ν -SQuIDS: A toolbox for neutrino oscillation experiments.

Carlos Alberto Argüelles Delgado^{a,b}, Jordi Salvado Serrat^{a,b}

^aDepartment of Physics, University of Wisconsin, Madison, WI 53706, USA

Abstract

The Neutrino Simple Quantum Integro-Differential equation Solver (ν -SQuIDS) is a C++ code based on SQuIDS that propagates an ensemble of neutrinos through a given media, e.g. Sun, Earth, Vacuum, etc, while considering, in a consistent way, the effect of neutrino oscillations, with coherent matter interactions, and non coherent interactions. The code has been design to be accurate and flexible, while at the same time maintain excellent performance. Furthermore, the user can easily change the neutrino oscillation parameters, propagation medium, and easily incorporate new oscillation physics.

Keywords: Neutrino oscillation, phenomenology, collective neutrino behavior, numerical techniques

Contents

1	Introduction Theory				
2					
3	Des	$\operatorname{criptio}$	on of the code		
	3.1	Body	& Track		
		3.1.1	Vacuum		
		3.1.2	ConstantDensity		
		3.1.3	VariableDensity		
		3.1.4	Earth		
		3.1.5	EarthAtm		
		3.1.6	Sun		
	3.2	Neutri	inoCrossSections		
		3.2.1	Constructors and Initializing Functions		
		3.2.2			

^b Wisconsin IceCube Particle Astrophysics Center, Madison, WI 53703, USA

The code can be found in http://icecube.wisc.edu/~carguelles/nusquids

Email addresses: arguelles@wisc.edu (Carlos Alberto Argüelles Delgado), salvadoserra@wisc.edu (Jordi Salvado Serrat)

	3.3	TauDecaySpectra	10
		3.3.1 Constructors and Initializing Functions	10
		3.3.2 Functions	11
	3.4	nuSQUIDS	11
		3.4.1 Constructors and Initializing Functions	12
		3.4.2 Functions	12
4	Exa	mples and benchmarks	15
	4.1	Single energy mode	15
		4.1.1 Example 1	15
		4.1.2 Example 2	16
		4.1.3 Example 3	16
	4.2	Multiple energy mode	17
		4.2.1 Example 1	18
		4.2.2 Example 2	19
5	Con	aclusions & Acknowledgements	20

1. Introduction

During the last decades a plethora of evidence that neutrinos change flavor as they propagate macroscopical distances due to the nonalignment of their mass and flavor eigenstates has accumulated from solar, atmospheric, accelerator, and reactor experiments (Mohapatra et al., 2005; de Gouvea et al., 2013). Thanks to these remarkable experimental results and related theoretical calculations the neutrino-mass induced flavor oscillation paradigm (Pontecorvo, 1968; Gribov and Pontecorvo, 1969; Fukugita and Yanagida, 2003; Akhmedov, 1999; Balantekin and Haxton, 2013; Gonzalez-Garcia and Maltoni, 2008) has been firmly established and the three mixing angles, that parametrize the lepton mixing matrix, together with the two square mass differences have been measured to good precision (Gonzalez-Garcia et al., 2014). Furthermore, it is the task of on going and future experiments to determine the neutrino mass ordering and the Dirac CP-violating phase (Hewett et al., 2012).

Even though most of the data can be explained in the standard three neutrino framework some puzzling anomalies still remain (Aguilar et al., 2001; Mention et al., 2011; MiniBooNE-Collaboration et al., 2012). These may be explained by introducing new neutrino states, often refer to as light sterile neutrinos (Kopp et al., 2013; Conrad et al., 2013; Abazajian et al., 2012). Moreover, other new physics scenarios such as the possibility of non standard interactions (Kopp and Welter, 2014; Pospelov, 2011) have also been considered. Also the interplay between cosmology and neutrino oscillation has been widely studied in the literature as well as a playground for new physics scenarios (Archidiacono et al., 2014; Dasgupta and Kopp, 2013; Saviano et al., 2014). In general, neutrino physics constitutes an excellent prove for fundamental physics (Hewett et al., 2012).

It is clear that a tool is needed to correctly solve the neutrino propagation problem, i.e. vacuum oscillation and matter induced interactions (Blennow and Smirnov, 2013). In particular, when only neutrino oscillations are important softwares such as GLobes (Huber et al.,

2007) or nuCRAFT (Wallraff and Wiebusch, 2014) have been made public. Unfortunately, when non coherent interactions are important they are no longer applicable. Furthermore, the ability to allow the user to incorporate new physics models and/or to change the propagation medium is also limited. nuSQuIDS pretends to address all of these problems by providing a highly customizable package while at the same time remain numerically efficient.

The rest of the paper is organized as follows: in section 2 we review neutrino oscillation theory and establish notation; in section 3 we describe the code; in section 4 we exemplify the code and test it in benchmark scenarios. Finally, section 5 presents concluding remarks.

2. Theory

We can represent the neutrino state using the density matrix formalism, e.g. in the weak-interaction flavor eigenstate basis $\{|\nu_{\alpha}\rangle\}$ it can be written as

$$\rho = \sum_{\alpha} \phi_{\alpha} |\nu_{\alpha}\rangle \langle \nu_{\alpha}| \tag{1}$$

where ϕ_{α} specifies the flavor content. Another important basis are the mass eigenstates $\{|\nu_i\rangle\}$, which are related to the former by

$$|\nu_{\alpha}\rangle = \sum_{i} U_{\alpha i} |\nu_{i}\rangle \tag{2}$$

where U is the unitary lepton mixing matrix. It is customary to parametrize the mixing matrix U with mixing angles, $\{\theta_{ij}\}$, and CP phases, $\{\delta_{ij}\}$; for example when considering the standard three flavor paradigm the following parametrization is often used

$$U = \begin{pmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta_{13}} \\ -s_{12}c_{23} - c_{12}s_{23}s_{13}e^{i\delta_{13}} & c_{12}c_{23} - s_{12}s_{23}s_{13}e^{i\delta_{13}} & s_{23}c_{13} \\ s_{12}s_{23} - c_{12}c_{23}s_{13}e^{i\delta_{13}} & -c_{12}s_{23} - s_{12}c_{23}s_{13}e^{i\delta_{13}} & c_{23}c_{13} \end{pmatrix},$$
(3)

where $c_{ij} = \cos \theta_{ij}$, $s_{ij} = \sin \theta_{ij}$. In the three flavor scenario we use the aforementioned parametrization (with values from (Gonzalez-Garcia et al., 2014)) and when more flavors are considered we used the prescription given in Argüelles Delgado and Salvado Serra (2014). Furthermore, the neutrino ensemble evolution is described by the following quantum Liouville equation

$$\frac{\partial \rho(E)}{\partial x} = -i[H(E, x), \rho(E)] . \tag{4}$$

In general we can always split the Hamiltonian, H, into a time dependent and independent parts. In particular, for neutrino oscillations the following splitting is convenient

