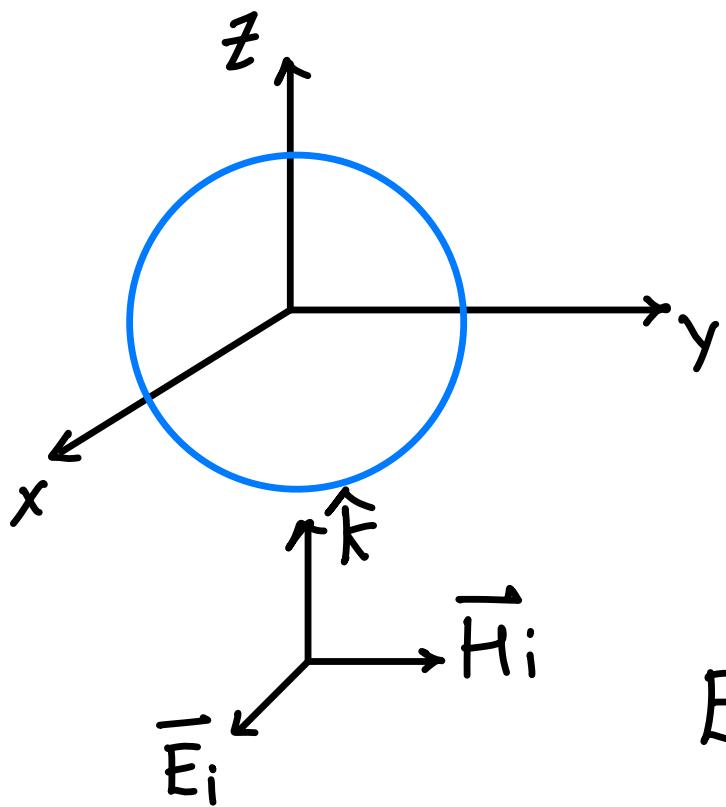




Plane Wave Scattering from Spheres.



$$\vec{E}^i = E_0 \hat{x} e^{ikz}$$

$$\vec{H}^i = \frac{E_0}{\eta} \hat{y} e^{ikz}$$

$$\vec{E}^i = E_r \hat{r} + E_\theta \hat{\theta} + E_\phi \hat{\phi}$$

$$E_r^i = \vec{E}^i \cdot \hat{r}$$

$$= E_0 \cos\phi \sin\theta e^{ikr \cos\theta}$$

$$= E_0 \frac{\cos\phi}{-ikr} \frac{d}{d\theta} e^{ikr \cos\theta}$$

Expanding plane wave in sph. modes.

$$e^{ikr \cos\theta} = \sum_{n=1}^{\infty} [i]^n (2n+1) j_n(kr) P_n(\cos\theta).$$

(Sec 8-2-1)

$$E_r^i = \frac{E_0 \cos\phi}{-i(kr)^2} \sum_{n=1}^{\infty} [i]^n (2n+1) \underbrace{j_n(kr)}_{kr j_n(kr)} \underbrace{P_n'(\cos\theta)}_{\frac{d}{d\theta} P_n^0 = P_n'}$$

$$kr j_n(kr) \frac{d}{d\theta} P_n^0 = P_n'$$

$$E_r = -\frac{1}{i\omega\mu\epsilon} \left(\frac{\partial^2}{\partial r^2} + k^2 \right) A_r \xrightarrow{\text{Spherical Gauge.}}$$

