

# Jürgen Schaffner-Bielich: Compact Stars

## 4.3 Scaling Solutions for Compact Stars



# Basic Idea

- scaling arguments = reasonable combination of dimensionful parameters

① make TOV dimensionless

② solve dimensionless TOV

③ rescale with  $\sqrt{\varepsilon_0}$  → physical solution



# Making TOV Dimensionless

- $\frac{dP}{dr} = -G \frac{m_r}{r^2} \varepsilon \left(1 + \frac{P}{\varepsilon}\right) \left(1 + \frac{4\pi r^3 P}{m_r}\right) \left(1 - \frac{2Gm_r}{r}\right)^{-1}$

$P = \varepsilon_0 P'$	$\varepsilon = \varepsilon_0 \varepsilon'$
$r = ar'$	$m_r = bm'_r$

- $\frac{\varepsilon_0 dP'}{adr'} = -G \frac{bm'_r}{(ar')^2} \varepsilon_0 \varepsilon' \left(1 + \frac{P'}{\varepsilon'}\right) \left(1 + \frac{4\pi (ar')^3 \varepsilon_0 P'}{bm'_r}\right) \left(1 - \frac{2Gbm'_r}{ar'}\right)^{-1}$

$Gb = a$	$a^3 \varepsilon_0 = b$
$a = (G\varepsilon_0)^{-\frac{1}{2}}$	$b = (G^3 \varepsilon_0)^{-\frac{1}{2}}$

$$\frac{dP'}{dr'} = -G \frac{m'_r}{r'^2} \varepsilon' \left(1 + \frac{P'}{\varepsilon'}\right) \left(1 + \frac{4\pi r'^3 P'}{m'_r}\right) \left(1 - \frac{2m'_r}{r'}\right)^{-1}$$

$$m_r = \frac{m'_r}{\sqrt{G^3 \varepsilon_0}} \sim \frac{1}{\sqrt{\varepsilon_0}} \quad r = \frac{r}{\sqrt{G\varepsilon_0}} \sim \frac{1}{\sqrt{\varepsilon_0}}$$



# Making TOV Dimensionless

$Gb = a$	$a^3 \varepsilon_0 = b$
$a = (G\varepsilon_0)^{-\frac{1}{2}}$	$b = (G^3 \varepsilon_0)^{-\frac{1}{2}}$

$$\frac{dP'}{dr'} = -G \frac{m'_r}{r'^2} \varepsilon' \left(1 + \frac{P'}{\varepsilon'}\right) \left(1 + \frac{4\pi r'^3 P'}{m'_r}\right) \left(1 - \frac{2m'_r}{r'}\right)^{-1}$$

$$m_r = \frac{m'_r}{\sqrt{G^3 \varepsilon_0}} \sim \frac{1}{\sqrt{\varepsilon_0}} \qquad r = \frac{r}{\sqrt{G \varepsilon_0}} \sim \frac{1}{\sqrt{\varepsilon_0}}$$

$$M_{max} = \frac{M'}{\sqrt{G^3 \varepsilon_0}} = M' \frac{m_p^3}{\sqrt{\varepsilon_0}}$$

$$R_{crit} = \frac{R'}{\sqrt{G \varepsilon_0}} = R' \frac{m_p}{\sqrt{\varepsilon_0}}$$



# Example: Free Massive Fermi Gas

- dimensionful quantity: fermion mass

$$P' = \frac{P'}{\varepsilon_0} = \frac{P'}{m_f^4}$$

$$\varepsilon' = \frac{\varepsilon}{\varepsilon_0} = \frac{\varepsilon'}{m_f^4}$$

- $\rightarrow \varepsilon_0 = m_f^4$

$$M_{max} = M' \frac{m_p^3}{\sqrt{\varepsilon_0}} = M' \frac{m_p^3}{m_f^2}$$

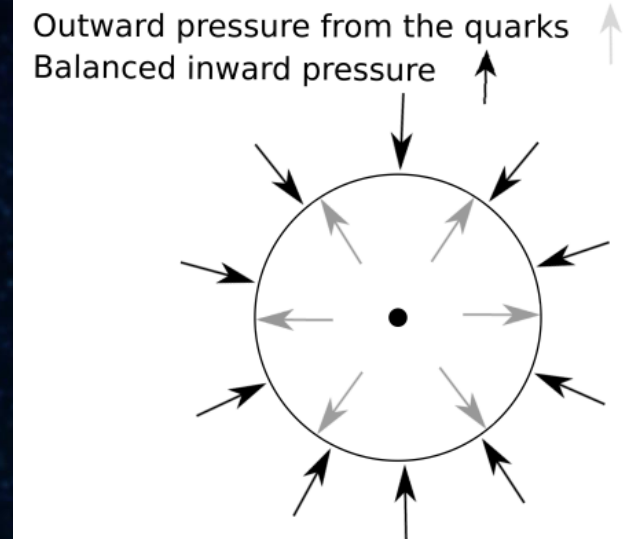
$$R_{crit} = R' \frac{m_p}{\sqrt{\varepsilon_0}} = R' \frac{m_p}{m_f^2}$$

