

Excited States and Nonadiabatic Dynamics CyberTraining School/Workshop 2022

Alexey Akimov

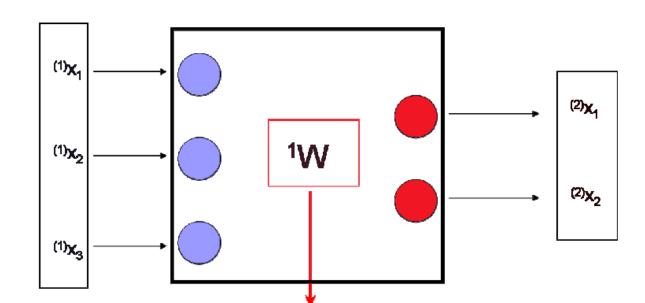
University at Buffalo, SUNY



Fundamentals of ML with Libra

Basics of the ANN





- The expected output should be scaled
- The inputs should be close to 0.0, where the transfer function slope is large

$$\begin{pmatrix} {}^{(2)}NET_1 \\ {}^{(2)}NET_2 \end{pmatrix} = \begin{pmatrix} {}^{(1)}w_{11} & {}^{(1)}w_{12} & {}^{(1)}w_{13} \\ {}^{(1)}w_{21} & {}^{(1)}w_{22} & {}^{(1)}w_{23} \end{pmatrix} \begin{pmatrix} {}^{(1)}x_1 \\ {}^{(1)}x_2 \\ {}^{(1)}x_3 \end{pmatrix} + \begin{pmatrix} {}^{(1)}b_1 \\ {}^{(1)}b_2 \\ {}^{(1)}b_3 \end{pmatrix}$$

$$f(^{(2)}NET_1)$$

$$f(^{(2)}NET_2)$$

$$f(x) = \tanh(\alpha x)$$

$$f(x) = \frac{1}{1 + \exp(-\alpha x)}$$

Function back-propagation



$$y_i^{(L)} = f\left(\sum_j W_{ij}^{(L)} y_j^{(L-1)} + B_i^{(L)}\right) \iff y^{(L)} = f(W^{(L)} y^{(L-1)} + B^{(L)})$$
$$a^{(L)} = W^{(L)} y^{(L-1)} + B^{(L)}$$

$$\frac{\partial y_{i}^{(L)}}{\partial W_{ab}^{(L)}} = f'\left(a_{i}^{(L)}\right) \sum_{j} \left(\frac{\partial W_{ij}^{(L)}}{\partial W_{ab}^{(L)}} y_{j}^{(L-1)}\right) = f'\left(a_{i}^{(L)}\right) \sum_{j} \left(\delta_{ia}\delta_{jb}y_{j}^{(L-1)}\right) = \delta_{ia}f'\left(a_{i}^{(L)}\right) y_{b}^{(L-1)}$$

$$\frac{\partial y_{i}^{(L)}}{\partial B_{a}^{(L)}} = f'\left(a_{i}^{(L)}\right) \frac{\partial B_{i}^{(L)}}{\partial B_{a}^{(L)}} = \delta_{ia}f'\left(a_{i}^{(L)}\right)$$

$$E = \frac{1}{2} \sum_{i} \left(y_i^{(NL)} - t_i \right)^2 \leftrightarrow E = \frac{1}{2} \left(\mathbf{y}^{(NL)} - \mathbf{t} \right)^T \left(\mathbf{y}^{(NL)} - \mathbf{t} \right)$$

