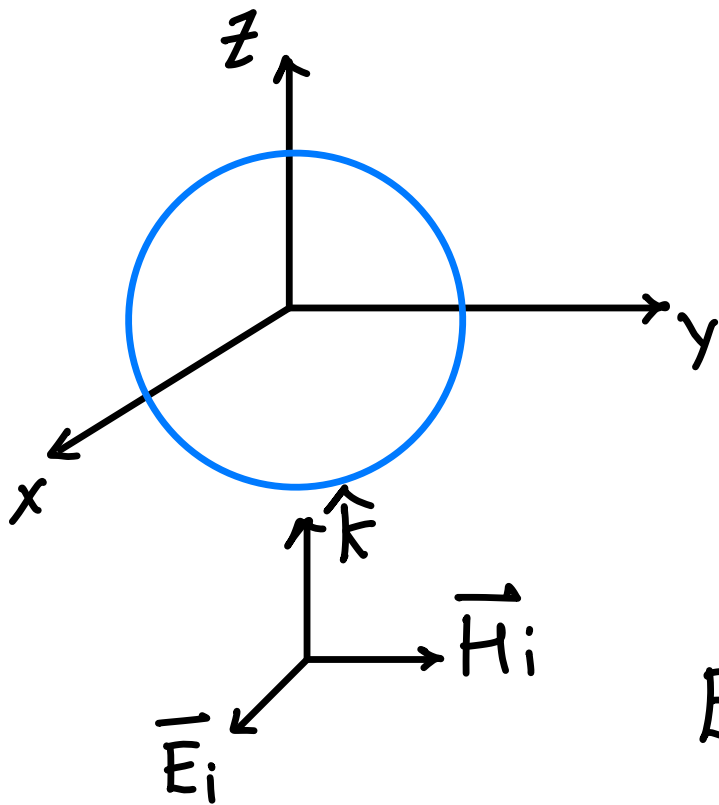




Plane Wave Scattering from Spheres.



$$\vec{E}^i = E_0 \hat{x} e^{ikz}$$

$$\vec{H}^i = \frac{E_0}{\eta} \hat{y} e^{ikz}$$

$$\vec{E}^i = E_r \hat{r} + E_\theta \hat{\theta} + E_\phi \hat{\phi}$$

$$E_r^i = \vec{E}^i \cdot \hat{r}$$

$$= E_0 \cos\phi \sin\theta e^{ikr \cos\theta}$$

$$= E_0 \frac{\cos\phi}{-ikr} \frac{d}{d\theta} e^{ikr \cos\theta}$$

Expanding plane wave in sph. modes.

$$e^{ikr \cos\theta} = \sum_{n=1}^{\infty} (i)^n (2n+1) j_n(kr) P_n(\cos\theta).$$

(Sec 8-2.1)

$$E_r^i = \frac{E_0 \cos\phi}{-i(kr)^2} \sum_{n=1}^{\infty} (i)^n (2n+1) \hat{j}_n(kr) \underbrace{P_n'(\cos\theta)}_{\substack{\text{red arrow} \\ kr j_n(kr) \frac{d}{d\theta} P_n^0 = P_n'}}$$

$$E_r = -\frac{1}{i\omega\mu\epsilon} \left(\frac{\partial^2}{\partial r^2} + k^2 \right) A_r \quad \text{Spherical Gauge.}$$