We set $c = \hbar = 1$.

$$H(E,x) = H_0(E) + H_1(E,x)$$
 (5a)

$$H_0(E) = \frac{1}{2E} \operatorname{diag}(0, \Delta m_{21}^2, \Delta m_{31}^2, ..., \Delta m_{i1}^2, ..., \Delta m_{n1}^2)$$
(5b)

$$H_1(E, x) = \sqrt{2}G_F U^{\dagger} \operatorname{diag}(N_e(x) - N_n(x)/2, -N_n(x)/2, -N_n(x)/2, 0, ..., 0)U$$
 (5c)

where n is the number of neutrino flavors, G_F is the Fermi constant, Δm_{i1}^2 are the neutrino mass splittings, and, finally, $N_e(x)$ and $N_n(x)$ are the electron and nucleon number densities along the neutrino path. On writing these equations we have used the convention that the first three flavor eigenstates corresponds to ν_e , ν_μ , and ν_τ , while the rest are assumed to be sterile neutrinos. Furthermore, H_0 arrises from the neutrino kinetic term, where as H_1 incorporates the matter potential, i.e. coherent forward scattering interactions (Mikheev and Smirnov, 1985, 1986; Wolfenstein, 1978). Given this splitting it is convenient to change to the so called *interaction picture*, generated by H_0 , defined by the following transformation for a given operator O

$$O \to \bar{O}(x) = exp(-iH_0x)Oexp(iH_0x)$$
 (6)

and then the evolution equation is

$$\frac{\partial \bar{\rho}(E)}{\partial x} = -i[\bar{H}_1(E, x), \bar{\rho}(E)] . \tag{7}$$

So far we have only incorporated neutrino oscillation and matter effects through coherent interactions, but we now wish to extend this formalism to incorporate non coherent interactions and collective neutrino behavior. This problem has been extensively discussed in the literature, in particular in (Duan et al., 2010; Strack and Burrows, 2005; Zhang and Burrows, 2013) and (Cirelli et al., 2005; Blennow et al., 2008; Argüelles Delgado and Kopp, 2012), for definiteness we follow the formalism given in (Gonzalez-Garcia et al., 2005). In what follows we will suppress the *bar* symbol and assume that all operators, unless specified, are on the interaction basis.

$$\frac{\partial \rho(E,x)}{\partial x} = -i[H_1(E,x), \bar{\rho}(E)] - \{\Gamma(E,x), \rho\} + F(\Omega, L; E, x)$$
(8a)

$$\frac{\partial l(E,x)}{\partial x} = -\gamma(E,x)l + G(\Omega;E,x)$$
(8b)

where we have introduced $\Omega = \{\rho(E)\}$ and $L = \{l(E)\} = \{e(E), \mu(E), \tau(E)\}$ the neutrino and lepton ensembles respectively. Γ incorporates the effect of attenuation due to non coherent interactions in neutrinos. The F term contains the interactions between the neutrino collective and the leptons, similarly the G term incorporates the effects of neutrinos into leptons. In most scenarios the e and μ leptons lose energy too fast to contribute significantly into the latter neutrino flux, thus we shall only consider the τ leptons since they

have a very short decay time (Halzen and Saltzberg, 1998). Thus, we write the right hand side terms of Eq. (8) explicitly as follows

$$\Gamma(E, x) = \sum_{\alpha} \frac{\Pi_{\alpha}}{2\lambda_{\text{total}}^{\alpha}}$$
(9a)

$$\gamma(E, x) = -\frac{1}{\lambda_{\text{dec}}^{\tau}(E, x)} \tag{9b}$$

$$F(\Omega, L; E, x) = \int_{E}^{\infty} d\tilde{E} \sum_{\alpha} \frac{1}{2} \left\{ \frac{\Pi_{\alpha}}{\lambda_{\text{NC}}^{\alpha}(\tilde{E}, x)}, \rho(\tilde{E}, x) \right\} \frac{\partial N_{\text{NC}}(\tilde{E}, E)}{\partial E}$$

$$+ \int_{E}^{\infty} d\tilde{E} \frac{1}{\lambda^{\tau}(\tilde{E}, x)} \tau(\tilde{E}, x) \frac{\partial N_{\text{dec}}(\tilde{E}, E)}{\partial E} \Pi_{\tau}$$
(9c)

$$+ \operatorname{Br}_{\operatorname{lep}} \int_{E}^{\infty} d\tilde{E} \frac{1}{\lambda^{\tau}(\tilde{E}, x)} \tilde{\tau}(\tilde{E}, x) \frac{\partial \tilde{N}_{\operatorname{dec}}(\tilde{E}, E)}{\partial E} \Pi_{\tau}$$

$$G(\Omega; E, x) = \int_{E}^{\infty} d\tilde{E} \frac{1}{\lambda^{\tau}_{CC}(\tilde{E}, x)} \operatorname{Tr} \left[\Pi_{\tau}, \rho(\tilde{E}, x) \right] \frac{\partial N_{\operatorname{CC}}(\tilde{E}, E)}{\partial E}$$
(9d)

where we have introduced the flavor projectors Π_{α} and the sums run over the active neutrino flavors. Furthermore, $\lambda_{\rm CC}$, $\lambda_{\rm NC}$, and $\lambda_{\rm total} = \lambda_{\rm CC} + \lambda_{\rm NC}$ are the charge, neutral and total neutrino interactions lengths respectively. More over, $\frac{\partial N}{\partial E}$ represent the outgoing neutrino (or τ lepton) spectral distribution for charge, neutral neutrino interactions and τ decay. Finally, λ_{dec}^{τ} is the τ decay length, which is assumed to be much smaller than the relevant neutrino oscillation and interaction scales.

