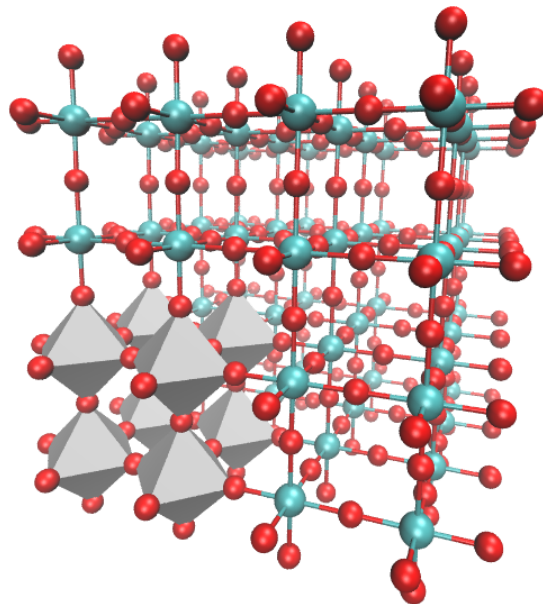




# Slater Determinant Energy



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# Slater Determinants

## Electronic energy of Slater determinant

- Represent wavefunction as a Slater determinant and compute expectation value

$$\langle x \rangle = \frac{\int_{-\infty}^{\infty} \Psi^*(x) \hat{x} \Psi(x) dx}{\int_{-\infty}^{\infty} \Psi^*(x) \Psi(x) dx} \rightarrow E = \frac{\langle \Psi_{SD} | \hat{H} | \Psi_{SD} \rangle}{\langle \Psi_{SD} | \Psi_{SD} \rangle}$$

- Assuming we know the form of the molecular orbitals (which generally we don't)
- Useful relation for the upcoming derivation:

$$\iint f(x)g(y)dx dy = \int f(x)dx \int g(y)dy$$

# Slater Determinants

## Overlap of Slater determinants

- $\langle \Psi_{SD} | \Psi_{SD} \rangle$  is the overlap of two Slater determinants (inner product)

$$\langle \Psi_{SD} | \Psi_{SD} \rangle = \left( \left\langle \prod_{i=1}^{N_{\text{elec}}} \chi_i \left| \hat{A}^\dagger \sqrt{N_{\text{elec}}!} \right. \right\rangle \left( \sqrt{N_{\text{elec}}!} \hat{A} \left| \prod_{j=1}^{N_{\text{elec}}} \chi_j \right. \right) \right)$$

$$\langle \Psi_{SD} | \Psi_{SD} \rangle = N_{\text{elec}}! \left\langle \prod_{i=1}^{N_{\text{elec}}} \chi_i \left| \hat{A} \prod_{j=1}^{N_{\text{elec}}} \chi_j \right. \right\rangle = \sum_{\pi \in S_N} \epsilon_\pi \left\langle \prod_{i=1}^{N_{\text{elec}}} \chi_i \left| \hat{\pi} \prod_{j=1}^{N_{\text{elec}}} \chi_j \right. \right\rangle$$

$$\langle \Psi_{SD} | \Psi_{SD} \rangle = \langle \chi_1 \chi_2 \dots \chi_{N_{\text{elec}}} | \chi_1 \chi_2 \dots \chi_{N_{\text{elec}}} \rangle + \epsilon \langle \chi_1 \chi_2 \dots \chi_{N_{\text{elec}}} | \chi_2 \chi_1 \dots \chi_{N_{\text{elec}}} \rangle + \dots$$

$$\langle \Psi_{SD} | \Psi_{SD} \rangle = \langle \chi_1 | \chi_1 \rangle \langle \chi_2 | \chi_2 \rangle \dots \langle \chi_{N_{\text{elec}}} | \chi_{N_{\text{elec}}} \rangle + \epsilon_{12} \langle \chi_1 | \chi_2 \rangle \langle \chi_2 | \chi_1 \rangle \dots \langle \chi_{N_{\text{elec}}} | \chi_{N_{\text{elec}}} \rangle + \dots$$

$$\langle \Psi_{SD} | \Psi_{SD} \rangle = \delta_{11} \delta_{22} \dots \delta_{N_{\text{elec}}} N_{\text{elec}} + \epsilon_{12} \delta_{12} \delta_{21} \dots \delta_{N_{\text{elec}}} N_{\text{elec}} + \dots$$

$$\langle \Psi_{SD} | \Psi_{SD} \rangle = 1 + 0 + \dots = 1 \quad \text{Only one term survives if ket and bra equal}$$

