

Simulating phase transition bubble hydrodynamics with general equations of state

Tampere Cosmology Meeting 2023

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11.-12.5.2023

Outline

Introduction to first-order phase transitions

Phase transition bubble hydrodynamics

Equations of state

Simulation with PTtools

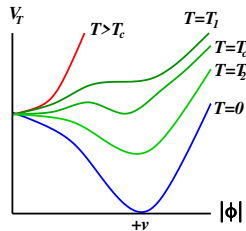
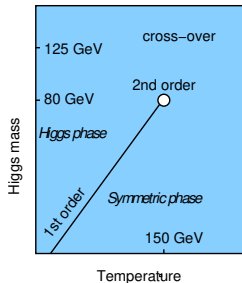
Conclusion

First-order phase transitions

- If there is a potential barrier between the local and the global minimum, the phase transition is of first order
- The Standard Model Higgs transition is a cross-over, but various extensions convert it to first order
- Klein-Gordon equation with a friction term

$$\square\phi - V_T'(\phi) = -\eta_T(\phi)u^\mu\partial_\mu\phi \quad (1)$$

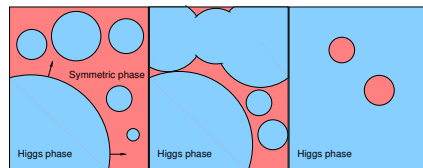
→ Constant wall speed



Hindmarsh et al., 2021

Bubble nucleation

- Spontaneous tunneling to the new phase → bubble nucleation and expansion
- Bubble expansion is governed by relativistic hydrodynamics → sound → GWs
- "Trying to determine the properties of a fluid in a water kettle based on the sound of boiling"
 - Listening to the sound indirectly through GWs
- Energy-momentum conservation gives rise to
 - Wave equation
 - Bubble wall junction conditions
 - Continuity equations, aka. hydrodynamic equations



Dimensionality of the problem

- Self-similarity
 - Friction results in a constant wall speed v_{wall}
 - As the bubble expands, its shape stays the same
 - → Time-independent solution
- Spherical symmetry
- 3+1 dimensional problem reduces to time-independent 1D

Wave equation

- Constant background space-time \rightarrow energy-momentum conservation $\nabla_\mu T^{\mu\nu} = 0$
- Energy-momentum tensor of an ideal fluid

$$T_f^{\mu\nu} = (e + p)u^\mu u^\nu + pg^{\mu\nu} \quad (2)$$

- For a one-dimensional flow in Cartesian coordinates

$$\partial_t [(e + pv^2)\gamma^2] + \partial_x [(e + p)\gamma^2 v] = 0, \quad (3)$$

$$\partial_t [(e + p)\gamma^2 v] + \partial_x [ev^2 + p]\gamma^2 = 0 \quad (4)$$

- First-order perturbation \rightarrow wave equation with speed of sound

$$\partial_t^2(\delta e) - \frac{\delta p}{\delta e} \partial_x^2(\delta e) = 0 \quad c_s^2 \equiv \frac{dp}{de} = \frac{dp/dT}{de/dT} \quad (5)$$

Phase boundary

- Energy-momentum conservation

$$\nabla_\mu T^{\mu\nu} = 0 \quad \Rightarrow \quad \partial_z T^{zz} = \partial_z T^{z0} = 0 \quad (6)$$

- Inserting ideal fluid $T_f^{\mu\nu} = (e + p)u^\mu u^\nu + pg^{\mu\nu} \rightarrow$ junction conditions

$$w_- \tilde{\gamma}_-^2 \tilde{v}_- = w_+ \tilde{\gamma}_+^2 \tilde{v}_+ \quad (7)$$

$$w_- \tilde{\gamma}_-^2 \tilde{v}_-^2 + p_- = w_+ \tilde{\gamma}_+^2 \tilde{v}_+^2 + p_+ \quad (8)$$