3. Description of the code

 ν -SQuIDS is a C++ code build using the SQuIDS class and framework (Argüelles Delgado and Salvado Serra, 2014). It is designed to propagate neutrinos through media while taking into account flavor oscillations and non coherent interactions. The program can run in two modes: single or multiple energy.

In the *single* energy mode only a fix neutrino energy is considered and no collective neutrino effects are implemented. In this case, only equation (7) is relevant for the neutrino propagation, and the nuSQUIDS class implements SQUIDS::HO as in equation (5b) and SQUIDS::HI as given in equation (5c).

In the *multiple* energy mode a statistical ensemble of neutrinos is considered. The ensemble is described by means of a set of SQUIDS::SU_vector *objects* located at fixed *energy nodes*, which can be either linearly or logarithmically spaced over the energy region under consideration. Besides defining SQUIDS::HO and SQUIDS::HI, as in the *sin-gle* energy mode, the following functions are also defined: SQUIDS::GammaRho by equation (9a), SQUIDS::InteractionsRho as in equation (9c), and SQUIDS::InteractionsScalar as equation (9d). Furthermore, τ regeneration is included under the approximation that the

neutrino interaction scale is much larger than the τ decay length by periodically reinjecting the τ flux decay products to the neutrino statistical ensemble.

Even though the nuSQUIDS class implements all the necessary differential equations we must still specify the neutrino propagation environment, its trajectory, the relevant cross sections, and - if τ regeneration is considered - the properties of τ decay. Thus, for example, when interactions are considered the nuSQUIDS instance will automatically construct appropriate NeutrinoCrossSections and TauDecaySpectra objects to evaluate cross sections and τ physics respectively. On the other hand, the user must explicitly specify the neutrino propagation medium and trajectory through relevant specialization of Body and Body::Track classes.

Finally, nuSQUIDS provides a set functions to evaluate the neutrino ensemble flavor and mass composition as well as the capability to store the calculation in an HDF5 (Folk et al., 1999) file for later study.

3.1. Body & Track

Body and Body::Track are virtual C++ classes which are used to represent the environment where the neutrino propages (Body) and the neutrino path inside it (Body::Track). Along this document we will sometime use the short hand Track for Body::Track, since its clear that the former is meaningless without its corresponding environment, and often we will refer to it as the *neutrino trajectory*. The Body class has two important *virtual* functions which are evaluated along the neutrino trajectory

• Density

```
virtual double density(std::shared_ptr<Track>);
```

Returns density at a give Track position in gr/cm³.

• Electron fraction

```
virtual double ye(std::shared_ptr<Track>);
```

Returns the electron fraction at a give Track position.

Furthermore, the Track object has the following members and functions

• Basic members

```
double x;
double xini;
double xned;
```

x represent the current position along the neutrino path, while xini and xend are the initial and final position in natural units.

• Functions

```
double GetX(void);
double GetInitialX(void);
double GetFinalX(void);
```

These functions return x, xini, and xend respectively.

Since Body and Track are abstract classes they themselves do not perform any task, but rather their specializations specify the real neutrino propagation environment and how it relates to its trajectory. ν -SQuIDS implements the most common used environments and trajectory configurations, but the *user* is free (and encouraged) to create new *classes* in order to extend ν -SQuIDS applicability.

The Body classes specializations implemented in ν -SQuIDS are the following: Vacuum, ConstantDensity, VariableDensity, Earth, EarthAtm, and Sun.

3.1.1. Vacuum

• Vacuum

```
Vacuum(void);
```

Initializes a Vacuum environment.

• Vacuum::Track

```
Vacuum::Track(double xini,double xend);
```

Initialize the corresponding Track setting the initial (xini) and final (xend) neutrino position in eV^{-1} .

3.1.2. ConstantDensity

• ConstantDensity

```
ConstantDensity(double rho, double ye);
```

Initializes a ConstantDensity environment with constant density rho, in gr/cm^3 , and electron fraction ye.

• ConstantDensity::Track

```
ConstantDensity::Track(double xini, double xend);
```

Initialize the corresponding Track setting the initial (xini) and final (xend) neutrino position in eV^{-1} .

3.1.3. VariableDensity

• VariableDensity

```
VariableDensity(std::vector < double > x, std::vector < double > density,
std::vector < double > ye);
```

Initializes a VariableDensity environment given three equal size arrays specifying the density and electron fraction at given positions. An object will be created that interpolates using gsl_spline (Gough, 2009) along the x array to get the density and electron fraction as continuous functions.

• VariableDensity::Track

```
VariableDensity::Track(double xini,double xend);
```

Initialize the corresponding Track setting the initial (xini) and final (xend) neutrino position in eV^{-1} .

3.1.4. Earth

• Earth

```
Earth(void);
```

Initializes an Earth environment as defined by the PREM (Dziewonski and Anderson, 1981).

```
Earth(string filepath);
```

Initializes an Earth environment as defined by a table given in the file specified by filepath. The table should have three columns: radius (where 0 is center and 1 is surface), density (gr/cm³), and y_e (dimensionless). gsl_spline (Gough, 2009) is used to interpolate ρ and y_e as a function of radius to the earth center.

• Earth::Track

```
Earth::Track(double xini,double xend, double L);
```

Initialize the corresponding Track setting the initial (xini) and final (xend) neutrino position along a baseline L.

3.1.5. EarthAtm

• EarthAtm

```
EarthAtm(void);
```

Initializes an EarthAtm environment as defined by the PREM (Dziewonski and Anderson, 1981).

```
EarthAtm(string filepath);
```

Initializes an EarthAtm environment as defined by a table given in the file specified by filepath. The table should have three columns: radius (where 0 is center and 1 is surface), density (gr/cm³), and ye (dimensionless). gsl_spline (Gough, 2009) is used to interpolate ρ and y_e as a function of radius to the earth center.

• EarthAtm::Track

```
EarthAtm::Track(double phi);
```

Initialize the corresponding Track by specifying the zenith angle in radians.

3.1.6. Sun

• Sun

```
Sun(void);
```

Initializes an Sun environment as defined by the *Standard Solar Model* (Bahcall et al., 2005).

• Sun::Track

```
Sun::Track(double xini, double xend);
```

Initialize the corresponding Track by the initial position in the sun xini and xend along the solar radius.