# Slater Determinants

## Hamiltonian matrix element of Slater determinants

$$E = \frac{\langle \Psi_{SD} | \hat{H} | \Psi_{SD} \rangle}{\langle \Psi_{SD} | \Psi_{SD} \rangle} = \langle \Psi_{SD} | \hat{H} | \Psi_{SD} \rangle$$

$$\hat{H}_{\text{elec}} = - \sum_{i=1}^{N_{\text{elec}}} \frac{1}{2} \nabla_{\mathbf{r}_i}^2 - \sum_I^{N_{\text{atoms}}} \sum_{i=1}^{N_{\text{elec}}} \frac{Z_I}{|\mathbf{r}_i - \mathbf{R}_I|} + \sum_{i=1}^{N_{\text{elec}}} \sum_{j=i+1}^{N_{\text{elec}}} \frac{1}{|\mathbf{r}_i - \mathbf{r}_j|}$$

$$\hat{H}_{\text{elec}} = \sum_{i=1}^{N_{\text{elec}}} \hat{h}_i + \sum_{i=1}^{N_{\text{elec}}} \sum_{j=1}^{N_{\text{elec}}} \frac{1}{|\mathbf{r}_i - \mathbf{r}_j|}$$

$$E_{\text{elec}} = \left\langle \Psi_{SD} \left| \sum_i \hat{h}_i \right| \Psi_{SD} \right\rangle + \left\langle \Psi_{SD} \left| \sum_{i < j} \frac{1}{|\mathbf{r}_{ij}|} \right| \Psi_{SD} \right\rangle$$

$$E_{\text{elec}} = E_{\text{k}} + E_{\text{ne}} + E_{\text{ee}} = E_0 + E_1$$

# Slater Determinants

## 1) One-electron Energy

$$E_0 = \left\langle \Psi_{SD} \left| \sum_i \hat{h}_i \right| \Psi_{SD} \right\rangle = \sum_i \left( \left\langle \prod_k \chi_k \left| \hat{A}^\dagger \sqrt{N_{\text{elec}}} \right. \right\rangle \hat{h}_i \left( \sqrt{N_{\text{elec}}} \left| \hat{A} \prod_j \chi_j \right\rangle \right) \right)$$

$$E_0 = \sum_i N_{\text{elec}}! \left\langle \prod_k \chi_k \left| \hat{h}_i \right| \hat{A} \prod_j \chi_j \right\rangle = \sum_i \left\langle \prod_k \chi_k \left| \hat{h}_i \right| \sum_{\pi \in S_N} \epsilon_\pi \hat{\pi} \prod_j \chi_j \right\rangle$$

$$E_0 = \sum_i \sum_{\pi \in S_N} \epsilon_\pi \left\langle \prod_k \chi_k \left| \hat{h}_i \right| \hat{\pi} \prod_j \chi_j \right\rangle$$

$$E_0 = \langle \chi_1 \chi_2 \dots | \hat{h}_1 | \chi_1 \chi_2 \dots \rangle - \langle \chi_1 \chi_2 \dots | \hat{h}_1 | \chi_2 \chi_1 \dots \rangle + \dots$$

$$E_0 = \langle \chi_1 | \hat{h}_1 | \chi_1 \rangle \langle \chi_2 | \chi_2 \rangle \dots - \langle \chi_1 | \hat{h}_1 | \chi_2 \rangle \langle \chi_2 | \chi_1 \rangle \dots + \dots$$

$$E_0 = \langle \chi_1 | \hat{h}_1 | \chi_1 \rangle + \langle \chi_2 | \hat{h}_2 | \chi_2 \rangle + \dots = \sum_i \langle \chi_i | \hat{h}_i | \chi_i \rangle$$



