## Assignment 4

Q1) Prove the following commutation relations,

$$[\hat{L}_i, \hat{r}^2] = 0, \quad [\hat{L}_i, \hat{\vec{r}} \cdot \hat{\vec{p}}] = 0$$
 (1)

- Q2) (a) Find the matrix representation of  $L_x, L_y, L_y$  for angular momentum l=1 and l=2.
- (b) For the angular momentum 1 (l=1) system, write down the matrix representation of the following Hamiltonian in simultaneous eigenfunctions of  $\hat{L}^2$  and  $\hat{L}_z$ ,

$$\hat{H} = -\hat{L}_x + \frac{\hat{L}_z^2}{2I} \tag{2}$$

Q3) A particle in a spherically symmetric potential is in a state described by the wave packet

$$\psi(x, y, z) = C\left(xy + yz + zx\right)e^{-\alpha r^2} \tag{3}$$

- (a) Write the wave-function in terms of the Spherical Harmonics  $Y_l^m(\theta,\phi)$  with appropriate l and m.
- (b) Find out the normalization constant C?
- (c) Show that the above wave function is an eigenfunction of  $\hat{L}^2$  with l=2.
- (d) What are the relative probabilities for m = 2, 1, 0, -1, -2?
- Q4) Let  $\hat{\mathbf{n}}$  be a unit vector in a direction specified by the polar angles  $(\theta, \phi)$ . Show that the component of the angular momentum in the direction  $\hat{\mathbf{n}}$  is

$$\hat{L}_n = \sin\theta \cos\phi \hat{L}_x + \sin\theta \sin\phi \hat{L}_y + \cos\theta \hat{L}_z \tag{4}$$

If the system is in simultaneous eigenstates of  $L^2$  and  $\hat{L}_z$  belonging to the eigenvalues  $l(l+1)\hbar^2$  and  $m\hbar$ ,

- (a) what are the expectation values of  $L_n$  and  $L_n^2$ ?
- Q5) a) Prove the following,

$$\mathbf{L}^2 = \mathbf{x}^2 \mathbf{p}^2 - (\mathbf{x} \cdot \mathbf{p})^2 + i\hbar \,\mathbf{x} \cdot \mathbf{p} \tag{5}$$

where x and p are position and momentum vector operators.

(b) Using the position representation for the above  $\langle \mathbf{x}' | \mathbf{L}^2 | \alpha \rangle$ , prove the following

$$\frac{1}{2m} \langle \mathbf{x}' | \mathbf{p}^2 | \alpha \rangle = -\left(\frac{\hbar^2}{2m}\right) \left(\frac{\partial^2}{\partial r^2} \langle \mathbf{x}' | \alpha \rangle + \frac{2}{r} \frac{\partial}{\partial r} \langle \mathbf{x}' | \alpha \rangle - \frac{1}{\hbar^2 r^2} \langle \mathbf{x}' | \mathbf{L}^2 | \alpha \rangle\right)$$
(6)

Q6) a) Prove the following relation using the commutation and anti-commutation relation of Pauli matrices

$$\sigma_i \sigma_k = \delta_{ik} I + i \epsilon_{ikl} \sigma_l \tag{7}$$

(b) Using the above relation obtain the following identity

$$(\vec{a} \cdot \vec{\sigma})(\vec{b} \cdot \vec{\sigma}) = (\vec{a} \cdot \vec{b})I + i(\vec{a} \times \vec{b}) \cdot \vec{\sigma} \tag{8}$$

where  $\vec{a}$  and  $\vec{b}$  are any arbitrary vectors.

