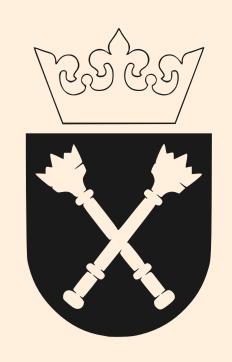
# Fluctuation-dissipation relations for thermodynamic distillation processes



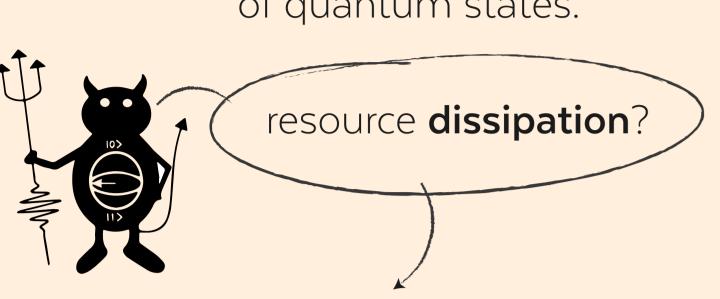
## Tanmoy Biswas<sup>1</sup>, A. de Oliveira Junior<sup>2</sup>, Michał Horodecki<sup>1</sup>, Kamil Korzekwa<sup>2</sup>

International Centre for Theory of Quantum Technologies, University of Gdansk, Wita Stwosza 63, 80-308 Gdansk, Poland<sup>1</sup> Faculty of Physics, Astronomy and Applied Computer Science, Jagiellonian University, 30-348 Krakow, Poland<sup>2</sup>

#### Background

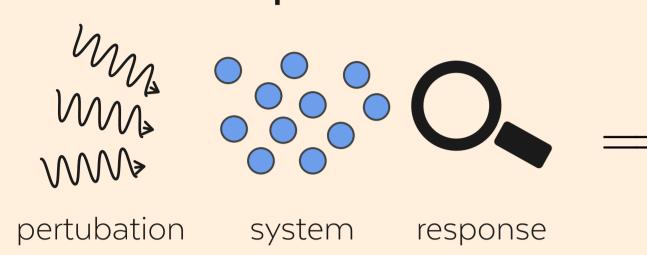
Goal: study resource dissipation and characterise optimal state transformation protocols.

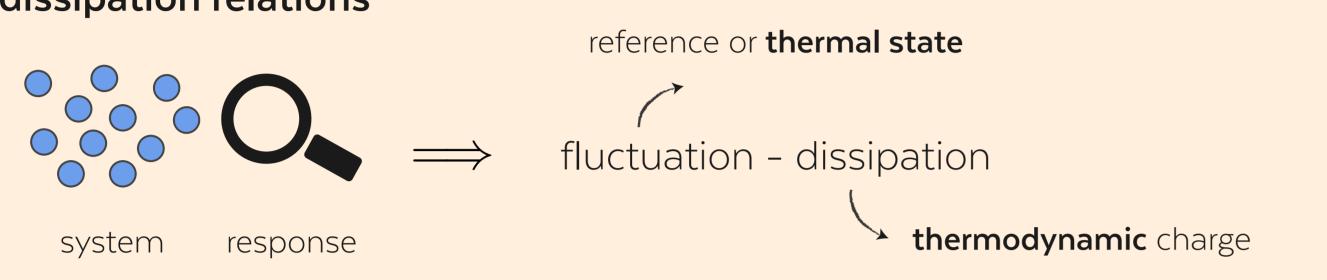
Framework: resource-theoretic ordering of quantum states.



resource theories in a nutshell Coherence Asymmetry Entanglement Non-local Coherent Asymetric Separable Symmetric

