

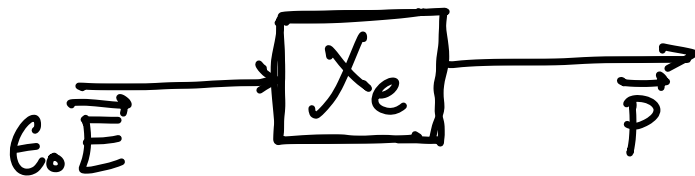


# EM03 - Dielectric Models

## Material Dispersion

$$\vec{P} = \epsilon_0 \chi_e \vec{E}$$

A better model of  $\chi_e$  is to say the response is that of a linear system.



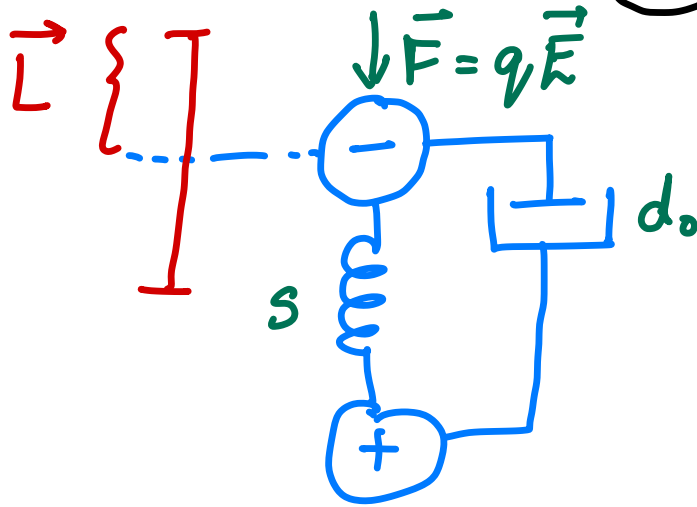
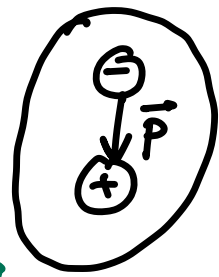
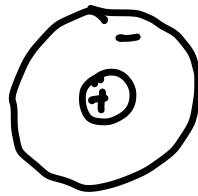
$$\vec{P} = \chi_e * \epsilon_0 \vec{E}$$

$$\vec{P}(t) = \epsilon_0 \int_{-\infty}^t \chi_e(t-\tau) \vec{E}(\tau) d\tau$$

$$\vec{P}(\omega) = \epsilon_0 \underbrace{\chi_e(\omega)}_? \vec{E}(\omega)$$

$$\vec{M}(\omega) = \mu_0 \chi_m(\omega) \vec{H}(\omega)$$

# Lorentz Model



$$m \frac{d^2 \vec{L}}{dt^2} + d_0 \frac{d\vec{L}}{dt} + s\vec{L} = q \vec{E}$$

$$\vec{E}(t) = \text{Re} \{ \tilde{E} e^{-i\omega t} \}$$

$$\vec{L}(t) = \text{Re} \{ \tilde{L} e^{-i\omega t} \}$$

$$\tilde{L} = \frac{\left(\frac{q}{m}\right) \tilde{E}}{\omega_0^2 - \omega^2 - i\gamma\omega}$$

resonant freq  
 $\omega_0^2 = \frac{s}{m}$

damping factor  
 $\gamma = \frac{d_0}{m}$

$$\vec{P} = Nq \vec{L} = \underbrace{\frac{\omega_p^2}{\omega_0^2 - \omega^2 - i\gamma\omega}}_{\chi_e} \epsilon_0 \vec{E}$$

plasma freq.

