# PRISMS-PF: Precipitate Evolution

#### 1 Variational formulation

The total free energy of the system (neglecting boundary terms) is of the form,

$$\Pi(c, \eta_1, \eta_2, \eta_3, \epsilon) = \int_{\Omega} f(c, \eta_1, \eta_2, \eta_3, \epsilon) \ dV$$
 (1)

where c is the concentration of the  $\beta$  phase,  $\eta_p$  are the structural order parameters and  $\varepsilon$  is the small strain tensor. f, the free energy density is given by

$$f(c, \eta_1, \eta_2, \eta_3, \epsilon) = f_{chem}(c, \eta_1, \eta_2, \eta_3) + f_{grad}(\eta_1, \eta_2, \eta_3) + f_{elastic}(c, \eta_1, \eta_2, \eta_3, \epsilon)$$
(2)

where

$$f_{chem}(c, \eta_1, \eta_2, \eta_3) = f_{\alpha}(c) \left( 1 - H(\eta_1) - H(\eta_2) - H(\eta_3) \right) + f_{\beta}(c) \left( H(\eta_1) + H(\eta_2) + H(\eta_3) \right) + W f_{Landau}(\eta_1, \eta_2, \eta_3)$$
(3)

$$f_{grad}(\eta_1, \eta_2, \eta_3) = \frac{1}{2} \sum_{p=1}^{3} \kappa_{ij}^{\eta_p} \eta_{p,i} \eta_{p,j}$$
(4)

$$f_{elastic}(c, \eta_1, \eta_2, \eta_3, \boldsymbol{\epsilon}) = \frac{1}{2} \boldsymbol{C}_{ijkl}(\eta_1, \eta_2, \eta_3) \left( \boldsymbol{\varepsilon}_{ij} - \boldsymbol{\varepsilon}_{ij}^0(c, \eta_1, \eta_2, \eta_3) \right) \left( \boldsymbol{\varepsilon}_{kl} - \boldsymbol{\varepsilon}_{kl}^0(c, \eta_1, \eta_2, \eta_3) \right)$$
(5)

$$\varepsilon^{0}(c, \eta_{1}, \eta_{2}, \eta_{3}) = H(\eta_{1})\varepsilon_{\eta_{1}}^{0}(c) + H(\eta_{2})\varepsilon_{\eta_{2}}^{0}(c) + H(\eta_{3})\varepsilon_{\eta_{3}}^{0}(c)$$
(6)

$$C(\eta_1, \eta_2, \eta_3) = H(\eta_1)C_{\eta_1} + H(\eta_2)C_{\eta_2} + H(\eta_3)C_{\eta_3} + (1 - H(\eta_1) - H(\eta_2) - H(\eta_3))C_{\alpha}$$
(7)

Here  $\varepsilon_{\eta_p}^0$  are the composition dependent stress free strain transformation tensor corresponding to each structural order parameter.

# 2 Required inputs

- $f_{\alpha}(c), f_{\beta}(c)$  Homogeneous chemical free energy of the components of the binary system, example form given in Appendix I
- $f_{Landau}(\eta_1, \eta_2, \eta_3)$  Landau free energy term that controls the interfacial energy and prevents precipitates with different orientation varients from overlapping, example form given in Appendix I
- W Barrier height for the Landau free energy term, used to control the thickness of the interface
- $H(\eta_p)$  Interpolation function for connecting the  $\alpha$  phase and the  $p^{th}$  orientation variant of the  $\beta$  phase, example form given in Appendix I
- $\kappa^{\eta_p}$  gradient penalty tensor for the  $p^{th}$  orientation variant of the  $\beta$  phase
- $C_{\eta_p}$  fourth order elasticity tensor (or its equivalent second order Voigt representation) for the  $p^{th}$  orientation variant of the  $\beta$  phase
- $C_{\alpha}$  fourth order elasticity tensor (or its equivalent second order Voigt representation) for the  $\alpha$  phase
- $\varepsilon_{n_p}^0$  stress free strain transformation tensor for the  $p^{th}$  orientation variant of the  $\beta$  phase

