

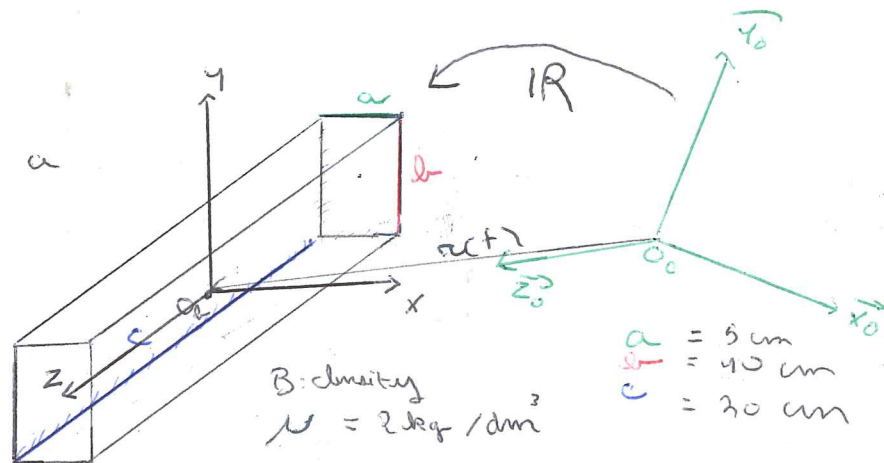
①

Term Project for TEK4040 Mathematical Modelling of Dynamic systems

November 2022

1.

Let us consider a brick B with affine space B



with center of mass b such that

$$B: \mathcal{F}_B = \{ O_b, \vec{x}, \vec{y}, \vec{z} \} = \{ b \}$$

Introducing inertial frame with affine space R

$$R: \mathcal{F}_R = \{ O_0, \vec{x}_0, \vec{y}_0, \vec{z}_0 \} = \{ 0 \}$$

introducing position vector $r(t) = \overrightarrow{O_0 O_b}(t)$ and Rotation operator IR between $\{0\} \rightarrow \{b\}$

$$\text{such that } IR = IR_{ob} = \underbrace{IR_{ob}^0}_{\text{rotational operator}} = \underbrace{IR_{ob}^b}_{\text{DCM}} = R_{ob}^0$$

introducing position vector $P_C^b = \begin{bmatrix} p_1 \\ p_2 \\ p_3 \end{bmatrix} = \begin{bmatrix} x \\ y \\ z \end{bmatrix}$

$$\text{such that } P_C^b = \begin{bmatrix} \vec{O_b C} \cdot \vec{x} \\ \vec{O_b C} \cdot \vec{y} \\ \vec{O_b C} \cdot \vec{z} \end{bmatrix}$$

$$J_b = \iiint_{C \in B} \begin{bmatrix} p_2^2 + p_3^2 & -p_1 p_2 & -p_1 p_3 \\ -p_2 p_1 & p_1^2 + p_3^2 & -p_2 p_3 \\ -p_3 p_1 & -p_3 p_2 & p_1^2 + p_2^2 \end{bmatrix} \rho(C) dx dy dz = \begin{bmatrix} J_{xx} & J_{xy} & J_{xz} \\ J_{xy} & J_{yy} & J_{yz} \\ J_{xz} & J_{yz} & J_{zz} \end{bmatrix}$$

it can be shown that product of inertia are null because of symmetry:

$$\textcircled{2} \quad \forall i \neq j$$

$$J_{ij}^{\text{e}} = - \iiint \rho_i \rho_j \, N \, dp_i \, dp_j \, dp_k$$

$$L_i = \int dp_i$$

$i \in B$

$$= \int \left[\int \left[\int \rho_i \rho_j \, N \, dp_j \right] dp_i \right] dp_k$$

$$= \iint \rho_i \left[\left[\frac{p_j^2}{2} \right]_{-L_{jk}/2}^{L_{jk}/2} \right] dp_i \, dp_k$$