$$\omega_p^2 = \frac{Nq^2}{m\epsilon_0}$$

$$\chi_e = \frac{\omega_p^2}{\omega_0^2 - \omega^2 - i\gamma\omega}$$

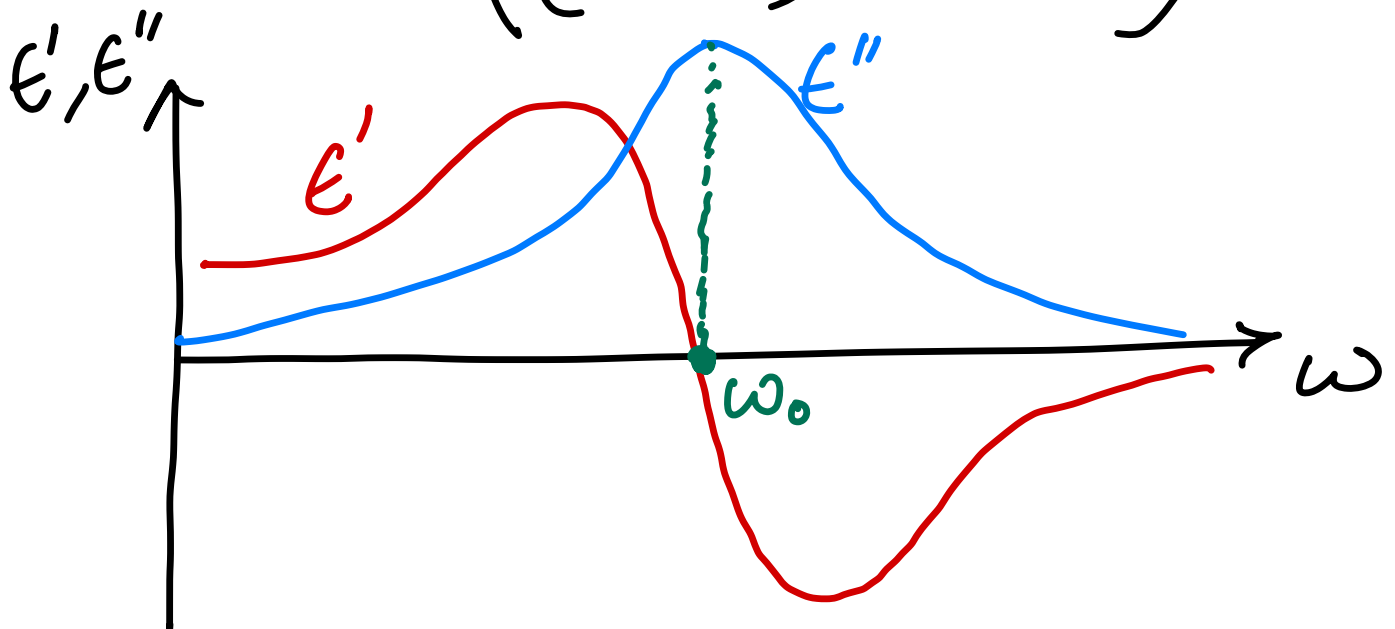
$$\epsilon = \epsilon_0 (1 + \chi_e)$$

$$= \epsilon_0 \left( 1 + \frac{\omega_p^2}{\omega_0^2 - \omega^2 - i\gamma\omega} \right)$$

$$\epsilon = \epsilon' + i\epsilon''$$

$$\epsilon' = \epsilon_0 \left( 1 + \frac{\omega_p^2 (\omega_0^2 - \omega^2)}{(\omega_0^2 - \omega^2)^2 + \gamma^2 \omega^2} \right)$$

$$\epsilon'' = \epsilon_0 \left( \frac{\omega_p^2 \gamma \omega}{(\omega_0^2 - \omega^2)^2 + \gamma^2 \omega^2} \right)$$



## Approximations

$$\underline{\omega \ll \omega_0}$$

$$\epsilon' \approx \epsilon_0 \left( 1 + \frac{\omega_p^2}{\omega^2} \right) \rightarrow \text{independent of freq.}$$

$$\epsilon'' \approx \epsilon_0 \left( \frac{\gamma \omega_p^2 \omega}{\omega_0^4} \right) \rightarrow \text{linear with freq.}$$

$$\underline{\text{Loss tangent: } \frac{\epsilon''}{\epsilon'}}$$

$$\underline{\omega \gg \omega_0}$$

$$\epsilon' \approx \epsilon_0 \left( 1 - \frac{\omega_p^2}{\omega^2} \right) \quad \epsilon'' \approx \epsilon_0 \left( \frac{\gamma \omega_p^2}{\omega^3} \right)$$

$$\text{As } \omega \rightarrow \infty, \quad \epsilon' = \epsilon_0 \quad ; \quad \epsilon'' = 0$$

$\Rightarrow$  Material disappears!

