

Commissioning and First Data Analysis of the Mainz Radius Experiment

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The Mainz Radius Experiment

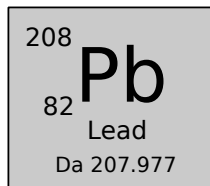
MREX

The Mainz Radius Experiment is an experimental campaign at the nuclear physics institute of Mainz, with the aim of investigating the properties of nuclear matter with imbalance in the number of protons and neutrons.

Objective

Determination of the neutron spacial density for ^{208}Pb nucleus, through the elastic electron scattering. From an accurate determination of the neutron spacial distribution, the *Neutron Skin Thickness* of ^{208}Pb is measured.

The results of the experiment will be valuable to constrain the Equation of State (EOS) of nuclear matter. It has also implications for the structure of neutron star.



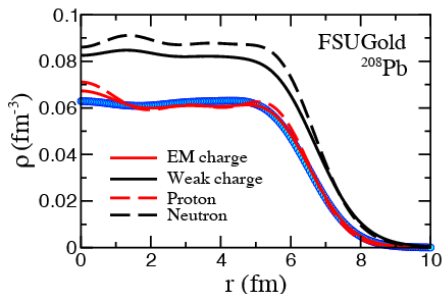
MREX and Neutron Skin Thickness

Definition

The neutron skin thickness is defined as the difference between rms radius of neutron and proton spacial distributions:

$$\delta r_{np} = \sqrt{\langle r_n^2 \rangle} - \sqrt{\langle r_p^2 \rangle}, \quad (1)$$

In neutron rich nuclei, the spacial distribution of neutrons is more extended than proton spacial distribution. Theoretical models link the Neutron skin thickness of heavy nuclei, such as ^{208}Pb , with the **slope of the symmetry energy L** .



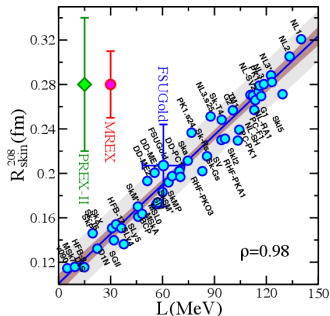
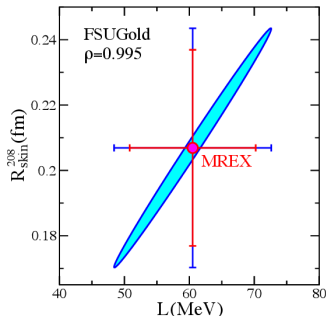
Symmetry Energy

The symmetry energy $S(\rho)$ is the key component of the equation of state which controls the neutron skin thickness. $S(\rho)$ quantifies the change in energy related to the neutron-proton asymmetry.

$$\epsilon(\rho, \alpha) = \epsilon_{SNM}(\rho) + \alpha^2 S(\rho) + O(\alpha^4)$$

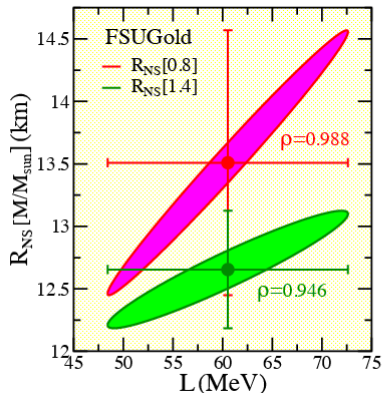
$$\epsilon(\rho) = J + L \cdot \left(\frac{\rho - \rho_0}{3\rho_0} \right) + \dots \quad (2)$$

Cosa è l'energia di simmetria, come è collegata con il comportamento della materia nucleare di neutroni



Neutron Skin and Neutron Star Radius

The slope of the symmetry energy is related to both the neutron skin of lead and neutron star radius. The radius of the neutron star is determined from the Tolman–Oppenheimer–Volkoff (TOV) equation. Giving the pressure P_c at the center of the star, the radius R can be determined. But for neutron star, the pressure at the center is strongly related to the pressure of pure neutron matter, in large part determined by L .



