

PRISMS-PF Application Formulation: dendriticSolidification

This example application implements a simple model of dendritic solidification based on the CHiMaD Benchmark Problem 3, itself based on the model given in the following article:
“Multiscale Finite-Difference-Diffusion-Monte-Carlo Method for Simulating Dendritic Solidification” by M. Plapp and A. Karma, *Journal of Computational Physics*, 165, 592-619 (2000)

1 Governing Equations

Consider a free energy density given by:

$$\Pi = \int_{\Omega} \left[\frac{1}{2} W^2(\hat{n}) |\nabla \phi|^2 + f(\phi, u) \right] dV \quad (1)$$

where ϕ is an order parameter for the solid phase and u is the dimensionless temperature:

$$u = \frac{T - T_m}{L/c_p} \quad (2)$$

for temperature T , melting temperature T_m , latent heat L , and specific heat c_p . The free energy density, $f(\phi, u)$ is given by a double-well potential:

$$f(\phi, u) = -\frac{1}{2}\phi^2 + \frac{1}{4}\phi^4 + \lambda u \phi \left(1 - \frac{2}{3}\phi^2 + \frac{1}{5}\phi^4 \right) \quad (3)$$

where λ is a dimensionless coupling constant. The gradient energy coefficient, W , is given by

$$W = W_0 [1 + \epsilon_m \cos[m(\theta - \theta_0)]] \quad (4)$$

where, W_0 , ϵ_m , and θ_0 are constants and θ is the in-plane azimuthal angle, where $\tan(\theta) = \frac{n_y}{n_x}$, where n_y and n_x are components of the normal vector $\hat{n} = \frac{\nabla \phi}{|\nabla \phi|}$.

The evolution equations are:

$$\frac{\partial u}{\partial t} = -\frac{\delta \Pi}{\delta \phi} = D \nabla^2 u + \frac{1}{2} \tau(\hat{n}) \frac{\partial \phi}{\partial t} \quad (5)$$

$$\tau(\hat{n}) \frac{\partial \phi}{\partial t} = [\phi - \lambda(1 - \phi^2)](1 - \phi^2) + \nabla \cdot [W(\hat{n})^2 \nabla \phi] + \frac{\partial}{\partial x} \left[|\nabla \phi|^2 W(\hat{n}) \frac{\partial W(\hat{n})}{\partial \left(\frac{\partial \phi}{\partial x} \right)} \right] \quad (6)$$

where

$$\tau(\hat{n}) = \tau_0 [1 + \epsilon_m \cos[m(\theta - \theta_0)]] \quad (7)$$

$$D = \frac{0.6267 \lambda W_0^2}{\tau_0} \quad (8)$$

2 Model Constants

$$\begin{aligned}
W_0 \\
\tau_0 \\
D \\
T_0 \\
\Delta = \frac{T_m - T_0}{L/c_p}
\end{aligned}$$

3 Time Discretization

Considering forward Euler explicit time stepping, we have the time discretized kinetics equation:

$$\eta^{n+1} = \eta^n - \Delta t M_\eta ((f_{\beta,c}^n - f_{\alpha,c}^n) H_{,\eta}^n - \kappa \Delta \eta^n) \quad (9)$$

$$c^{n+1} = c^n + \Delta t M_\eta \nabla \cdot (\nabla (f_{\alpha,c}^n (1 - H^n) + f_{\beta,c}^n H^n)) \quad (10)$$

4 Weak Formulation

In the weak formulation, considering an arbitrary variation w , the above equations can be expressed as residual equations:

$$\int_{\Omega} w \eta^{n+1} dV = \int_{\Omega} w \eta^n - w \Delta t M_\eta ((f_{\beta,c}^n - f_{\alpha,c}^n) H_{,\eta}^n - \kappa \Delta \eta^n) dV \quad (11)$$

$$= \int_{\Omega} w \left(\underbrace{\eta^n - \Delta t M_\eta ((f_{\beta,c}^n - f_{\alpha,c}^n) H_{,\eta}^n)}_{r_\eta} \right) + \nabla w \cdot \underbrace{(-\Delta t M_\eta \kappa \nabla \eta^n)}_{r_{\eta x}} dV \quad (12)$$

and

$$\int_{\Omega} w c^{n+1} dV = \int_{\Omega} w c^n + w \Delta t M_c \nabla \cdot (\nabla (f_{\alpha,c}^n (1 - H^n) + f_{\beta,c}^n H^n)) dV \quad (13)$$

$$= \int_{\Omega} w \underbrace{c^n}_{r_c} + \nabla w \cdot \underbrace{(-\Delta t M_c) [(f_{\alpha,cc}^n (1 - H^n) + f_{\beta,cc}^n H^n) \nabla c + ((f_{\beta,c}^n - f_{\alpha,c}^n) H_{,\eta}^n \nabla \eta)]}_{r_{cx}} dV \quad (14)$$

The above values of r_η , $r_{\eta x}$, r_c and r_{cx} are used to define the residuals in the following parameters file:
applications/coupledCahnHilliardAllenCahn/parameters.h