In addition, to drive the kinetics, we need:

- M mobility value for the concentration field
- L mobility value for the structural order parameter field

### 3 Variational treatment

From the variational derivatives given in Appendix II, we obtain the chemical potentials for the concentration and the structural order parameters:

$$\mu_{c} = f_{\alpha,c} \left( 1 - H(\eta_{1}) - H(\eta_{2}) - H(\eta_{3}) \right) + f_{\beta,c} \left( H(\eta_{1}) + H(\eta_{2}) + H(\eta_{3}) \right) + \boldsymbol{C}_{ijkl} \left( -\boldsymbol{\varepsilon}_{ij,c}^{0} \right) \left( \boldsymbol{\varepsilon}_{kl} - \boldsymbol{\varepsilon}_{kl}^{0} \right)$$
(8)  
$$\mu_{\eta_{p}} = \left( f_{\beta} - f_{\alpha} \right) H(\eta_{p})_{,\eta_{p}} + W f_{Landau,\eta_{p}} - \boldsymbol{\kappa}_{ij}^{\eta_{p}} \eta_{p,ij} + \boldsymbol{C}_{ijkl} \left( -\boldsymbol{\varepsilon}_{ij,\eta_{p}}^{0} \right) \left( \boldsymbol{\varepsilon}_{kl} - \boldsymbol{\varepsilon}_{kl}^{0} \right) + \frac{1}{2} \boldsymbol{C}_{ijkl,\eta_{p}} \left( \boldsymbol{\varepsilon}_{ij} - \boldsymbol{\varepsilon}_{ij}^{0} \right) \left( \boldsymbol{\varepsilon}_{kl} - \boldsymbol{\varepsilon}_{kl}^{0} \right)$$
(9)

#### 4 Kinetics

Now the PDE for Cahn-Hilliard dynamics is given by:

$$\frac{\partial c}{\partial t} = \nabla \cdot (M \nabla \mu_c) \tag{10}$$

and the PDE for Allen-Cahn dynamics is given by:

$$\frac{\partial \eta_p}{\partial t} = -L\mu_{\eta_p} \tag{11}$$

where M and L are the constant mobilities.

#### 5 Mechanics

Considering variations on the displacement u of the from  $u + \epsilon w$ , we have

$$\delta_u \Pi = \int_{\Omega} \nabla w : C(\eta_1, \eta_2, \eta_3) : \left( \varepsilon - \varepsilon^0(c, \eta_1, \eta_2, \eta_3) \right) dV = 0$$
 (12)

(13)

where  $\sigma = C(\eta_1, \eta_2, \eta_3) : (\varepsilon - \varepsilon^0(c, \eta_1, \eta_2, \eta_3))$  is the stress tensor.

Now consider

$$R = \int_{\Omega} \nabla w : C(\eta_1, \eta_2, \eta_3) : \left( \varepsilon - \varepsilon^0(c, \eta_1, \eta_2, \eta_3) \right) dV = 0$$
(14)

We solve for R=0 using a gradient scheme which involves the following linearization:

$$R\mid_{u} + \frac{\partial R}{\partial u} \Delta u = 0 \tag{15}$$

$$\Rightarrow \frac{\partial R}{\partial u} \Delta u = -R \mid_{u} \tag{16}$$

This is the linear system Ax = b which we solve implicitly using the Conjugate Gradient scheme. For clarity, here in the left hand side (LHS)  $A = \frac{\partial R}{\partial u}$ ,  $x = \Delta u$  and the right hand side (RHS) is  $b = -R \mid_{u}$ .