- By defining new variables $\theta = \frac{1}{4}(e - 3p)$, $\alpha_+ \equiv \frac{4}{3} \frac{\theta_+ (w_+) - \theta_- (w_-)}{w_+}$, $r = \frac{w_+}{w_-}$

$$\tilde{v}_+ = \frac{1}{2(1 + \alpha_+)} \left[\left(\frac{1}{3\tilde{v}_-} + \tilde{v}_- \right) \pm \sqrt{\left(\frac{1}{3\tilde{v}_-} - \tilde{v}_- \right)^2 + 4\alpha_+^2 + \frac{8}{3}\alpha_+} \right], \quad (9)$$

$$\tilde{v}_- = \frac{1}{2} \left[\left((1 + \alpha_+) \tilde{v}_+ + \frac{1 - 3\alpha_+}{3\tilde{v}_+} \right) \pm \sqrt{\left((1 + \alpha_+) \tilde{v}_+ + \frac{1 - 3\alpha_+}{3\tilde{v}_+} \right)^2 - \frac{4}{3}} \right]. \quad (10)$$

Hydrodynamic equations

- Energy-momentum conservation $\nabla_\mu T^{\mu\nu} = 0$
- Projection \rightarrow hydrodynamic equations away from the phase boundary

$$0 = u_\mu \partial_\nu T^{\mu\nu} = -\partial_\mu (w u^\mu) + u^\mu \partial_\mu p, \quad (11)$$

$$0 = \bar{u}_\mu \partial_\nu T^{\mu\nu} = w \bar{u}^\nu u^\mu \partial_\mu u_\nu + \bar{u}^\mu \partial_\mu p. \quad (12)$$

- Using self-similarity $\xi = \frac{r}{t}$ and parametrising

$$\frac{d\xi}{d\tau} = \xi [(\xi - v)^2 - c_s^2(1 - \xi v)^2], \quad (13)$$

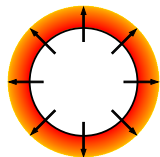
$$\frac{dv}{d\tau} = 2v c_s^2(1 - v^2)(1 - \xi v), \quad (14)$$

$$\frac{dw}{d\tau} = w \left(1 + \frac{1}{c_s^2}\right) \gamma^2 \mu \frac{dv}{d\tau}. \quad (15)$$

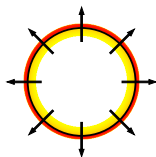
- Note that c_s^2 is computed using user-provided functions

Fluid shells

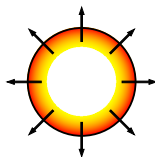
- Three types of solutions, determined by
 - Wall velocity v_{wall}
 - Transition strength α_n
 - Speed of sound $c_s(T, \phi)$



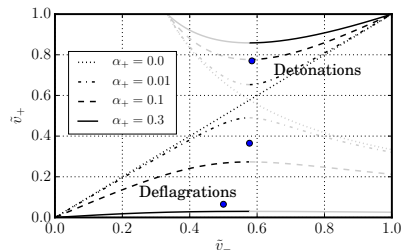
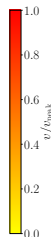
subsonic deflagration
 $v_w \leq c_s$



supersonic deflagration
 $c_s < v_w < c_J$



detonation
 $c_J \leq v_w$



Black = bubble wall / phase boundary

General equation of state

- Equation of state for an ultrarelativistic plasma with multiple degrees of freedom

$$p(T, \phi) = \frac{\pi^2}{90} g_p(T, \phi) T^4 - V_0(\phi) \quad (16)$$

- The rest can be deduced with thermodynamics
 - Entropy density $s = \frac{dp}{dT}$
 - Enthalpy density $w = Ts$
 - Energy density $e = w - p$
 - Sound speed $c_s = \sqrt{\frac{dp}{de}}$

Simple models

- Bag model: equation of state with constant degrees of freedom

$$g_{\pm} = \frac{90}{\pi^2} a_{\pm} \quad \Rightarrow \quad p_{\pm} = a_{\pm} T^4 - V_{\pm}, \quad c_s^2 \equiv \frac{dp}{de} = \frac{1}{3} \quad (17)$$