} because for all $i \in B$ p_j is centered

$$= 0$$

$$= 0$$

$$\forall i, \quad \begin{matrix} i \neq c \\ k \neq c \end{matrix}; \quad i \neq k$$

$$J_{ic}^{\text{e}} = \iiint \rho \left[p_i^2 + p_k^2 \right] \rho \, dp_i \, dp_j \, dp_k$$

$$= \rho L_i \iint \left[p_j^2 + p_k^2 \right] dp_j \, dp_k$$

$$= \rho L_i \int \left[\frac{p_j^3}{3} + p_k^2 p_j \right]_{-L_{jk}/2}^{L_{jk}/2} dp_k$$

$$= \rho L_i \int p_k^2 L_j + \frac{L_j^3}{12} dp_k$$

$$= \rho L_i L_j \left[\frac{p_k^3}{3} + \frac{L_j^2}{12} p_k \right]_{-L_k/2}^{L_k/2}$$

$$= \rho L_i L_j L_k \left[\frac{L_j^2}{12} + \frac{L_k^2}{12} \right]$$

$$= m \left[\frac{L_j^2}{12} + \frac{L_k^2}{12} \right]$$

$$m = 2 \, \text{e-06} \, \text{kg}$$

$$J_{11}^{\text{e}} \approx 8,3 \, \text{e-09} \, \text{kg} \cdot \text{m}^2$$

$$J_{22}^{\text{e}} \approx 7,083 \, \text{e-09} \, \text{kg} \cdot \text{m}^2$$

$$J_{33}^{\text{e}} \approx 2,083 \, \text{e-09} \, \text{kg} \cdot \text{m}^2$$

③ 2) Given constant energy J , the inertia ellipsoid is defined w.r.t for all $\underline{x}^T = \begin{bmatrix} p_1 \\ p_2 \\ p_3 \end{bmatrix}$ such that

$$\frac{1}{2} \underline{x}^T J \underline{x} = J$$

$$\Leftrightarrow \frac{1}{2} [p_1, p_2, p_3] \begin{bmatrix} J_{11} & 0 & 0 \\ 0 & J_{22} & 0 \\ 0 & 0 & J_{33} \end{bmatrix} \begin{bmatrix} p_1 \\ p_2 \\ p_3 \end{bmatrix} = J$$

$$\Leftrightarrow \frac{J_{11} p_1^2 + J_{22} p_2^2 + J_{33} p_3^2}{2} = J$$

$$\Leftrightarrow \frac{p_1^2}{\frac{2J}{J_{11}}} + \frac{p_2^2}{\frac{2J}{J_{22}}} + \frac{p_3^2}{\frac{2J}{J_{33}}} = 1$$

\hookrightarrow this is an ellipsoid with half axis $a_i = \sqrt{\frac{2J}{J_{ii}}}$

Given free torque motion the kinetic rotational energy ^{ellipsoid} is derived from Euler equations for all (i, j, k) , $i \neq j \neq k$, $k \equiv j+1 [3]$ (using $J = \begin{bmatrix} J_{11} & J_{12} & J_{13} \\ J_{21} & J_{22} & J_{23} \\ J_{31} & J_{32} & J_{33} \end{bmatrix}$)

$$m_i = 0 = J_{ii} \dot{\omega}_i + \omega_j \omega_k (J_{kk} - J_{jj})$$

$$\Leftrightarrow 0 = J_{ii} \dot{\omega}_i \omega_i + \omega_i \omega_j \omega_k (J_{kk} - J_{jj})$$

$$\Rightarrow \sum_{\substack{i=0 \\ (i \neq j \neq k \\ k \equiv j+1 [3])}}^2 J_{ii} \dot{\omega}_i \omega_i + \sum_{\substack{i=0 \\ (i \neq j \neq k \\ k \equiv j+1 [3])}}^2 \omega_i \omega_j \omega_k (J_{kk} - J_{jj}) = 0$$

$$\Rightarrow \sum_{i=0}^2 J_{ii} \dot{\omega}_i \omega_i = 0$$

$$\Rightarrow \int_{t_0}^t \sum_{i=0}^2 J_{ii} \dot{\omega}_i \omega_i dt = K$$

$$\sum_{i=0}^2 \left[\frac{J_{ii} \omega_i^2}{2} \right] = K$$

$$\sum_{i=0}^2 \frac{\omega_i^2}{\frac{2K}{J_{ii}}} = 1$$

ellipsoid with half axis $a_i = \sqrt{\frac{2K}{J_{ii}}}$

not needed for the answer

④ Given free torque motion, the angular momentum ellipsoid is derived from

$$\frac{d\vec{h}}{dt} = \vec{0}$$

$$\Rightarrow \|\vec{h}\| = h_0 \quad (\text{constant length})$$

$$\vec{h} = \vec{J} \vec{\omega} = \begin{bmatrix} J_{xx} \omega_x \\ J_{yy} \omega_y \\ J_{zz} \omega_z \end{bmatrix}$$

