

Commissioning and First Data Analysis of the Mainz Radius Experiment

Adriano Del Vincio 562946

Relatori: prof. Francesco Forti, Prof. Concettina Sfienti

Università di Pisa

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UNIVERSITÀ DI PISA

The Mainz Radius Experiment

MREX

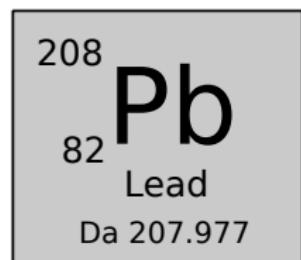
The Mainz Radius Experiment is an experimental campaign at the nuclear physics institute of Mainz, with the aim of investigating the properties of nuclear matter with imbalance in the number of protons and neutrons.

Objective

Determination of the neutron spacial density for ^{208}Pb nucleus, through the elastic electron scattering. From an accurate determination of the neutron spacial distribution, the *Neutron Skin Thickness* of ^{208}Pb is measured.

The neutron skin thickness is defined as the difference between rms radius of neutron and proton spacial distributions:

$$\delta r_{np} = \sqrt{\langle r_n^2 \rangle} - \sqrt{\langle r_p^2 \rangle}, \quad (1)$$



MREX and Neutron Skin Thickness

The results of the experiment will be valuable to constrain the Equation of State (**EOS**) of nuclear matter. The equation of state plays an important role for the description of the **neutron star structure**, specifically the determination of the radius.

In neutron rich nuclei, the spacial distribution of neutrons is more extended than proton spacial distribution. Theoretical models link the Neutron skin thickness of heavy nuclei, such as ^{208}Pb , with the **slope of the symmetry energy L**.

The neutron skin thickness is given by the competition of two terms of the nuclear binding energy: the surface term and the symmetry energy.

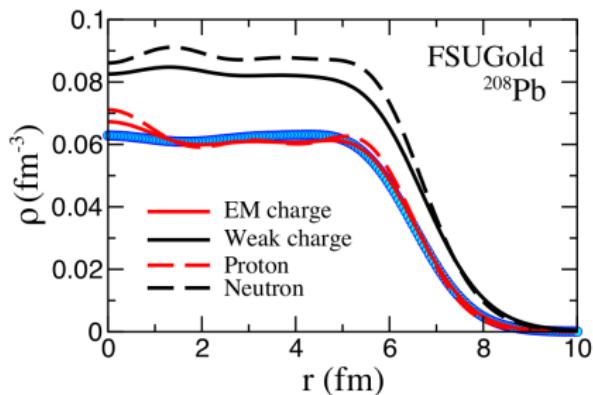
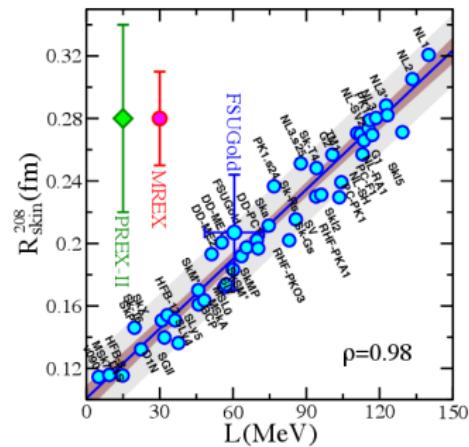
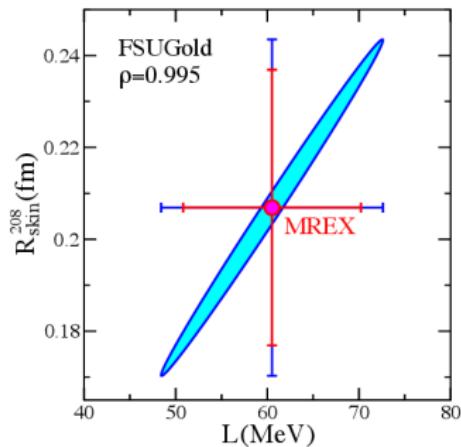


Figure: Neutron (black line) and proton (red line) spacial distributions as predicted by FSUGold model. The blue dots represents the experimental data of ^{208}Pb charge density.

Symmetry Energy

The symmetry energy $S(\rho)$ is the key component of the equation of state which controls the neutron skin thickness. $S(\rho)$ quantifies the change in energy related to the neutron-proton asymmetry.



$$\begin{aligned}\epsilon(\rho, \alpha) &= \epsilon_{SNM}(\rho) + \alpha^2 S(\rho) + O(\alpha^4) \\ S(\rho) &= J + L \cdot \left(\frac{\rho - \rho_0}{3\rho_0} \right) + \dots\end{aligned}\tag{2}$$

Neutron Skin and Neutron Star Radius

The slope of the symmetry energy is related to both the **neutron skin** of lead and **neutron star radius**. The radius of the neutron star is determined from Tolman-Oppenheimer-Volkoff (TOV) equation. Giving the pressure P_c at the center of the star, the radius R can be determined. But for neutron star, the pressure at the center is strongly related to the **pressure of pure neutron matter**, in large part determined by L .

$$P = \frac{1}{3} \rho_0 L \quad (3)$$

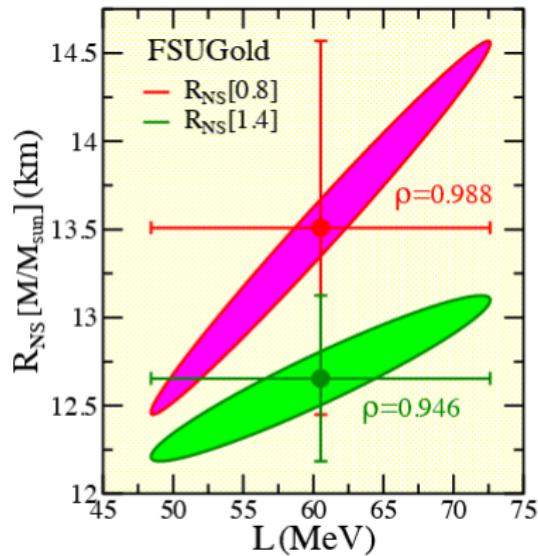


Figure: Covariance ellipses displaying the correlation between the neutron star radius and the slope of the symmetry energy, as predicted by FSUGold model.