3.2. NeutrinoCrossSections

This object contains neutrino cross section information used when considering neutrino interactions. The given cross sections are a pQCD deep inelastic calculation using the CTEQ6 parton distributions functions on an isoscalar target and are valid for $E_{\nu} > O(10 \text{GeV})$. Furthermore, the ν_e and ν_{μ} cross sections are the same, where as the ν_{τ} includes the τ final state mass suppression; the same holds true for the respective antineutrino cross sections. The cross sections are loaded from tables included in nuSQUIDS/data/generate/. The following quantities are given

- Total neutrino (antineutrino) cross section: $\sigma_{\alpha}^{CC/NC}(E_{\nu})$
- Single differential neutrino (antineutrino) cross section: $\frac{d\sigma_{\alpha}^{CC}}{dE_{\nu}}(E_{\nu}, E_{lep}), \frac{d\sigma_{\alpha}^{NC}}{dE_{\nu}}(E_{\nu}, \tilde{E}_{\nu})$
- 3.2.1. Constructors and Initializing Functions
 - Standard void constructor.

```
NeutrinoCrossSections();
```

• Constructor and initializing function with memory reservation.

```
NeutrinoCrossSections(double Emin, double Emax, int Esize);
void Init(double, double, int);
```

This constructor and initialization functions calculate and store the cross sections on logarithmically spaced energy nodes between Emin and Emax with Esize divisions. For the total cross section gsl_spline (Gough, 2009) is used to interpolate in neutrino energy, where as for the differential cross section simple bilinear interpolation has been implemented.

3.2.2. Functions

The following functions assume that the τ and $\bar{\tau}$ have the same decay distribution.

• Total neutrino cross sections

```
double sigma_CC(int e1,int flv,int nt);
double sigma_NC(int e1,int flv,int nt);
```

sigma_CC (sigma_NC) returns the total deep inelastic neutrino charge (neutral) current cross section at an energy node e1 with neutrino flavor specified by flv (0 = e, 1 = μ , 2 = τ), and nt toggles between neutrinos (0) and antineutrinos (1).

• Neutrino single differential cross sections

```
double dsde_CC(int e1,int e2,int flv, int nt);
double dsde_NC(int e1,int e2,int flv, int nt);
```

dsde_CC (dsde_NC) returns the neutrino single differential charge (neutral) current cross section between energy nodes e1 and e2. Furthermore, the neutrino flavor specified by flv (0 = e, 1 = μ , 2 = τ), and nt toggles between neutrinos (0) and antineutrinos (1).

3.3. TauDecaySpectra

This object contains the information about the τ decay into leptons and hadrons. The formulas implemented in this class were taken from (Dutta et al., 2000) and are implemented in natural units. It is only used when tau regeneration is considered and it returns the following quantities on the energy nodes

$$\frac{dN_{dec}^{lep/had}(E_{\tau}, E_{\nu})}{dE_{\nu}}, \frac{d\tilde{N}_{dec}^{lep/had}(E_{\tau}, E_{\nu})}{dE_{\nu}}$$
(10)

i.e. the neutrino and antineutrino spectral distributions from τ lepton (hadron) decay mode.

3.3.1. Constructors and Initializing Functions

• Standard void constructor.

```
TauDecaySpectra();
```

• Constructor and initializing function with memory reservation.

```
TauDecaySpectra(double Emin, double Emax, int Esize);
void Init(double, double, int);
```

This constructor and initialization functions calculate and store the τ decay spectra on logarithmically spaced energy nodes between Emin and Emax with Esize divisions.

3.3.2. Functions

The following functions assume that the τ and $\bar{\tau}$ have the same decay distribution.

• (Anti)Neutrino spectra with respect to neutrino energy

```
double dNdEnu_All(int e1,int e2);
```

Returns neutrino decay spectra evaluated between energy nodes e1 and e2 when τ decays into leptons or hadrons.

```
double dNdEnu_Lep(int e1,int e2);
```

Returns neutrino decay spectra evaluated between energy nodes e1 and e2 when τ decays into leptons.

• (Anti)Neutrino spectra with respect to τ energy

```
double dNdEle_All(int e1,int e2);
```

Returns neutrino decay spectra evaluated between energy nodes e1 and e2 when τ decays into leptons or hadrons. with respect to the initial τ energy.

```
double dNdEle_Lep(int e1,int e2);
```

Returns neutrino decay spectra evaluated between energy nodes e1 and e2 when τ decays into leptons. with respect to the initial τ energy.

3.4. nuSQUIDS

This object is an specialization of the SQUIDS class (Argüelles Delgado and Salvado Serra, 2014) that implements the differential equations as described in Sec. 2. In particular, it is used to specify the propagation Body and its associated Track. Moreover it uses the NeutrinoCrossSections and TauDecaySpectra in order to evaluate the neutrino cross sections and τ decay spectra; the latter is only used then τ regeneration is enabled. Furthermore, through the SQUIDS::Set function it enables the user to modify the neutrino oscillation parameters as well as the differential equation numerical precision. Finally, it also has the capability to create and read HDF5 files that store the program results and configuration.

3.4.1. Constructors and Initializing Functions

• Standard void constructor.

```
nuSQUIDS();
```

• Single energy mode constructor.

```
nuSQUIDS(int numneu, string NT);
Init(int, string);
```

This constructor and initialization function initializes nuSQUIDS in the *single energy mode*. numneu specifies the number of neutrino flavors which can go from 2 to 6, while NT can be set to "neutrino" or "antineutrino".

• Multiple energy mode constructor.

This constructor and initialization function initializes nuSQUIDS in the *multiple energy* mode. Emin, Emax and Esize define the energy nodes to be used spaced in either logaritmic or linear scales depending on the value of elogscale (true or false). Furthermore, numneu specifies the number of neutrino flavors which can go from 2 to 6, while NT can be set to "neutrino" or "antineutrino" or "both". Finally, iinteraction toggles the neutrino interactions on (true) and off (false).

• Constructing from a ν -SQuIDS-HDF5 file

```
nuSQUIDS(string filepath);
Init(string);
```

This constructor and initialization function initializes nuSQUIDS from a previously generated ν -SQuIDS-HDF5 file. The result nuSQUIDS object will be given in *single* or *multiple* energy mode depending on the HDF5 file configuration. filepath must specify the full path of the HDF5 file.

3.4.2. Functions

Brief description of the nuSQUIDS public functions given in alphabetical order.

• Flavor composition evaluator (single energy mode)

```
double EvalFlavor(int flv);
```

Returns the content a given neutrino flavor specified by flv (0 = e, 1 = μ , 2 = τ , ...). This function can only be use in the *single energy mode*.