$$(\vec{h})^T (\vec{h}) = h_0^2 = J_{xx}^2 \omega_x^2 + J_{yy}^2 \omega_y^2 + J_{zz}^2 \omega_z^2$$

$$\Leftrightarrow 1 = \frac{\omega_x^2}{\frac{h_0^2}{J_{xx}^2}} + \frac{\omega_y^2}{\frac{h_0^2}{J_{yy}^2}} + \frac{\omega_z^2}{\frac{h_0^2}{J_{zz}^2}}$$

\hookrightarrow ellipsoid with half axes $\frac{h_0}{J_{ii}}$, $i=1,2,3$

We will perform rotation along axis "near" \vec{x} , \vec{y} , \vec{z} , with the same kinetic energy as a pure rotation along \vec{z} of periode $T=1s$.

$$\hookrightarrow \vec{\omega}_{ref_z} = \begin{bmatrix} 0 \\ 0 \\ \tilde{\omega}_z \end{bmatrix} \quad \tilde{\omega}_z = \frac{2\pi}{T} = 2\pi$$

$$\hookrightarrow K = \frac{1}{2} \tilde{\omega}_z^2 J_{zz} \approx 4,14e-08 \quad \frac{\text{kg} \cdot \text{m}^2 \cdot \text{rad}^2 \cdot \text{s}^{-2}}{\text{N} \cdot \text{m} \cdot \text{s}^{-4} \cdot \text{rad}^2}$$

We can find $\tilde{\omega}_x = \sqrt{\frac{2K}{J_{xx}}} \approx 3,1416 \text{ rad/s}$
 similarly $\tilde{\omega}_y = \sqrt{\frac{2K}{J_{yy}}} \approx 3,1075 \text{ rad/s}$

such that $\vec{\omega}_{ref_x} = \begin{bmatrix} \tilde{\omega}_x \\ 0 \\ 0 \end{bmatrix}$ and $\vec{\omega}_{ref_y} = \begin{bmatrix} 0 \\ \tilde{\omega}_y \\ 0 \end{bmatrix}$
 gives same kinetic energy as $\vec{\omega}_{ref_z}$

⑤ We will choose

$$\underline{\omega}_{\text{near } z}^{\text{idr}} = \begin{bmatrix} \omega_{x, \text{near } z} \\ 0 \\ \omega_z \end{bmatrix}$$

$$\omega_{\text{near } z} = \frac{\tilde{\omega}_x}{10}$$

$$\underline{\omega}_{\text{near } y}^{\text{idr}} = \begin{bmatrix} 0 \\ \omega_y \\ \omega_{z, \text{near } y} \end{bmatrix}$$

$$\omega_{z, \text{near } y} = \omega_{z, \text{near } x} = \frac{\tilde{\omega}_z}{10}$$

$$\underline{\omega}_{\text{near } x}^{\text{idr}} = \begin{bmatrix} \omega_x \\ 0 \\ \omega_{z, \text{near } x} \end{bmatrix}$$

such that the kinetic Energy induced is equal to K_0

it can be shown that

$$\begin{cases} \omega_z = \tilde{\omega}_z \sqrt{1 - \frac{\omega_{\text{near } z}^2}{\tilde{\omega}_x^2}} \\ \omega_y = \tilde{\omega}_y \sqrt{1 - \frac{1}{100}} \\ \omega_x = \tilde{\omega}_x \sqrt{1 - \frac{1}{100}} \end{cases}$$