# Slater Determinants

## 2) Two-electron Energy

$$E_1 = \left\langle \Psi_{SD} \left| \sum_{i < j} r_{ij}^{-1} \right| \Psi_{SD} \right\rangle = \sum_{i < j} \left( \left\langle \prod_k \chi_k \left| \hat{A}^\dagger \sqrt{N_{\text{elec}}}! \right. \right) r_{ij}^{-1} \left( \sqrt{N_{\text{elec}}}! \left| \hat{A} \prod_l \chi_l \right. \right) \right)$$

$$E_1 = \sum_{i < j} N_{\text{elec}}! \left\langle \prod_k \chi_k \left| r_{ij}^{-1} \right| \hat{A} \prod_l \chi_l \right\rangle = \sum_{i < j} \left\langle \prod_k \chi_k \left| r_{ij}^{-1} \right| \sum_{\pi \in S_N} \epsilon_\pi \hat{\pi} \prod_l \chi_l \right\rangle$$

$$E_1 = \sum_{i < j} \sum_{\pi \in S_N} \epsilon_\pi \left\langle \prod_k \chi_k \left| r_{ij}^{-1} \right| \hat{\pi} \prod_l \chi_l \right\rangle$$

$$E_1 = \langle \chi_1 \chi_2 \dots | r_{12}^{-1} | \chi_1 \chi_2 \dots \rangle - \langle \chi_1 \chi_2 \dots | r_{12}^{-1} | \chi_2 \chi_1 \dots \rangle + \dots$$

$$E_1 = \langle \chi_1 \chi_2 | r_{12}^{-1} | \chi_1 \chi_2 \rangle \langle \chi_3 | \chi_3 \rangle \dots - \langle \chi_1 \chi_2 | r_{12}^{-1} | \chi_2 \chi_1 \rangle \langle \chi_3 | \chi_3 \rangle \dots + \dots$$

$$E_1 = \langle \chi_1 \chi_2 | r_{12}^{-1} | \chi_1 \chi_2 \rangle - \langle \chi_1 \chi_2 | r_{12}^{-1} | \chi_2 \chi_1 \rangle + \dots$$

$$E_1 = \sum_{i < j} \langle \chi_i \chi_j | r_{ij}^{-1} | \chi_i \chi_j \rangle - \langle \chi_i \chi_j | r_{ij}^{-1} | \chi_j \chi_i \rangle$$

← Coulomb integral

← Exchange integral

# Slater Determinants

## 2) Two-electron Energy

$$\begin{aligned}
 E_1 &= \sum_{i < j} \langle \chi_i \chi_j | r_{ij}^{-1} | \chi_i \chi_j \rangle - \langle \chi_i \chi_j | r_{ij}^{-1} | \chi_j \chi_i \rangle \\
 &= \frac{1}{2} \sum_{ij} \langle \chi_i \chi_j | r_{ij}^{-1} | \chi_i \chi_j \rangle - \langle \chi_i \chi_j | r_{ij}^{-1} | \chi_j \chi_i \rangle \\
 &= \frac{1}{2} \sum_{ij} \langle \chi_i \chi_j | r_{ij}^{-1} (1 - \hat{\pi}_{12}) | \chi_i \chi_j \rangle
 \end{aligned}$$

## Electronic energy of Slater determinant

$$E = \langle \Psi_{SD} | \hat{H} | \Psi_{SD} \rangle = \sum_i \langle \chi_i | \hat{h} | \chi_i \rangle + \frac{1}{2} \sum_{ij} \langle \chi_i \chi_j | r_{ij}^{-1} (1 - \hat{\pi}_{12}) | \chi_i \chi_j \rangle$$



