

Technical note: Estimating relative entropies of path ensembles

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We consider the problem of estimating the relative Kolmogorov-Sinai entropies of path (trajectory) ensembles aided by transition path sampling biased by trajectory activity. We loosely follow the development and terminology in [1].

PRELIMINARIES

Consider a system of N particles, with a point in phase space for this system denoted by x . Let X denote a trajectory of the system observed at discrete intervals for a fixed total observation time t_{obs} . For convenience, we write $X \equiv \{x_0, x_1, \dots, x_T\}$, where the subscripted x_t denote these discrete observation times indexed by $t = 0, 1, \dots, T$.

The equilibrium dynamics of this system is given by the probability functional (sometimes termed a “path ensemble”)

$$P_0[A] = Z_0^{-1} q_0[X] \quad (1)$$

$$Z_0 = \int dX q_0[X] \quad (2)$$

where Z_0 is a normalizing constant, or partition function over trajectories, the integral is taken over the entire phase space, dX is an appropriate path differential or measure, and $q_0[X] \geq 0$ is an unnormalized probability density.

For simplicity, we suppose here that the system evolves under Hamiltonian dynamics with initial phase space points x_0 canonical distributed according to an equilibrium NVT ensemble:

$$q_0[X] = e^{-\beta H(x_0)} \quad (3)$$

where β denotes the inverse temperature $1/k_B T$ and $H(x)$ the Hamiltonian for the system. This assumption is not essential to the subsequent analysis and suggested sampling scheme; other dynamical models (e.g. Langevin, Brownian) or ensembles (e.g. NPT) can be treated by substituting alternative definitions of $q_0[X]$ that will be related to the exponentiated path action.

Define an order parameter related to the dynamical activity of the system

$$K[X] \equiv \sum_{t=1}^T \|x_t - x_{t-1}\|^2 \quad (4)$$

We consider modified path ensembles $P_s[X]$ parameter-

ized by the conjugate field variable s :

$$P_s[X] = Z_s^{-1} q_s[X] \quad (5)$$

$$Z_s = \int dX q_s[X] \quad (6)$$

$$q_s[X] = q_0[X] e^{-s K[X]} \quad (7)$$

When $s < 0$, trajectories with large activity $K[X]$ will be more frequently sampled than in the equilibrium ensemble $P_0[X]$. Similarly, when $s > 0$, these trajectories will be suppressed. Choosing $s = 0$ recovers equilibrium dynamics.

KOLMOGOROV-SINAI TRAJECTORY ENTROPY

Suppose we are interested in the Kolmogorov-Sinai entropy of trajectories for some arbitrary path ensemble $P[X] = Z^{-1} q[X]$:

$$S\{P[X]\} \equiv - \int dX P[X] \ln P[X] \quad (8)$$

$$= - \langle \ln P[X] \rangle \quad (9)$$

$$= - \langle \ln q[X] \rangle + \ln Z \quad (10)$$

While it is difficult to compute the *absolute* entropy for a given path ensemble due to the difficulty of determining the absolute normalization constant Z , we are often only interested in the *relative entropy* of two path ensembles. The relative entropy of the path ensembles $P_s[X]$ at different values of s , for example, is:

$$\begin{aligned} \Delta S(s_1, s_2) &\equiv S\{P_2[X]\} - S\{P_1[X]\} \\ &= - \langle \ln P_{s_2}[X] \rangle_{s_2} + \langle \ln P_{s_1}[X] \rangle_{s_1} \\ &= - \langle \ln q_{s_2}[X] \rangle_{s_2} + \langle \ln q_{s_1}[X] \rangle_{s_1} + \ln(Z_{s_2}/Z_{s_1}) \end{aligned} \quad (11)$$

This quantity involves only expectations of known functions $q_s[X]$ and ratios of normalization constants Z_{s_2}/Z_{s_1} , which are readily computed by standard simulation approaches.

For convenience, we define the functional

$$U_s[X] \equiv - \ln q_s[X] = - \ln q_0[X] + s K[X] \quad (12)$$

which, for the case of Hamiltonian dynamics, reduces to

$$U_s[X] = \beta H(x_0) + s K[X] \quad (13)$$

This functional is analogous to an *internal energy* with contribution from an external field parameterized by s .