\Rightarrow

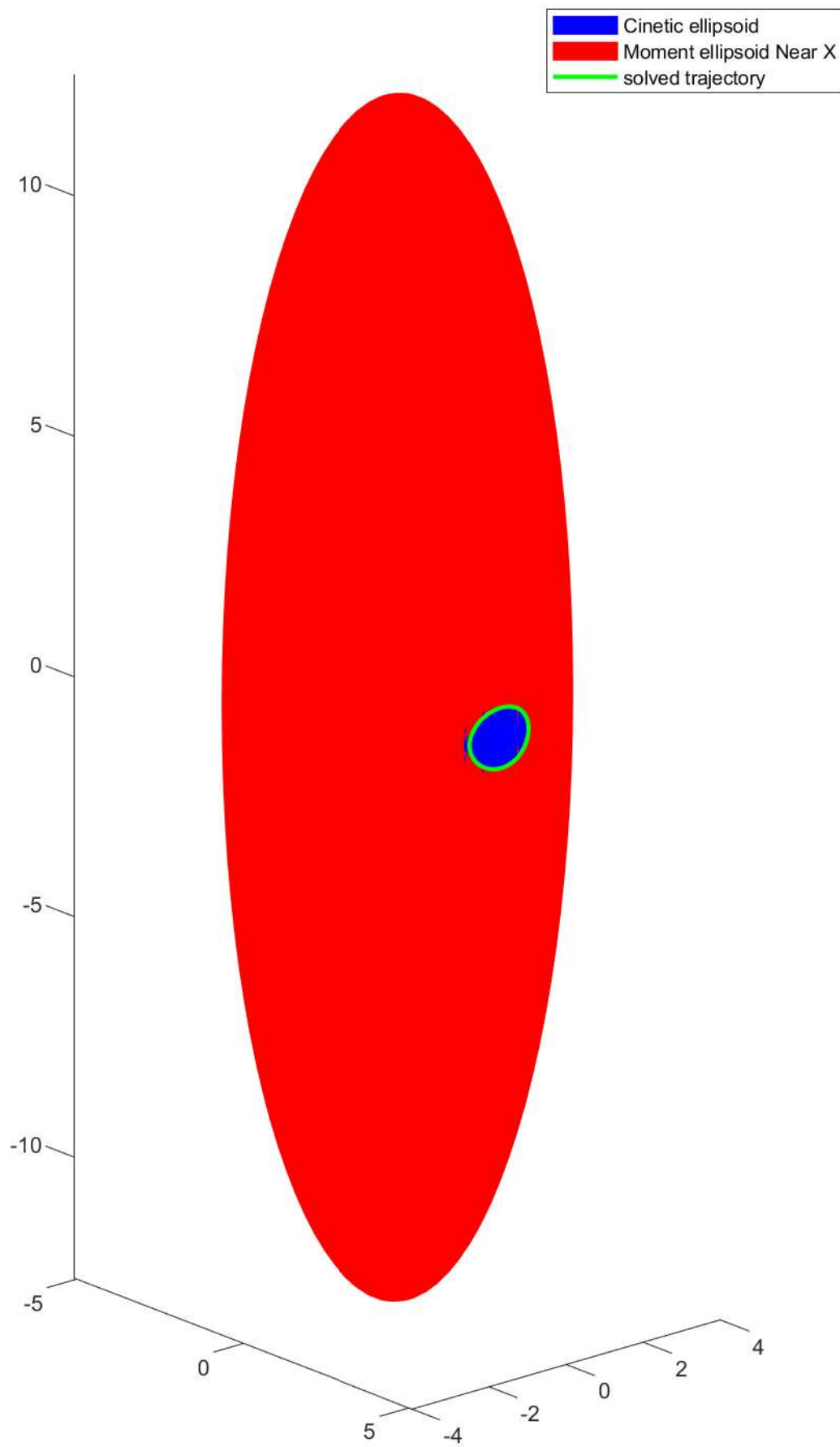
$$\begin{cases} \omega_z \approx 6.2517 \text{ rad/s} \\ \omega_y \approx 3.3905 \text{ rad/s} \\ \omega_x \approx 3.1258 \text{ rad/s} \end{cases} \quad (\text{ellipsoid available after } 3)$$