#### 6 Time discretization

Using forward Euler explicit time stepping, equations ?? and ?? become:

$$c^{n+1} = c^n + \Delta t [\nabla \cdot (M \nabla \mu_c)] \tag{17}$$

$$\eta_p^{n+1} = \eta_p^n - \Delta t L \mu_{\eta_p} \tag{18}$$

### 7 Weak formulation

Writing equations  $\ref{eq:condition}$  and  $\ref{eq:condition}$  in the weak form, with the arbitrary variation given by w yields:

$$\int_{\Omega} wc^{n+1}dV = \int_{\Omega} wc^n + w\Delta t [\nabla \cdot (M\nabla \mu_c)]dV$$
(19)

$$\int_{\Omega} w \eta_p^{n+1} dV = \int_{\Omega} w \eta_p^n - w \Delta t L \mu_{\eta_p} dV \tag{20}$$

The gradient of  $\mu_c$  is:

$$\nabla \mu_{c} = \nabla c \left[ f_{\alpha,cc} + \sum_{p=1}^{3} H(\eta_{p}) (f_{\beta,cc} - f_{\alpha,cc}) \right] + \sum_{p=1}^{3} \nabla \eta_{p} H(\eta_{p})_{,\eta_{p}} (f_{\beta,c} - f_{\alpha,c})$$

$$+ \left[ \sum_{p=1}^{3} (C_{ijkl}^{\eta_{p}} - C_{ijkl}^{\alpha}) \nabla \eta_{p} H(\eta_{p})_{,\eta_{p}} \right] (-\epsilon_{ij,c}^{0}) (\epsilon_{ij} - \epsilon_{ij}^{0})$$

$$- C_{ijkl} \left[ \sum_{p=1}^{3} H(\eta_{p})_{,\eta_{p}} \epsilon_{ij,c}^{0\eta_{p}} \nabla \eta_{p} + H(\eta_{p}) \epsilon_{ij,cc}^{0\eta_{p}} \nabla c \right] (\epsilon_{kl} - \epsilon_{kl}^{0})$$

$$+ C_{ijkl} (-\epsilon_{ij,c}^{0}) \left[ \nabla \epsilon_{ij} - \left( \sum_{p=1}^{3} H(\eta_{p})_{,\eta_{p}} \epsilon_{kl}^{0\eta_{p}} \nabla \eta_{p} + H(\eta_{p}) \epsilon_{kl,c}^{0\eta_{p}} \nabla c \right) \right]$$

$$(21)$$

Applying the divergence theorem to equation ??, one can derive the residual terms  $r_c$  and  $r_{cx}$ :

$$\int_{\Omega} w c^{n+1} dV = \int_{\Omega} w \underbrace{c^n}_{r_c} + \nabla w \cdot (\underbrace{-\Delta t M \nabla \mu_c}_{r}) dV$$
 (22)

Expanding  $\mu_{\eta_p}$  in equation ?? and applying the divergence theorem yields the residual terms  $r_{\eta_p}$  and  $r_{\eta_p x}$ :

$$\int_{\Omega} w \eta_{p}^{n+1} dV = \int_{\Omega} w \left\{ \underbrace{\eta_{p}^{n} - \Delta t L \left[ (f_{\beta} - f_{\alpha}) H(\eta_{p}^{n})_{,\eta_{p}} + W f_{Landau,\eta_{p}} - C_{ijkl} \left( H(\eta_{p})_{,\eta_{p}} \epsilon_{ij}^{0\eta_{p}} \right) \left( \epsilon_{kl} - \epsilon_{kl}^{0} \right) \right.} + \underbrace{\frac{1}{2} \left[ \left[ (C_{ijkl}^{\eta_{p}} - C_{ijkl}^{\alpha}) H(\eta_{p})_{,\eta_{p}} \right] \left( \epsilon_{ij} - \epsilon_{ij}^{0} \right) \left( \epsilon_{kl} - \epsilon_{kl}^{0} \right) \right] \right\}}_{r_{\eta_{p}} cont.} + \nabla w \cdot \underbrace{\left( -\Delta t L \kappa_{ij}^{\eta_{p}} \eta_{p,i}^{n} \right) dV}_{r_{\eta_{p},x}} dV$$
(23)