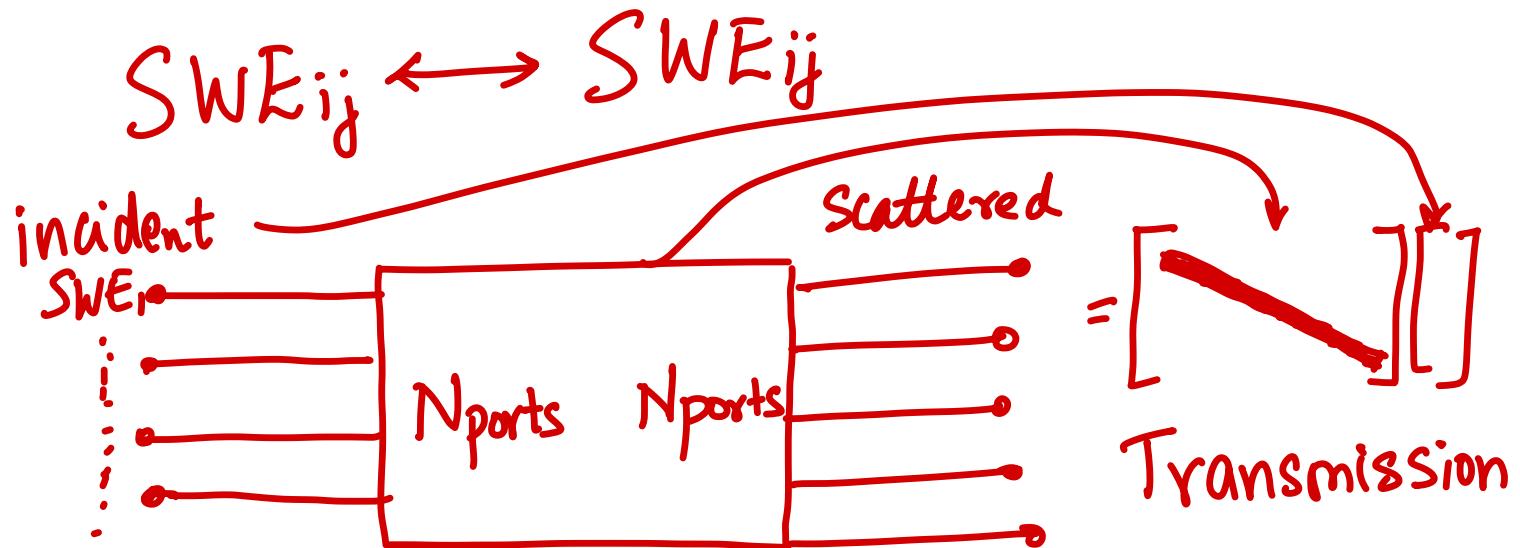
$$A_r^i = \frac{E_0}{\omega} \cos\phi \sum_{n=1}^{\infty} \frac{(i)^n (2n+1)}{n(n+1)} \hat{j}_n(kr) P_n'(\cos\theta).$$

$$\frac{d^2}{dr^2} \hat{j}_n(kr) + k^2 \hat{j}_n(kr) = \frac{n(n+1)}{r^2} \hat{j}_n(kr)$$

$$A_{mr}^i = \frac{E_0}{\omega\eta} \sin\phi \sum_{n=1}^{\infty} \frac{(i)^n (2n+1)}{n(n+1)} \hat{j}_n(kr) P_n'(\cos\theta).$$

SWE_{ij}

> SWF are Orthogonal !



$$A_r^s = \frac{E_0}{\omega} \cos \phi \sum_{n=1}^{\infty} a_n \hat{h}_n^{(0)}(kr) P_n'(\cos \theta)$$

$$A_{mr}^s = \frac{E_0}{\omega \eta} \sin \phi \sum_{n=1}^{\infty} b_n \hat{h}_n^{(1)}(kr) P_n'(\cos \theta)$$

We apply BC to find a_n, b_n .

$$\left. E_\theta^i + E_\theta^s \right|_{r=a} = 0 ; \left. E_\phi^i + E_\phi^s \right|_{r=a} = 0 .$$

$$a_n = - \frac{(i)^n (2n+1)}{n(n+1)} \frac{\hat{j}_n'(ka)}{\hat{h}_n^{(1)'}(ka)}$$

$$b_n = - \frac{(i)^n (2n+1)}{n(n+1)} \frac{\hat{j}_n(ka)}{\hat{h}_n^{(1)}(ka)}$$