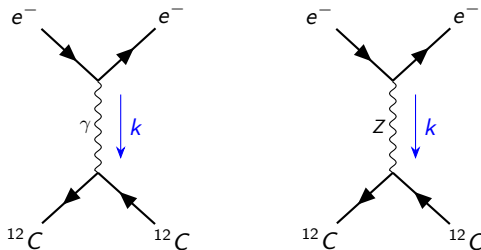
Parity Violating Asymmetry

Measurement of the Neutron Spatial Distribution of Lead

The determination of the neutron spatial distribution relies on the electron nucleus elastic scattering experiment, planned at the future MESA accelerator, in Mainz. The neutron spatial distribution is measured via the parity violating scattering, where longitudinal polarized electrons scatter from a fixed lead target. The key quantity is the asymmetry in the cross-section due to the different polarization state of the beam,

$$A_{pv} = \frac{\sigma_R - \sigma_L}{\sigma_R + \sigma_L} \quad (3)$$

The asymmetry is due to the interference between two Feynmann diagrams.



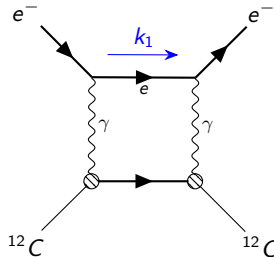
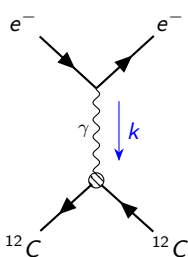
Transverse Asymmetry

Argomento principale della tesi: misura dell'asimmetria trasversa, fondo sistematico di A_{pv} da determinare. Introduzione alla fisica del processo

The transverse asymmetry is defined as the ratio between the sum and the difference of the elastic cross section for the two different polarized electrons:

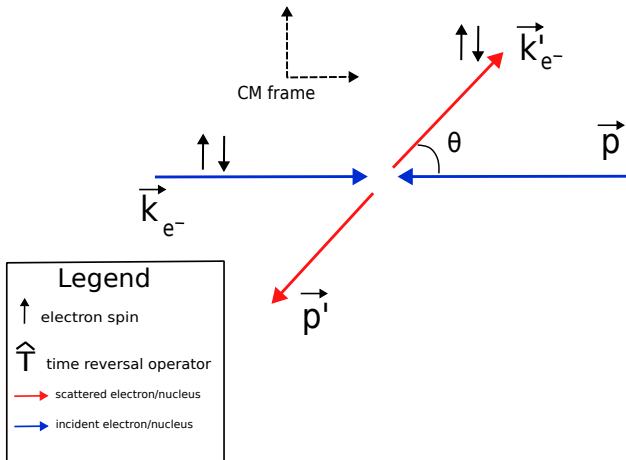
$$A_{transverse} = \frac{\sigma_{\uparrow} - \sigma_{\downarrow}}{\sigma_{\uparrow} + \sigma_{\downarrow}}$$

Before moving on to the experimental details, we identify the kinematics of the experiment. For the beam normal spin asymmetry, the electrons are polarized in the normal plane identified by the $\frac{\vec{k}' \wedge \vec{k}}{|\vec{k}| |\vec{k}'|}$



Description of the Process

The transverse asymmetry arises considering the time reversal operator. The time reversal operator reverses all the momenta and the spin direction.



Scattering Process

The incident beam is made by 570 MeV electrons, that are polarized along the transverse axes (\uparrow and \downarrow). The physical quantity to measure is the asymmetry between the number of scattered electrons, due to the change of the polarity:

$$asym = \frac{N_+ - N_-}{N_+ + N_-} (\text{expected} \sim \pm 20 \text{ ppm}, Q = 0,2 \text{ GeVc}^{-1}) \quad (4)$$

It's possible to obtain a final formula for the transverse Asymmetry, writing the Amplitude of the 1-loop diagram, considering the elastic intermediate state and the inelastic intermediate state (whose contribution is higher):

$$A_n \simeq C_0 \log \frac{Q^2}{m_e^2 c^2} \frac{F_{Compton}(Q^2)}{F_{ch}(Q^2)} \quad (5)$$

MENZIONARE PREX!!!

During the last beam-time, several measurements were performed at Mainz Mikrotron MAMI. The last data acquisition campaign had the following goals:

- Test the new data acquisition system, developed for the new setup with a low rate signals ($\simeq 1$ MHz).
- Measure the transverse asymmetry A_n of ^{12}C .
- Measure the expected rates on ^{208}Pb target, in anticipation of the future measurement of A_n for lead.
- Long term goal: acquire more knowledge on the systematic effects that the transverse asymmetry has on the measurement of the Parity-violating asymmetry.