- Constant sound speed model

$$g_{p\pm} = \frac{90}{\pi^2} a_{\pm} \left(\frac{T}{T_0} \right)^{\mu_{\pm}-4} \quad (18)$$

$$\Rightarrow \quad p_{\pm} = a_{\pm} \left(\frac{T}{T_0} \right)^{\mu_{\pm}-4} T^4 - V_{\pm} \approx a_{\pm} T^{\mu_{\pm}} - V_{\pm} \quad (19)$$

$$c_{s\pm}^2 = \frac{1}{\mu_{\pm} - 1} \quad (20)$$

Equation of state from an arbitrary particle physics model

- Non-constant degrees of freedom \rightarrow varying sound speed $c_s(T, \phi)$
- The equation of state can be constructed from $V_0(\phi)$, $g_p(T, \phi)$ and preferably also $g_e(T, \phi)$ or $g_s(T, \phi)$

$$g_p = 4g_s - 3g_e \quad (21)$$

$$e(T, \phi) = \frac{\pi^2}{30} g_e(T, \phi) T^4 + V_0(\phi) \quad (22)$$

$$p(T, \phi) = \frac{\pi^2}{90} g_p(T, \phi) T^4 - V_0(\phi) \quad (23)$$

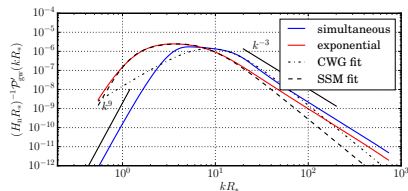
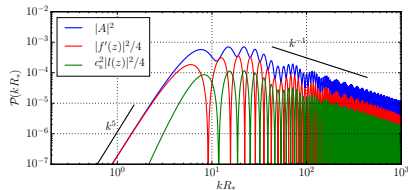
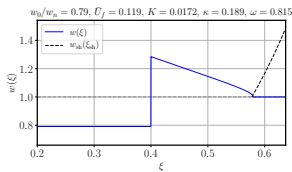
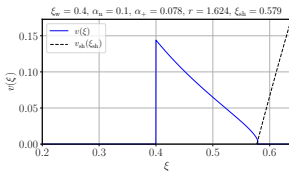
$$s(T, \phi) = \frac{2\pi^2}{45} g_s(T, \phi) T^3 \quad (24)$$

PTtools

- Python-based simulation framework for GW power spectra of first-order phase transitions
 - Based on Numpy, SciPy (solvers & ODE integration) & Numba
- Input
 - Equation of state, either directly as $p(T, \phi)$ (and $e(T, \phi)$ or $s(T, \phi)$), or as degrees of freedom: $g_p(T, \phi)$ (and $g_e(T, \phi)$ or $g_s(T, \phi)$), and potential $V_0(\phi)$
 - Bag model, constant sound speed model and the Standard Model with adjustable phases are already included
 - User provides the model as a Python class. The g_{eff} functions can be arbitrary, e.g. interpolated from data.
 - Bubble wall speed v_{wall}
 - Transition strength parameter α_n
 - Transition rate parameter β
- Output: gravitational wave power spectrum in wavenumber units of mean bubble spacing
- Use case: estimate the likelihood of various Standard Model extensions with future LISA data in the 2030s

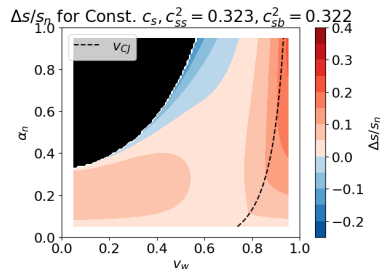
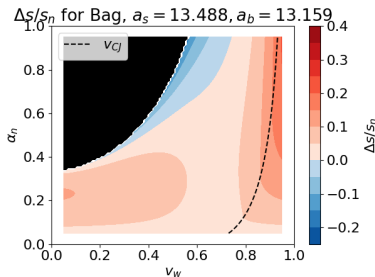
Gravitational wave production

- Sound Shell Model (Hindmarsh, 2018)
- Velocity profile of a single bubble $v_{ip}(\xi) + \text{sine transform}$
→ velocity power spectrum
- Convolution with the nucleation rate function
→ Overall velocity power spectrum
- Convolution with a Green's function → GW power spectrum