Gradients w.r.t the weights and biases of the layer NL



$$\frac{\partial E}{\partial W_{ab}^{(NL)}} = \sum_{i} \left(y_i^{(NL)} - t_i \right) f' \left(a_i^{(NL)} \right) \sum_{j} y_j^{(NL-1)} \delta_{ai} \delta_{bj} = \left(y_a^{(NL)} - t_a \right) f' \left(a_a^{(NL)} \right) y_b^{(NL-1)}$$

$$\frac{\partial E}{\partial B_a^{(NL)}} = \sum_{i} \left(y_i^{(NL)} - t_i \right) f' \left(a_i^{(NL)} \right) \delta_{ia} = \left(y_a^{(NL)} - t_a \right) f' \left(a_a^{(NL)} \right)$$

$$d_a^{(NL)} = \left(y_a^{(NL)} - t_a\right) f'\left(a_a^{(NL)}\right)$$

Introducing matrix $\pmb{ au}^{(NL)}$ - the diagonal matrix, with $\tau_{ii}^{(NL)}=f'\left(a_i^{(NL)}\right)$ $\pmb{d}^{(NL)}=\pmb{ au}^{(NL)}(\pmb{v}^{(NL)}-\pmb{t})$

$$\frac{\partial E}{\partial W_{ab}^{(NL)}} = d_a^{(NL)} y_b^{(NL-1)} \longleftrightarrow \frac{\partial E}{\partial \boldsymbol{W}^{(NL)}} = (\boldsymbol{d}^{(NL)})^T \boldsymbol{y}^{(NL-1)}$$

$$\frac{\partial E}{\partial B_a^{(NL)}} = d_a^{(NL)} \longleftrightarrow \frac{\partial E}{\partial \boldsymbol{B}^{(NL)}} = \boldsymbol{d}^{(NL)}$$

In general



$$\frac{\partial E}{\partial W_{ab}^{(NL-1)}} = d_a^{(NL-1)} y_b^{(NL-2)} \iff \frac{\partial E}{\partial \boldsymbol{W}^{(NL-1)}} = (\boldsymbol{d}^{(NL-1)})^T \boldsymbol{y}^{(NL-2)}$$

$$\frac{\partial E}{\partial B_a^{(NL-1)}} = d_a^{(NL-1)} \iff \frac{\partial E}{\partial \boldsymbol{B}^{(NL-1)}} = \boldsymbol{d}^{(NL-1)}$$

Derivatives of the ANN w.r.t. Inputs



$$y_i^{(L)} = f\left(\sum_j W_{ij}^{(L)} y_j^{(L-1)} + B_i^{(L)}\right) \iff y^{(L)} = f(W^{(L)} y^{(L-1)} + B^{(L)})$$
$$a^{(L)} = W^{(L)} y^{(L-1)} + B^{(L)}$$

$$\begin{split} &\frac{\partial y_{i}^{(L)}}{\partial y_{a}^{(0)}} = f'\left(a_{i}^{(L)}\right) \sum_{j} \left(W_{ij}^{(L)} \frac{\partial y_{j}^{(L-1)}}{\partial y_{a}^{(0)}}\right) \\ &= f'\left(a_{i}^{(L)}\right) \sum_{j} \left(W_{ij}^{(L)} f'\left(a_{j}^{(L-1)}\right) \sum_{k} W_{jk}^{(L-1)} \frac{\partial y_{k}^{(L-2)}}{\partial y_{a}^{(0)}}\right) \\ &= \sum_{i,k} f'\left(a_{i}^{(L)}\right) W_{ij}^{(L)} f'\left(a_{j}^{(L-1)}\right) W_{jk}^{(L-1)} \frac{\partial y_{k}^{(L-2)}}{\partial y_{a}^{(0)}} = \cdots \end{split}$$

$$G^{(L)} = \tau^{(L)} W^{(L)} G^{(L-1)}$$

More specifically,

$$\mathbf{G}^{(NL)} = \mathbf{\tau}^{(NL)} \mathbf{W}^{(NL)} \mathbf{\tau}^{(NL-1)} \mathbf{W}^{(NL-1)} \mathbf{\tau}^{(NL-2)} \mathbf{W}^{(NL-2)} \dots \mathbf{\tau}^{(1)} \mathbf{W}^{(1)}$$