Structure of the event

The Data are divided in a series of events (80 ms), that correspond to 4 sequential sub-event. For each sub-event there is a precise polarization of the Beam. For each sub-event all the scattering electrons are counted.

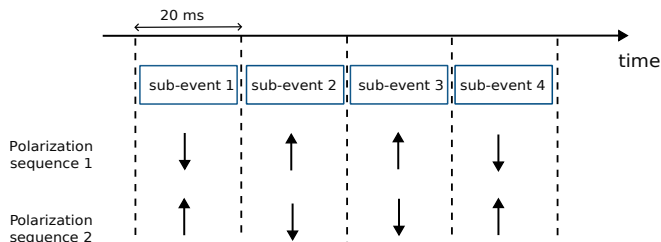
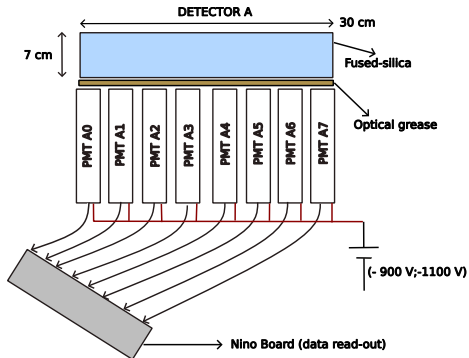
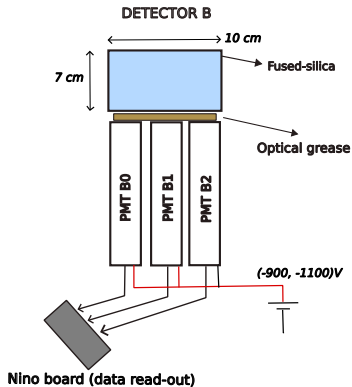


Figure: Event sequence: all the particle

Schema dell'acceleratore e descrizione breve del suo funzionamento

Immagine A1.

Detectors



Nino board, immagine + spiegazione di come sono acquisiti i segnali.

False Asymmetries

The counts of the pmts can be slightly different due to the variation of the position of the beam on the target, the variations of the incident angles, the uncertain associated with the energy and the current of the beam. All this quantity can influence the asymmetry measured by the pmts, considering also that the expected asymmetry is in the order of ten part per million, and small asymmetry introduced by fluctuations of the beam parameters are not negligible:

$$Asym = A_{physical} \cdot P + \delta_I + A_x \delta x + A_y \delta y + A_{\theta_x} \delta \theta_x + A_{\theta_y} \delta \theta_y + A_E \delta E \quad (6)$$

Descrizione dei principi di funzionamento dei monitors di MAMI.

Voltage to Frequency Converter

Breve descrizione di come funzionano i voltage to frequency converter

General Scheme of the Experiment

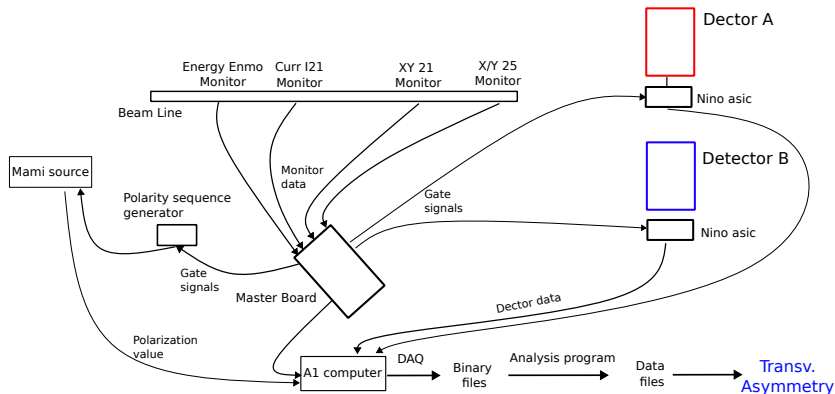


Figure: Scheme of the experiment.

Calibration of the Beam Parameters

Beam Position

Beam Energy

Auto-Calibration procedure

Analysis on Carbon Target

Modello lineare tra asimmetria e beam parameters, discutere differenza corrente e altri parametri del fascio.

Discutere la rilevante perdita di polarizzazione che è avvenuta ed il modo in cui si sono identificati questi dati.