• Flavor composition evaluator (multiple energy mode)

```
double EvalFlavorAtNode(int flv,int ie,int rho = 0);
double EvalFlavor(int flv,double Enu,int rho = 0);
```

EvalFlavorAtNode returns the content a given neutrino flavor specified by flv (0 = e, 1 = μ , 2 = τ , ...) at an energy node ie. Furthermore, EvalFlavor returns the approximate content of a given flavor for a specific neutrino energy Enu by interpolating in the interaction basis. In each function, when considering NT = "both", the parameter rho toggles between neutrino (0) and antineutrino (1).

• Mass composition evaluator (single energy mode)

```
double EvalMass(int eig);
```

Returns the content a given neutrino mass eigenstate specified by eig $(0 = \nu_1, 1 = \nu_2, 2 = \nu_3, \ldots)$. This function can only be use in the *single energy mode*.

• Mass composition evaluator (multiple energy mode)

```
double EvalMassAtNode(int eig,int ie,int rho = 0);
double EvalMass(int eig,double Enu,int rho = 0);
```

EvalMassAtNode returns the content a given neutrino mass eigenstate specified by eig $(0 = \nu_1, 1 = \nu_2, 2 = \nu_3, \ldots)$ at an energy node ie. Furthermore, EvalMass returns the approximate content of a given mass eigenstate for a specific neutrino energy Enu by interpolating in the interaction basis. In each function, when considering NT = "both", the parameter rho toggles between neutrino (0) and antineutrino (1).

• Evolve ν state

```
void EvolveState(void);
```

Once the neutrino propagation problem has been setup this function evolves the neutrino state (in either *single* energy or *multiple* energy mode) from its initial position to its final position.

• Get energy nodes values

```
std::vector < double > GetERange(void);
```

Returns a vector containing the energy nodes positions given in natural units, i.e. eV.

• Get number of energy nodes

```
int GetNumE(void);
```

Returns the number of energy nodes.

• Get number neutrino flavors

```
int GetNumNeu(void);
```

Returns the number of neutrino flavors.

• Get Hamiltonian at current position

Returns the SU_vector that represents the (anti)neutrino Hamiltonian at a position specified by a Track and given neutrino energy Enu (in eV). Furthermore, rho specifies whether the neutrino or antineutrino Hamiltonian is returned.

• Get Body

```
std::shared_ptr <Body > GetBody(void);
```

Returns the Body instance currently stored in the nuSQUIDS object.

• Get Track

```
std::shared_ptr<Track> GetTrack(void);
```

Returns the Track instance currently stored in the nuSQUIDS object.

• Set Body

```
void Set_Body(std::shared_ptr <Body>);
```

Sets the Body instance in which the neutrino propagation will take place.

• Set Track

```
void Set_Track(std::shared_ptr<Track>);
```

Sets the Track instance which describes the neutrino propagation inside a given Body.

• Set the initial state

```
void Set_initial_state(array1D, string basis);
void Set_initial_state(array2D, string basis);
void Set_initial_state(array3D, string basis);
```

Set_initial_state sets the initial neutrino (and antineutrino) state. The states can be specified for the single and multiple energy modes can be specified using the different C++ signatures.

- array1D state: Can only be used in single energy mode and is defined by state $[\alpha] = \phi_{\alpha}$ where α is a flavor or mass eigenstate index.
- array2D state: Can only be used in multiple energy mode and is defined by state[ei][α] = $\phi_{\alpha}(E[\text{ei}])$, i.e. the flavor (mass) eigenstate composition at a given energy node ei.

- array3D state: Can only be used in multiple energy mode and is defined by state[ei][ρ][α] = $\phi_{\alpha}^{\rho}(E[ei])$, i.e. the flavor (mass) eigenstate composition at a given energy node ei, and where $\rho = 0 \equiv$ neutrino and $\rho = 0 \equiv$ antineutrino.
- Write the state into an HDF5 file.

```
void WriteStateHDF5(string);
```

Writes the current nuSQUIDS configuration and state into an HDF5 file for later use.

• Read from an HDF5 file.

```
void ReadStateHDF5(string);
```

Reads an previously generated HDF5 file and sets the nuSQUIDS object accordingly, i.e. it configures it and loads the saved state.

4. Examples and benchmarks

4.1. Single energy mode

In this section we illustrate the usage of the *single* energy mode. A full written example can be found in nuSQUIDS/examples/osc_example.cpp; here only code snippets will be shown.

4.1.1. Example 1

In this first example, shown in Listing 4.1.1, we will define a nuSQUIDS instance, set the mixing angles, and then calculate $P(\nu_{\mu} \to \nu_{e})$ for a 500 km baseline experiment in the Earth.

```
// create a nuSQUIDS object setting number of flavors = 3
   nuSQUIDS nus(3,"neutrino");
3
  // We can change the mixing angles from their defaults
4
   nus.Set("th12",0.563942);
  nus.Set("th13",0.154085);
   nus.Set("th23",0.785398);
   // square mass differences
   nus.Set("dm21sq",7.65e-05);
10
   nus.Set("dm31sq",0.00247);
11
   // CP phase
12
   nus.Set("delta1",0.0);
13
14
   // We must set the neutrino energy
   nus.Set_E(10.0*nus.units.GeV);
15
16
  // For this example lets consider a long baseline experiment
17
  double baseline = 500.0*nus.units.km;
18
19 // create a body object, in this case the Earth
```

```
20
  |std::shared_ptr<Earth> earth = std::make_shared<Earth>();
  // create a trajectory on the body. In this case the Earth
  // trajectory is given by three quantities: Initial position,
  // Final position, and Baseline.
   std::shared_ptr<Earth::Track> earth_track = std::make_shared<Earth</pre>
      :: Track > (0.0, baseline, baseline);
   // Then we set the body and trajectory in the nuSQUIDS object.
25
26
   nus.Set_Body(earth);
27
   nus.Set_Track(earth_track);
28
29
   // We must specify the initial neutrino flavor composition
   // In this case we will start with a pure nu_mu state.
30
   std::vector<double> ini_state{0,1,0};
31
32
  nus.Set_initial_state(ini_state, "flavor");
33
34
  // Now that we have all the pieces in place we can tell the
35
   // nuSQUIDS object to evolve the given state.
  nus.EvolveState();
36
37
38
   // Then we can evaluate, for example, the nu_e content
  nus.EvalFlavor(0);
```

