一元傅里叶变换的应用

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1 方程的频率方法(常微分方程)

$$\mathcal{F}D^{\alpha} = (i\omega)^{\alpha}\mathcal{F},\tag{1.1}$$

一维时,

$$(f')^{\wedge}(\omega) = (i\omega)\hat{f}(\omega) \tag{1.2}$$

其中, $D^{\alpha} = D_1^{\alpha_1} D_2^{\alpha_2} \cdots D_n^{\alpha_n}$, $D_j = \frac{\partial}{\partial x_j}$, $X = (x_1, \dots, x_n)$, $X^{\alpha} = x_1^{\alpha_1} \cdots x_n^{\alpha_n}$

例题 1.1.

$$\int_0^{+\infty} g(w)\sin wx dw = f(x) \tag{1.3}$$

求积分方程(1.3)的解 g(w)

其中

$$f(x) = \begin{cases} \frac{\pi}{2} \frac{e^{-x}}{x}, & x > 0\\ 0, & x \le 0 \end{cases}$$
 (1.4)

解 1.1. 方法一: 用正弦积分变换(略), 见教材 P99 例 1.6

方法二 假设 g 为奇函数,则:

$$g(w) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \hat{g}(w)e^{-iwx}dw = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \hat{g}(w)(-i)\sin wxdw$$
$$= \frac{-i}{\sqrt{2\pi}} \int_{0}^{\infty} g(w)\sin wxdw,$$
 (1.5)

由上可得

$$\frac{\sqrt{2\pi}}{-2i}\hat{g}(x) = f(x) \tag{1.6}$$

又 g 为奇函数, 有 $\hat{g}(-w) = -\hat{g}(w)$, 故方程变为:

$$\frac{\sqrt{2\pi}}{-2i}\hat{g}(t) = f(t) \tag{1.7}$$

最后解得 g(w) 为奇函数, 仍记为 f.

$$f(t) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(w)e^{-iwt}dw = \frac{-2i}{\sqrt{2\pi}} \int_{0}^{+\infty} f(w)\sin wt dw$$

$$= \frac{-2i}{\sqrt{2\pi}} \int_{0}^{\pi} \frac{\pi}{2}\sin wt \sin wt dw$$

$$= \frac{-2i}{\sqrt{2\pi}} \cdot \frac{\pi}{4} \int_{0}^{\pi} [\cos(1-t)w - \cos(1+t)w]dw$$

$$= \frac{-2i}{\sqrt{2\pi}} \cdot \frac{\pi}{4} \left(\frac{\sin(1-t)\pi}{1-t} + \frac{\sin(1+t)\pi}{1+t}\right)$$

$$= \frac{-2i}{\sqrt{2\pi}} \cdot \frac{\pi}{4} \cdot 2 \cdot \frac{\sin \pi t}{1-t^2} = \frac{-2i}{\sqrt{2\pi}} \cdot \frac{\pi}{2} \cdot \frac{\sin \pi t}{1-t^2}$$

$$(1.8)$$

对(1.7)两边做傅里叶变换

$$\frac{\sqrt{2\pi}}{-2i}\hat{g}(t) = f(t) \tag{1.9}$$

利用做两次傅里叶变换是镜像对称变换的性质

$$\mathcal{F}^2\{f(t)\} = \mathcal{F}\{\mathcal{F}\{f(t)\}\} = f(-t) \tag{1.10}$$

$$\frac{\sqrt{2\pi}}{-2i}\hat{g}(-t) = \hat{f}(t) \tag{1.11}$$

而且 g 为奇函数可知

$$g(-t) = \frac{-2i}{\sqrt{2\pi}} \cdot \frac{\pi}{2} \cdot \frac{\sin \pi t}{1 - t^2}$$
 (1.12)

$$g(t) = \frac{\sin \pi t}{1 - t^2} \tag{1.13}$$

例题 1.2. 积分方程解

$$g(t) = h(t) + \int_{-\infty}^{\infty} f(t)g(t-x)dx$$
(1.14)

h、f已知,且g、h、f的傅里叶变换存在。

解 1.2. 由傅里叶变换的卷积公式:

$$\mathcal{F}\{f * g\} = \sqrt{2\pi} \cdot \hat{f} \cdot \hat{g} \tag{1.15}$$

原方程可化为:

$$\hat{g}(w) = \hat{h}(w) + \hat{f}(w) \cdot \hat{g}(w) \cdot \sqrt{2\pi}$$
(1.16)

解得:

$$\hat{g}(w) = \frac{\hat{h}(w)}{1 - \sqrt{2\pi} \cdot \hat{f}(w)} \tag{1.17}$$

因此:

$$g(t) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \frac{\hat{h}(w)}{1 - \sqrt{2\pi} \cdot \hat{f}(w)} e^{iwt} dw$$
 (1.18)