"Ultraviolet transparency."

# Conducting Media (Drude Model)

$$\vec{J} = Nq \frac{d\vec{L}}{dt}$$

$$= Nq \frac{d}{dt} \left( \frac{q\vec{E}}{\omega_0^2 - \omega^2 - i\omega\gamma} \cdot \frac{1}{m} \right)$$

$$\vec{J} = \frac{-i\omega q^2 N/m}{\cancel{\omega_0^2} - \omega^2 - i\omega\gamma} \vec{E}$$

$$\vec{J} = \frac{Nq^2}{\underbrace{m(\gamma - i\omega)}_{\sigma}} \vec{E}$$

At  $\omega \ll \gamma$

$$\vec{J} \approx \underbrace{\frac{Nq^2}{m\gamma}}_{\sigma} \vec{E}$$

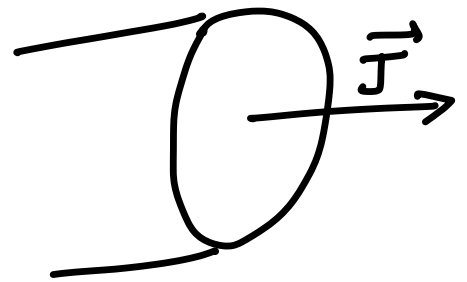
$$\Rightarrow \boxed{\vec{J} = \sigma \vec{E}}$$

Ohm's Law.

$\rho = \frac{1}{\sigma} \rightarrow$  resistivity

$$\int \vec{E} \cdot d\vec{l} = \int \rho \vec{J} \cdot d\vec{l}$$

$$V = \rho \int \frac{I}{A} \hat{n} \cdot d\vec{l} \hat{n} = \frac{\rho I}{A} L = \underline{\underline{IR}}$$



$$S = 0$$

$$\Rightarrow \omega_0 = 0$$

Drude model.

Conductivity

## Relationship b/w $\sigma$ & $\epsilon$

$$\epsilon = \epsilon_0 \left( 1 - \frac{\omega_p^2}{\omega^2 + i\gamma\omega} \right) \quad \text{when } \omega_0 \rightarrow 0$$

$$\Rightarrow \epsilon_r' = 1 - \frac{\omega_p^2}{\omega^2 + \gamma^2} \quad \epsilon_r'' = \frac{\omega_p^2 \gamma}{\omega(\omega^2 + \gamma^2)}$$

When  $\omega \ll \gamma$

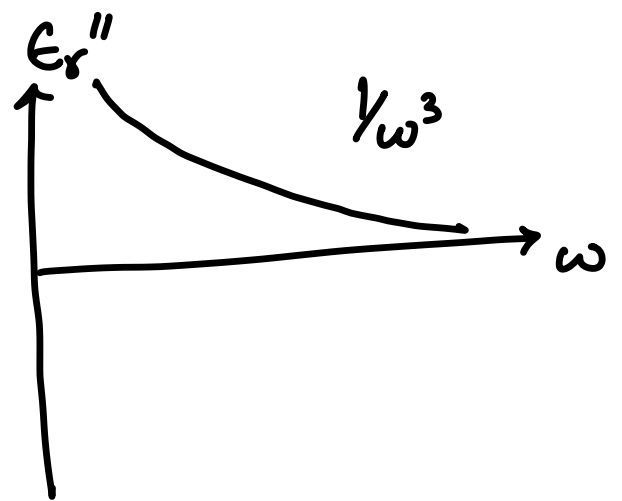
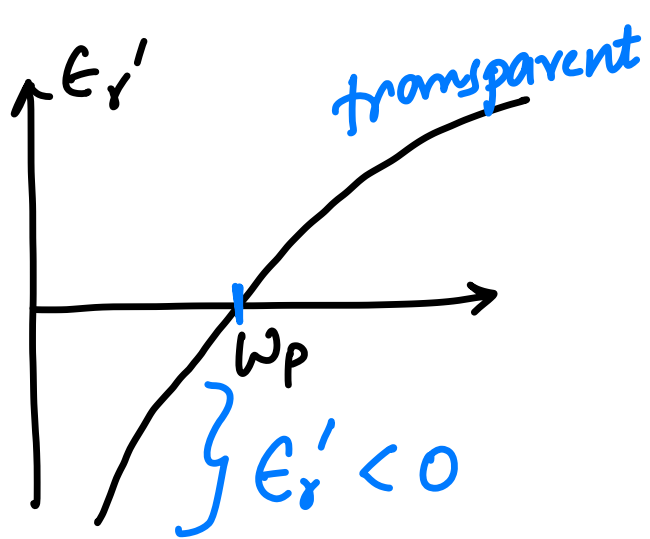
$$\epsilon_r'' = \frac{Nq^2}{m\epsilon_0\omega\gamma} = \frac{\sigma}{\omega\epsilon_0}$$