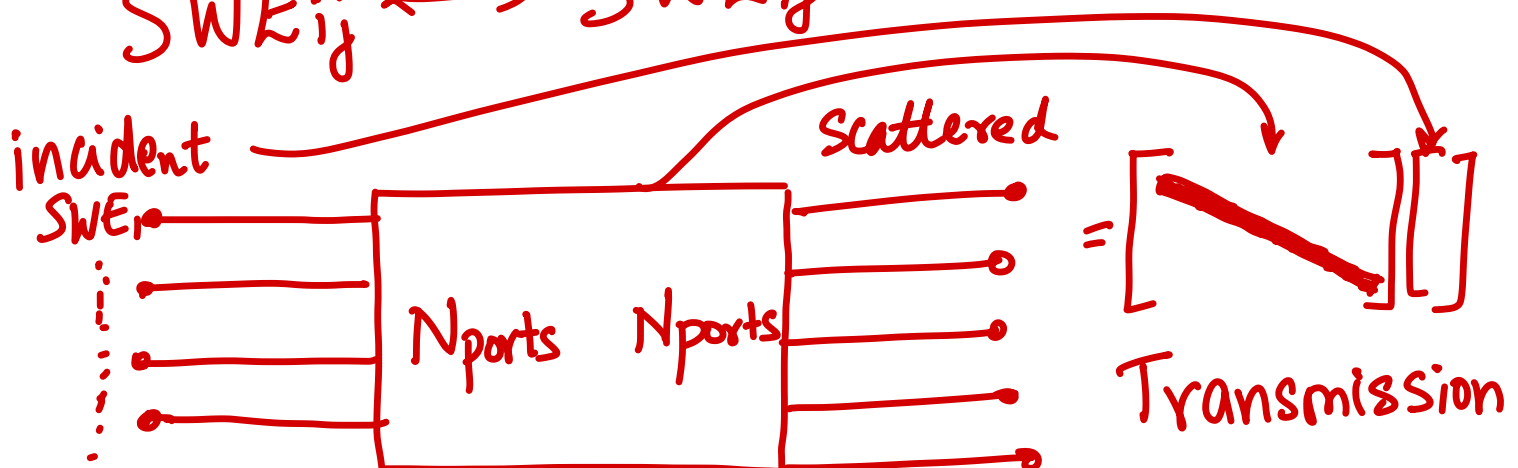
$$A_r^i = \frac{E_0}{\omega} \cos\phi \sum_{n=1}^{\infty} \frac{(i)^n (2n+1)}{n(n+1)} \hat{j}_n(kr) P_n'(\cos\theta).$$

$$\frac{d^2}{dr^2} \hat{j}_n(kr) + k^2 \hat{j}_n(kr) = \frac{n(n+1)}{r^2} \hat{j}_n(kr)$$

$$A_{mr}^i = \frac{E_0}{\omega\eta} \sin\phi \sum_{n=1}^{\infty} \frac{(i)^n (2n+1)}{n(n+1)} \underbrace{\hat{j}_n(kr) P_n'(\cos\theta)}_{SWE_{ij}}$$

> SWF are orthogonal!

$$SWE_{ij} \leftrightarrow SWE_{ij}$$



$$A_r^s = \frac{E_0}{\omega} \cos \phi \sum_{n=1}^{\infty} a_n \hat{h}_n^{(1)'}(kr) P_n'(\cos \theta)$$

$$A_{m\phi}^s = \frac{E_0}{\omega \eta} \sin \phi \sum_{n=1}^{\infty} b_n \hat{h}_n^{(1)}(kr) P_n'(\cos \theta)$$

We apply BC to find a_n, b_n .

$$E_\theta^i + E_\theta^s \Big|_{r=a} = 0 ; E_\phi^i + E_\phi^s \Big|_{r=a} = 0 .$$

$$a_n = - \frac{(i)^n (2n+1)}{n(n+1)} \frac{\hat{j}_n'(ka)}{\hat{h}_n^{(1)'}(ka)}$$

$$b_n = - \frac{(i)^n (2n+1)}{n(n+1)} \frac{\hat{j}_n(ka)}{\hat{h}_n^{(1)}(ka)}$$

\vec{E}^s, \vec{H}^s can be derived from $A_r^s, A_{m\phi}^s$
(Eq: 8.139 Pg 549)

Far-Field

$$\lim_{kr \rightarrow \infty} \hat{h}_n^{(1)'}(kr) \approx (-i)^n e^{ikr}$$

$$E_{\theta}^s = -i E_0 \cos \phi \frac{e^{ikr}}{kr} \sum_{n=1}^{\infty} (-i)^n \left(a_n \sin \theta P_n'(\cos \theta) - b_n \frac{P_n'(\cos \theta)}{\sin \theta} \right)$$

$$E_{\phi}^s = -i E_0 \sin \phi \frac{e^{ikr}}{kr} \sum_{n=1}^{\infty} (-i)^n \left(a_n \frac{P_n'(\cos \theta)}{\sin \theta} - b_n \sin \theta P_n''(\cos \theta) \right)$$