例题 1.3. 常微分非齐次线性积分方程:

$$y'' - y = -f \tag{1.19}$$

其中 f 为已知函数

解 1.3. 对两边取傅里叶变换:

$$(iw)^2 \hat{y}(w) - \hat{y}(w) = -\hat{f}(w)$$
(1.20)

解得:

$$\hat{y}(w) = \frac{\hat{f}(w)}{1 + w^2} \tag{1.21}$$

因此:

$$y(x) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \frac{\hat{f}(w)}{1+w^2} e^{iwx} dw$$
 (1.22)

将 $\frac{\hat{f}(w)}{1+w^2}$ 视为 $\hat{f}(w)$ 与 $\frac{1}{1+w^2}$ 的乘积。

由卷积定理:

$$(f * h)(w) = \sqrt{2\pi} \hat{f}(w) \cdot \hat{h}(w), \quad \sharp + \hat{h}(w) = \frac{1}{1 + w^2}.$$
 (1.23)

因此:

$$\hat{y}(w) = \sqrt{2\pi} \hat{f}(w) \cdot \hat{g}(w), \quad \sharp \, \hat{g}(w) = \frac{1}{\sqrt{2\pi}} \cdot \frac{1}{1 + w^2}. \tag{1.24}$$

所以:

$$y(t) = (f * g)(t), \quad \sharp + g(w) = \frac{1}{2}e^{-|t|}.$$
 (1.25)

因此:

$$y(t) = \frac{1}{2} \int_{-\infty}^{\infty} f(t)e^{-|t-x|} dx.$$
 (1.26)

例题 1.4. 求解积分方程:

$$ax'(t) + bx(t) + c \int_0^t x(t)dt = h(t),$$
 (1.27)

其中 $a,b,c \in \mathbb{R}$, h 已知

解 1.4. 关键是通过傅里叶变换求解。设 G'=g,且 G 有傅里叶变换。对两边取傅里叶变换:

$$(iw)\hat{G}(w) = \hat{g}(w) \Rightarrow \hat{G}(w) = \frac{\hat{g}(w)}{iw}.$$
(1.28)

利用此公式,原方程可变换为:

$$a(iw)\hat{x}(w) + b\hat{x}(w) + c\frac{\hat{x}(w)}{iw} = h(w)$$
 (1.29)

解得:

$$\hat{x}(w) = \frac{h(w)}{iaw + b + \frac{c}{iw}}. (1.30)$$

PDE 的傅里叶变换 2

-维波动方程初值问题 2.1

例题 2.1. 求解一维波动方程初值问题:

$$\left(\frac{\partial^2 u}{\partial t^2} = \frac{\partial^2 u}{\partial x^2}, \quad x \in \mathbb{R}, t > 0 \right)$$
(2.1a)

$$u|_{t=0} = \varphi_0(x) \tag{2.1b}$$

$$\begin{cases} \frac{\partial^2 u}{\partial t^2} = \frac{\partial^2 u}{\partial x^2}, & x \in \mathbb{R}, t > 0 \\ u|_{t=0} = \varphi_0(x) & (2.1a) \\ \frac{\partial u}{\partial t}\Big|_{t=0} = \varphi_1(x) & (2.1c) \end{cases}$$

解 2.1. 对二元函数 u(x,t) 的 x 变量作傅里叶变换,记记之为 V(w,t)。

则:

$$V(w,t) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} u(x,t)e^{-iwx}dx$$
 (2.2)

于是:

$$\frac{\partial V}{\partial t}(w,t) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \frac{\partial u}{\partial t}(x,t) e^{-iwx} dx = \mathcal{F}_1\left(\frac{\partial u}{\partial t}(x,t)\right)$$
 (2.3)

有以下 6 个公式

$$\mathcal{F}_1\left(\frac{\partial u}{\partial x}\right)(w,t) = \frac{\partial}{\partial t}V(w,t) \quad (\partial w \, \, \mathbf{h} \, \, \mathbf{\ddot{x}} \, \, \mathbf{\dot{y}}) \tag{2.4}$$

$$\mathcal{F}_1\left(\frac{\partial u}{\partial x}\right)(w,t) = (iw)V(w,t) \tag{2.5}$$

$$\mathcal{F}_1\left(\frac{\partial^2 u}{\partial x^2}\right)(w,t) = \frac{\partial^2}{\partial t^2}V(w,t) \tag{2.6}$$

$$\mathcal{F}_1\left(\frac{\partial^2 u}{\partial x^2}\right)(w,t) = (iw)^2 V(w,t) \tag{2.7}$$

$$\mathcal{F}_1(\sin x)(w) = \sqrt{\frac{\pi}{2}}i\left[\delta(w+1) - \delta(w-1)\right]$$
(2.8)

$$\mathcal{F}_1(\cos x)(w) = \sqrt{\frac{\pi}{2}} \left[\delta(w+1) + \delta(w-1) \right]$$
 (2.9)

$$\mathcal{F}_1(\cos x)(w) = \sqrt{\frac{\pi}{2}} \left[\delta(w+1) + \delta(w-1) \right]$$
 (2.10)