\vec{E}^s, \vec{H}^s can be derived from A_r^s, A_{mr}^s
 (Eq: 8.139 Pg 549)

Far-Field

$$\lim_{kr \rightarrow \infty} \hat{h}_n^{(1)'}(kr) \simeq (-i)^n e^{ikr}$$

$$E_\theta^s = -i E_0 \cos \phi \frac{e^{ikr}}{kr} \sum_{n=1}^{\infty} (-i)^n \left(a_n \sin \theta P_n'(\cos \theta) - b_n \frac{P_n'(\cos \theta)}{\sin \theta} \right)$$

$$E_\phi^s = -i E_0 \sin \phi \frac{e^{ikr}}{kr} \sum_{n=1}^{\infty} (-i)^n \left(a_n \frac{P_n'(\cos \theta)}{\sin \theta} - b_n \sin \theta P_n''(\cos \theta) \right)$$

Radar cross section (RCS)

$$\sigma_c(\theta, \phi) = \lim_{r \rightarrow \infty} 4\pi r^2 \frac{|E_\theta^s|^2}{|E_0|^2}$$

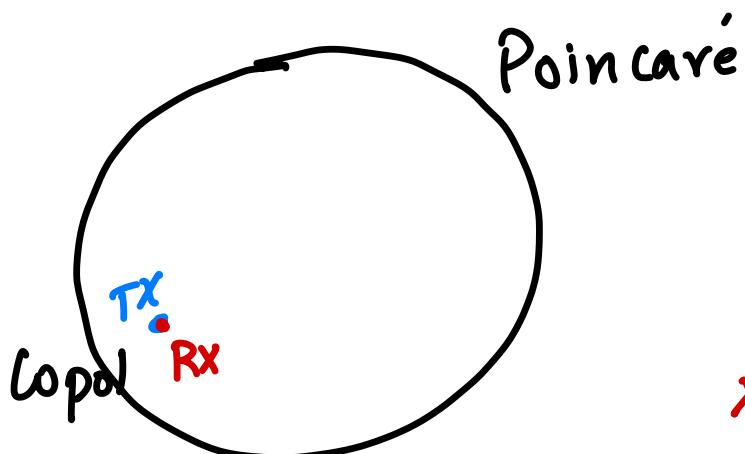
monostatic RCS : $\frac{Tx}{Rx} <$

bistatic RCS : $Tx <$



Y_{Rx}

$$\mathcal{T}(\theta_T, \phi_T; \theta_R, \phi_R)$$

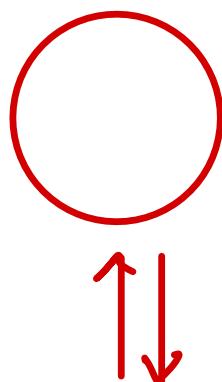


$$\mathcal{T}_c(\theta, \phi) = \frac{\lambda^2 \cos^2 \phi}{\pi} \left| \sum_{n=1}^{\infty} (-i)^n \left(a_n \sin \theta P_n''(\cos \theta) - b_n \frac{P_n'(\cos \theta)}{\sin \theta} \right) \right|^2$$

$$\mathcal{T}_x(\theta, \phi) = \dots$$

$$\theta = \pi, \text{ co-pol} : \phi = 0$$

$$\phi = \pi/2$$



$$\mathcal{T}_c = \frac{\lambda^2}{4\pi} \left| \sum_{n=1}^{\infty} (i)^n n(n+1) (a_n - b_n) \right|^2$$

$$\mathcal{T}_x = 0$$

Low-freq approx.

$$a \ll \lambda ; a < 0.1\lambda$$

$$\Gamma_c^{LF} = \frac{9\lambda^2}{4\pi} (ka)^6$$

$$\boxed{\Gamma_c^{LF} \propto \frac{1}{\lambda^4}}$$

High freq approx

$a \gg \lambda \Rightarrow$ Physical Optics (PO)

$$\Gamma_c^{PO} = \pi a^2$$