Entropy conservation

- 10 000 bubbles with varying $v_{\text{wall}}, \alpha_n$
- Top-left region ruled out by $\alpha_+ > \frac{1}{3}$
- Unphysical region (blue): decrease in total entropy



Summary

- A phase transition converts energy from the scalar field to kinetic and thermal energy → Gravitational wave production
- The phase transition proceeds by bubble nucleation
- The hydrodynamics is based on energy-momentum conservation
- Solving a bubble starts with the equation of state
 - Equation solving & ODE integration → fluid velocity profile
- Sound shell model converts the velocity profile to GW power spectrum
 - Sine transform & convolutions → GW power spectrum
- PTtools has been updated to support arbitrary equations of state
- PTtools will be published soon

Thank you!

Sources

- M. Hindmarsh, M. Lüben, J. Lumma, and M. Pauly, Phase transitions in the early universe, SciPost Phys. Lect. Notes, Feb. 2021, doi: 10.21468/SciPostPhysLectNotes.24.
- Hindmarsh, Mark, and Mulham Hijazi. Gravitational Waves from First Order Cosmological Phase Transitions in the Sound Shell Model. Journal of Cosmology and Astroparticle Physics 2019, Dec. 2019, doi: 10.1088/1475-7516/2019/12/062.
- Caprini et al. Detecting Gravitational Waves from Cosmological Phase Transitions with LISA: An Update. Journal of Cosmology and Astroparticle Physics 2020, Mar. 2020, doi: 10.1088/1475-7516/2020/03/024.
- Espinosa et al. Energy Budget of Cosmological First-Order Phase Transitions. Journal of Cosmology and Astroparticle Physics 2010, Jun. 2010, doi: 10.1088/1475-7516/2010/06/028.
- Giese et al., Model-Independent Energy Budget for LISA. Journal of Cosmology and Astroparticle Physics, Jan. 2021, doi: 10.1088/1475-7516/2021/01/072.
- Borsanyi, Sz, Z. Fodor, K. H. Kampert, S. D. Katz, T. Kawanai, T. G. Kovacs, S. W. Mages, et al. Lattice QCD for Cosmology., Jun. 2016, ArXiv: 1606.07494

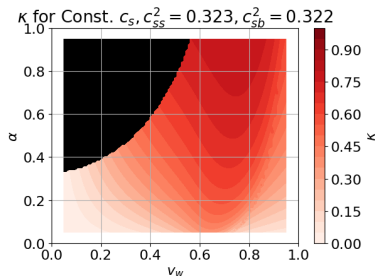
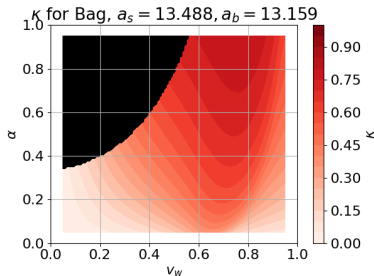
Energy redistribution

- Efficiency factor κ : ratio of kinetic energy to total energy released

$$\kappa = \frac{e_K}{|\Delta e_\theta|} \quad (25)$$

$$e_K = 4\pi \int_0^{\xi_{\max}} d\xi \xi^2 w \gamma^2 v^2 \quad (26)$$

$$\Delta e_\theta = 4\pi \int_0^{\xi_{\max}} d\xi \xi^2 (\theta - \theta_n) \quad (27)$$



Speedup example with Numba

- Add JIT decorator
- Replace unsupported features with simpler code or split the unsupported parts to another function
- (Restructure the function to make the possible parallelism explicit)

`@numba.njit(parallel=True)`

```
def sin_transform(t: np.ndarray, f: np.ndarray, freq: np.ndarray) -> np.ndarray:  
    integral = np.zeros_like(freq)  
    for i in numba.prange(freq.size):  
        integrand = f * np.sin(freq[i] * t)  
        integral[i] = np.trapz(integrand, t)  
    return integral
```