$$\boxed{\epsilon_r'' = \frac{\sigma}{\omega\epsilon_0}}$$

$$\boxed{\epsilon'' = \frac{\sigma}{\omega}}$$

$\omega \gg \gamma$

$$\epsilon_r' = 1 - \frac{\omega_p^2}{\omega^2} \quad \epsilon_r'' = \frac{\omega_p^2 \gamma}{\omega^3}$$



$\Rightarrow$  reflection.

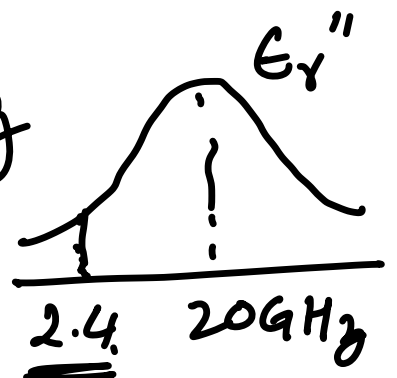
$\omega_p$  for most metals  $\approx 10^{15} - 10^{16}$  Hz.  
 $\Rightarrow$  high UV

$\Rightarrow$  Al & Silver are reflective.

Cu & Gold?  $\rightarrow$  They have resonances around blue-green region of the spectrum.

$\Rightarrow$  yellow-red are reflective.

$>$  Water has a resonance at  $\approx 20$  GHz.  
 Microwave oven  $\approx 2.4$  GHz





>  $O_2 \rightarrow \underline{60\text{ GHz}}$ . {Secure comm.}

> Radiative cooling. 

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$$\vec{J} = \sigma \vec{E} \quad \text{PEC } \sigma = \infty \Rightarrow \rho = 0$$
$$\Rightarrow \vec{E} = \vec{J} \rho = 0.$$

$$\nabla \cdot \vec{J} = -\frac{\partial \rho}{\partial t} = \nabla \cdot \sigma \vec{E} = \sigma \nabla \cdot \vec{E} = \frac{\sigma \rho}{\epsilon}$$