# Spin Independent Eqns.

Integrate out spin

$$E = \langle \Psi_{SD} | \hat{H} | \Psi_{SD} \rangle = \sum_i \langle \chi_i | \hat{h} | \chi_i \rangle + \frac{1}{2} \sum_{ij} \langle \chi_i \chi_j | r_{ij}^{-1} (1 - \hat{\pi}_{12}) | \chi_i \chi_j \rangle$$

$$\iint f(x)g(y)dx dy = \int f(x)dx \int g(y)dy$$
$$\begin{aligned} |\chi_i(\mathbf{x})\rangle &= |\psi_i(\mathbf{r})\sigma_i(\omega)\rangle \\ \langle \alpha | \alpha \rangle &= \langle \beta | \beta \rangle = 1 \\ \langle \alpha | \beta \rangle &= \langle \beta | \alpha \rangle = 0 \end{aligned}$$

## 1) One-electron integral

$$\begin{aligned} \langle \chi_1(\mathbf{x}_1) | \hat{h}_1(\mathbf{r}_1) | \chi_1(\mathbf{x}_1) \rangle &= \langle \psi_1(\mathbf{r}_1)\sigma_1(\omega_1) | \hat{h}(\mathbf{r}_1) | \psi_1(\mathbf{r}_1)\sigma_1(\omega_1) \rangle \\ &= \langle \psi_1(\mathbf{r}_1) | \hat{h}_1(\mathbf{r}_1) | \psi_1(\mathbf{r}_1) \rangle \langle \sigma_1(\omega_1) | \sigma_1(\omega_1) \rangle \\ &= \langle \psi_1(\mathbf{r}_1) | \hat{h}_1(\mathbf{r}_1) | \psi_1(\mathbf{r}_1) \rangle \end{aligned}$$





# Spin Independent Eqns.

## 2) Coulomb integral

$$\begin{aligned} \left\langle \chi_1(\mathbf{x}_1)\chi_2(\mathbf{x}_2) \left| \frac{1}{r_{12}}(\mathbf{r}_1, \mathbf{r}_2) \right| \chi_1(\mathbf{x}_1)\chi_2(\mathbf{x}_2) \right\rangle &= \\ \left\langle \psi_1(\mathbf{r}_1)\sigma_1(\omega_1)\psi_2(\mathbf{r}_2)\sigma_2(\omega_2) \left| \frac{1}{r_{12}}(\mathbf{r}_1, \mathbf{r}_2) \right| \psi_1(\mathbf{r}_1)\sigma_1(\omega_1)\psi_2(\mathbf{r}_2)\sigma_2(\omega_2) \right\rangle &= \\ \left\langle \psi_1(\mathbf{r}_1)\psi_2(\mathbf{r}_2) \left| \frac{1}{r_{12}}(\mathbf{r}_1, \mathbf{r}_2) \right| \psi_1(\mathbf{r}_1)\psi_2(\mathbf{r}_2) \right\rangle \left\langle \sigma_1(\omega_1)\sigma_2(\omega_2) \left| \sigma_1(\omega_1)\sigma_2(\omega_2) \right\rangle \right. &= \\ \left\langle \psi_1(\mathbf{r}_1)\psi_2(\mathbf{r}_2) \left| \frac{1}{r_{12}}(\mathbf{r}_1, \mathbf{r}_2) \right| \psi_1(\mathbf{r}_1)\psi_2(\mathbf{r}_2) \right\rangle \left\langle \sigma_1(\omega_1) \left| \sigma_1(\omega_1) \right\rangle \right\rangle \left\langle \sigma_2(\omega_2) \left| \sigma_2(\omega_2) \right\rangle \right\rangle &= \\ \left\langle \psi_1(\mathbf{r}_1)\psi_2(\mathbf{r}_2) \left| \frac{1}{r_{12}}(\mathbf{r}_1, \mathbf{r}_2) \right| \psi_1(\mathbf{r}_1)\psi_2(\mathbf{r}_2) \right\rangle & \end{aligned}$$



