# Derivations inspired from Lukin notes, Ran et al. [1], Perplexity and Bankim's code.

## Aditya

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## 1 Two-Level Atom Coupled to Light: Lindblad Equation Solution

## 1.1 System Definition

The Hamiltonian in the rotating wave approximation (RWA):

$$H = \frac{\hbar\Delta}{2}\sigma_z + \hbar g(\sigma_+ + \sigma_-) \tag{1}$$

Where:

- $\Delta = \omega_0 \omega_L$  is the detuning
- *g* is the coupling strength
- $\sigma_z, \sigma_+, \sigma_-$  are Pauli matrices

Lindblad operator for spontaneous emission:

$$L = \sqrt{\gamma}\sigma_{-} \tag{2}$$

This hamiltonian assumes that the states are at  $\frac{-\hbar\Delta}{2}$  and  $\frac{\hbar\Delta}{2}$ . Some references derive results by adding constant  $\frac{\hbar\Delta}{2}$  so that the ground state comes at 0 energy and the upper state at  $\hbar\Delta$ .

### 1.2 Lindblad Equation

$$\frac{d\rho}{dt} = -\frac{i}{\hbar}[H, \rho] + \gamma \left(\sigma_{-}\rho\sigma_{+} - \frac{1}{2}\{\sigma_{+}\sigma_{-}, \rho\}\right)$$
(3)

#### 1.3 **Expanded Equations**

Density matrix:

$$\rho = \begin{pmatrix} \rho_{ee} & \rho_{eg} \\ \rho_{ge} & \rho_{gg} \end{pmatrix}$$
(4)

Coupled differential equations:

$$\frac{d\rho_{ee}}{dt} = -\gamma \rho_{ee} + ig(\rho_{eg} - \rho_{ge})$$

$$\frac{d\rho_{gg}}{dt} = \gamma \rho_{ee} - ig(\rho_{eg} - \rho_{ge})$$
(6)

$$\frac{d\rho_{gg}}{dt} = \gamma \rho_{ee} - ig(\rho_{eg} - \rho_{ge}) \tag{6}$$

$$\frac{d\rho_{eg}}{dt} = -\left(\frac{\gamma}{2} + i\Delta\right)\rho_{eg} + ig(\rho_{ee} - \rho_{gg}) \tag{7}$$

$$\frac{d\rho_{ge}}{dt} = -\left(\frac{\gamma}{2} - i\Delta\right)\rho_{ge} - ig(\rho_{ee} - \rho_{gg}) \tag{8}$$

## **Steady-State Solution**

Setting time derivatives to zero:

$$0 = -\gamma \rho_{ee} + ig(\rho_{eg} - \rho_{ge}) \tag{9}$$

$$0 = \gamma \rho_{ee} - ig(\rho_{eg} - \rho_{ge}) \tag{10}$$

$$0 = -\left(\frac{\gamma}{2} + i\Delta\right)\rho_{eg} + ig(\rho_{ee} - \rho_{gg}) \tag{11}$$

$$0 = -\left(\frac{\gamma}{2} - i\Delta\right)\rho_{ge} - ig(\rho_{ee} - \rho_{gg}) \tag{12}$$

#### 1.5 Solution

Steady-state populations:

$$\rho_{ee} = \frac{g^2}{2g^2 + \frac{\gamma^2}{4} + \Delta^2} \tag{13}$$

$$\rho_{gg} = 1 - \rho_{ee} \tag{14}$$

Coherences:

$$\rho_{eg} = \rho_{ge}^* = \frac{ig(\frac{\gamma}{2} - i\Delta)}{2g^2 + \frac{\gamma^2}{4} + \Delta^2}$$
 (15)

Usually in our QNLO group we tend to call  $\gamma' = \frac{\gamma}{2}$ , so watchout for different conventions. Here in this derivation I have not used our group convention.

To learn trick to write  $\chi$  for any level refer to [1]. For Susceptibility, we work in low probe power regime, hence we work with linear susceptibility which allows us to use following equations:

$$\alpha = k \operatorname{Im}(\chi(0)) \tag{16}$$

$$n = 1 + \frac{\text{Re}(\chi(0))}{2} \tag{17}$$

Note that these formulae are only valid when the slowly varying envelope approximation holds, i.e.

$$|\operatorname{Im}[\chi(0)]|, |\operatorname{Re}[\chi(0)]| \ll 1.$$

Now lets understand how to include OD in these calculations.

$$\chi = \frac{OD \times \text{Coherence}}{\Omega_{probe}} \tag{18}$$

Transmission = 
$$e^{-kl \times \text{Im}(\chi)}$$
 (19)

$$Phase = Re(\chi) \tag{20}$$

## 1.6 Involving OD and susceptibity by using Beer Lambert's law

We start with the Beer-Lambert law:

$$I = I_0 e^{-\alpha l} \tag{21}$$

where  $\alpha$  is the absorption coefficient.