Back-prop with momentum



http://page.mi.fu-berlin.de/rojas/neural/chapter/K8.pdf

https://arxiv.org/pdf/1711.05101.pdf

http://citeseerx.ist.psu.edu/viewdoc/summary?doi=10.1.1.17.1332

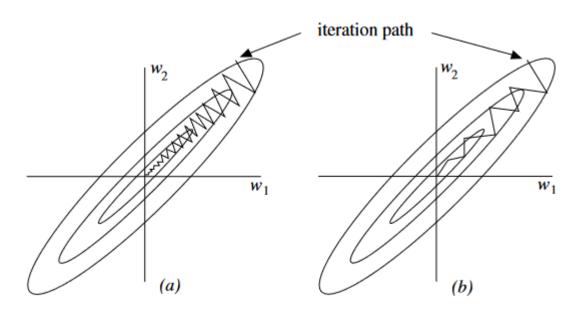


Fig. 8.1. Backpropagation without (a) or with (b) momentum term

$$\Delta w_k(i) = -\gamma \frac{\partial E}{\partial w_k} + \alpha \Delta w_k(i-1),$$

Rprop



$$\gamma_i^{(k+1)} = \begin{cases} \min(\gamma_i^{(k)} u, \gamma_{max}) & \text{if } \nabla_i E^{(k)} \cdot \nabla_i E^{(k-1)} > 0 \\ \max(\gamma_i^{(k)} d, \gamma_{min}) & \text{if } \nabla_i E^{(k)} \cdot \nabla_i E^{(k-1)} < 0 \\ \gamma_i^{(k)} & \text{otherwise,} \end{cases}$$

where the constants u and d satisfy u > 1 and d < 1, as usual. When $\nabla_i E^{(k)} \cdot \nabla_i E^{(k-1)} \ge 0$ the weight updates are given by

$$\Delta^{(k)}w_i = -\gamma_i^{(k)}\operatorname{sgn}(\nabla_i E^{(k)}),$$

following rule:

$$\Delta_{ij}^{(t)} := \begin{cases} \eta^{+} \cdot \Delta_{ij}^{(t-1)} &, & \text{if } \frac{\partial E}{\partial w_{ij}}^{(t-1)} \cdot \frac{\partial E}{\partial w_{ij}}^{(t)} > 0 \\ \eta^{-} \cdot \Delta_{ij}^{(t-1)} &, & \text{if } \frac{\partial E}{\partial w_{ij}}^{(t-1)} \cdot \frac{\partial E}{\partial w_{ij}}^{(t)} < 0 \\ \Delta_{ij}^{(t-1)} &, & \text{else} \end{cases},$$

$$(1)$$

where $0 < \eta^- < 1 < \eta^+$. If the partial deriva-



Fundamentals of DVR with Libra

Wavefunction is discretized on a grid



Wavefunction is discretized on a grid

$$\langle r|\Psi\rangle = \Psi(r,t) = \sum_{\substack{i \in grid, \\ a}} \Psi_a(r_i,t) \delta(r-r_i) |a\rangle \qquad \text{PSI_dia} = \left\{ \begin{pmatrix} \Psi_0(r_0) \\ \dots \\ \Psi_{N-1}(r_0) \end{pmatrix}, \begin{pmatrix} \Psi_0(r_1) \\ \dots \\ \Psi_{N-1}(r_1) \end{pmatrix}, \dots, \begin{pmatrix} \Psi_0(r_{Npts-1}) \\ \dots \\ \Psi_{N-1}(r_{Npts-1}) \end{pmatrix} \right\}$$

In Libra, any N-dimensional grid is "linearized" this way via a mapping function

This could be thought of as using the basis of grid-point functions $|i,a\rangle$: $\langle r|\Psi\rangle = \delta(r-r_i)|a\rangle$

$$\langle \Psi | \Psi \rangle = \sum_{\mathbf{a}, b, i, j} \int dr \Psi_a^*(r_i) \Psi_b(r_j) \delta(r - r_i) \delta(r - r_j) \langle a | b \rangle = \Delta r \sum_{\mathbf{a}, i} \Psi_a^*(r_i) \Psi_a(r_i)$$

