is the atomic charge for atom A.

The damping and scaling parameters in the dispersion model are:

$$a1 = 0.52$$
 |  $a2 = 5.0$  |  $s6 = 1.0$  |  $s8 = 2.7$  |  $s9 = 5.0$ 

 $C_9^{ABC}$  is the triple-dipole constant<sup>1</sup>:

$$C_9^{ABC} = \frac{3}{\pi} \int_0^\infty \alpha^A(i\omega) \alpha^B(i\omega) \alpha^C(i\omega) d\omega$$
 (23)

The three-body contribution is typically < 5-10% of  $E_{disp}$ , so it is small enough that we can reasonably approximate the coefficients by a geometric mean as<sup>1</sup>:

$$C_9^{ABC} \approx -\sqrt{C_6^{AB}C_6^{AC}C_6^{BC}} \tag{24}$$

 $\theta_{ABC}$  is the angle between the two edges going from B to the other two atoms.  $\theta_{BCA}$  is the angle between the edges going from C to the other two and so on.

 $<sup>^1</sup> https://www.researchgate.net/publication/43347348\_A\_Consistent\_and\_Accurate\_Ab\_Initio\_Parametrization\_of\_Density\_FD\_for\_the\_94\_Elements\_H-Pu$ 

BJ = Becke-Johnson

$$f_n^{damp,BJ}(R_{AB}) = \frac{R_{AB}^n}{R_{AB}^n + (a_1 \cdot R_{AB}^{crit} + a_2)^6}$$
 (25)

$$R_{AB}^{crit} = \sqrt{\frac{C_8^{AB}}{C_6^{AB}}} \tag{26}$$

$$f_9^{damp,zero}(R_{AB}, R_{AC}, R_{BC}) = \left(1 + 6\left(\sqrt{\frac{R_{AB}^{crit}R_{BC}^{crit}R_{CA}^{crit}}{R_{AB}R_{BC}R_{CA}}}\right)^{16}\right)^{-1}$$
(27)