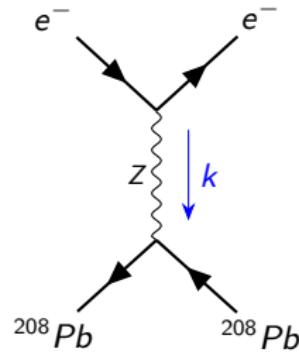
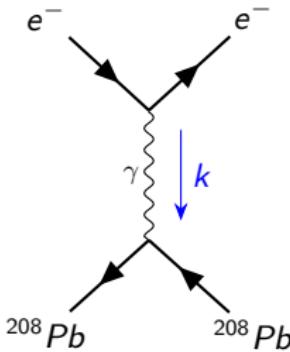
Parity Violating Asymmetry

Measurement of the Neutron Spacial Distribution of Lead

The determination of the neutron spacial distribution relies on the electron nucleus elastic scattering experiment, planned at the future MESA accelerator, in Mainz. The neutron spacial distribution is measured via the parity violating scattering, where longitudinal polarized electrons scatter from a fixed lead target. The key quantity is the asymmetry in the cross-section due to the different polarization state of the beam,

$$A_{pv} = \frac{\sigma_R - \sigma_L}{\sigma_R + \sigma_L} \quad (4)$$

The asymmetry is due to the interference between two Feynmann diagrams.

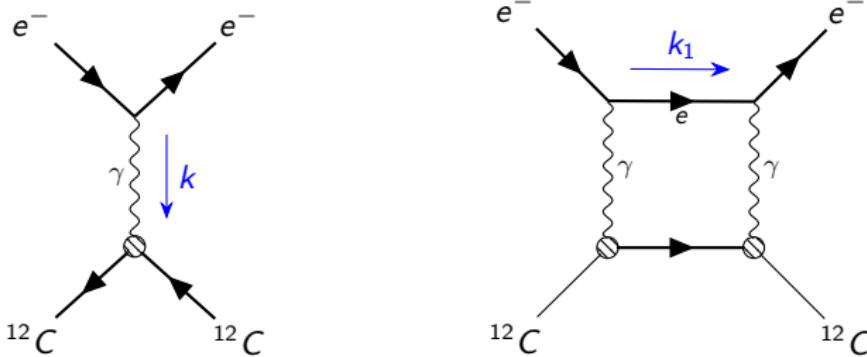


Transverse Asymmetry

The transverse asymmetry is defined as the ratio between the sum and the difference of the elastic cross section for the two different polarized electrons:

$$A_{transverse} = \frac{\sigma_{\uparrow} - \sigma_{\downarrow}}{\sigma_{\uparrow} + \sigma_{\downarrow}}$$

Before moving on to the experimental details, we identify the kinematics of the experiment. For the beam normal single spin asymmetry, the electrons are polarized in the normal plane identified by the $\frac{\vec{k}' \wedge \vec{k}}{|\vec{k}| |\vec{k}'|}$



Scattering Process

The incident beam is made by 570 MeV electrons, that are polarized along the transverse axes (\uparrow and \downarrow). The physical quantity to measure is the asymmetry between the number of scattered electrons, due to the change of the polarity:

$$asym = \frac{N_+ - N_-}{N_+ + N_-} \text{ (expected } \sim +/- 20 \text{ ppm, } Q = 0.2 \text{ GeVc}^{-1}) \quad (5)$$

The prediction for the transverse asymmetry is given by the formula below:

$$A_n \simeq C_0 \log \frac{Q^2}{m_e^2 c^2} \frac{F_{Compton}(Q^2)}{F_{ch}(Q^2)} \quad (6)$$

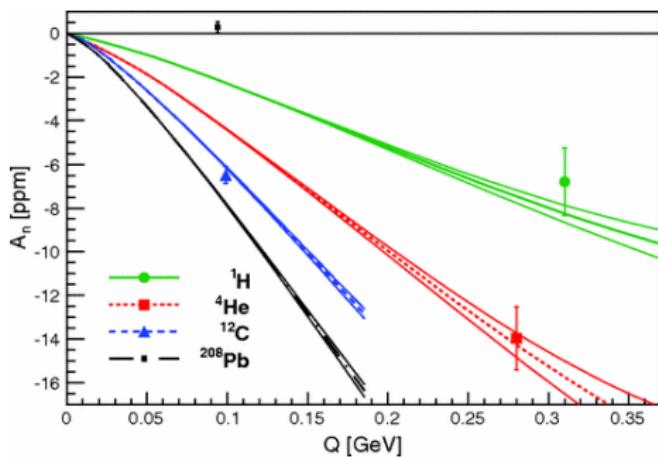


Figure: Asymmetry versus transferred momentum for selected nuclei.

Beam-time 29/11/2022 - 5/12/2022

Beam Normal Single Spin Asymmetry at MAMI

During the last beam-time, several measurements were performed at Mainz Mikrotron MAMI. The last data acquisition campaign had the following goals:

- Test the new **data acquisition system**, developed for the new setup with a low rate signals ($\simeq 1 \text{ MHz}$).
- Measure the **transverse asymmetry** A_n of ^{12}C .
- Measure the expected **rates** on ^{208}Pb target, in anticipation of the future measurement of A_n for lead.
- Long term goal: acquire more knowledge on the **systematic effects** that the transverse asymmetry has on the measurement of the Parity-violating asymmetry.

Structure of the event

$$A_{mes.} = \frac{(N_1 + N_4) - (N_2 + N_3)}{(N_1 + N_4) + (N_2 + N_3)} \quad (7)$$

The Data are divided in a series of events (80 ms), that correspond to 4 sequential sub-event. For each sub-event there is a precise polarization of the Beam. For each sub-event all the scattering electrons are counted.