# 8 Appendix I: Example functions for $f_{\alpha}$ , $f_{\beta}$ , $f_{Landau}$ , $H(\eta_p)$

$$f_{\alpha}(c) = A_{2,\alpha}c^2 + A_{1,\alpha}c + A_{0,\alpha} \tag{24}$$

$$f_{\beta}(c) = A_{2,\beta}c^2 + A_{1,\beta}c + A_{0,\beta} \tag{25}$$

$$f_{Landau}(\eta_1, \eta_2, \eta_3) = (\eta_1^2 + \eta_2^2 + \eta_3^2) - 2(\eta_1^3 + \eta_2^3 + \eta_3^3) + (\eta_1^4 + \eta_2^4 + \eta_3^4) + 5(\eta_1^2 \eta_2^2 + \eta_2^2 \eta_3^2 + \eta_1^2 \eta_3^2) + 5(\eta_1^2 \eta_2^2 \eta_3^2)$$
(26)

$$H(\eta_p) = 3\eta_p^2 - 2\eta_p^3 \tag{27}$$

## 9 Appendix II: Variational Derivatives

Variational derivative of  $\Pi$  with respect to  $\eta_p$  (where  $\eta_q$  and  $\eta_r$  correspond to the structural order parameters for the other two orientational variants):

$$\delta_{\eta_p} \Pi = \frac{d}{d\alpha} \left[ \int_{\Omega} f_{chem}(c, \eta_p + \alpha w, \eta_q, \eta_r) + f_{grad}(\eta_p + \alpha w, \eta_q, \eta_r) + f_{el}(c, \eta_p + \alpha w, \eta_q, \eta_r, \epsilon) dV \right]_{\alpha = 0}$$
(28)

Breaking up each of these terms yields:

$$\frac{d}{d\alpha} \left[ f_{chem}(c, \eta_p + \alpha w, \eta_q, \eta_r) \right]_{\alpha=0} = f_{\alpha}(c) \left[ -\frac{\partial H(\eta_p + \alpha w)}{\partial (\eta_p + \alpha w)} \frac{\partial (\eta_p + \alpha w)}{\partial \alpha} \right]_{\alpha=0} 
+ f_{\beta}(c) \left[ \frac{\partial H(\eta_p + \alpha w)}{\partial (\eta_p + \alpha w)} \frac{\partial (\eta_p + \alpha w)}{\partial \alpha} \right]_{\alpha=0} 
+ W \left[ \frac{\partial f_{Landau}(\eta_p + \alpha w, \eta_q, \eta_r)}{\partial (\eta_p + \alpha w)} \frac{\partial (\eta_p + \alpha w)}{\partial \alpha} \right]_{\alpha=0} 
= f_{\alpha}(c) \left[ -\frac{\partial H(\eta_p)}{\partial \eta_p} w \right] + f_{\beta}(c) \left[ \frac{\partial H(\eta_p)}{\partial \eta_p} w \right] + W \left[ \frac{\partial f_{Landau}(\eta_p, \eta_q, \eta_r)}{\partial \eta_p} w \right]$$
(29)

$$\frac{d}{d\alpha} \left[ f_{grad}(\eta_p + \alpha w, \eta_q, \eta_r) \right]_{\alpha=0} = \frac{1}{2} \left[ \kappa_{ij}^{\eta_p} (\eta_p + \alpha w)_{,i} (\eta_p + \alpha w)_{,j} + \kappa_{ij}^{\eta_q} (\eta_q)_{,i} (\eta_q)_{,j} + \kappa_{ij}^{\eta_r} (\eta_r)_{,i} (\eta_r)_{,j} \right]_{\alpha=0} \\
= \kappa_{ij} w_{,i} \eta_{p,j} \tag{30}$$