Beam Parameters Correlation

Variance of the Asymmetry Data

The statistical error of the measured asymmetries is now computed:

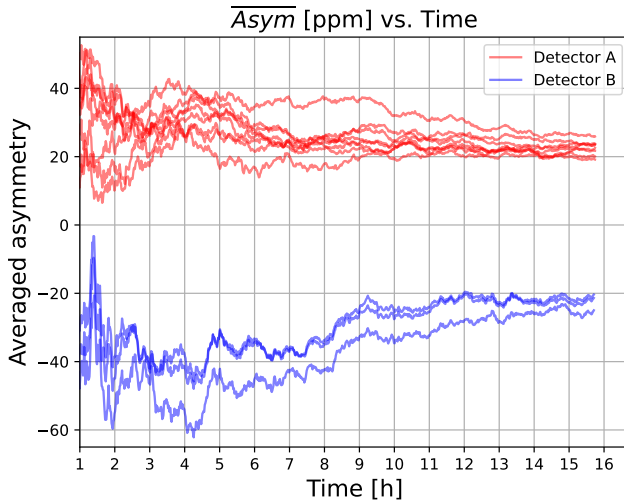
$$\text{Var}[A_{\text{asym}}] = \text{Var}\left[\frac{N_{\uparrow} - N_{\downarrow}}{N_{\uparrow} + N_{\downarrow}}\right] \simeq \frac{\text{Var}[N_{\uparrow} - N_{\downarrow}]}{(N_{\uparrow} + N_{\downarrow})^2}$$
$$\frac{2\text{Var}[N]}{4N^2} = \frac{1}{2N} \quad \sigma = \frac{1}{\sqrt{2N}}$$

Where it is supposed that the PMTs counts are normal distributed, with μ equal to σ^2 .

The rms associated to the sample mean decreases as the $\sqrt{N_{\text{measure}}}$.

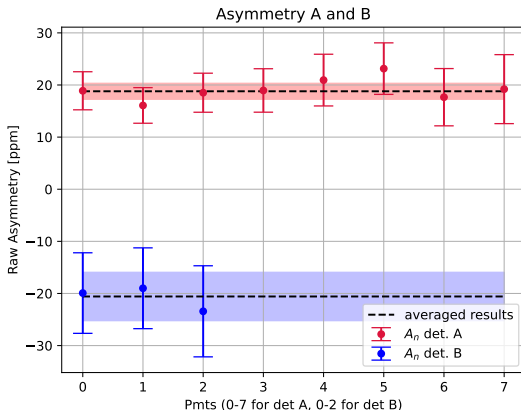
Considering $5 \cdot 10^5$ events and $\mu = 40000$ counts per PMT (similar to what was measured for detector A) we obtain an error of $\simeq 5\text{ppm}$.

Here a plot about the trend of the asymmetry as the data increases. The band is the error computed as showed in the previous slide, centered around the values of $+20\text{ppm}$ for detector A and -20ppm for detector B.



Results

For each pmt, we present the raw values of the asymmetry, obtained by subtracting the Raw current asymmetry, that is roughly -1.11 ppm, then we compute the averaged values:



Combining the result of each pmt, assuming all the asymmetries measured are independent of each other, we obtain the following quantities for Beam normal single spin asymmetries:

$$\hat{A} = \frac{\sum_i A_i \frac{1}{w_i}}{\frac{1}{w_i}} \quad w_i = \frac{1}{\sigma_i^2}$$

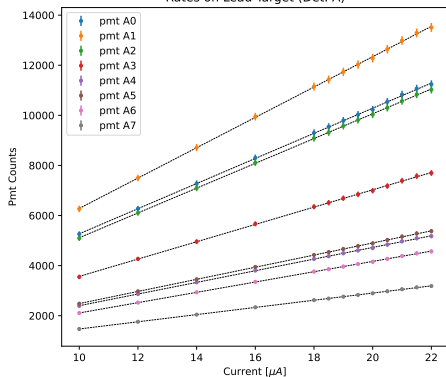
We obtain the following:

$$A_A = (23.1 \pm 1.7)ppm \quad A_B = (-21 \pm 5)ppm \quad (7)$$

Reversing the sign of the asymmetry for detB we notice that the two measurement are consistent, and this show the good behaviour of the electronic setup used for the experiment.

Rates on Lead

Rates on Lead Target (Det. A)



Rates on lead Target (Det. B)