# Spin Independent Eqns.

## 3) Exchange integral

$$\begin{aligned} & \left\langle \chi_1(\mathbf{x}_1)\chi_2(\mathbf{x}_2) \left| \frac{1}{r_{12}}(\mathbf{r}_1, \mathbf{r}_2)\hat{\pi}_{12} \right| \chi_1(\mathbf{x}_1)\chi_2(\mathbf{x}_2) \right\rangle \quad \text{Note the change of indices in ket functions} \\ & = \left\langle \chi_1(\mathbf{x}_1)\chi_2(\mathbf{x}_2) \left| \frac{1}{r_{12}}(\mathbf{r}_1, \mathbf{r}_2) \right| \chi_1(\mathbf{x}_2)\chi_2(\mathbf{x}_1) \right\rangle \\ & = \left\langle \psi_1(\mathbf{r}_1)\sigma_1(\omega_1)\psi_2(\mathbf{r}_2)\sigma_2(\omega_2) \left| \frac{1}{r_{12}}(\mathbf{r}_1, \mathbf{r}_2) \right| \psi_1(\mathbf{r}_2)\sigma_1(\omega_2)\psi_2(\mathbf{r}_1)\sigma_2(\omega_1) \right\rangle \\ & = \left\langle \psi_1(\mathbf{r}_1)\psi_2(\mathbf{r}_2) \left| \frac{1}{r_{12}}(\mathbf{r}_1, \mathbf{r}_2) \right| \psi_1(\mathbf{r}_2)\psi_2(\mathbf{r}_1) \right\rangle \langle \sigma_1(\omega_1)\sigma_2(\omega_2) | \sigma_1(\omega_2)\sigma_2(\omega_1) \rangle \\ & = \left\langle \psi_1(\mathbf{r}_1)\psi_2(\mathbf{r}_2) \left| \frac{1}{r_{12}}(\mathbf{r}_1, \mathbf{r}_2) \right| \psi_1(\mathbf{r}_2)\psi_2(\mathbf{r}_1) \right\rangle \langle \sigma_1(\omega_1) | \sigma_2(\omega_1) \rangle \langle \sigma_2(\omega_2) | \sigma_1(\omega_2) \rangle \\ & = \left\langle \psi_1(\mathbf{r}_1)\psi_2(\mathbf{r}_2) \left| \frac{1}{r_{12}}(\mathbf{r}_1, \mathbf{r}_2) \right| \psi_1(\mathbf{r}_2)\psi_2(\mathbf{r}_1) \right\rangle \delta_{\sigma_1\sigma_2} \delta_{\sigma_2\sigma_1} \end{aligned}$$

Exchange is zero between electrons of different spins!

# Integral Notation

## Physicist's notation

$$\langle i | \hat{h} | j \rangle = \int \chi_i^*(\mathbf{x}_1) \hat{h} \chi_j(\mathbf{x}_1) d\mathbf{x}_1$$

$$\langle ij | r_{ij}^{-1} | kl \rangle = \langle ij | kl \rangle = \iint \chi_i^*(\mathbf{x}_1) \chi_j^*(\mathbf{x}_2) r_{12}^{-1} \chi_k(\mathbf{x}_1) \chi_l(\mathbf{x}_2) d\mathbf{x}_1 d\mathbf{x}_2$$

$$\langle ij || kl \rangle = \langle ij | kl \rangle - \langle ij | lk \rangle$$

Core:  $h_{ii} = \langle i | \hat{h} | i \rangle$

Coulomb:  $J_{ij} = \langle ij | ij \rangle$

Exchange:  $K_{ij} = \langle ij | ji \rangle$

# Integral Notation

## Chemist's notation

- Spin orbitals

$$[i|\hat{h}|j] = \int \chi_i^*(\mathbf{x}_1) \hat{h} \chi_j(\mathbf{x}_1) d\mathbf{x}_1$$

$$[ij|kl] = \iint \chi_i^*(\mathbf{x}_1) \chi_j(\mathbf{x}_1) r_{12}^{-1} \chi_k^*(\mathbf{x}_2) \chi_l(\mathbf{x}_2) d\mathbf{x}_1 d\mathbf{x}_2 = \langle ik|jl \rangle$$