4.1.2. Example 2

Even though we recommend to use the *multiple* energy mode when considering a range of energies this can also be performed in the single energy mode. Lets assume that a set of energies is provided in std::vector<double> energy_range, then the code snip in Listing 4.1.2 presents how this could be implemented (results of this can be seen in Fig. 4.1.2)

```
nuSQUIDS nus(3,"neutrino");
   double phi = acos(-1.0);
   std::vector<double> ini_state{0,1,0};
   nus.Set_Body(std::make_shared < EarthAtm >());
   nus.Set_Track(std::make_shared < EarthAtm::Track > (phi));
   for( double enu : energy_range ){
6
7
           nus.Set_E(enu);
8
           nus.Set_initial_state(ini_state, "flavor");
9
           nus.EvolveState();
10
           std::cout << nus.EvalFlavor(1) << std::endl;</pre>
11 | }
```

In the above example we could had rather fix the energy and reset the Track in order to explore the probability as a function of baseline.

4.1.3. Example 3

We can also use the custom environments provided, such as ConstantDensity and VariableDensity. For example in case we want to consider $\rho = 13 \mathrm{gr/cm^3}$ and $y_e = 0.5$ we

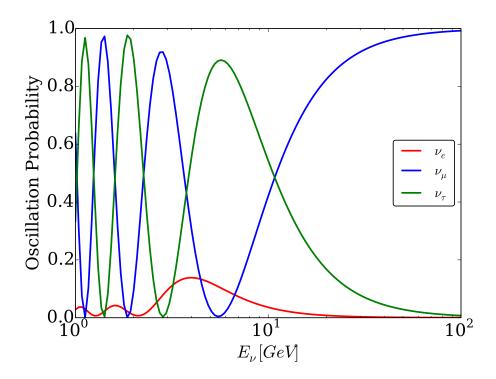


Figure 4.1: Single mode neutrino oscillation probability in a constant density matter environment.

can use the code snipped in Listing 4.1.3, which result is shown in Fig. 4.1.3,

```
nuSQUIDS nus(3,"neutrino");
1
2
   double phi = acos(-1.0);
3
   std::vector<double> ini_state{0,1,0};
   nus.Set_Body(std::make_shared < ConstantDensity > (13.0,0.5));
4
   nus.Set_Track(std::make_shared < ConstantDensity::Track > (0.0,1000.0*
      nus.units.km));
6
   for( double enu : energy_range ){
7
           nus.Set_E(enu);
8
           nus.Set_initial_state(ini_state, "flavor");
9
           nus.EvolveState();
10
           std::cout << nus.EvalFlavor(0) << std::endl;</pre>
11 | }
```

4.2. Multiple energy mode

As explained in the text the *multiple* energy mode can consider a *set* of neutrino energies. The simplest thing that can be done is to

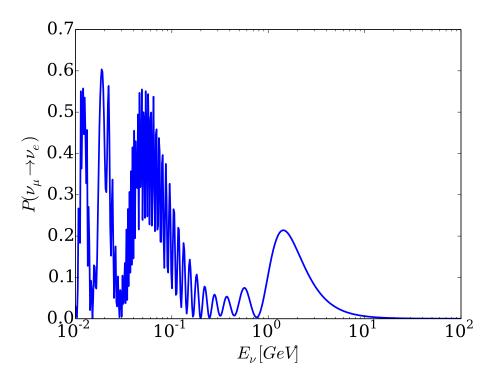


Figure 4.2: Single mode neutrino oscillation probability in a constant density matter environment.

4.2.1. Example 1

In this mode we can actually input a neutrino flux and propagate it in its environment. For example we could propagate $\phi(E) = \phi_0 E^{-1} \Pi_{\mu}$, i.e. power law pure ν_{μ} flux, through the Earth. In Listing 4.2.1 we shall consider the this case and further demonstrate the how to consider more neutrino flavors. The results of the following code are shown in Figure 4.2.1.

```
1
   nuSQUIDS nus(1.,1.e2,200,4,"neutrino",true,false);
2
3
   double phi = acos(-1.);
   std::shared_ptr<EarthAtm> earth_atm = std::make_shared<EarthAtm>();
   std::shared_ptr<EarthAtm::Track> track_atm = std::make_shared<</pre>
      EarthAtm::Track>(phi);
6
7
   // set sterile parameters while leaving the rest at default
   nus.Set("dm41sq", 3.0);
9
   nus.Set("th24", 0.7);
10
   nus.Set("rel_error", 1.0e-9);
11
12
   nus.Set("abs_error", 1.0e-9);
13
14 // construct the initial state
```

```
15
   vector < double > E_range = nus.GetERange();
16
   array2D inistate(E_range.size());
17
   double NO = 1.0e18;
   for ( int i = 0 ; i < inistate.size(); i++){</pre>
18
19
          inistate[i].resize(3);
20
         for ( int k = 0; k < 3; k ++){
21
            // initialze muon state
22
            inistate[i][k] = (k == 1) ? NO*pow(E_range[i],-2) : 0.0;
         }
23
24
   // set the initial state
25
26
   nus.Set_initial_state(inistate, "flavor");
27
28
   nus.EvolveState();
29
   // save the result
30 | nus.WriteStateHDF5("./mulene_example1.hdf5");
```

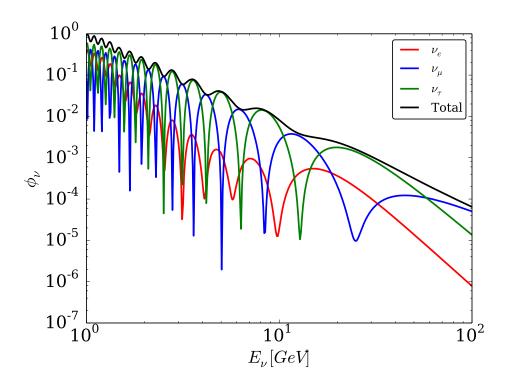


Figure 4.3: Single mode neutrino oscillation probability in a constant density matter environment.

4.2.2. Example 2

We can further consider the propagation of $\phi(E) = \phi_0 E^{-1}(\Pi_\mu + \Pi_\tau)$, again through the Earth, but considering τ regeneration.