- Compactness  $C = \frac{GM}{R}$
- $C_{max} = \frac{M'}{R'} = \frac{0.384}{3.367} = 0.11$



# Example: Relativistic Fermi Gas with Vacuum Term

- relativistic particles:  $P = \frac{1}{3}\varepsilon \rightarrow P = \frac{1}{3}(\varepsilon - \varepsilon_0)$
- vanishing  $P$  at nonvanishing  $\varepsilon$
- MIT bag model:
  - $P = P_{free} - B \leftrightarrow \varepsilon = \varepsilon_{free} + B$
  - relativistic massless gas of particles:  $P_{free} = \frac{1}{3}\varepsilon_{free}$
  - $P = \frac{1}{3}\varepsilon_{free} - B = \frac{1}{3}(\varepsilon - B) - B = \frac{1}{3}(\varepsilon - 4B)$   
 $\rightarrow \varepsilon_0 = 4B$
  - $C = 0.271$  ( $C_{free Fermion} = 0.11$ )





# Limiting EOS from causality

- stiffest possible EOS (Zel'dovich):  $P = \varepsilon$

- $c_{sound}^2 = \frac{\partial P}{\partial \varepsilon} = 1$
- maximum  $P$  for given  $\varepsilon$
- unstable TOV solutions

$$P = P_f + (\varepsilon - \varepsilon_f)$$

- know low-density EOS for  $0 < P < P_f$
- stiffest possible EOS for higher densities must be Zel'dovich EOS
  - maximum  $P \leftrightarrow$  maximum  $M$

- Rhoades-Ruffini mass limit:  $M_{max} = M' \frac{1}{\sqrt{G^3 \varepsilon_f}} = 4.2 M_{solar} \sqrt{\frac{\varepsilon_{nm}}{\varepsilon_f}}$



# Selfbound Linear EOS

- selfbound stars:
  - vanishing  $P$  at nonvanishing  $\varepsilon$
  - vanishing  $P$  at surface
  - do not need gravity to be stabilized
- $P = s(\varepsilon - \varepsilon_0)$ 
  - $c_{sound}^2 = \frac{\partial P}{\partial \varepsilon} = s$
- redshift factor:  $1 + z = \frac{1}{\sqrt{\frac{1}{1-2C}}}$



# Selfbound Linear EOS

s=1/3 MIT bag model	$M_{max}$ $= 2.57 M_{solar} \sqrt{\frac{\epsilon_{nm}}{\epsilon_0}}$	$R_{crit}$ $= 14 \text{ km} \sqrt{\frac{\epsilon_{nm}}{\epsilon_0}}$	$C_{max}$ $= 0.271$	$Z_{max}$ $= 0.478$	$\frac{\epsilon_{max}}{\epsilon_0} = 4.81$	$\frac{\epsilon_{max}(2M_{solar})}{\epsilon_{nm}} = 7.91$
s=2/3	$M_{max}$ $= 3.64 M_{solar} \sqrt{\frac{\epsilon_{nm}}{\epsilon_0}}$	$R_{crit}$ $= 16.4 \text{ km} \sqrt{\frac{\epsilon_{nm}}{\epsilon_0}}$	$C_{max}$ $= 0.328$	$Z_{max}$ $= 0.705$	$\frac{\epsilon_{max}}{\epsilon_0} = 3.54$	$\frac{\epsilon_{max}(2M_{solar})}{\epsilon_{nm}} = 11.8$
s=3/3 Fel'dovic h	$M_{max}$ $= 4.23 M_{solar} \sqrt{\frac{\epsilon_{nm}}{\epsilon_0}}$	$R_{crit}$ $= 17.6 \text{ km} \sqrt{\frac{\epsilon_{nm}}{\epsilon_0}}$	$C_{max}$ $= 0.354$	$Z_{max}$ $= 0.851$	$\frac{\epsilon_{max}}{\epsilon_0} = 3.03$	$\frac{\epsilon_{max}(2M_{solar})}{\epsilon_{nm}} = 13.5$