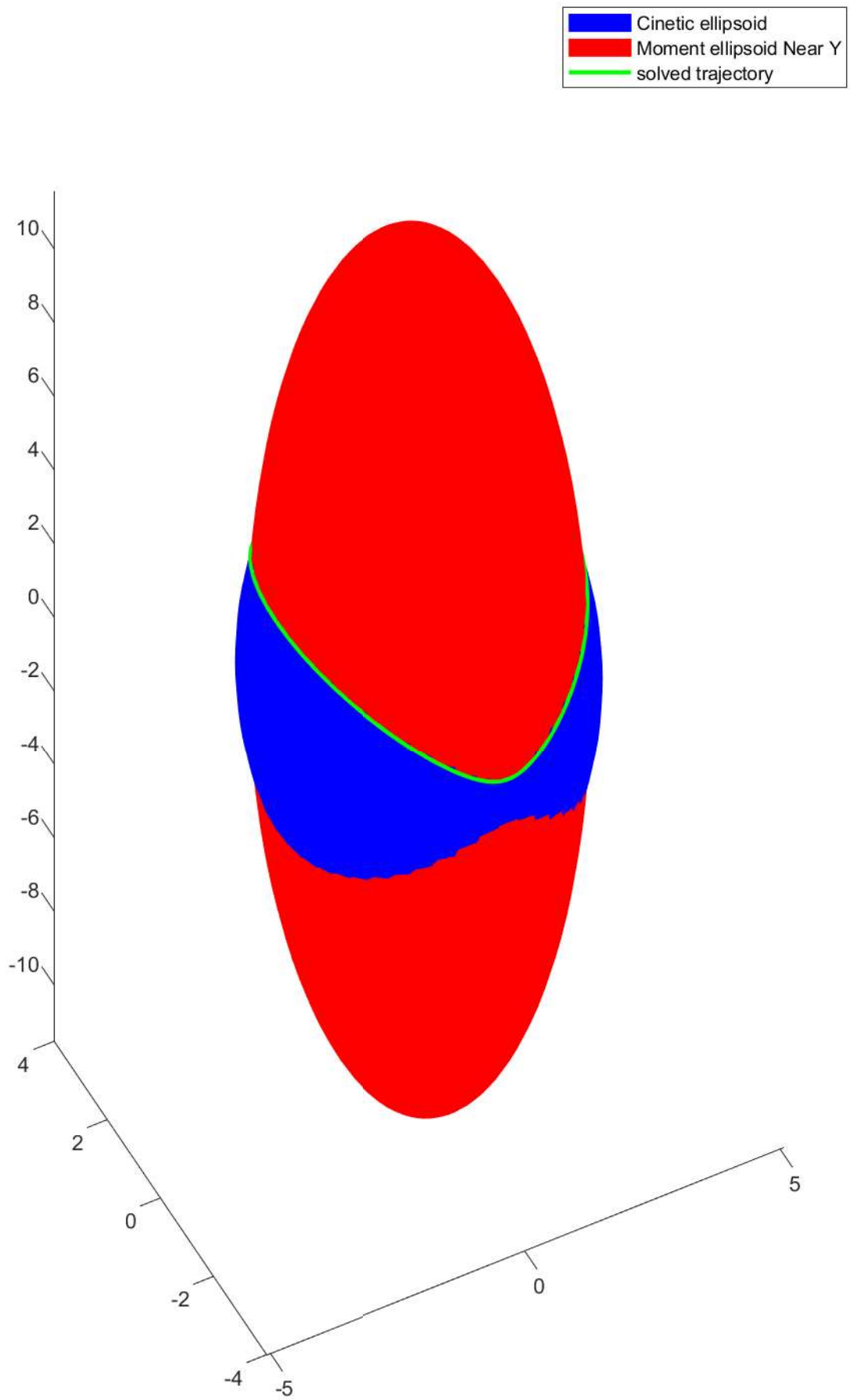
3) $\underline{\omega} = \begin{bmatrix} \omega_x \\ \omega_y \\ \omega_z \end{bmatrix}$ $A(\underline{\omega}) = \begin{bmatrix} 0 & \frac{1}{2\sqrt{J_{xx}}} (J_{yy} - J_{zz}) \omega_z & \frac{1}{2\sqrt{J_{xx}}} (J_{yy} - J_{zz}) \omega_y \\ \frac{1}{2\sqrt{J_{yy}}} (J_{zz} - J_{xx}) \omega_z & 0 & \frac{1}{2\sqrt{J_{yy}}} (J_{zz} - J_{xx}) \omega_x \\ \frac{1}{2\sqrt{J_{zz}}} (J_{xx} - J_{yy}) \omega_y & \frac{1}{2\sqrt{J_{zz}}} (J_{xx} - J_{yy}) \omega_x & 0 \end{bmatrix}$