Matrix elements of operators

$$\langle \Psi | \hat{A} | \Psi \rangle = \sum_{a,b,i,j} \int dr \Psi_a^*(r_i) \Psi_b(r_j) \delta(r - r_i) \delta(r - r_j) A_{ab}(r) = \Delta r \sum_{a,b,i} \Psi_a^*(r_i) A_{ab}(r_i) \Psi_a(r_i)$$

Momentum representation



Real-space (coordinate) wavefunction

 $\psi_a(\mathbf{r},t) = \int \tilde{\psi}_a(\mathbf{k},t) e^{2\pi i \mathbf{r} \mathbf{k}} d\mathbf{k}$

Reciprocal-space (momentum) wavefunction

$$\tilde{\psi}_i(\mathbf{k},t) = \int \psi_i(\mathbf{r},t)e^{-2\pi i \mathbf{r} \mathbf{k}}d\mathbf{r}$$

$$\begin{split} & \left| \psi_{i}(x) \right| \left(-i \frac{\partial}{\partial x} \right)^{n} \left| \psi_{j}(x) \right\rangle = \sum_{i,j} \int dx \left(\int \tilde{\psi}_{i}(k) e^{2\pi i x k} dk \right) \left(-i \frac{\partial}{\partial x} \right)^{n} \left(\int \tilde{\psi}_{j}(k') e^{2\pi i x k'} dk' \right) \\ &= (-i)^{n} \sum_{i,j} \int dx \left(\int \tilde{\psi}_{i}(k) e^{2\pi i x k} dk \right)^{*} \left((2\pi i)^{n} \int k'^{n} \tilde{\psi}_{j}(k') e^{2\pi i x k'} dk' \right) \\ &= (2\pi)^{n} \sum_{i,j} \int dx dk dk' \tilde{\psi}_{i}^{*}(k) e^{-2\pi i x k} (k')^{n} \tilde{\psi}_{j}(k') e^{2\pi i x k'} \\ &= (2\pi)^{n} \sum_{i,j} \int dk dk' \tilde{\psi}_{i}^{*}(k) \delta(k - k') (k')^{n} \tilde{\psi}_{j}(k') = (2\pi)^{n} \sum_{i,j} \int dk \tilde{\psi}_{i}^{*}(k) k^{n} \tilde{\psi}_{j}(k) \\ &\rightarrow (2\pi)^{n} \Delta k \sum_{i,j,m} \tilde{\psi}_{i}^{*}(k_{m}) k_{m}^{n} \tilde{\psi}_{j}(k_{m}) \end{split}$$

Solution of the TD-SE



$$i\hbar \frac{\partial \Psi(r,t)}{\partial t} = \widehat{H}\Psi(r,t) = (\widehat{T} + \widehat{V})\Psi(r,t)$$

Finite difference evaluation of the derivatives

$$\partial_t \Psi_i(\boldsymbol{r}_n, t_m) = \frac{1}{2\Delta t} \left[\Psi_i(\boldsymbol{r}_n, t_{m+1}) - \Psi_i(\boldsymbol{r}_n, t_{m-1}) \right]$$
$$\nabla_{\boldsymbol{r}_{\alpha}} \Psi_i(\boldsymbol{x}_n, t_m) = \frac{1}{2\Delta r_{\alpha}} \left[\Psi_i(\boldsymbol{r}_{\alpha, n+1}, t_m) - \Psi_i(\boldsymbol{r}_{\alpha, n-1}, t_m) \right]$$