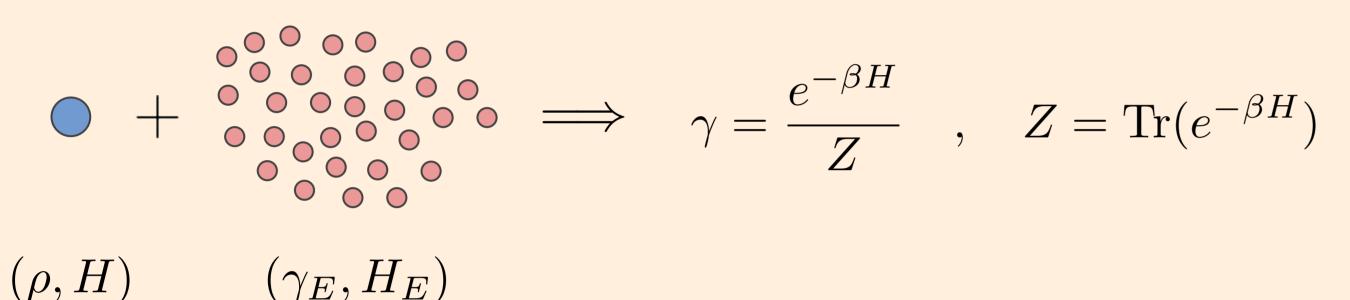
Fluctuation-dissipation relations





## Resource theory of thermodynamics

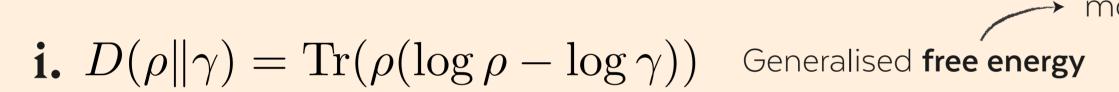
Indentifying the set of thermodynamically-free states



Thermodynamic transformations are modelled by thermal operations

$$\mathcal{E}(\rho) = \mathrm{Tr}_E(U(\rho \otimes \gamma_E)U^\dagger)$$
 with  $[U, H \otimes \mathbb{1}_E + \mathbb{1}_E \otimes H_E] = 0$ 

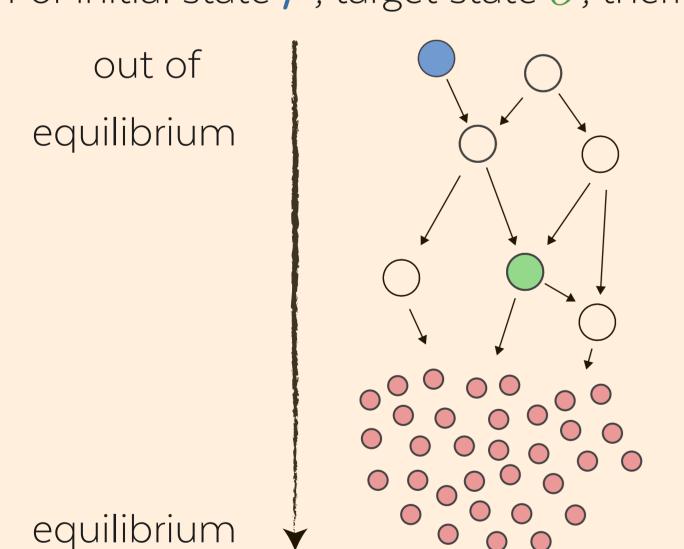
Relevant information-theoretic quantities



ii. 
$$V(\rho\|\gamma) = \operatorname{Tr}\left(\rho\left(\log\rho - \log\gamma - D(\rho\|\gamma)\right)^2\right)$$
 Fluctuations

### Thermodynamic distillation process

For initial state  $\rho$ , target state  $\sigma$ , thermal bath  $\beta \Longrightarrow \mathcal{E}(\rho) = \sigma$ .

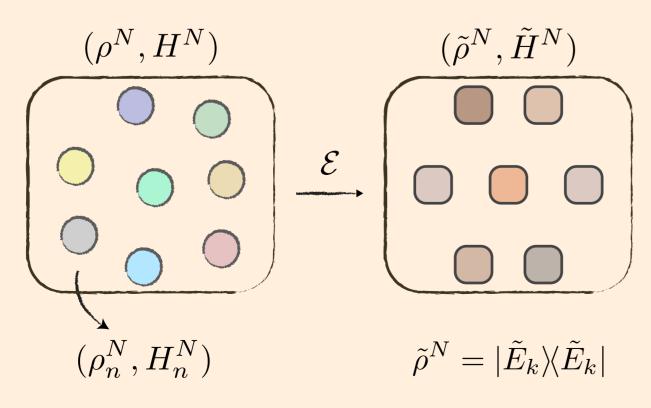


- General answer not known beyond the simplest qubit case.
  - Phys. Rev. X 5, 021001 (2015) Nat.Commun. 6, 7689 (2015)
- For energy-incoherent states. Nat.Commun. 4, 2059 (2013)

 $\mathcal{E}(\rho) = \sigma : \mathbf{p} \succ^{\beta} \mathbf{q}$ 

 $\epsilon$  -approximate **interconversion** problem:

- Approximate interconversion problem with finite system. Quantum, vol. 2, p.108, (2018)
- target state final state
- For general states.
- Dissipation of resources.
- Different subsystems.
- **Optimal** processes.
- $\epsilon$  -approximate thermodynamic distillation process:



Relevant thermo information-theoretic quantities

Free energy difference

iii. 
$$\Delta F^N := \frac{1}{\beta} \left( \sum_{n=1}^N D(\rho_n^N \| \gamma_n^N) - D(\tilde{\rho}^N \| \tilde{\gamma}^N) \right)$$
 iv.  $\sigma^2(F^N) := \frac{1}{\beta^2} \sum_{n=1}^N V(\rho_n^N \| \gamma_n^N)$ 

Free energy **fluctuations** 

#### Results

Theorem. For a distillation setting with energy incoherent or pure initial states, the transformation error of the approximate distillation process in the asymptotic limit is given by

$$\lim_{N \to \infty} \epsilon_N = \lim_{N \to \infty} \Phi \left( -\frac{\Delta F^N}{\sigma(F^N)} \right)$$

Moreover, for any N there exist an approximate distillation process with enegy incoherent initial state with the transformation error bounded by

$$\epsilon_N \le \Phi\left(-\frac{\Delta F^N}{\sigma(F^N)}\right) + \frac{C\kappa^3(F^N)}{\sigma^3(F^N)}$$

The amout of free energy dissipated in the above setting satisfies

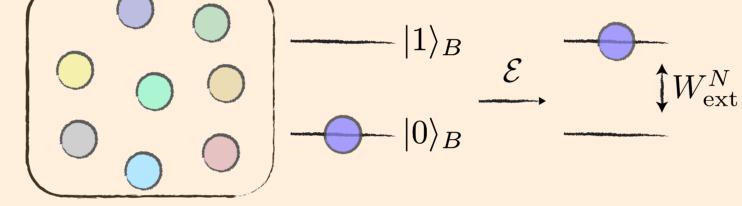
$$F_{\rm diss}^{\rm tot} = a(\epsilon) \, \sigma^{\rm tot}(F) \quad \text{with } a(\epsilon) = -\Phi^{-1}(\epsilon)(1-\epsilon) + \frac{\exp\left(\frac{-(\Phi^{-1}(\epsilon))^2}{2}\right)}{\sqrt{2\pi}}$$

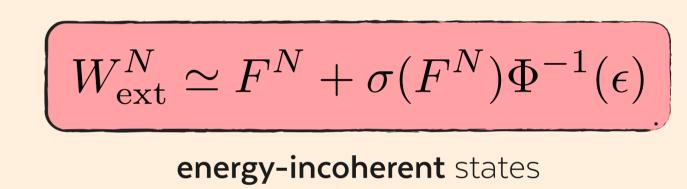
with  $\Phi^{-1}$  being the inverse of the Gaussian cumulative distribution

## Aplications

Optimal work extraction  $(\rho^N, H^N)$ 

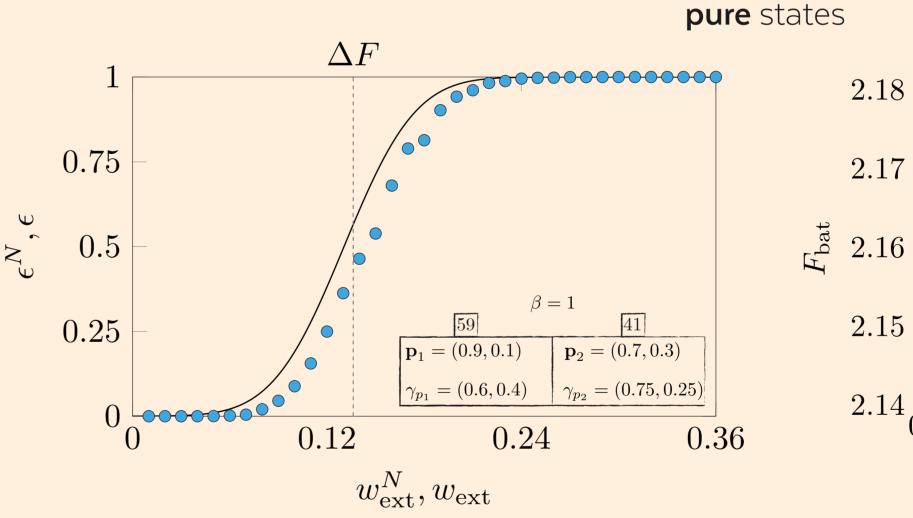
Applying the above Theorem, the extracted work for a given error is given by:

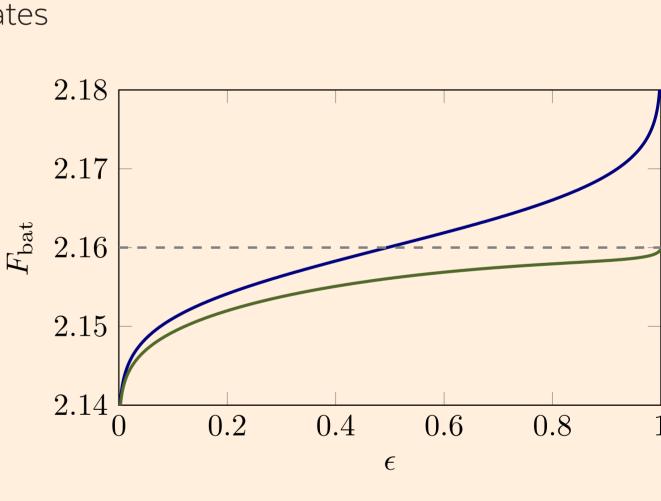




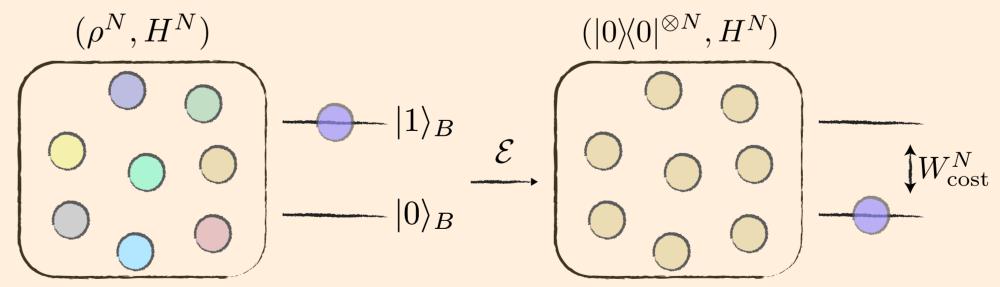
The optimal amount of work extracted from N pure quantum systems up to second-order is given by:

$$W_{\rm ext} \simeq N \left( \langle H \rangle_{\psi} + \frac{\log Z}{\beta} + \frac{\langle H^2 \rangle_{\psi} - \langle H \rangle_{\psi}^2}{\sqrt{N}} \Phi^{-1}(\epsilon) \right)$$





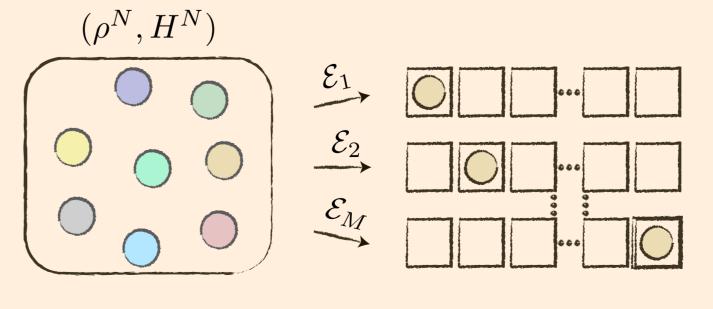
Optimal information erasure



The above Theorem also provides the second-order asymptotics for the cost of erasure:

$$W_{\text{cost}}^N \simeq \frac{S(\rho^N)}{\beta} - \sigma(F^N)\Phi^{-1}(\epsilon).$$

Optimal thermodynamically-free communication rate



The optimal number of messages that can be encoded into in a thermodynamically-free way:

I This result is valid for either a pure or incoherent state!

$$R(\rho^{\otimes N}, \epsilon_{\mathrm{d}}) \simeq D(\rho \| \gamma) + \frac{\sqrt{V(\rho \| \gamma)}}{\sqrt{N}} \Phi^{-1}(\epsilon_{\mathrm{d}}).$$

AOJ acknowledge financial support by the Foundation for Polish Science through TEAM-NET project (contract no. POIR.04.04.00- 00-17C1/18-00).