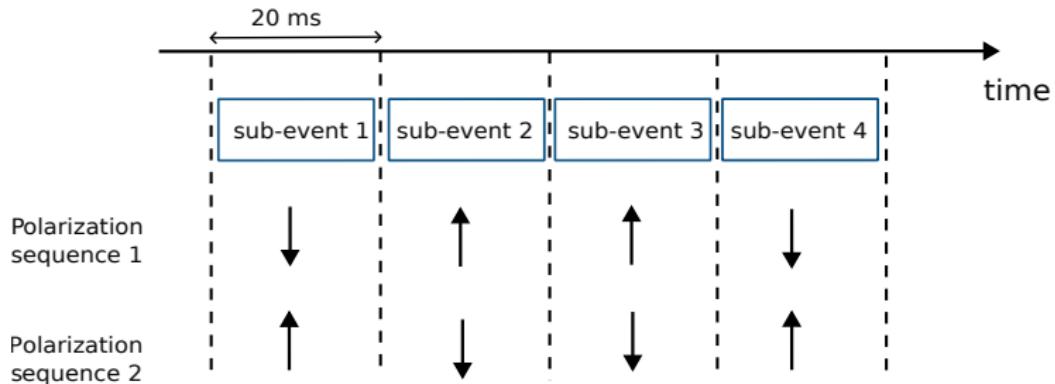


Figure: Event sequence

MAMI Electron Accelerator

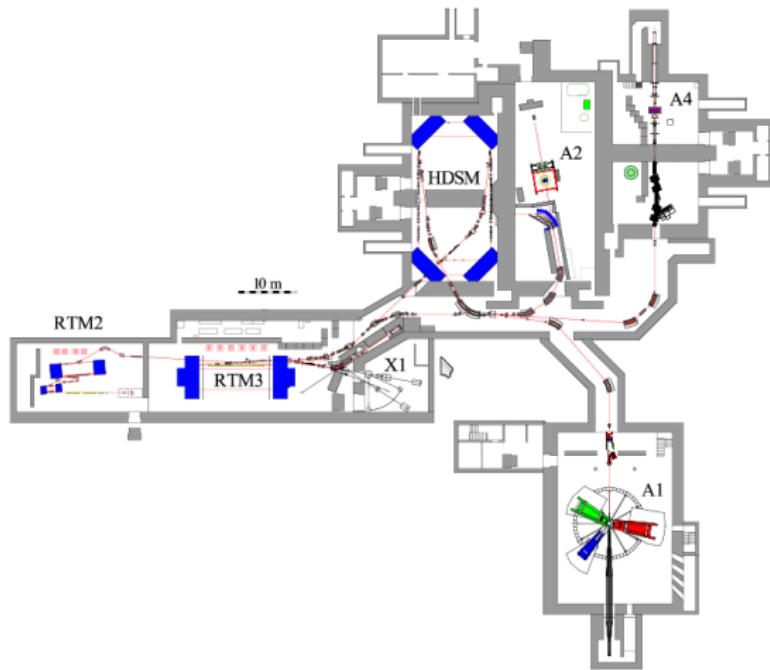
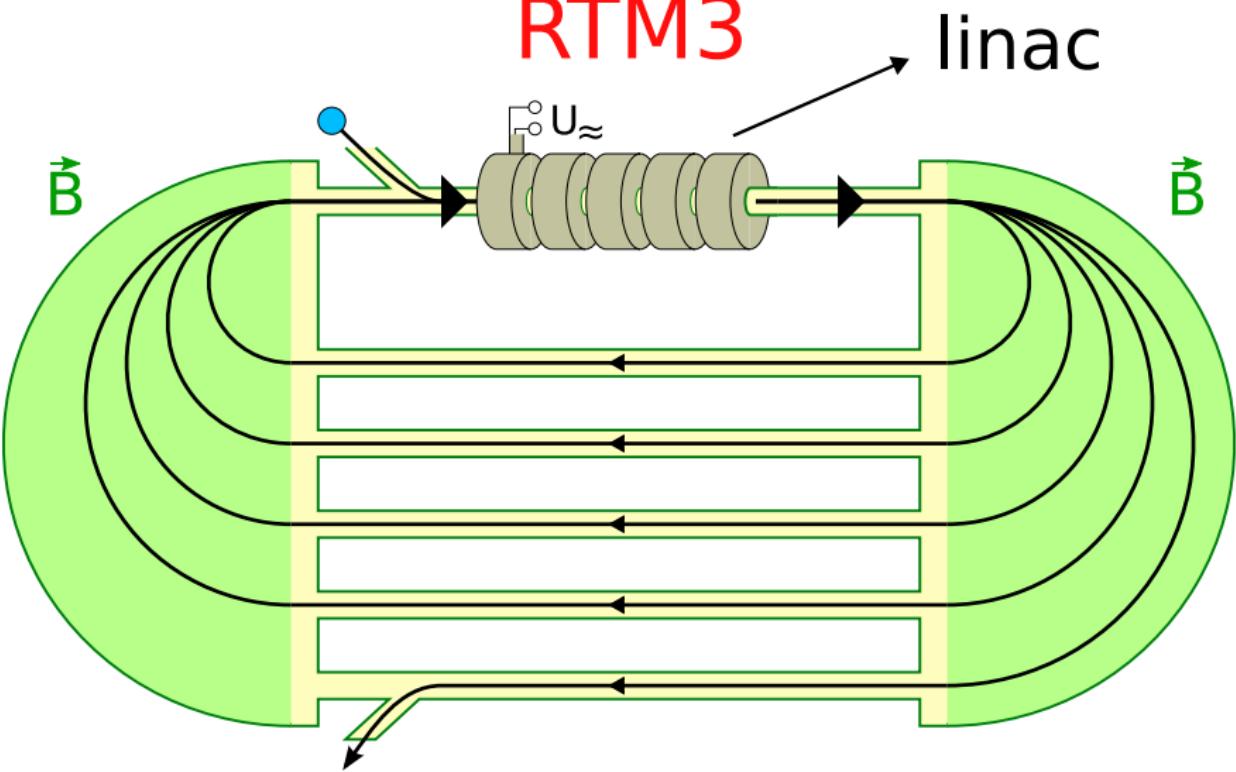


Figure: Scheme of MAMI accelerator. RTM2 is the first acceleration stage, where the electrons are extracted from the source. RTM3 corresponds to the final acceleration stage of the experiment (570 MeV), A1 is the experiment hall where the experiment described is the thesis is set up.

Race Track Microtron Scheme

RTM3



RTM3

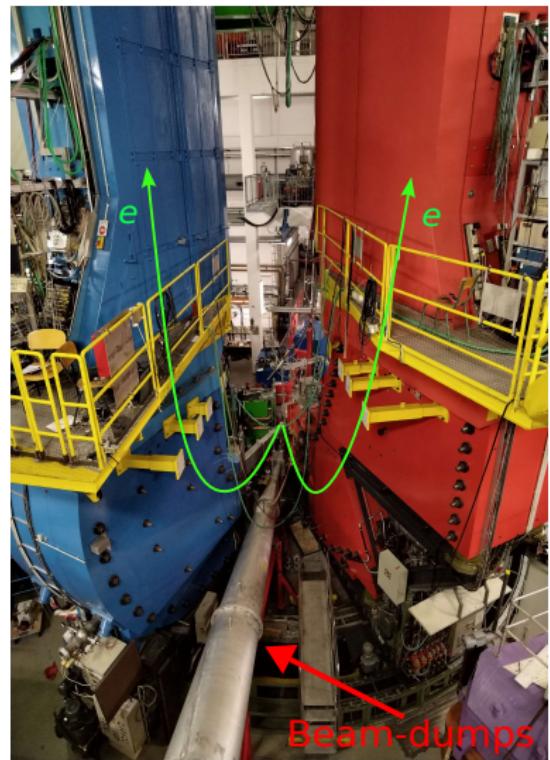
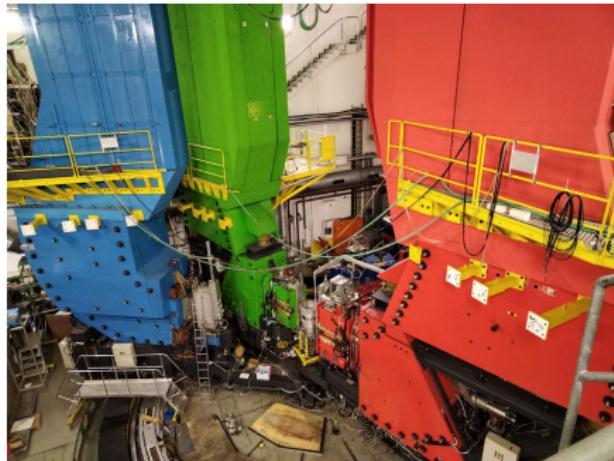
Race Track Microtron 3 of MAMI



A1 Experimental Hall

MAMI Electron Accelerator

In MAMI A1, three large spectrometers are positioned on a rail track around the scattering chamber. For the experiment, only the red and blue spectrometers were used.