$$\frac{d}{d\alpha} \left[ f_{el}(c, \eta_p + \alpha w, \eta_q, \eta_r, \epsilon) \right]_{\alpha=0} = \frac{1}{2} \left[ \frac{\partial C_{ijkl}(\eta_p + \alpha w, \eta_q, \eta_r)}{\partial (\eta_p + \alpha w)} \frac{\partial (\eta_p + \alpha w)}{\partial \alpha} \right. \\
\left. \cdot \left( \epsilon_{ij} - \epsilon_{ij}^0(c, \eta_p + \alpha w, \eta_q, \eta_r) \right) \left( \epsilon_{kl} - \epsilon_{kl}^0(c, \eta_p + \alpha w, \eta_q, \eta_r) \right) \right. \\
\left. + C_{ijkl}(\eta_p + \alpha w, \eta_q, \eta_r) \left( - \frac{\partial \epsilon_{ij}^0(c, \eta_p + \alpha w, \eta_q, \eta_r)}{\partial (\eta_p + \alpha w)} \frac{\partial (\eta_p + \alpha w)}{\partial \alpha} \right) \right. \\
\left. \cdot \left( \epsilon_{kl} - \epsilon_{kl}^0(c, \eta_p + \alpha w, \eta_q, \eta_r) \right) \right. \\
\left. + C_{ijkl}(\eta_p + \alpha w, \eta_q, \eta_r) \left( \epsilon_{ij} - \epsilon_{ij}^0(c, \eta_p + \alpha w, \eta_q, \eta_r) \right) \right. \\
\left. \cdot \left( - \frac{\partial \epsilon_{kl}^0(c, \eta_p + \alpha w, \eta_q, \eta_r)}{\partial (\eta_p + \alpha w)} \frac{\partial (\eta_p + \alpha w)}{\partial \alpha} \right) \right]_{\alpha=0} \right. \\
= \frac{1}{2} \left[ \frac{\partial C_{ijkl}(\eta_p, \eta_q, \eta_r)}{\partial \eta_p} w \left( \epsilon_{ij} - \epsilon_{ij}^0(c, \eta_p, \eta_q, \eta_r) \right) \left( \epsilon_{kl} - \epsilon_{kl}^0(c, \eta_p, \eta_q, \eta_r) \right) \right. \\
\left. + C_{ijkl}(\eta_p, \eta_q, \eta_r) \left( - \frac{\partial \epsilon_{ij}^0(c, \eta_p, \eta_q, \eta_r)}{\partial \eta_p} w \right) \left( \epsilon_{kl} - \epsilon_{kl}^0(c, \eta_p, \eta_q, \eta_r) \right) \right.$$
(31)

Putting the terms back together yields:

$$\delta_{\eta_{p}}\Pi = \int_{\Omega} f_{\alpha}(c) \left[ -\frac{\partial H(\eta_{p})}{\partial \eta_{p}} w \right] + f_{\beta}(c) \left[ \frac{\partial H(\eta_{p})}{\partial \eta_{p}} w \right] + W \left[ \frac{\partial f_{Landau}(\eta_{p}, \eta_{q}, \eta_{r})}{\partial \eta_{p}} w \right] 
+ \kappa_{ij} w_{,i} \eta_{p,j} 
+ \frac{1}{2} \left[ \frac{\partial C_{ijkl}(\eta_{p}, \eta_{q}, \eta_{r})}{\partial (\eta_{p})} w \left( \epsilon_{ij} - \epsilon_{ij}^{0}(c, \eta_{p}, \eta_{q}, \eta_{r}) \right) \left( \epsilon_{kl} - \epsilon_{kl}^{0}(c, \eta_{p}, \eta_{q}, \eta_{r}) \right) \right] 
+ C_{ijkl}(\eta_{p}, \eta_{q}, \eta_{r}) \left( -\frac{\partial \epsilon_{ij}^{0}(c, \eta_{p}, \eta_{q}, \eta_{r})}{\partial \eta_{p}} w \right) \left( \epsilon_{kl} - \epsilon_{kl}^{0}(c, \eta_{p}, \eta_{q}, \eta_{r}) \right) dV$$
(32)