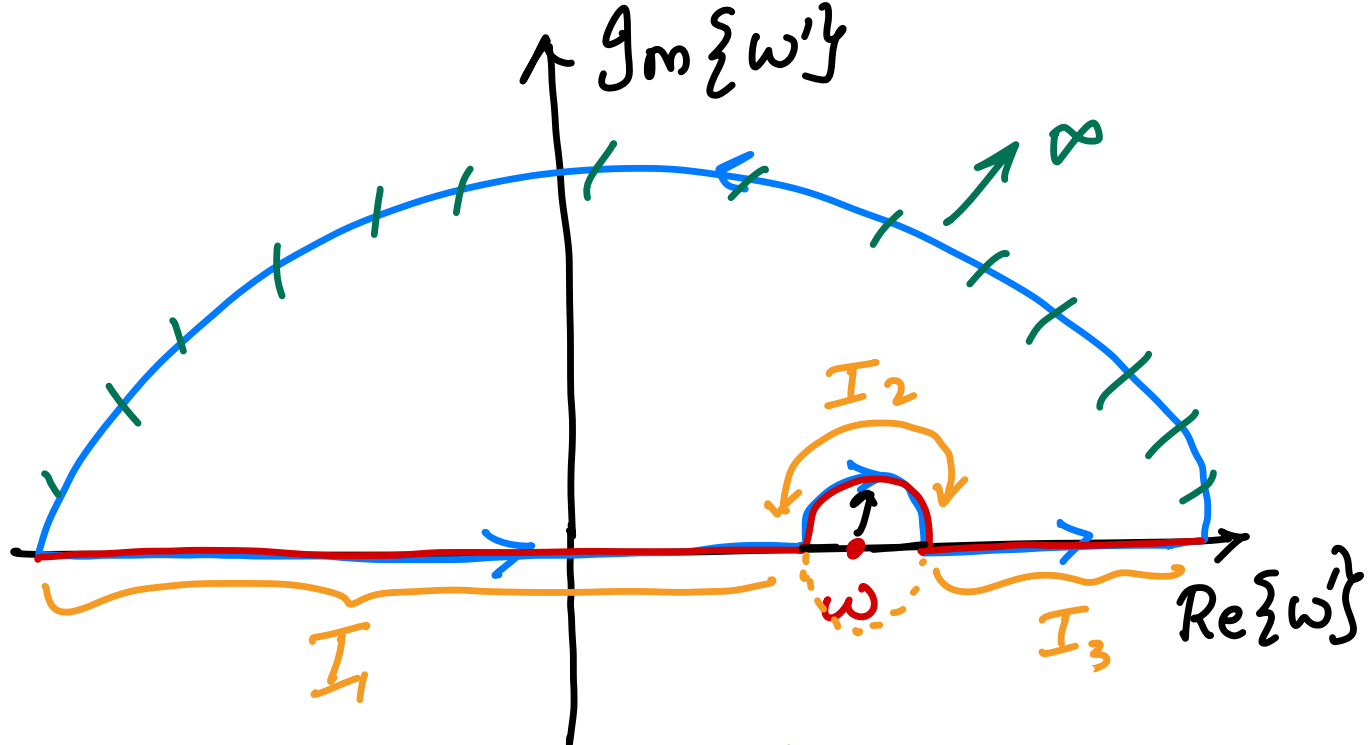
$$\Rightarrow \frac{\partial \rho}{\partial t} + \frac{\sigma \rho}{\epsilon} = 0 \Rightarrow \boxed{\rho(t) = \rho_0 e^{-\frac{\sigma}{\epsilon} t}}.$$

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## Kramers Kronig Relations

> Causality, Stability, linearity.

" $\chi_e(\omega)$  is "analytic" in the upper half complex  $\omega$  plane."



$$\oint_C \frac{\chi_e(w')}{w' - w} dw' = 0$$

$$\int_C \frac{\chi_e(w')}{w' - w} dw' = 0$$

$$I_1 + I_2 + I_3 = 0$$

$$\Rightarrow I_1 + I_3 = -I_2$$

$$\int_{-\infty}^{\infty} \frac{\chi_e(w')}{w' - w} dw' =$$

$$I_2 = \frac{1}{2} \{ \text{Residue} \}$$

$$\int_{-\infty}^{\infty} \frac{\chi_e(w')}{w' - w} dw' = \underbrace{i\pi \chi_e(w)}_{\text{Residue}/2}$$

$$\chi_e'(\omega) = \frac{1}{\pi} \int_{-\infty}^{\infty} \frac{\chi_e''(\omega')}{\omega' - \omega} d\omega'$$

$$\chi_e''(\omega) = -\frac{1}{\pi} \int_{-\infty}^{\infty} \frac{\chi_e'(\omega')}{\omega' - \omega} d\omega' \quad \leftarrow \text{Hilbert.}$$

▶  $\epsilon' \rightarrow$  propagation of phase }  
 $\epsilon'' \rightarrow$  loss or absorption. }

> Phase retrieval!

> KK retrieval!

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