- Spatial orbitals

$$(i|\hat{h}|j) = \int \psi_i^*(\mathbf{r}_1) \hat{h} \psi_j(\mathbf{r}_1) d\mathbf{r}_1$$

$$(ij|kl) = \iint \psi_i^*(\mathbf{r}_1) \psi_j(\mathbf{r}_1) r_{12}^{-1} \psi_k^*(\mathbf{r}_2) \psi_l(\mathbf{r}_2) d\mathbf{r}_1 d\mathbf{r}_2$$

$$\langle ij||kl \rangle = (ij|kl) - (il|kj)$$

$$\text{Core: } h_{ii} = (i|\hat{h}|i)$$

$$\text{Coulomb: } J_{ij} = (ii|jj)$$

$$\text{Exchange: } K_{ij} = (ij|ji)$$

# Physical Interpretation

## Interpretation of terms in SD energy

- In the chemists spin orbital notation the energy is

$$E_{SD} = \sum_i^{N_{\text{elec}}} h_{ii} + \sum_{i>j}^{N_{\text{elec}}} [ii|jj] - [ij|ji]$$

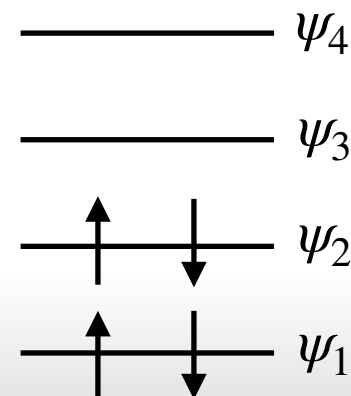
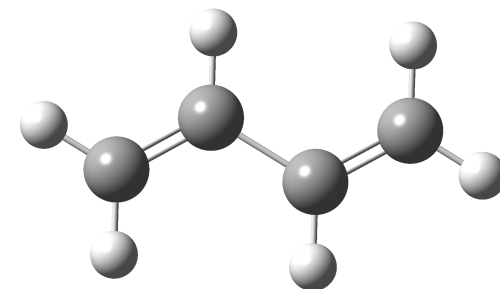
- Not all these terms survive spin integration

$$E_{SD} = h_{11} + h_{\bar{1}\bar{1}} + h_{22} + h_{\bar{2}\bar{2}} + J_{1\bar{1}} + J_{12} + J_{1\bar{2}} + J_{\bar{1}2} + J_{\bar{1}\bar{2}} + J_{2\bar{2}} - K_{12} - K_{\bar{1}\bar{2}}$$

- If spatial parts of each MO are equal for  $\alpha$  and  $\beta$  and  $N_{\alpha}=N_{\beta}$  - **Restricted** wavefunction

$$E_{SD} = 2(h_{11} + h_{22}) + J_{11} + 4J_{12} + J_{22} - 2K_{12}$$

$$E_{SD} = 2 \sum_i^{N_{\text{elec}}/2} h_{ii} + \sum_{ij}^{N_{\text{elec}}/2} 2J_{ij} - K_{ij}$$



# Physical Interpretation

## Interpretation of terms in SD energy

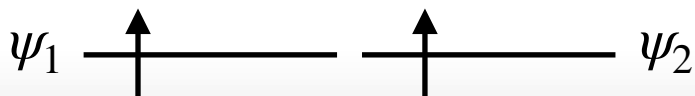
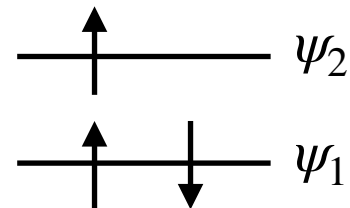
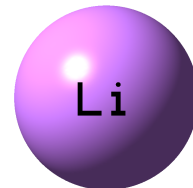
- If  $N_\alpha \neq N_\beta$  but spatial parts of each MO are equal for  $\alpha$  and  $\beta$  - **Restricted open-shell** wavefunction

$$E_{SD} = 2h_{11} + h_{22} + J_{11} + 2J_{12} - K_{12}$$

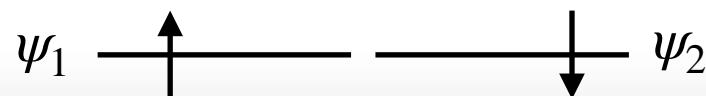
- If spatial parts of any MO are different for  $\alpha$  and  $\beta$  (possibly  $N_\alpha \neq N_\beta$ ) - **Unrestricted** wavefunction

$$E_{SD} = h_{11} + h_{\bar{1}\bar{1}} + h_{22} + J_{1\bar{1}} + J_{12} + J_{\bar{1}\bar{2}} - K_{12}$$