Variational derivative of  $\Pi$  with respect to c:

$$\delta_c \Pi = \frac{d}{d\alpha} \left[ \int_{\Omega} f_{chem}(c + \alpha w, \eta_p, \eta_q, \eta_r) + f_{grad}(\eta_p, \eta_q, \eta_r) + f_{el}(c + \alpha w, \eta_p, \eta_q, \eta_r, \epsilon) dV \right]_{\alpha = 0}$$
(33)

Breaking up each of these terms yields:

$$\frac{d}{d\alpha} \left[ f_{chem}(c + \alpha w, \eta_p, \eta_q, \eta_r) \right]_{\alpha=0} = \left[ \frac{\partial f_{\alpha}(c + \alpha w)}{\partial (c + \alpha w)} \frac{\partial (c + \alpha w)}{\partial \alpha} \left( 1 - \sum_{p=1}^{3} H(\eta_p) \right) + W \frac{\partial f_{Landau}(\eta_p, \eta_q, \eta_r)}{\partial (c + \alpha w)} \frac{\partial (c + \alpha w)}{\partial \alpha} \right]_{\alpha=0} \\
= \frac{\partial f_{\alpha}(c)}{\partial c} w \left( 1 - \sum_{p=1}^{3} H(\eta_p) \right) + \frac{\partial f_{\beta}(c)}{\partial c} w \left( \sum_{p=1}^{3} H(\eta_p) \right) \tag{34}$$

$$\frac{d}{d\alpha} \left[ f_{grad}(\eta_p, \eta_q, \eta_r) \right]_{\alpha=0} = 0 \tag{35}$$

$$\frac{d}{d\alpha} \left[ f_{el}(c + \alpha w, \eta_p, \eta_q, \eta_r, \epsilon) \right]_{\alpha=0} = \frac{1}{2} C_{ijkl}(\eta_p, \eta_q, \eta_r) \left[ -\frac{\partial \epsilon_{ij}^0(c + \alpha w, \eta_p, \eta_q, \eta_r)}{\partial (c + \alpha w)} \frac{\partial (c + \alpha w)}{\partial \alpha} \left( \epsilon_{kl} - \epsilon_{kl}^0(c + \alpha w, \eta_p, \eta_q, \eta_r) \right) - \left( \epsilon_{ij} - \epsilon_{ij}^0(c + \alpha w, \eta_p, \eta_q, \eta_r) \right) \frac{\partial \epsilon_{ij}^0(c + \alpha w, \eta_p, \eta_q, \eta_r)}{\partial (c + \alpha w)} \frac{\partial (c + \alpha w)}{\partial \alpha} \right]_{\alpha=0}$$

$$= -C_{ijkl}(\eta_p, \eta_q, \eta_r) \frac{\partial \epsilon_{ij}^0(c, \eta_p, \eta_q, \eta_r)}{\partial c} w \left( \epsilon_{kl} - \epsilon_{kl}^0(c + \alpha w, \eta_p, \eta_q, \eta_r) \right)$$
(36)

Putting the terms back together yields:

$$\delta_{c}\Pi = \int_{\Omega} \frac{\partial f_{\alpha}(c)}{\partial c} w \left( 1 - \sum_{p=1}^{3} H(\eta_{p}) \right) + \frac{\partial f_{\beta}(c)}{\partial c} w \left( \sum_{p=1}^{3} H(\eta_{p}) \right)$$

$$- C_{ijkl}(\eta_{p}, \eta_{q}, \eta_{r}) \frac{\partial \epsilon_{ij}^{0}(c, \eta_{p}, \eta_{q}, \eta_{r})}{\partial c} w \left( \epsilon_{kl} - \epsilon_{kl}^{0}(c + \alpha w, \eta_{p}, \eta_{q}, \eta_{r}) \right) dV$$

$$(37)$$