Cherenkov Detectors

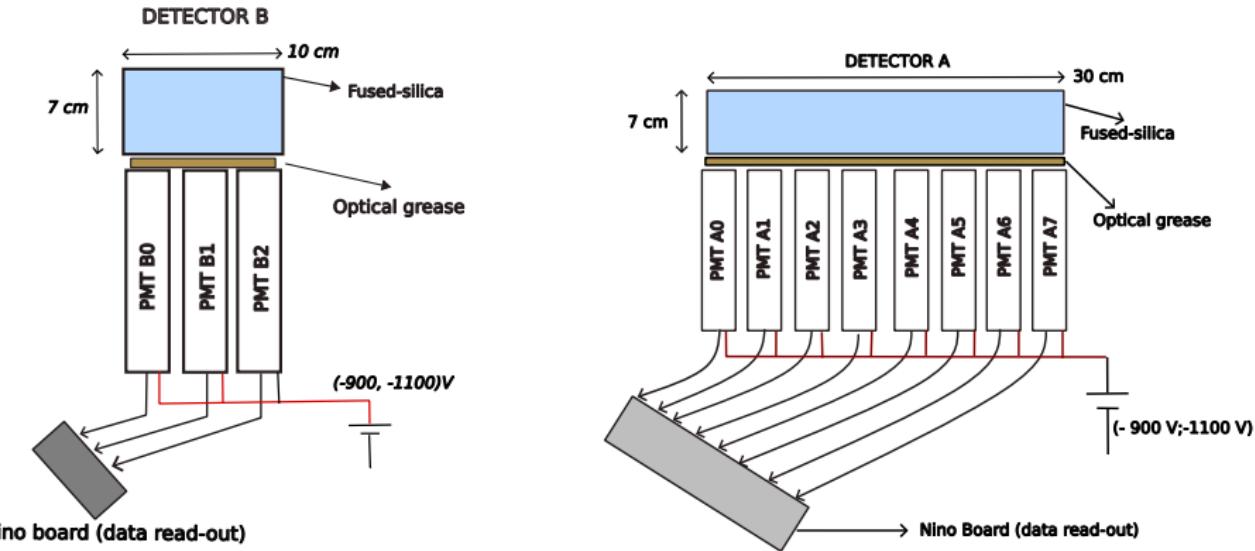


Figure: Scheme of detector B (on the left) and detector A (on the right).

NINO Asic Board

Data Acquisition Electronics

The NINO board is the DAQ system that counts the input signals coming from the Cherenkov detectors.

- 32 input channels.
- Attenuation circuit for each input channel.
- Common current Threshold for 4 adjacent channels.



False Asymmetries

The number of scattered electrons could depend on the beam parameters. Variation of the beam parameters can influence the measured asymmetries. The important beam parameters are:

- **X, Y** positions of the beam spot on the target.
- beam current **I**.
- beam energy **E**.
- incident angles of the beam θ_x, θ_y

Considering that the expected asymmetry is expected to be $\simeq 20\text{ ppm}$, small fluctuation of the beam correlated to the polarization can produce false asymmetries. The effects can be summarized by the model:

$$\text{Asym} = A_{\text{physical}} \cdot P + \delta_I + A_x \delta x + A_y \delta y + A_{\theta_x} \delta \theta_x + A_{\theta_y} \delta \theta_y + A_E \delta E \quad (8)$$

MAMI Beam Monitors

At MAMI the beam parameters, as **X,Y** transverse positions, **Energy** and **Current** are measured by resonant cavities.

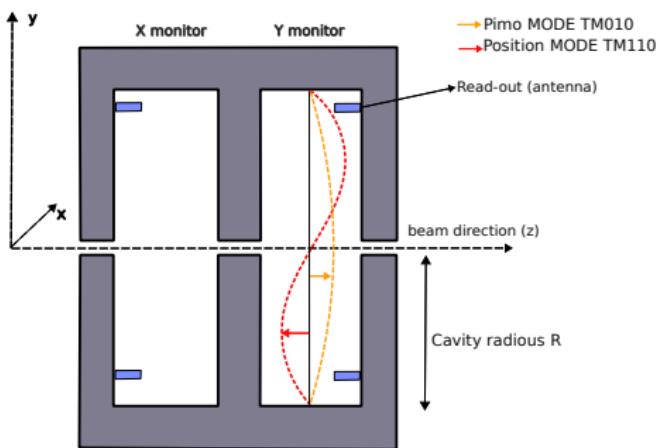
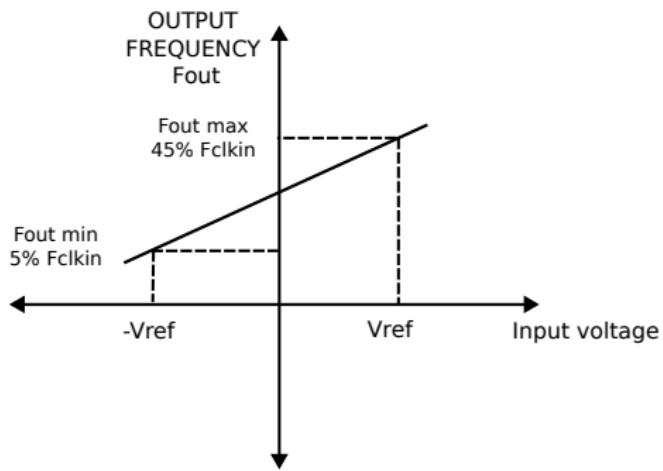


Figure: Scheme of the resonant cavity in MAMI, with the different electromagnetic modes excited by the beam passage.

The beam monitors are read-out with the voltage-to-frequency converters **VFC** devices, which produces an output signal whose frequency is proportional to the amplitude of the beam monitor signals.