- Origin of Hund's rules (predict high spin)



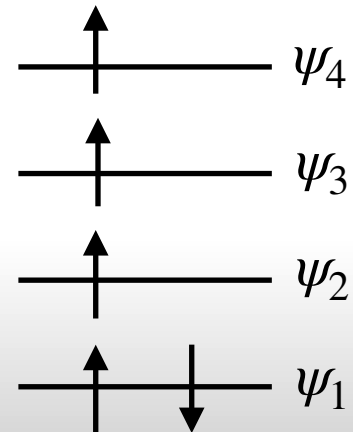
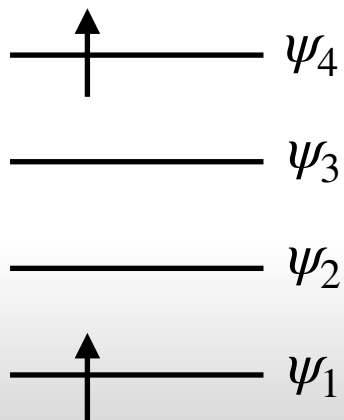
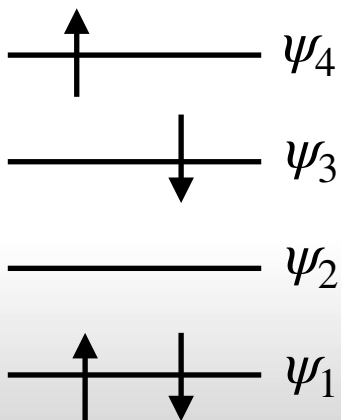
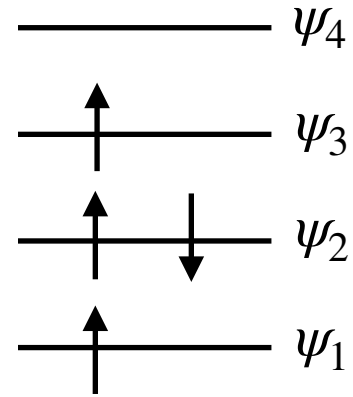
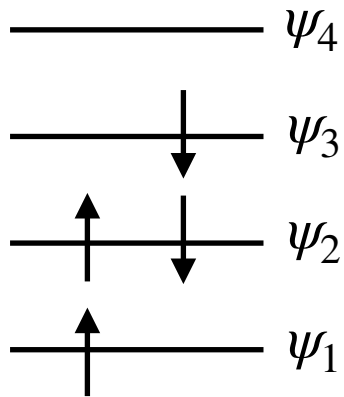
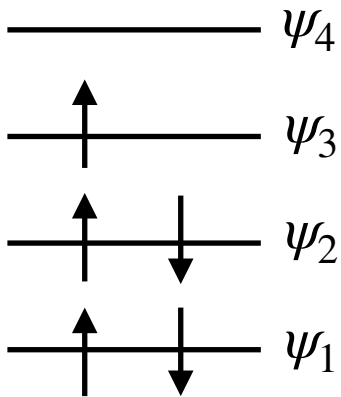
$$E_{SD} = h_{11} + h_{22} + J_{12} - K_{12}$$



$$E_{SD} = h_{11} + h_{\bar{2}\bar{2}} + J_{1\bar{2}}$$



# Physical Interpretation





# Summary

- Need to build in wavefunction antisymmetry manually - Slater Determinant
- Can derive energy of Slater determinant wavefunction given set of known spatial orbitals
- **The problem is, we still don't know the actual form of the orbitals**