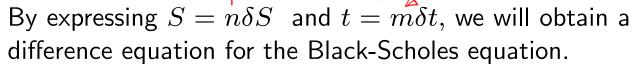
Stability Analysis



$$V(S,t) = V(n\delta S, m\delta t) = V_n^m.$$

 $\delta S = \frac{S^*}{N}$ where $S^* \gg E$ is a suitably large value of S ; $\delta t = \frac{T}{M}.$ Taking N and M steps for S and t respectively, so

$$S = n\delta S$$
 $0 \le n \le N$
 $t = m\delta t$ $0 \le m \le M$.

$$\begin{cases} \frac{V_n^m - V_n^{m+1}}{\delta t} + \frac{1}{2}n^2\sigma^2\left(V_{n-1}^m - 2V_n^m + V_{n+1}^m\right) + \\ \frac{1}{2}(r-D)n\left(V_{n+1}^m - V_{n-1}^m\right) - rV_n^m = 0 \end{cases}$$

and rearrange to obtain a *forward marching* scheme in time

$$V_{n}^{m+1} = V_{n}^{m} + \delta t \left(\frac{1}{2} n^{2} \sigma^{2} \left(V_{n-1}^{m} - 2V_{n}^{m} + V_{n+1}^{m} \right) \right)$$

$$+ \delta t \left(\frac{1}{2} (r - D) n \left(V_{n+1}^{m} - V_{n-1}^{m} \right) - r V_{n}^{m} \right)$$

$$\equiv F \left(V_{n-1}^{m}, V_{n}^{m}, V_{n+1}^{m} \right)$$

$$\vdots$$

Now for the RHS collect coefficients of each variable term V, to get

$$V_n^{m+1} = \alpha_n V_{n-1}^m + \beta_n V_n^m + \gamma_n V_{n+1}^m \tag{1}$$

where

$$\alpha_{n} = \frac{1}{2} \left(n^{2} \sigma^{2} - n (r - D) \right) \delta t,$$

$$\beta_{n} = 1 - \left(r + n^{2} \sigma^{2} \right) \delta t,$$

$$\gamma_{n} = \frac{1}{2} \left(n^{2} \sigma^{2} + n (r - D) \right) \delta t$$

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Fourier Stability (Von Neumann's) Method

A method is called step-wise unstable if for a fixed grid (i.e. δt , δS constant) there exists an initial perturbation which "blows up" as $t\to\infty$, i.e. as we march in time. Here in a forward marching scheme. The question we wish to answer is "do small errors propagate along the grid and grow exponentially?".

Assume an initial disturbance which is proportional to $\exp(in\omega)$. We therefore study the propagation of perturbations created at any given point in time.

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If \hat{V}_n^m is an approximation to the exact solution V_n^m then for exact exact

$$\widehat{V}_n^m = V_n^m + E_n^m$$

where $E_n^{\ m}$ is the associated error. Then $E_n^{\ m}$ also satisfies the difference equation (4) to give

$$E_n^{m+1} = \alpha_n E_{n-1}^m + \beta_n E_n^m + \gamma_n E_{n+1}^m.$$

Put
$$E_n^m = \overline{a}^m \exp(in\omega) \tag{3}$$

which is oscillatory of amplitude \overline{a} and frequency ω . Substituting (3) into (1) gives

$$\overline{a}^{n+1}e^{(in\omega)} = \alpha_n \overline{a}^{n}e^{i(n-1)\omega} + \beta_n \overline{a}^{n}e^{in\omega} + \gamma_n \overline{a}^{n}e^{i(n+1)\omega}$$
 which becomes (
$$\underbrace{a}_{n} + \underbrace{a}_{n} = \underbrace{a}_{n} - \underbrace{a}_{n} = \underbrace{a}_{n$$

$$\overline{a} = \alpha_n e^{-i\omega} + \beta_n + \gamma_n e^{i\omega}.$$

Now stability criteria arises from the balancing of the time dependency and diffusion terms, so that

$$\left(\frac{\partial V}{\partial t} + \frac{1}{2}\sigma^2 S^2 \frac{\partial^2 V}{\partial S^2}\right) = 0$$

From (2) we take the following contributions

$$\alpha_n = \frac{1}{2}n^2\sigma^2\delta t, \ \beta_n = 1 - n^2\sigma^2\delta t, \ \gamma_n = \frac{1}{2}n^2\sigma^2\delta t$$

$$SLX = \frac{ix - ix}{2i}$$

$$\overline{a} = \frac{2}{2}n^2\sigma^2\delta t \left(\underline{e^{i\omega} + e^{-i\omega}}\right) + 1 - n^2\sigma^2\delta t$$

$$\alpha = n^2 \sigma^2 \delta t \left(\cos \omega - 1\right) + 1. \quad \cos 2 \zeta = \cos^2 3 \zeta - \sin^2 3 \zeta$$

we have

$$\overline{a} = 1 - 2n^2 \sigma^2 \sin^2 \frac{\omega}{2} \delta t$$

$$= 2 \cos^2 \delta c - 1$$

For stability \overline{a} must be bounded, i.e. $(|\overline{a}| < 1)$

$$\left|1 - 2n^2\sigma^2\sin^2\frac{\omega}{2}\delta t\right| < 1$$

which upon simplifying we find is

so
$$\delta t \sim O\left(N^{-2}\right)$$
.

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. $\delta t < \frac{1}{\sigma^2 N^2}$ stable (4)

$$\omega = 1 - 2sh \frac{\omega}{2}$$

$$\cos w - 1 = -2 \sin^2 w$$