$$\nabla_{\boldsymbol{r}_{\alpha}}^{2}\Psi_{i}(x_{n},t_{m}) = \frac{1}{4\Delta\boldsymbol{r}_{\alpha}^{2}}\left[\Psi_{i}(\boldsymbol{r}_{\alpha,n+2},t_{m}) - \Psi_{i}(\boldsymbol{r}_{n},t_{m}) - \left[\Psi_{i}(\boldsymbol{r}_{n},t_{m}) - \Psi_{i}(\boldsymbol{r}_{\alpha,n-2},t_{m})\right]\right] = \frac{1}{4\Delta\boldsymbol{r}_{\alpha}^{2}}\left[\Psi_{i}(\boldsymbol{r}_{\alpha,n+2},t_{m}) - 2\Psi_{i}(\boldsymbol{r}_{n},t_{m}) + \Psi_{i}(\boldsymbol{r}_{\alpha,n-2},t_{m})\right]$$

Solution of the TD-SE



$$i\hbar \frac{\partial |\Psi(t)\rangle}{\partial t} = \hat{H}|\Psi(t)\rangle = (\hat{T} + \hat{V})|\Psi(t)\rangle$$

Split-operator method (Kosloff & Kosloff)

$$|\Psi(t+\Delta t)\rangle = \exp\left(-\frac{i\Delta t}{\hbar}\widehat{H}\right)|\Psi(t)\rangle = \exp\left(-\frac{i\Delta t}{\hbar}\big(\widehat{T}+\widehat{V}\big)\right)|\Psi(t)\rangle \approx \exp\left(-\frac{i\Delta t}{2\hbar}\widehat{V}\right)\exp\left(-\frac{i\Delta t}{\hbar}\widehat{T}\right)\exp\left(-\frac{i\Delta t}{2\hbar}\widehat{V}\right)|\Psi(t)\rangle$$

$$\begin{split} &\Psi_{a}(r_{i},t') = \langle r_{i},a|\exp\left(-\frac{i\Delta t}{2\hbar}\hat{V}\right)|\Psi(t)\rangle = \langle r_{i}|\exp\left(-\frac{i\Delta t}{2\hbar}\hat{V}\right)\sum_{j,b}|r_{j},b\rangle\langle r_{j},b|\Psi(t)\rangle = \sum_{j,b}\left\langle r_{i},a\right|\exp\left(-\frac{i\Delta t}{2\hbar}\hat{V}\right)\left|r_{j},b\rangle\langle r_{j},b|\Psi(t)\rangle \\ &= \sum_{j,b}\left\langle a\right|\exp\left(-\frac{i\Delta t}{2\hbar}V\left(r_{i}\right)\right)\left|b\right\rangle\delta_{ij}\Psi_{b}(r_{j},t) = \sum_{b}\left[\exp\left(-\frac{i\Delta t}{2\hbar}V\left(r_{i}\right)\right)\right]_{ab}\Psi_{b}(r_{i},t) \end{split}$$

$$\begin{split} \widetilde{\Psi}_{\mathbf{a}}(k_{i},t'') &= \langle k_{i},a| \exp\left(-\frac{i\Delta t}{\hbar}\widehat{T}\right) |\Psi(t)\rangle = \langle k_{i},a| \exp\left(-\frac{i\Delta t}{\hbar}\widehat{T}\right) \sum_{j,b} |k_{j},b\rangle \langle k_{j},b| \Psi(t)\rangle \\ &= \sum_{j,b} \left\langle k_{i},a \right| \exp\left(-\frac{i\Delta t}{2\hbar}\widehat{T}\right) \left|k_{j},b\right\rangle \langle k_{j},b| \Psi(t)\rangle = \sum_{j,b} \exp\left(-\frac{i\Delta t}{2\hbar}\frac{k_{i}^{2}}{2m}\right) \delta_{ij}\delta_{ab}\widetilde{\Psi}_{b}(t) = \exp\left(-\frac{i\Delta t}{2\hbar}\frac{k_{i}^{2}}{2m}\right) \widetilde{\Psi}_{a}(t) \end{split}$$



Fundamentals of HEOM with Libra

See here

https://compchem-cybertraining.github.io/Cyber_Training_Workshop_2021/files/Jain-HEOM.pdf