

Math 228B - HW4

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15 March 2017

1 Problem No.1

1.1 Problem Description:

Write programs to solve the advection equation

$$u_t + au_x = 0,$$

on $[0, 1]$ with periodic boundary conditions using upwinding and Lax-Wendroff. For smooth solutions we expect upwinding to be first-order accurate and Lax-Wendroff to be second-order accurate, but it is not clear what accuracy to expect for nonsmooth solution.

1. Let $a = 1$ and solve the problem up to time $t = 1$. Perform a refinement study for both upwinding and Lax-Wendroff with $\Delta t = 0.9a\Delta x$ with a smooth initial condition. Compute the rate of converge in the 1-norm and 2-norm, and max-norm. Note that the exact solution at time $t = 1$ is the initial condition, and so computing the error is easy.
2. Repeat the previous problem with the discontinuous initial condition

$$u(x, 0) = \begin{cases} 1 & \text{if } |x - \frac{1}{2}| < \frac{1}{4} \\ 0 & \text{otherwise} \end{cases}$$

1.2 Solution:

Part 1: Under upwinding scheme, the advection equation is discretized as:

$$\frac{u_j^{n+1} - u_j^n}{\Delta t} + a \frac{u_j^n - u_{j-1}^n}{\Delta x} = 0$$
$$u_j^{n+1} = u_j^n - \frac{a\Delta t}{\Delta x}(u_j^n - u_{j-1}^n)$$

For which we can define *Courant* number ν as $\nu = a\Delta t/\Delta x$. For Lax-Wendroff scheme, the discretization would be

$$\frac{u_j^{n+1} - u_j^n}{\Delta t} + a \frac{u_{j+1}^n - u_{j-1}^n}{2\Delta x} - a^2 \Delta t \frac{u_j^{n-1} - 2u_j^n + u_{j+1}^n}{2\Delta x^2} = 0$$
$$u_j^{n+1} = u_j^n - \frac{a\Delta t}{2\Delta x}(u_{j+1}^n - u_{j-1}^n) + \frac{a^2 \Delta t^2}{2\Delta x^2}(u_{j-1}^n - 2u_j^n + u_{j+1}^n)$$

For smooth initial conditions we use $\cos(2.0 * x * PI) + \sin(6.0 * x * PI)/2.0$ as initial conditions which is smooth continuous function for the periodic domain $[0,1]$. The periodic boundary conditions mean that the value of the last grid point is equal to the value of the first grid point. We can consider this as solving over a circle.

Since the problem asking for solving such that $\Delta t = 0.9a\Delta x$ and $a = 1.0$, thus $\nu = 0.9$ and the numerical scheme is stable. This applies for upwinding and Lax-Wendroff.

The error was calculated for 1-norm as $\|e\|_1 = h \sum_j |e_j|$, where e_j is the absolute difference (error) between the numerical solution and exact solution at grid point j . For 2-norm, the error is $\|e\|_2 = h \sqrt{\sum_j |e_j|^2}$. For max-norm, the error becomes $\|e\|_{max} = \max_j |e_j|$.

The exact and numerical solutions for both the smooth and discontinuous initial conditions are shown in Figure 1.