Radar Cross Section (RCS)

$$\sigma_c(\theta, \phi) = \lim_{r \rightarrow \infty} 4\pi r^2 \frac{|E_{\theta}^s|^2}{|E_0|^2}$$

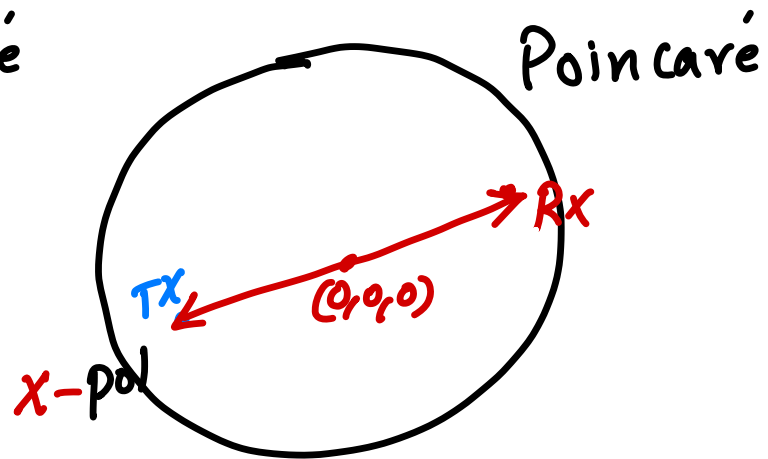
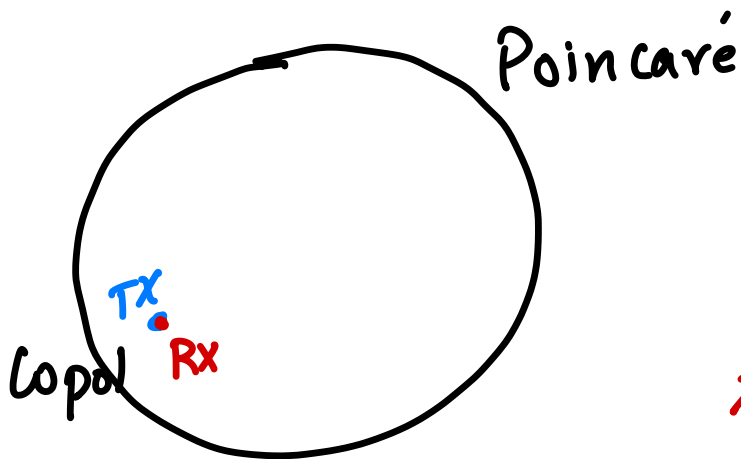
monostatic RCS : $T_x < R_x <$

bistatic RCS : $T_x <$

R_x



$$\sigma(\theta_T, \phi_T; \theta_R, \phi_R)$$

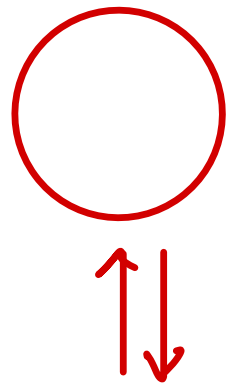


$$\sigma_c(\theta, \phi) = \frac{\lambda^2 \cos^2 \phi}{\pi} \left| \sum_{n=1}^{\infty} (-i)^n \left(a_n \sin \theta P_n'(\cos \theta) - b_n \frac{P_n'(\cos \theta)}{\sin \theta} \right) \right|^2$$

$$\sigma_x(\theta, \phi) = \dots$$

$$\theta = \pi, \quad \omega\text{-pol} : \phi = 0$$

$$\phi = \pi/2$$



$$\sigma_c = \frac{\lambda^2}{4\pi} \left| \sum_{n=1}^{\infty} (i)^n n(n+1) (a_n - b_n) \right|^2$$

$$\sigma_x = 0$$

Low-freq approx.

$$a \ll \lambda ; a < 0.1 \lambda$$

$$\sigma_c^{LF} = \frac{9\lambda^2}{4\pi} (ka)^6$$

$$\sigma_c^{LF} \propto \frac{1}{\lambda^4}$$

High freq approx

$a \gg \lambda \Rightarrow$ Physical Optics (PO)

$$\sigma_c^{PO} = \pi a^2$$