Q7) (a) Show that any  $2 \times 2$  matrix can be written in terms of Pauli matrices as,

$$A + \vec{B} \cdot \vec{\sigma} \tag{9}$$

- (b) What are the conditions that A and  $\vec{B}$  must satisfy if the matrix to be unitary. If it is to be Hermitian.
- Q8) In the Holstein-Primakoff representation, one can express the Spin matrices in terms of the bosonic creation  $\hat{a}^{\dagger}$  and annihilation operator  $\hat{a}$  as,

$$\hat{S}_z = \hbar (S - \hat{a}^{\dagger} \hat{a}) \tag{10}$$

$$\hat{S}_{+} = \hbar \sqrt{2S - a^{\dagger} \hat{a}} \hat{a} \tag{11}$$

$$\hat{S}_{-} = \hbar \hat{a}^{\dagger} \sqrt{2S - a^{\dagger} \hat{a}} \tag{12}$$

where the bosonic operators follows the commutation relation  $[\hat{a}, \hat{a}^{\dagger}] = 1$ .

(a) Using this representation show that, the Spin matrices follows commutation relation

$$[\hat{S}_i, \hat{S}_j] = i\hbar \epsilon_{ijk} \hat{S}_k \tag{13}$$

- (b) Also show that the total spin angular momentum  $\hat{S}^2 = S(S+1)\mathbb{I}$ .
- Q9) (a) Given an generalized Laguerre polynomial  $L^{\beta}_{\alpha}(x)$ , prove the following

$$\int_0^\infty e^{-x} x^{\beta} L_{\alpha}^{\beta}(x) L_{\alpha'}^{\beta}(x) dx = \frac{(\alpha + \beta)!}{\alpha!}$$
(14)

Hint! use the Generating function  $e^{-xz/(1-z)}/(1-z)^{\beta+1} = \sum_{\alpha=0}^{\infty} L_{\alpha}^{\beta}(x)z^n$  and  $(t+y)^r = \sum_{\alpha=0}^{\infty} {}^rC_{\alpha}t^{r-\alpha}y^{\alpha}$ .

(b) The radial part of the wave function of the Hydrogen atom is given by the Laguerre polynomial  $R_{nl}(r) = C(kr)^l e^{-kr/2} L_{n-l-1}^{2l+1}(kr)$ , where  $k=2/na_0$ . Calculate the normalization constant C using

$$\int_0^\infty r^2 [R_{nl}(r)]^2 dr = 1. \tag{15}$$

 $\textbf{Hint! Use the following recursion relation } xL_{\alpha}^{\beta} = (2\alpha+\beta+1)L_{\alpha}^{\beta} - (\alpha+\beta)L_{\alpha-1}^{\beta} - (\alpha+1)L_{\alpha+1}^{\beta}$ 

(c) Using the properties of Laguerre polynomial evaluate  $\langle r \rangle$  from

$$\langle r \rangle = \int_0^\infty r^2 [R_{nl}(r)]^2 dr. \tag{16}$$

Q10) Consider the Hamiltonian  $\hat{H}(\lambda)$  which depends on some parameter  $\lambda$ , follows the eigenvalue equation

$$\hat{H}(\lambda) |\psi_n(\lambda)\rangle = E_n(\lambda) |\psi_n(\lambda)\rangle \tag{17}$$

(a) If we can consider  $\lambda$  (which for example can be chosen as  $\hbar, m$  etc.) to be a continuous variable and perform derivative with respect to this parameter, show that

$$\frac{\partial E_n(\lambda)}{\partial \lambda} = \left\langle \frac{\partial \hat{H}(\lambda)}{\partial \lambda} \right\rangle \tag{18}$$

This relation is known as "Feynman-Hellmann" theorem.

(b) Consider the Hydrogen atom problem. Now using the Feynman-Hellmann theorem and by choosing the appropriate parameter  $\lambda$  show that,

$$\left\langle \frac{1}{r} \right\rangle_{nl} = \frac{Z}{a_0 n^2}, \quad \langle \hat{p}^2 \rangle = \frac{m^2 c^2 \alpha^2}{n^2} \tag{19}$$

(c) Using the above expectation values show that,  $\langle \hat{T} \rangle = -\frac{1}{2} \langle \hat{V}(r) \rangle$ , where  $\hat{T}$  is kinetic energy operator.