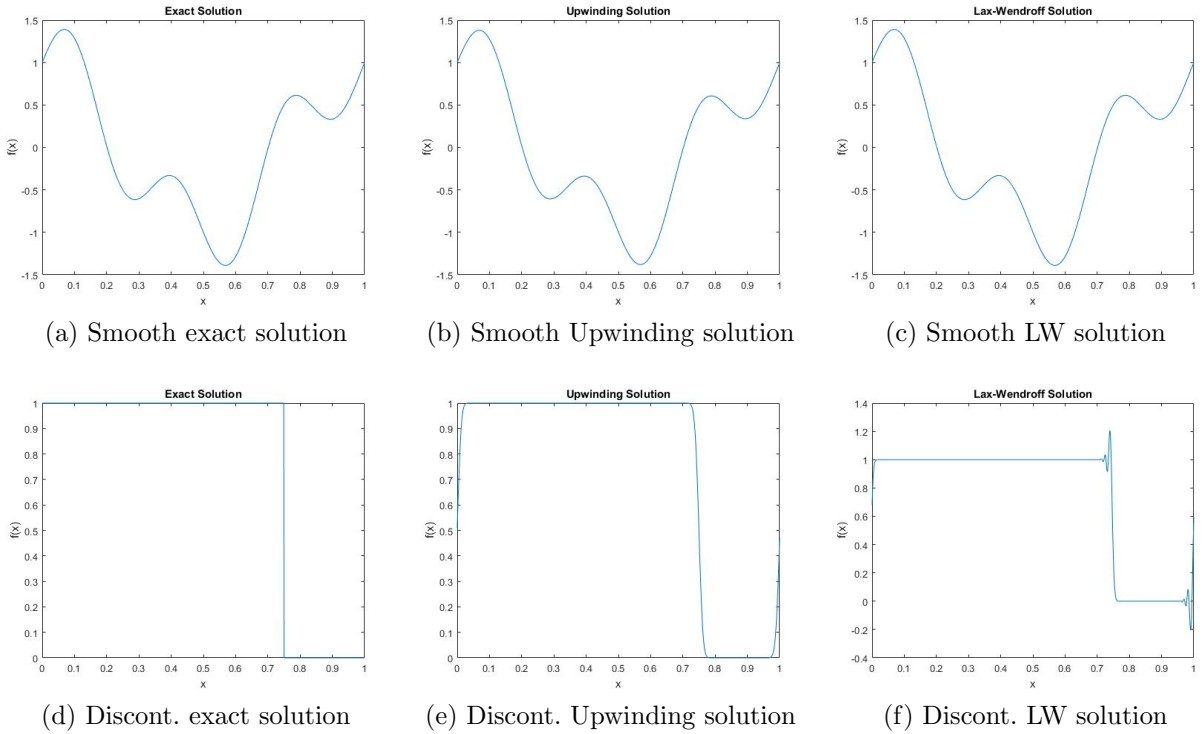


Figure 1: The exact solution (first column), upwinding solution (second column) and Lax-Wendroff solution (third column) for both smooth (first row) and discontinuous (second row) initial conditions for the 1d advection equation.