Now, let's express  $\alpha$  as:

$$\alpha = 2k \operatorname{Im}(n) \tag{22}$$

The refractive index n can be written as:

$$n = \sqrt{1 + \chi} \tag{23}$$

For weak  $\chi$  during Electromagnetically Induced Transparency (EIT), we approximate:

$$n \approx 1 + \frac{\chi}{2} \tag{24}$$

Thus, the imaginary part of n is:

$$Im(n) = \frac{Im(\chi)}{2} \tag{25}$$

Substituting this into the expression for  $\alpha$ , we get:

$$\alpha = k \operatorname{Im}(\chi) \tag{26}$$

Therefore, the intensity I can be written as:

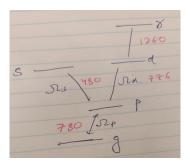


Figure 1: Five level system

$$I = I_0 e^{-kl\operatorname{Im}(\chi)} \tag{27}$$

So transmission can be plotted from here and phase is just the real part of  $\chi$ . The kl factor is normalized by putting  $\frac{1}{kl}$  in  $\chi$ . For example consider the  $\chi$  example from [1]

The expression for  $\chi_p$  is given by:

$$\chi_p = \frac{\alpha_0 \gamma}{\gamma - i\Delta_p + \frac{|\Omega_c|^2}{\gamma_{21} - i\Delta_{21}}} \tag{28}$$

, after comparing from other resources including Bankim's code and Lee's thesis I found that

$$\alpha_0 = \frac{i \, OD}{KL} \tag{29}$$

For five level system as shown in Fig.1, the analytical expression after following the trick in [1] for  $\chi$  is

$$\chi = \frac{i \cdot od \cdot \Gamma}{\Gamma - i(\Delta_p) + \frac{\Omega_s^2}{\gamma_{gs} - i(\Delta_p + \Delta_s)} + \frac{\Omega_d^2}{\gamma_{gd} - i(\Delta_p + \Delta_d) + \frac{\Omega_r^2}{\gamma_{gr} - i(\Delta_p + \Delta_d + \Delta_r)}}} +$$
(30)

We normalize medium length of medium by considering kl = 1

```
"""Calculate coherence using QuTiP and analytical methods."""
   # QuTiP calculation
   H = create_hamiltonian(Delta, g)
   c_ops = [np.sqrt(gamma) * qt.sigmam()]
   rho_ss = qt.steadystate(H, c_ops)
   rho_eg_qutip = rho_ss[0, 1]
   # Analytical calculation
   \label{eq:rho_eg_analytical} \mbox{ = -1j * g * ((gamma / 2) - 1j * Delta) / (2 * )}
       \hookrightarrow g**2 + (gamma**2) / 4 + Delta**2)
   return rho_eg_qutip, rho_eg_analytical
def plot_coherences(g, gamma, OD):
   # System parameters
   Delta_values = np.linspace(-5, 5, 100)
   # Arrays to store coherence values
   coherence_qutip = []
   coherence_analytical = []
   # Calculate coherences for each Delta
   for Delta in Delta_values:
       qutip_result , analytical_result = calculate_coherence(Delta
           \hookrightarrow , g, gamma)
       coherence_qutip.append(qutip_result)
       coherence_analytical.append(analytical_result)
   # Convert lists to numpy arrays for plotting
   coherence_qutip = np.array(coherence_qutip)
   coherence_analytical = np.array(coherence_analytical)
   chi_analytical = OD*coherence_analytical/g
   chi_numerical = OD*coherence_qutip/g
   {\tt transmission\_analytical = np.exp(-np.real(1j*chi\_analytical))}
   transmission_numerical = np.exp(-np.real(1j*chi_numerical))
   phase_analytical = np.real(chi_analytical)
   phase_numerical = np.imag(1j*chi_numerical)
   fig, axes = plt.subplots(1, 3, figsize=(15, 5))
   # Plot on each subplot
   axes[0].plot(Delta_values, transmission_analytical, label='
       \hookrightarrow Analytical transmission', color='blue')
    axes[0].plot(Delta_values, transmission_numerical, '--', label=
       axes[0].set_title("Transmission")
   axes[0].set_xlabel("Detuning Delta")
   axes[0].set_ylabel("Transmission")
   axes[0].legend()
   axes[1].plot(Delta_values, phase_analytical, label='Analytical
      axes[1].plot(Delta_values, phase_numerical, '--', label='QuTiP
       ⇔ Phase', color='red')
    axes[1].set_title("Phase")
```

```
axes[1].set_xlabel("Detuning Delta")
   axes[1].set_ylabel("Phase")
   axes[1].legend()
   axes[2].plot(Delta_values, abs(coherence_analytical), label='
       axes[2].plot(Delta_values, abs(coherence_qutip), '--', label='
      axes[2].set_title(r"coherence rho_12")
   axes[2].set_xlabel("Detuning Delta")
   axes[2].set_ylabel("coherence rho_12")
   axes[2].legend()
   plt.suptitle("Always consider weak probe limit that is very low
       \hookrightarrow g", fontsize=16)
   plt.tight_layout()
   plt.show()
# Create sliders
g_slider = FloatSlider(value=0.1, min=0.001, max=.5, step=0.1,
   → description='g:')
gamma_slider = FloatSlider(value=0.1, min=0.05, max=1.0, step=0.01,

    description = 'gamma:')

OD = FloatSlider(value=50, min=0, max=150, step=5, description='OD:
   → ')
# Create interactive plot
interactive_plot = interactive(plot_coherences, g=g_slider, gamma=
   → gamma_slider, OD=OD)
# Display the interactive plot
display(interactive_plot)
```

## References

[1] Ran Finkelstein and Ofer Firstenberg. "A practical guide to electromagnetically induced transparency in atomic vapor". In: *New Journal of Physics* 25.6 (2023), p. 063016. DOI: 10.1088/1367-2630/acbc40.