5. Conclusions & Acknowledgements

Abazajian, K. N., Acero, M. A., Agarwalla, S. K., Aguilar-Arevalo, A. A., Albright, C. H., Antusch, S., Argüelles, C. A., Balantekin, A. B., Barenboim, G., Barger, V., Bernardini, P., Bezrukov, F., Bjaelde, O. E., Bogacz, S. A., Bowden, N. S., Boyarsky, A., Bravar, A., Berguno, D. B., Brice, S. J., Bross, A. D., Caccianiga, B., Cavanna, F., Chun, E. J., Cleveland, B. T., Collin, A. P., Coloma, P., Conrad, J. M., Cribier, M., Cucoanes, A. S., D'Olivo, J. C., Das, S., de Gouvea, A., Derbin, A. V., Dharmapalan, R., Diaz, J. S., Ding, X. J., Djurcic, Z., Donini, A., Duchesneau, D., Ejiri, H., Elliott, S. R., Ernst, D. J., Esmaili, A., Evans, J. J., Fernandez-Martinez, E., Figueroa-Feliciano, E., Fleming, B. T., Formaggio, J. A., Franco, D., Gaffiot, J., Gandhi, R., Gao, Y., Garvey, G. T., Gavrin, V. N., Ghoshal, P., Gibin, D., Giunti, C., Gninenko, S. N., et al. (129 additional authors not shown), 04 2012. Light sterile neutrinos: A white paper.

URL http://arxiv.org/abs/1204.5379

Aguilar, A., Auerbach, L., Burman, R., Caldwell, D., Church, E., Cochran, A., Donahue, J., Fazely, A., Garvey, G., Gunasingha, R., Imlay, R., Louis, W., Majkic, R., Malik, A., Metcalf, W., Mills, G., Sandberg, V., Smith, D., Stancu, I., Sung, M., Tayloe, R., VanDalen, G., Vernon, W., Wadia, N., White, D., Yellin, S., 2001. Evidence for neutrino oscillations from the observation of electron anti-neutrinos in a muon anti-neutrino beam.

URL http://arxiv.org/abs/hep-ex/0104049

Akhmedov, E. K., 1999. Neutrino physics, 103–164.

Archidiacono, M., Fornengo, N., Gariazzo, S., Giunti, C., Hannestad, S., Laveder, M., 04 2014. Light sterile neutrinos after bicep-2.

URL http://arxiv.org/abs/1404.1794

Argüelles Delgado, C. A., Kopp, J., 2012. Sterile neutrinos and indirect dark matter searches in IceCube. JCAP 1207, 016.

Argüelles Delgado, C. A., Salvado Serra, J., 2014. A simple quantum integro differential solver (squids).

Bahcall, J. N., Serenelli, A. M., Basu, S., 2005. New solar opacities, abundances, helioseismology, and neutrino fluxes. The Astrophysical Journal Letters 621 (1), L85.

Balantekin, A. B., Haxton, W. C., 03 2013. Neutrino oscillations.

URL http://arxiv.org/abs/1303.2272

Blennow, M., Edsjo, J., Ohlsson, T., 2008. Neutrinos from WIMP annihilations using a full three-flavor Monte Carlo. JCAP 0801, 021.

Blennow, M., Smirnov, A. Y., 06 2013. Neutrino propagation in matter.

URL http://arxiv.org/abs/1306.2903

Cirelli, M., Fornengo, N., Montaruli, T., Sokalski, I., Strumia, A., Vissani, F., 2005. Spectra of neutrinos from dark matter annihilations.

URL http://arxiv.org/abs/hep-ph/0506298

Conrad, J. M., Louis, W. C., Shaevitz, M. H., 06 2013. The lsnd and miniboone oscillation searches at high m^2 .

URL http://arxiv.org/abs/1306.6494

Dasgupta, B., Kopp, J., 10 2013. A ménage à trois of ev-scale sterile neutrinos, cosmology, and structure formation.

URL http://arxiv.org/abs/1310.6337

de Gouvea, A., Pitts, K., Scholberg, K., Zeller, G., Alonso, J., Bernstein, A., Bishai, M., Elliott, S., Heeger, K., Hoffman, K., Huber, P., Kaufman, L., Kayser, B., Link, J., Lunardini, C., Monreal, B., Morfin, J., Robertson, H., Tayloe, R., Tolich, N., Abazajian, K., Akiri, T., Albright, C., Asaadi, J., Babu, K., Balantekin, A., Barbeau, P., Bass, M., Blake, A., Blondel, A., Blucher, E., Bowden, N., Brice, S., Bross, A., Carls, B., Cavanna, F., Choudhary, B., Coloma, P., Connolly, A., Conrad, J., Convery, M., Cooper, R., Cowen, D., da Motta, H., de Young, T., Lodovico, F. D., Diwan, M., Djurcic, Z., Dracos, M., Dodelson, S., Efremenko, Y., Ekelof, T., Feng, J., Fleming, B., Formaggio, J., Friedland, A., Fuller, G., Gallagher, H., Geer, S., Gilchriese, M., Goodman, M., Grant, D., Gratta, G., Hall, C., Halzen, F.,