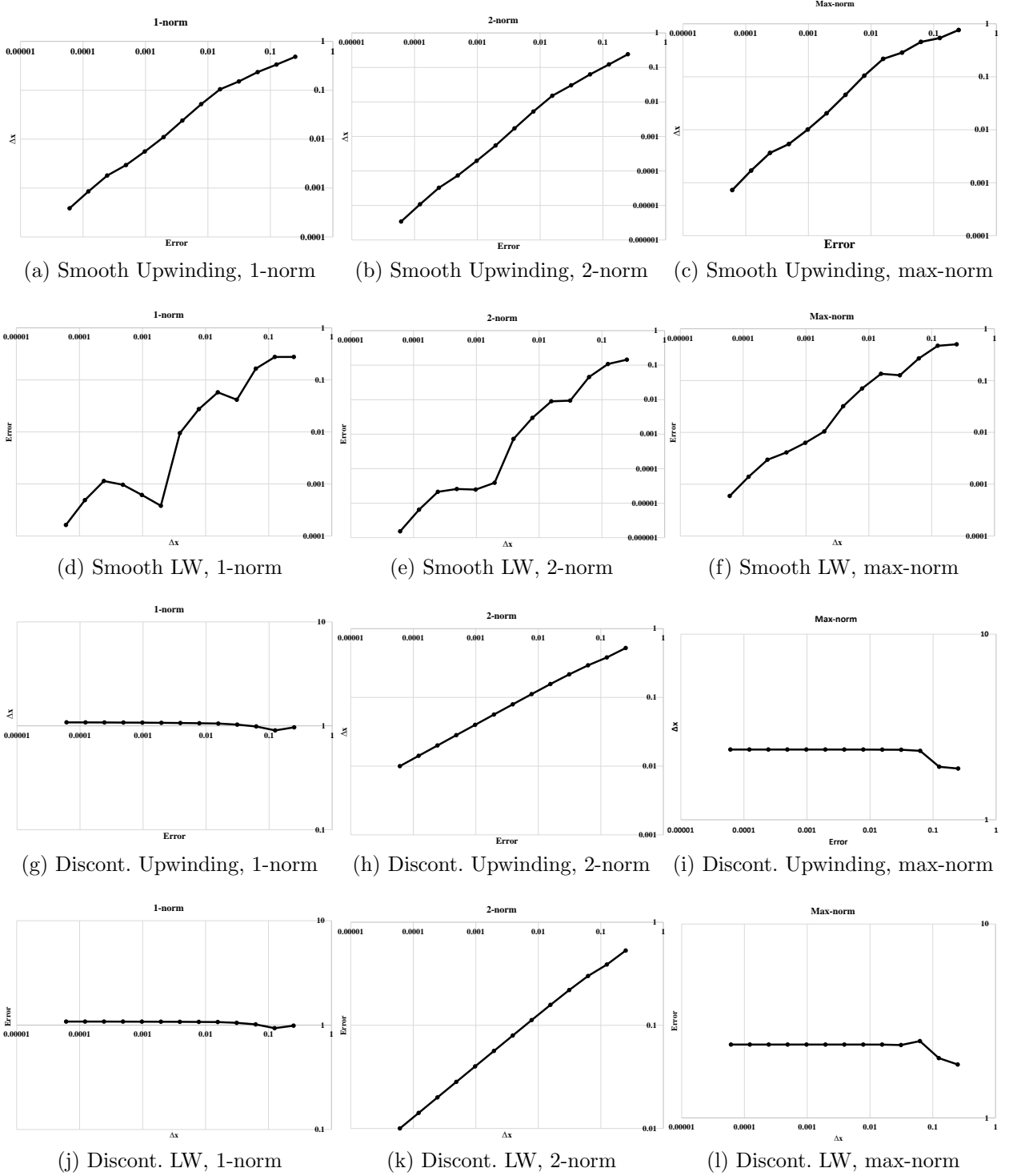


Figure 2: The rate convergence in 1-norm (first column), 2-norm (second column) and max-norm (third column) for smooth (first two rows) and discontinuous (last two rows) initial condition using upwinding and Lax-Wendroff scheme to solve the 1d advection equation.

2 Problem No.2

2.1 Problem Description:

For solving the heat equation we frequently use Crank-Nicolson, which is trapezoidal rule time integration with a second-order space discretization. The analogous scheme for the linear advection equation

is

$$u_j^{n+1} - u_j^n + \frac{\nu}{4}(u_{j+1}^n - u_{j-1}^n) + \frac{\nu}{4}(u_{j+1}^{n+1} - u_{j-1}^{n+1}) = 0,$$

where $\nu = a\Delta t/\Delta x$

1. Use von Neumann analysis to show that this scheme is unconditionally stable and that $\|u^n\|_2 = \|u^0\|_2$. This scheme is said to be non-dissipative - i.e. there is no amplitude error. This seems reasonable because this is a property of the PDE.
2. Solve the advection equation on the periodic domain $[0, 1]$ with the initial condition from problem 1 (part 2). Show the solution and comment on your results.
3. Compute the relative phase error as $\arg(g(\theta)/(-\nu\theta))$, where g is the amplification factor and $\theta = \zeta\Delta t$, and plot it for $\theta \in [0, \pi]$. How does the relative phase error and lack of amplitude error relate to the numerical solutions you observed in part 2.

2.2 Solution:

Part 1:

Appendix