原方程在频率域可转化为常微分方程问题

$$\frac{d^2V}{dt^2} = -w^2V$$
(2.11a)

$$V|_{t=0} = \sqrt{\frac{\pi}{2}} \left[\delta(w+1) + \delta(w-1) \right]$$
 (2.11b)

$$\begin{cases} \frac{d^2V}{dt^2} = -w^2V & (2.11a) \\ V|_{t=0} = \sqrt{\frac{\pi}{2}} \left[\delta(w+1) + \delta(w-1) \right] & (2.11b) \\ \frac{dV}{dt} \Big|_{t=0} = \sqrt{\frac{\pi}{2}} \cdot i \left[\delta(w+1) - \delta(w-1) \right] & (2.11c) \end{cases}$$

通解如下:

$$V(w,t) = C_1 \sin(wt) + C_2 \cos(wt)$$
 (2.12)

注记 2.1. 齐次方程特征方程为 $\lambda^2+w^2=0$,其解为 $\lambda=\pm iw$,对应的解为 $e^{\pm iwt}$,即 $\cos(wt)$ 和 $\sin(wt)$ 。

由(2.11b)

$$C_2 = \sqrt{\frac{\pi}{2}} \left[\delta(w+1) + \delta(w-1) \right]$$
 (2.13)

曲(2.11c)

$$C_1 = \sqrt{\frac{\pi}{2}} \cdot \frac{i}{w} \left[\delta(w+1) - \delta(w-1) \right]$$
 (2.14)

最后得

$$V(w,t) = \sqrt{\frac{\pi}{2}} \cdot \frac{i}{w} \left[\delta(w+1) - \delta(w-1) \right] \sin(wt)$$

$$+ \sqrt{\frac{\pi}{2}} \left[\delta(w+1) + \delta(w-1) \right] \cos(wt)$$

$$= \left(\sqrt{\frac{\pi}{2}} \cdot \frac{i}{w} \sin(wt) + \sqrt{\frac{\pi}{2}} \cos(wt) \right) \delta(w+1)$$

$$+ \left(-\sqrt{\frac{\pi}{2}} \cdot \frac{i}{w} \sin(wt) + \sqrt{\frac{\pi}{2}} \cos(wt) \right) \delta(w-1)$$
(2.15)

做傅里叶逆变换, 从频率域变回时间域

$$u(x,t) = \mathcal{F}_{1}^{-1}(V(\cdot,t))(x)$$

$$= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \left\{ \left[\sqrt{\frac{\pi}{2}} \cdot \frac{i}{w} \sin(wt) + \sqrt{\frac{\pi}{2}} \cos(wt) \right] \delta(w+1) + \left[-\sqrt{\frac{\pi}{2}} \cdot \frac{i}{w} \sin(wt) + \sqrt{\frac{\pi}{2}} \cos(wt) \right] \delta(w-1) \right\} e^{iwx} dw$$

$$= \frac{1}{\sqrt{2\pi}} \left\{ \left[\sqrt{\frac{\pi}{2}} \cdot \frac{i}{-1} \sin(-t) + \sqrt{\frac{\pi}{2}} \cos(-t) \right] e^{-ix} + \left[\sqrt{\frac{\pi}{2}} \cdot \frac{i}{1} \sin(t) + \sqrt{\frac{\pi}{2}} \cos(t) \right] e^{ix} \right\}$$

$$= \frac{1}{2} \left\{ [-i \sin(t) + \cos(t)] e^{-ix} + [i \sin(t) + \cos(t)] e^{ix} \right\}$$

$$= \frac{1}{2} \left\{ \cos(t) e^{-ix} + \cos(t) e^{ix} + i \sin(t) e^{ix} - i \sin(t) e^{-ix} \right\}$$

$$= \frac{1}{2} \left\{ \cos(t) \left(e^{ix} + e^{-ix} \right) + i \sin(t) \left(e^{ix} - e^{-ix} \right) \right\}$$

$$= \frac{1}{2} \left\{ 2 \cos(t) \cos(x) + 2i \sin(t) \sin(x) \right\}$$

$$= \cos(t - x)$$

一维热传导方程

例题 2.2. 求解一维热传导方程初值问题:

$$\begin{cases} \frac{\partial u}{\partial t} = a^2 \frac{\partial^2 u}{\partial x^2} + f(x, t), x \in \mathbb{R}, t > 0 \\ u|_{t=0} = \varphi(x) \end{cases}$$
 (2.17a)