Similarly, we define a quantity F_s that is analogous to a free energy:

$$F_s \equiv -\ln Z_s \quad (14)$$

With these definitions, the relative entropy can be written a form reminiscent of a restatement of the Helmholtz free energy:

$$\Delta S(s_1, s_2) = [\langle U_{s_2} \rangle_{s_2} - \langle U_{s_1} \rangle_{s_1}] - [F_{s_2} - F_{s_1}] \quad (15)$$

Alternatively, consider the computation of the relative entropy of the subset of trajectories in the unbiased path ensemble $P_0[X]$ for which $K[X] > K_*$. We define a modified path ensemble $P_*[X]$ such that

$$\begin{aligned} P_*[X] &= Z_*^{-1} q_*[X] \\ Z_*[X] &= \int dX q_*[X] \\ q_*[X] &= \chi_*[X] q_0[X] \end{aligned} \quad (16)$$

where $\chi_*[X]$ is an indicator function such that

$$\chi_*[X] = \begin{cases} 1 & \text{if } K[X] > K_* \\ 0 & \text{otherwise} \end{cases} \quad (17)$$

This gives the relative entropy

$$\begin{aligned} \Delta S_* &= S\{P_*[X]\} - S\{P_0[X]\} \\ &= -\langle \ln q_*[X] \rangle_* + \langle \ln q_0[X] \rangle_0 + \ln(Z_*/Z_0) \\ &= [\langle U_0 \rangle_* - \langle U_0 \rangle_0] - [F_* - F_0] \end{aligned} \quad (18)$$

The next section shows how to efficiently estimate these quantities from simulations such as those conducted in [1].

EFFICIENT ESTIMATION FROM SIMULATION

Suppose we have selected K values of the field s , denoted $\{s_1, \dots, s_K\}$, from which we have sampled N_k i.i.d. trajectories from each of the corresponding path ensembles $P_{s_k}[X]$:

$$X_{kn} \sim P_{s_k}[X], \quad n = 1, \dots, N_k \quad (19)$$

These samples can be generated efficiently from transition path sampling [2] simulations, potentially employing a replica-exchange scheme [3] among field values to ensure ergodic sampling from all selected values s_k . We will additionally use an indexing scheme where all $N = \sum_{k=1}^K N_k$ samples are denoted by X_n indexed by a single $n = 1, \dots, N$, where the order in which samples coming from different $P_{s_k}[X]$ are arranged is irrelevant.

We can optimally estimate path averages of the form

$$\langle \mathcal{F} \rangle \equiv \int dX P[X] \mathcal{F}[X] \quad (20)$$

and ratios of normalization constants Z_i/Z_j from trajectories sampled from multiple path ensembles (in the sense that the estimator is *asymptotically efficient*) using a recently-developed formalism [4] based on the optimal *extended bridge sampling* estimators developed in statistics [5–8] and recently applied to statistical mechanics [9].

To compute the optimal estimator, we first solve a set of K nonlinear equations self-consistently for the estimates $\hat{Z}_{s_1}, \dots, \hat{Z}_{s_K}$ of the normalizing constants Z_{s_1}, \dots, Z_{s_K} :

$$\begin{aligned} \hat{Z}_{s_i} &= \sum_{n=1}^N \frac{q_{s_i}[X_n]}{\sum_{k=1}^K N_k \hat{Z}_{s_k}^{-1} q_{s_k}[X_n]} \\ &= \sum_{n=1}^N \frac{q_0[X_n] e^{-s_i K[X_n]}}{\sum_{k=1}^K N_k \hat{Z}_{s_k}^{-1} q_0[X_n] e^{-s_k K[X_n]}} \\ &= \sum_{n=1}^N \left[\sum_{k=1}^K N_k \hat{Z}_{s_k}^{-1} e^{-(s_k - s_i) K[X_n]} \right]^{-1} \end{aligned} \quad (21)$$

Though any scheme to solve the K coupled nonlinear equations indexed by $i = 1, \dots, K$ can be used here, several efficient and numerically stable approaches are detailed in [9]. Note that the \hat{Z}_{s_i} are only determined up to a multiplicative constant, such that only ratios $\hat{Z}_{s_i}/\hat{Z}_{s_j}$ are meaningful.

For an arbitrary path ensemble $P[X] = Z^{-1} q[X]$, we can estimate the normalization constant Z up to the same normalization constant without further iteration by

$$\hat{Z} = \sum_{n=1}^N \frac{q[X_n]}{\sum_{k=1}^K N_k \hat{Z}_{s_k}^{-1} q_{s_k}[X_n]} \quad (22)$$

For some path ensemble $P_s[X]$ in particular, even if no samples are collected from this path ensemble, this reduces to

$$\hat{Z}_s = \sum_{n=1}^N \left[\sum_{k=1}^K N_k \hat{Z}_{s_k}^{-1} e^{-(s_k - s) K[X_n]} \right]^{-1} \quad (23)$$

Path expectations $\langle \mathcal{F} \rangle$ for a path function $\mathcal{F}[X]$ for an arbitrary path ensemble $P[X] = Z^{-1} q[X]$ can also be easily computed. Define a the normalized weight function

$$w_n = \hat{Z}^{-1} \frac{q[X_n]}{\sum_{k=1}^K N_k \hat{Z}_{s_k}^{-1} q_{s_k}[X_n]} \quad (24)$$

Note that, because of the large dynamic range that can be assumed by quantities like w_n and Z_s , it is more numerically stable to work with their logarithms.