General Scheme of the Experiment

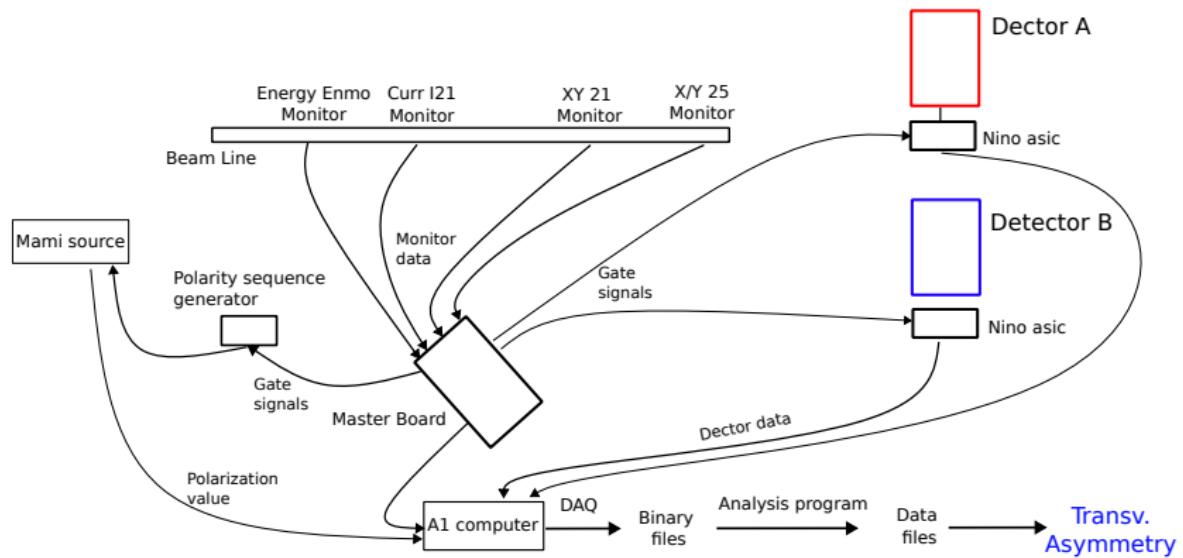


Figure: Scheme of the experiment.

Calibration of the Beam Parameters

working point of the PMTS

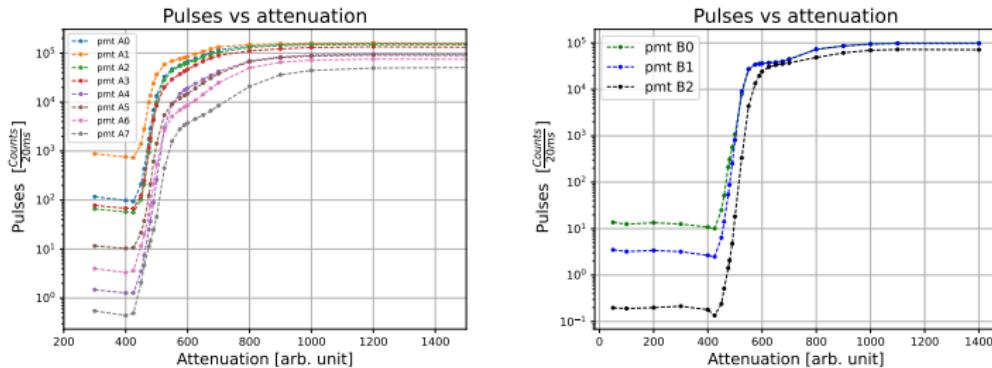
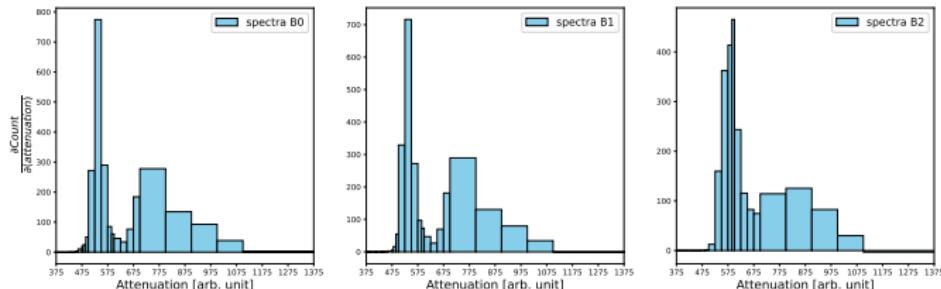
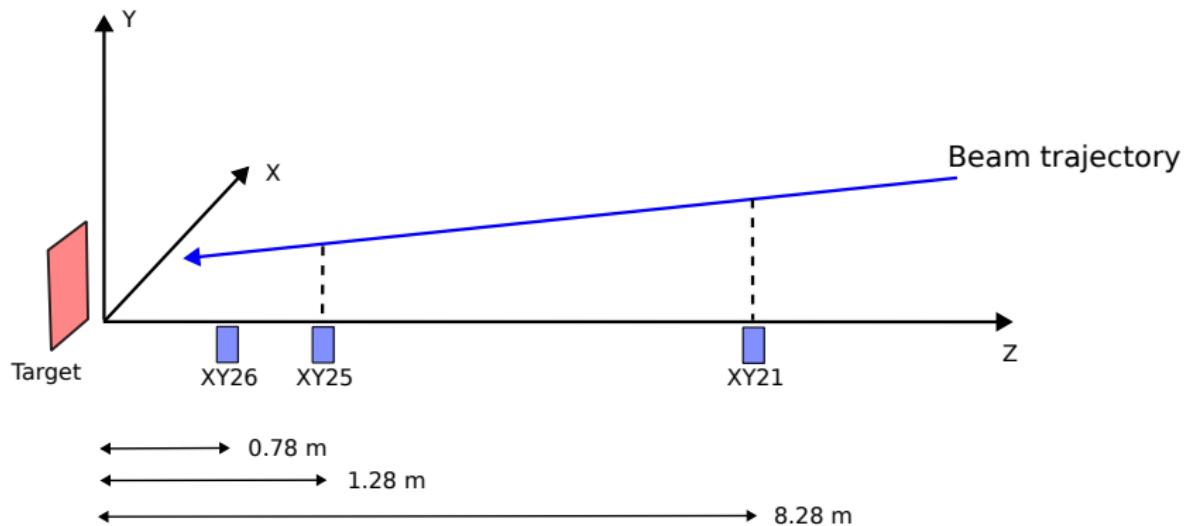


Figure: Counts versus attenuation value for PMTs of detector A (left plot) and B (right plot).



Beam Position

Computing the beam spot (X,Y) coordinates on the target



The beam is assumed following a straight path given by $x, y = m \cdot z + q_{x,y}$, the (X, Y) coordinates are given by $q_{x,y}$

Beam Energy and Beam Current

The conversion parameters from V to μA and keV are measured with a current and energy scan.

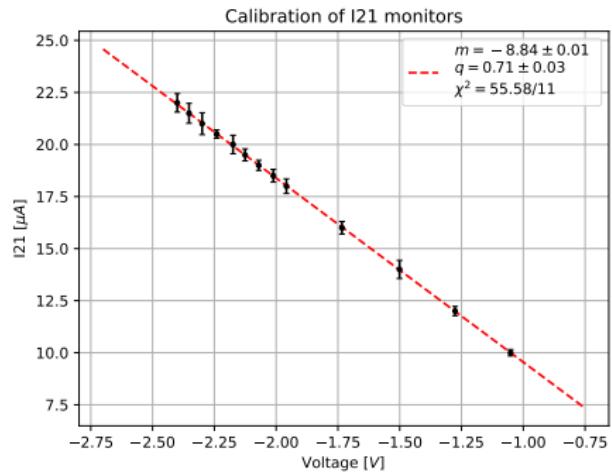


Figure: Current monitor I21. On the x axis the voltage signal of the PIMO monitor, on the y axis the nominal current of the beam.