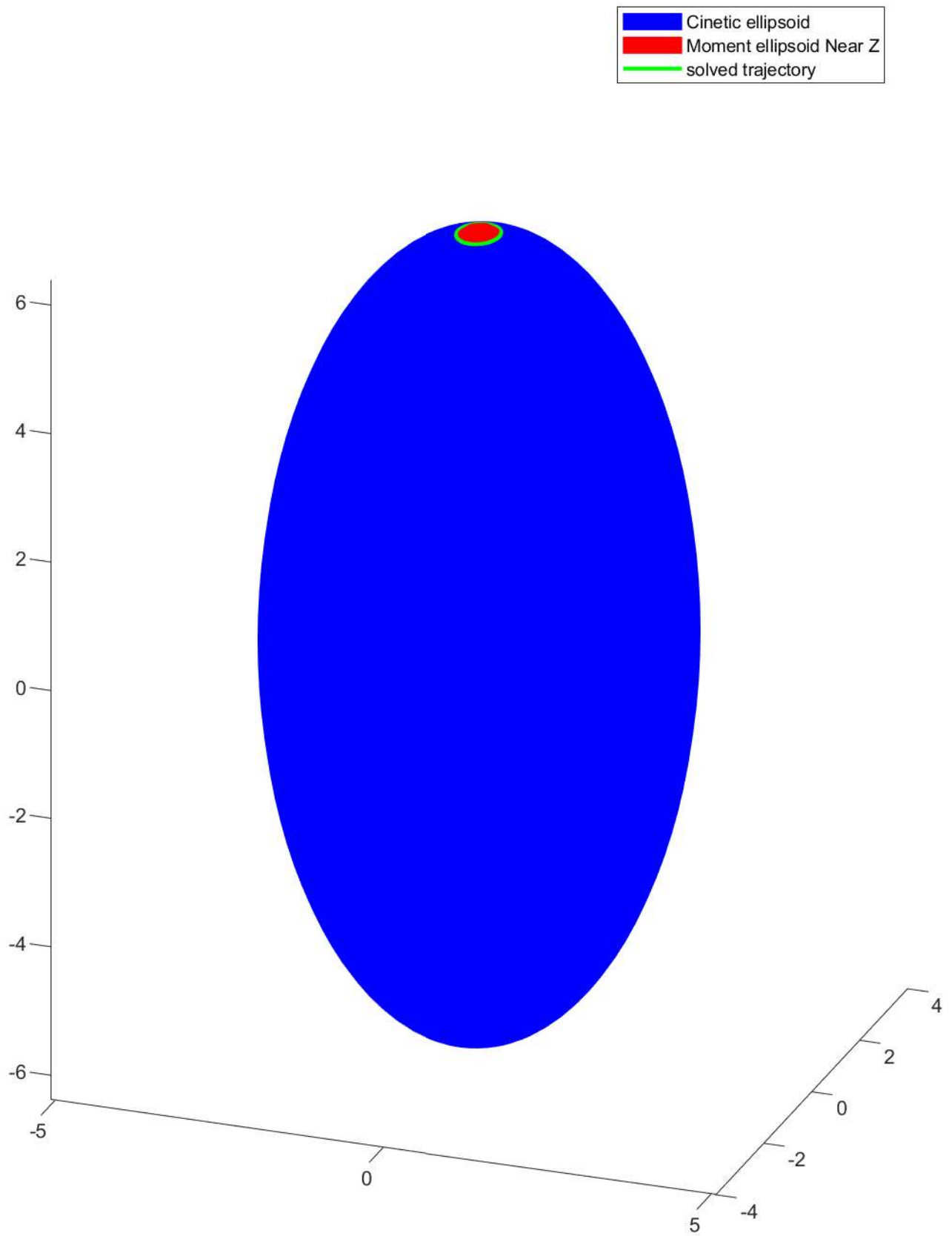
it can be shown that

(0) $\dot{\underline{\omega}} = A(\underline{\omega}) \underline{\omega}$ with euler equations

The ellipsoids with solutions from (0) are printed below







4) We have $\dot{w} = A(w) w \quad (1)$
 since $\dot{\theta} = D_{\theta}^{\circ}(\theta) w(2)$, we can find
 3-2-1
 euler
 angles $\theta(t)$ by resolving (1),
 then (2)

The corresponding Euler angles are to be found in simulation plot -

Video :

https://www.youtube.com/watch?v=8EGT9M3CTxQab_channel=Adrien

Source code :

<https://github.com/Adrien902/TEK4040-A1>