The path average estimate $\bar{\mathcal{F}} \approx \langle \mathcal{F} \rangle$ is then given by

$$\bar{\mathcal{F}} = \sum_{n=1}^N w_n \mathcal{F}[X_n] \quad (25)$$

For example, we estimate $\langle U_s \rangle_s$ by

$$\langle U_s \rangle_s \approx \sum_{n=1}^N w_{ns} U_s[X_n] \quad (26)$$

where the s -dependent weight w_{ns} has simple form:

$$w_{ns} = \hat{Z}_s^{-1} \left[\sum_{k=1}^K N_k \hat{Z}_{s_k}^{-1} e^{-(s_k-s)K[X_n]} \right]^{-1} \quad (27)$$

This allows us to estimate the relative entropy between two s -ensembles as

$$\Delta \hat{S}(s_1, s_2) = \sum_{n=1}^N (w_{ns_2} U_{s_2}[X_n] - w_{ns_1} U_{s_1}[X_n]) + \ln(\hat{Z}_{s_2}/\hat{Z}_{s_1}) \quad (28)$$

or the relative entropy of the subset of equilibrium trajectories with $K[X] > K_*$ for $s = 0$ by

$$\Delta S_* = \sum_{n=1}^N (w_{n*} - w_{n0}) U_0[X_n] + \ln(\hat{Z}_*/\hat{Z}_0) \quad (29)$$

where w_{n*} and Z_* contain contributions from only those trajectories with $K[X] > K_*$:

$$w_{n*} = \hat{Z}_*^{-1} \chi_*[X] \left[\sum_{k=1}^K N_k \hat{Z}_{s_k}^{-1} e^{-s_k K[X_n]} \right]^{-1} \quad (30)$$

$$Z_* = \sum_{n=1}^N \chi_*[X] \left[\sum_{k=1}^K N_k \hat{Z}_{s_k}^{-1} e^{-s_k K[X_n]} \right]^{-1} \quad (31)$$

Estimation of the statistical error in these estimates (through estimation of the asymptotic variance) is more involved, but mechanically straightforward; the procedure is described in detail and illustrated in [4]. The algebra is lengthy and unilluminating, but the Python `pymbar` package [available at <http://www.simtk.org/home/pymbar>] has a function to do this analysis automatically.

OTHER OBSERVABLES

(There is a slight change of notation here.)

Expected activity as a function of field s

We can estimate $\langle K \rangle_s$ as

$$\langle K \rangle_s \approx \sum_{n=1}^N w_n(s) K[X_n] \quad (32)$$

where

$$w_n(s) = \hat{Z}(s)^{-1} \left[\sum_{k=1}^K N_k \hat{Z}_k^{-1} e^{-(s_k-s)K[X_n]} \right]^{-1}$$

$$\hat{Z}(s) = \sum_{n=1}^N \left[\sum_{k=1}^K N_k \hat{Z}_k^{-1} e^{-(s_k-s)K[X_n]} \right]^{-1} \quad (33)$$

Susceptibility

We can also estimate the susceptibility, $\chi(s) \equiv -\partial \langle K \rangle_s / \partial s$, in either of two mathematically (but not numerically) equivalent ways. Differentiating our estimate of $\langle K \rangle_s$ produces an estimator expression

$$\chi(s) = \frac{\partial \langle K \rangle_s}{\partial s} \approx \sum_{n=1}^N \frac{\partial w_n}{\partial s}(s) K[X_n] \quad (34)$$

where the weight derivatives are given by

$$\frac{\partial w_n}{\partial s}(s) = w_n(s) (-K[X_n] + \langle K \rangle_s) \quad (35)$$

Alternatively, the mathematically equivalent fluctuation form of $\chi(s)$ can be used,

$$\chi(s) = \langle (K - \langle K \rangle_s) \rangle_s = \langle K^2 \rangle_s - \langle K \rangle_s^2 \quad (36)$$

and the two moments estimated separately. While this expression may run into numerical precision issues if $\langle K^2 \rangle$ and $\langle K \rangle^2$ are individually large but have a small difference, their uncertainties can be propagated to produce an uncertainty estimate for $\chi(s)$ (see Ref. [4]).

Probability density of K

The probability density $p(K; s)$ or free energy $f(K; s) = -\ln p(K; s)$ can also be estimated by use of the functional $\mathcal{F}[X] \equiv h(K[X] - K')$, where $h(x)$ here denotes a positive kernel localized around $x = 0$, normalized such that $\int dx h(x) = 1$, such as a top-hat function or Gaussian:

$$p(K'; s) \approx \langle h(K[X] - K') \rangle_s \quad (37)$$

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