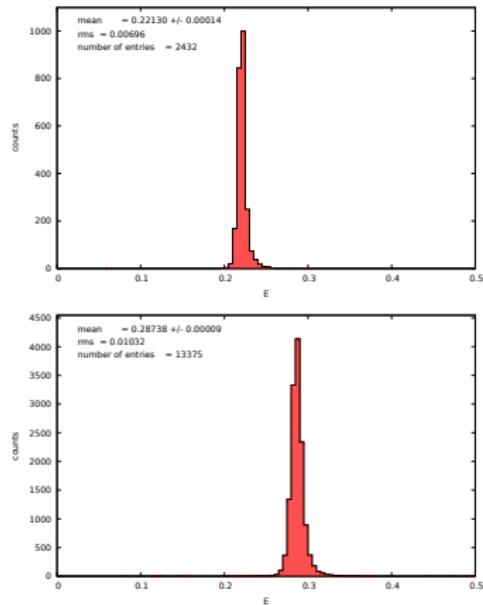
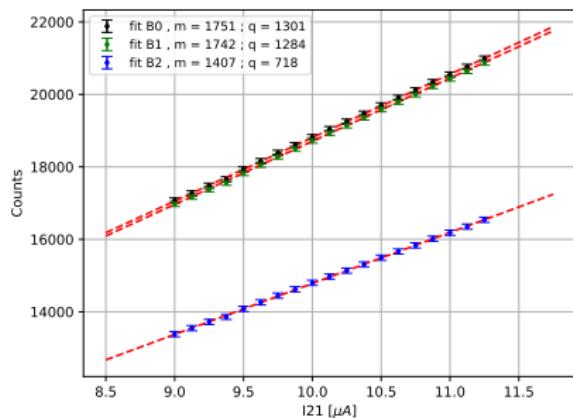
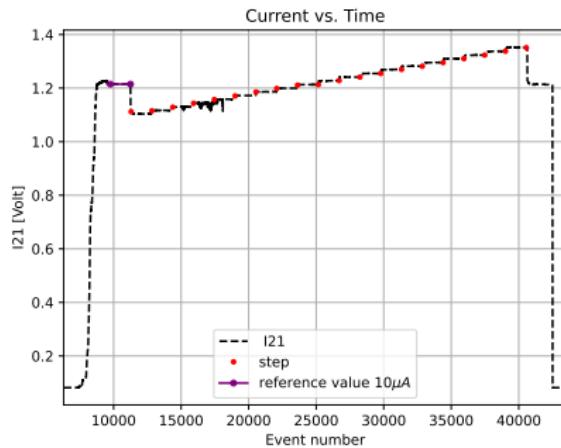


Figure: Calibration of the energy monitor. For the first histogram the current is 15 μA , 20 μA for the second one.

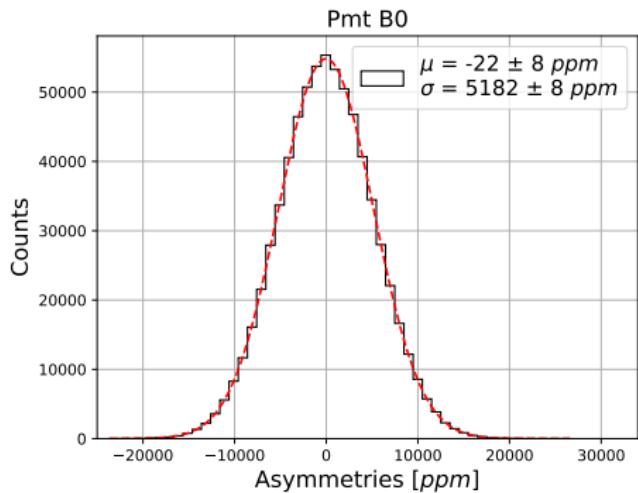
Auto-Calibration procedure

Every 3 hours of beam, a special operation mode of MAMI is set. The current is raised in small steps, from $9\text{ }\mu\text{A}$ to $11.125\text{ }\mu\text{A}$. This is used to check the linearity of the PMTs through time, and to measure the PMT offset, that is subtracted later from the data.



Analysis on Carbon Target

Each event corresponds to a single measurement of the asymmetry A . With this experiment, an amount of 1 million events have been collected. The asymmetry data are normal distributed:



$$\text{Var}[A_{\text{asym}}] = \text{Var}\left[\frac{N_{\uparrow} - N_{\downarrow}}{N_{\uparrow} + N_{\downarrow}}\right] \simeq \frac{\text{Var}[N_{\uparrow} - N_{\downarrow}]}{(N_{\uparrow} + N_{\downarrow})^2}$$
$$\frac{2\text{Var}[N]}{4N^2} = \frac{1}{2N} \quad \sigma = \frac{1}{\sqrt{2N}}$$

Where it is supposed that the PMTs counts are normal distributed, with μ equal to σ^2 . The rms associated to the sample mean decreases as the $\sqrt{N_{\text{measure}}}$. With the statistical error is:

$$\sigma = \frac{1}{\sqrt{N \cdot N_{\text{measure}}}} \quad (9)$$

Considering $5 \cdot 10^5$ events and $\mu = 20000$ counts per PMT (similar to what was measured for detector B)

- $\sigma = 8 \text{ ppm}$.

Model For Fitting the Data

The final model that describes the asymmetry takes care of the possible false asymmetries, assuming a **linear dependence** on the beam parameters.

$$A_{\text{tot}} = A_{\text{physical}} \cdot P + \delta_I + A_x \delta x + A_y \delta y + A_{\theta_x} \delta \theta_x + A_{\theta_y} \delta \theta_y + A_E \delta E \quad (10)$$

Each δq is computed as the mean difference between up and down polarized sub-events:

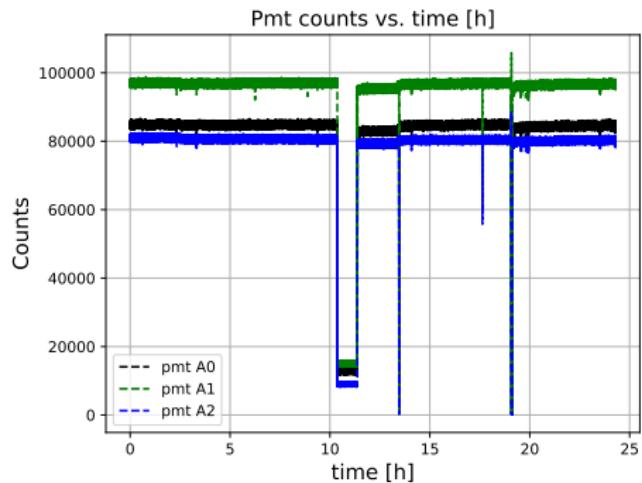
$$\delta q = \frac{q_{\uparrow,0} + q_{\uparrow,1}}{2} - \frac{q_{\downarrow,0} + q_{\downarrow,1}}{2} \quad (11)$$