- Harris, D., Heffner, M., et al. (100 additional authors not shown), 10 2013. Neutrinos. URL http://arxiv.org/abs/1310.4340
- Duan, H., Fuller, G. M., Qian, Y.-Z., 01 2010. Collective neutrino oscillations. URL http://arxiv.org/abs/1001.2799
- Dutta, S. I., Reno, M. H., Sarcevic, I., 2000. Tau neutrinos underground: Signals of muon-neutrino \rightarrow tau-neutrino oscillations with extragalactic neutrinos. Phys.Rev. D62, 123001.
- Dziewonski, A. M., Anderson, D. L., 1981. Preliminary reference earth model. Physics of the earth and planetary interiors 25 (4), 297–356.
- Folk, M., Cheng, A., Yates, K., 1999. Hdf5: A file format and i/o library for high performance computing applications. In: Proceedings of Supercomputing. Vol. 99.
- Fukugita, M., Yanagida, T., 2003. Physics of Neutrinos: And Applications to Astrophysics. Physics and astronomy online library. Springer.
 - URL http://books.google.fr/books?id=E18R8nC5QegC
- Gonzalez-Garcia, M., Halzen, F., Maltoni, M., 2005. Physics reach of high-energy and high-statistics icecube atmospheric neutrino data. Phys.Rev. D71, 093010.
- Gonzalez-Garcia, M., Maltoni, M., 2008. Phenomenology with Massive Neutrinos. Phys.Rept. 460, 1–129.
- Gonzalez-Garcia, M., Maltoni, M., Schwetz, T., 2014. Updated fit to three neutrino mixing: status of leptonic CP violation.
- Gough, B., 2009. GNU scientific library reference manual. Network Theory Ltd.
- Gribov, V., Pontecorvo, B., 1969. Neutrino astronomy and lepton charge. Phys.Lett. B28, 493.
- Halzen, F., Saltzberg, D., 1998. Tau neutrino appearance with a 1000 megaparsec baseline. URL http://arxiv.org/abs/hep-ph/9804354
- Hewett, J., Weerts, H., Brock, R., Butler, J., Casey, B., Collar, J., de Gouvea, A., Essig, R., Grossman, Y., Haxton, W., Jaros, J., Jung, C., Lu, Z., Pitts, K., Ligeti, Z., Patterson, J., Ramsey-Musolf, M., Ritchie, J., Roodman, A., Scholberg, K., Wagner, C., Zeller, G., Aefsky, S., Afanasev, A., Agashe, K., Albright, C., Alonso, J., Ankenbrandt, C., Aoki, M., Arguelles, C., Arkani-Hamed, N., Armendariz, J., Armendariz-Picon, C., Diaz, E. A., Asaadi, J., Asner, D., Babu, K., Bailey, K., Baker, O., Balantekin, B., Baller, B., Bass, M., Batell, B., Beacham, J., Behr, J., Berger, N., Bergevin, M., Berman, E., Bernstein, R., Bevan, A., Bishai, M., Blanke, M., Blessing, S., Blondel, A., Blum, T., Bock, G., Bodek, A., Bonvicini, G., Bossi, F., Boyce, J., Breedon, R., Breidenbach, M., Brice, S., Briere, R., Brodsky, S., et al. (403 additional authors not shown), 05 2012. Fundamental physics at the intensity frontier.
 URL http://arxiv.org/abs/1205.2671
- Huber, P., Kopp, J., Lindner, M., Rolinec, M., Winter, W., 2007. New features in the simulation of neutrino oscillation experiments with GLoBES 3.0: General Long Baseline Experiment Simulator. Comput. Phys. Commun. 177, 432–438.
- Kopp, J., Machado, P. A., Maltoni, M., Schwetz, T., 2013. Sterile neutrino oscillations: the global picture. Journal of High Energy Physics 2013 (5), 1–52.
- Kopp, J., Welter, J., 08 2014. The not-so-sterile 4th neutrino: Constraints on new gauge interactions from neutrino oscillation experiments.
 - URL http://arxiv.org/abs/1408.0289
- Mention, G., Fechner, M., Lasserre, T., Mueller, T. A., Lhuillier, D., Cribier, M., Letourneau, A., 01 2011. The reactor antineutrino anomaly.
 - URL http://arxiv.org/abs/1101.2755
- Mikheev, S., Smirnov, A. Y., 1985. Resonance Amplification of Oscillations in Matter and Spectroscopy of Solar Neutrinos. Sov.J.Nucl.Phys. 42, 913–917.
- Mikheev, S., Smirnov, A. Y., 1986. Resonant amplification of neutrino oscillations in matter and solar neutrino spectroscopy. Nuovo Cim. C9, 17–26.
- MiniBooNE-Collaboration, Aguilar-Arevalo, A. A., Brown, B. C., Bugel, L., Cheng, G., Church, E. D., Conrad, J. M., Dharmapalan, R., Djurcic, Z., Finley, D. A., Ford, R., Garcia, F. G., Garvey, G. T., Grange, J., Huelsnitz, W., Ignarra, C., Imlay, R., Johnson, R. A., Karagiorgi, G., Katori, T., Kobilarcik, T., Louis, W. C., Mariani, C., Marsh, W., Mills, G. B., Mirabal, J., Moore, C. D., Mousseau, J., Nienaber,

P., Osmanov, B., Pavlovic, Z., Perevalov, D., Polly, C. C., Ray, H., Roe, B. P., Russell, A. D., Shaevitz, M. H., Spitz, J., Stancu, I., Tayloe, R., de Water, R. G. V., White, D. H., Wickremasinghe, D. A., Zeller, G. P., Zimmerman, E. D., 07 2012. A combined \rightarrow_e and $^- \rightarrow_e^-$ oscillation analysis of the miniboone excesses.

URL http://arxiv.org/abs/1207.4809

Mohapatra, R., Antusch, S., Babu, K., Barenboim, G., Chen, M.-C., Davidson, S., de Gouvea, A., de Holanda, P., Dutta, B., Grossman, Y., Joshipura, A., Kayser, B., Kersten, J., Keum, Y., King, S., Langacker, P., Lindner, M., Loinaz, W., Masina, I., Mocioiu, I., Mohanty, S., Murayama, H., Pascoli, S., Petcov, S., Pilaftsis, A., Ramond, P., Ratz, M., Rodejohann, W., Shrock, R., Takeuchi, T., Underwood, T., Wolfenstein, L., 2005. Theory of neutrinos: A white paper.

URL http://arxiv.org/abs/hep-ph/0510213

Pontecorvo, B., 1968. Neutrino Experiments and the Problem of Conservation of Leptonic Charge. Sov.Phys.JETP 26, 984–988.

Pospelov, M., 03 2011. Neutrino physics with dark matter experiments and the signature of new baryonic neutral currents.

URL http://arxiv.org/abs/1103.3261

Saviano, N., Pisanti, O., Mangano, G., Mirizzi, A., 09 2014. Unveiling secret interactions among sterile neutrinos with big-bang nucleosynthesis.

URL http://arxiv.org/abs/1409.1680

Strack, P., Burrows, A., 2005. Generalized boltzmann formalism for oscillating neutrinos.

URL http://arxiv.org/abs/hep-ph/0504035

Wallraff, M., Wiebusch, C., 09 2014. Calculation of oscillation probabilities of atmospheric neutrinos using nucraft.

URL http://arxiv.org/abs/1409.1387

Wolfenstein, L., 1978. Neutrino Oscillations in Matter. Phys.Rev. D17, 2369–2374.

Zhang, Y., Burrows, A., 10 2013. Transport equations for oscillating neutrinos.

URL http://arxiv.org/abs/1310.2164