Thank you for listening



# NOTES 1

- free massive fermi gas
  - $P=?$
  - rescale with fermi mass<sup>4</sup>
  - $C=0,11$
- relativistic fermi gas with vacuum term
  - $P=1/3 e \rightarrow$  polytrope 1
  - rescale with offset  $P=1/3(e-e_0) \rightarrow$  pressure vanishes at nonvanishing  $e$
  - MIT bag model
    - gas of free relativistic quarks in bag stabilizes by outside pressure  $B \rightarrow P_{\text{tot}} = P_{\text{free}} - B$  and  $e = e_{\text{free}} + B$  (nonvanishing vacuum energy density contributions due to quark gluon condensates)
    - relativistic massless gas of particles:  $P_{\text{free}} = 1/3 e_{\text{free}} \rightarrow P = 1/3 e_{\text{free}} - B = 1/3 (e-B) - B = 1/3 (e-4B)$
    - $M_{\text{max}}, R_{\text{crit}} \rightarrow C=0.271 \rightarrow$  more compact than free fermion star



# NOTES 2

- limiting EOS from causality
  - stiffest possible EOS (Zel'dovich):  $P = \varepsilon$ 
    - $c_{sound}^2 = \frac{\partial P}{\partial \varepsilon} = 1$
    - maximum  $P$  for given  $\varepsilon$
    - unstable TOV solutions
- know low-density EOS for  $0 < P < P_f$
- stiffest possible EOS for higher densities must be Zel'dovich EOS
  - maximum  $P \leftrightarrow$  maximum  $M$
- Rhoades-Ruffini mass limit:  $M_{max} = M' \frac{1}{\sqrt{G^3 \varepsilon_f}} = 4.2 M_{solar} \sqrt{\frac{\varepsilon_{nm}}{\varepsilon_f}}$ 
  - would give  $M_{max} = 3 M_s$  for  $\varepsilon_f = \varepsilon_{nm}$  but we know  $4.2 M_s$  is true

$$P = P_f + (\varepsilon - \varepsilon_f)$$



# NOTES 3

- selfbound stars:
  - vanishing  $P$  at nonvanishing  $\varepsilon$
  - vanishing  $P$  at surface
  - do not need gravity to be stabilized

- selfbound linear eos → see table

MIT bag model $s=1/3$	$M_{max}$ $= 2.57 M_{solar} \sqrt{\frac{\varepsilon}{\varepsilon_0}}$	$R_{crit}$ $= 14 km \sqrt{\frac{\varepsilon_{nm}}{\varepsilon_0}}$	$C_{max} = 0.271$	$z_{max} = 0.478$
$s=1/3$	$M_{max}$ $= 3.64 M_{solar} \sqrt{\frac{\varepsilon}{\varepsilon_0}}$	$R_{crit}$ $= 16.4 km \sqrt{\frac{\varepsilon_{nm}}{\varepsilon_0}}$	$C_{max} = 0.328$	$z_{max} = 0.705$
$s=1/3$ Fel'dovich	$M_{max}$ $= 4.23 M_{solar} \sqrt{\frac{\varepsilon}{\varepsilon_0}}$	$R_{crit}$ $= 17.6 km \sqrt{\frac{\varepsilon_{nm}}{\varepsilon_0}}$	$C_{max} = 0.354$	$z_{max} = 0.851$



# Selfbound Linear EOS- whole table

s=1/3 MIT bag model	$M_{max}$ $= 2.57 M_{solar} \sqrt{\frac{\epsilon_{nm}}{\epsilon_0}}$	$R_{crit}$ $= 14 \text{ km} \sqrt{\frac{\epsilon_{nm}}{\epsilon_0}}$	$C_{max}$ $= 0.271$	$Z_{max}$ $= 0.478$	$\frac{\epsilon_{max}}{\epsilon_0} = 4.81$	$\frac{\epsilon_{max}(2M_{solar})}{\epsilon_{nm}} = 7.91$
s=2/3	$M_{max}$ $= 3.64 M_{solar} \sqrt{\frac{\epsilon_{nm}}{\epsilon_0}}$	$R_{crit}$ $= 16.4 \text{ km} \sqrt{\frac{\epsilon_{nm}}{\epsilon_0}}$	$C_{max}$ $= 0.328$	$Z_{max}$ $= 0.705$	$\frac{\epsilon_{max}}{\epsilon_0} = 3.54$	$\frac{\epsilon_{max}(2M_{solar})}{\epsilon_{nm}} = 11.8$
s=3/3 Fel'dovic h	$M_{max}$ $= 4.23 M_{solar} \sqrt{\frac{\epsilon_{nm}}{\epsilon_0}}$	$R_{crit}$ $= 17.6 \text{ km} \sqrt{\frac{\epsilon_{nm}}{\epsilon_0}}$	$C_{max}$ $= 0.354$	$Z_{max}$ $= 0.851$	$\frac{\epsilon_{max}}{\epsilon_0} = 3.03$	$\frac{\epsilon_{max}(2M_{solar})}{\epsilon_{nm}} = 13.5$