The current asymmetry is defined differently. Because the scattering rate is proportional to the beam current, the model contains the current asymmetry directly:

$$\delta I = \frac{I_{\uparrow} - I_{\downarrow}}{I_{\uparrow} + I_{\downarrow}} \quad (12)$$

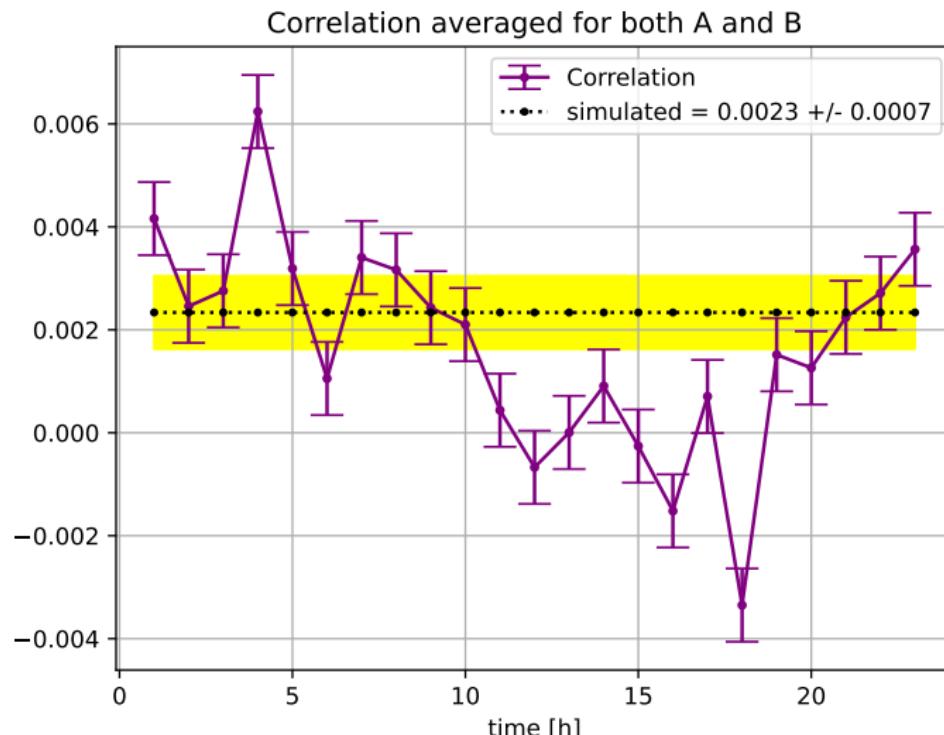
Data Selection

Before proceeding in the analysis, an accurate data selection was performed, to remove the outliers, as well as data affected by anomalous variation of the beam current.

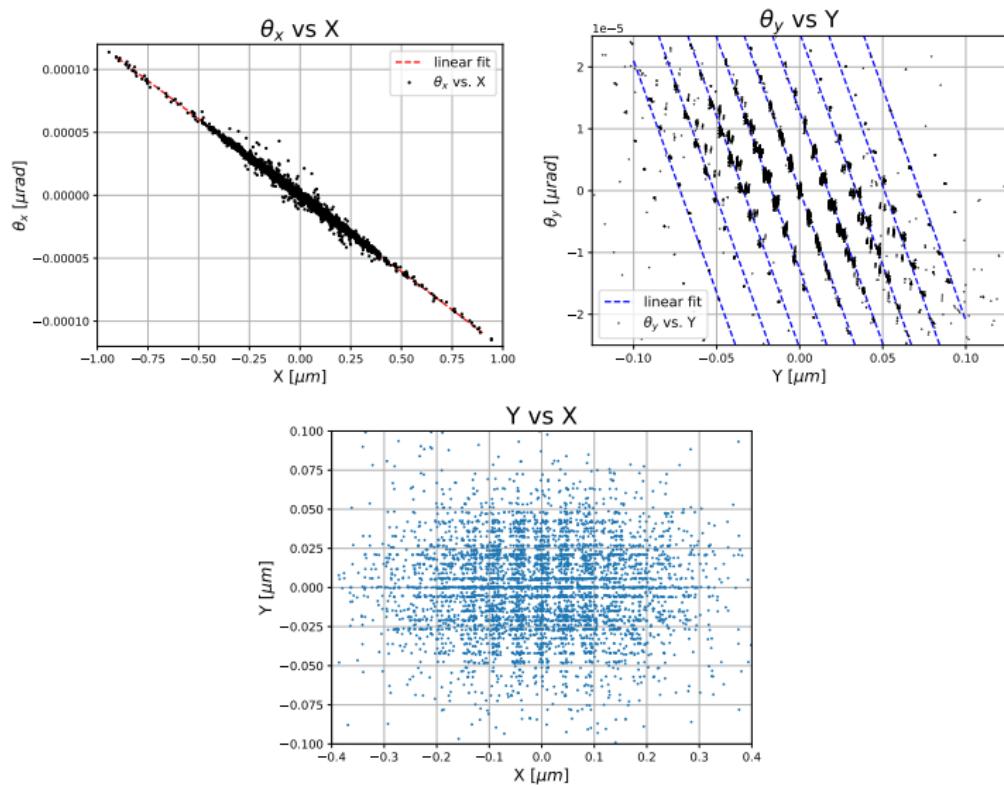


Polarization Loss

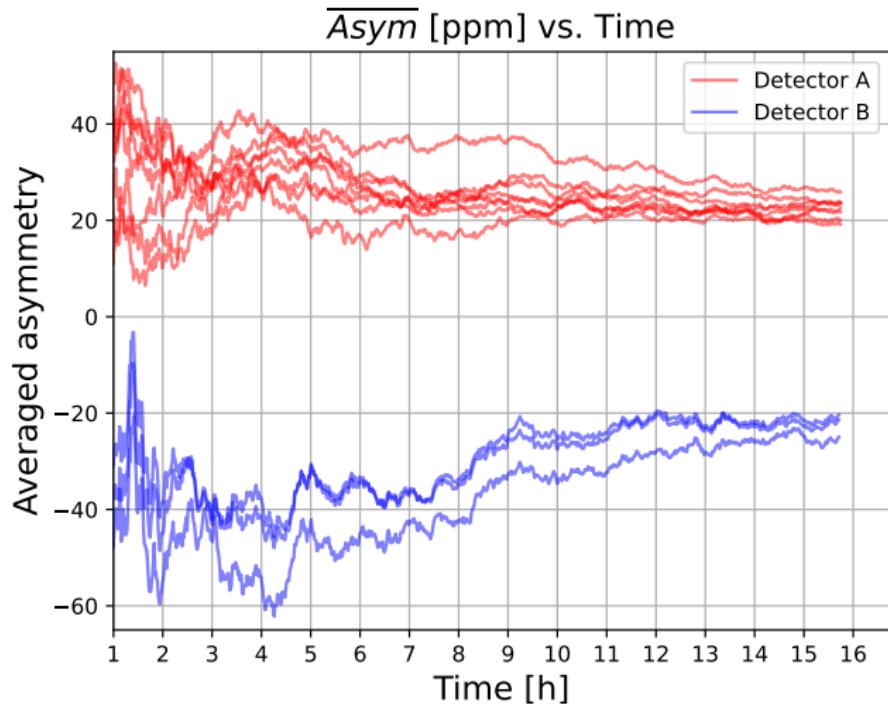
During the experiment, a relevant part of the data was affected by a loss of polarization of the beam:



Beam Parameters Correlation

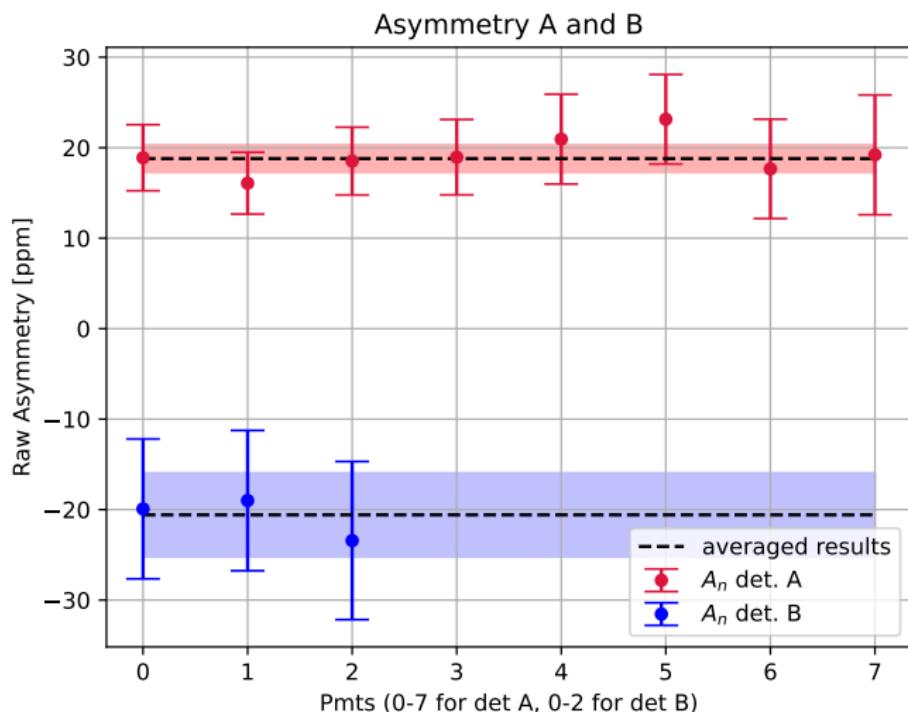


Here a plot about the trend of the asymmetry as the data increases. The band is the error computed as showed in the previous slide, centered around the values of $+20\text{ppm}$ for detector A and -20ppm for detector B.



Results

Final asymmetry result for each PMTs:



Results

The result of the analysis is obtained combining the measured asymmetries for PMTs.

$$\hat{A} = \frac{\sum_i A_i \frac{1}{w_i}}{\frac{1}{w_i}} \quad w_i = \frac{1}{\sigma_i^2}$$

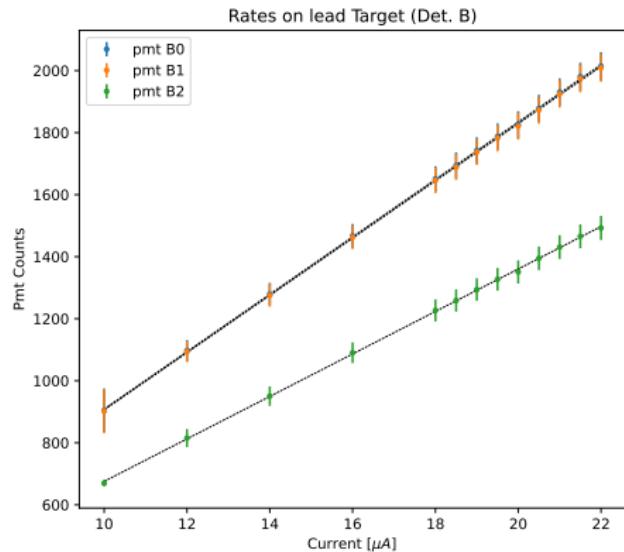
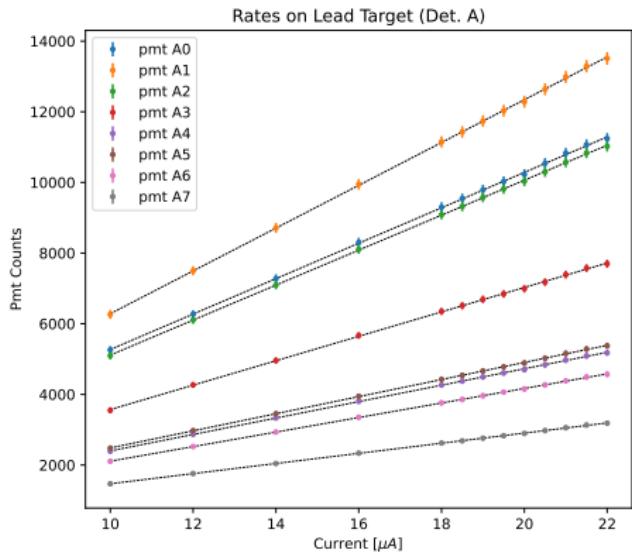
the final values of the Beam Normal Single Spin asymmetries are:

$$A_A = (23.1 \pm 1.7) \text{ ppm} \quad A_B = (-21 \pm 5) \text{ ppm} \quad (13)$$

With a transfer momentum $Q = 0.04 \text{ GeV}$. Reversing the sign of the asymmetry for Det. B we notice that the two measurement are consistent within 1σ .

Rates on Lead

The rates on lead target (1 mm thick) have been measured. The transverse asymmetry on lead will be the next step of MREX experiment. This measurement is mandatory to constrain the systematic error of PV experiment.



Since the rates are lower for lead compared to carbon, the statistics needed to measure A_n for ^{208}Pb is 23 times the statistic accumulated for carbon.

Ringraziamenti

GRAZIE A TUTTI PER L'ATTENZIONE