解 2.2. 令 $V(w,t) = \mathcal{F}_1(u(\cdot,t))(w)$. 则

$$\begin{cases}
\mathcal{F}_1\left(\frac{\partial u}{\partial t}(\cdot,t)\right) = \frac{\partial}{\partial t}(\mathcal{F}_1 u(w,t)) = \frac{\partial}{\partial t}V(w,t) = \frac{dV}{dt}(w,t) \\
\mathcal{F}_1\left(\frac{\partial^2 u}{\partial x^2}\right)(w) = (iw)^2 V(w,t) = -w^2 V(w,t)
\end{cases} (2.18a)$$

$$\mathcal{F}_1\left(\frac{\partial^2 u}{\partial x^2}\right)(w) = (iw)^2 V(w,t) = -w^2 V(w,t)$$
(2.18b)

记

$$\begin{cases} \hat{f}_1(w,t) = \mathcal{F}_1 f(\cdot,t)(w) \\ \hat{\varphi}(w) = (\mathcal{F}\varphi)(w) \end{cases}$$
(2.19a)

则原方程在频域表示为:

$$\begin{cases} \frac{dV}{dt} = -a^2 w^2 V + \hat{f}_1(w, t) \\ V|_{t=0} = \hat{\varphi}(w) \end{cases}$$
 (2.20a)

方程(2.20a)为一阶非齐次常微分方程。由常数变异法,可得解如下:

$$V(w,t) = \hat{\varphi}(w)e^{-a^2w^2t} + \int_0^t \hat{f}_1(w,\tau)e^{-a^2w^2(t-\tau)}d\tau$$
 (2.21)

$$\mathcal{F}^{-1}\left(e^{-a^2w^2t}\right)(x) = \sqrt{2\pi} \cdot \frac{1}{2a\sqrt{\pi t}}e^{-\frac{x^2}{4a^2t}}$$
 (2.22)

其中,

$$\mathcal{F}\left(e^{-\beta x^2}\right)(w) = \frac{\sqrt{\pi}}{\sqrt{\beta}}e^{-\frac{w^2}{4\beta}} \tag{2.23}$$

 $\diamondsuit \beta = a^2t$,则

$$\mathcal{F}^{-1}\left(e^{-a^2w^2t}\right)(x) = \frac{1}{a\sqrt{2\pi t}}e^{-\frac{x^2}{4a^2t}}$$
(2.24)

进一步地,

$$\mathcal{F}^{-1}\left(e^{-a^2(w-k)^2t}\right)(y) = \frac{1}{a\sqrt{2\pi(t-\tau)}}e^{-\frac{(y-x)^2}{4a^2(t-\tau)}}$$
(2.25)

最终得解:

$$u(x,t) = \left(\varphi(x) * \frac{1}{a\sqrt{2\pi t}} e^{-\frac{x^2}{4a^2t}}\right)(x)$$

$$+ \int_0^t \left(f(x,\tau) * \frac{1}{a\sqrt{2\pi(t-\tau)}} e^{-\frac{x^2}{4a^2(t-\tau)}}\right)(x)d\tau$$
(2.26)

2.3 附

例题 2.3. 求解微分方程初值问题:

$$\begin{cases} \frac{dy}{dx} = Ay + f(x), x \in \mathbb{R} \\ y(0) = C \end{cases}$$
 (2.27a)

解 2.3. 齐次方程 y' = Ay。解得:

$$\frac{y'}{y} = A \implies (\ln y)' = A \implies \ln y = Ax + C_1 \implies y = e^{C_1} e^{Ax}$$
 (2.28)

 $\mathbb{P} y = Ce^{Ax}$.

令 $y = C(x)e^{Ax}$ 。代入原方程:

$$y' = C'(x)e^{Ax} + C(x)e^{Ax}A$$
(2.29)

代入 y' = Ay + f(x):

$$C'(x)e^{Ax} + C(x)e^{Ax}A = AC(x)e^{Ax} + f(x)$$
(2.30)

解得:

$$C'(x) = f(x)e^{-Ax} (2.31)$$

因此:

$$C(x) = \int f(x)e^{-Ax}dx \tag{2.32}$$

特解:

$$C(x) = \int_0^x f(x)e^{-Ax}dx + C$$
 (2.33)

解为:

$$y = \int_0^x f(x)e^{-Ax}dx \cdot e^{Ax} + Ce^{Ax}$$
 (2.34)

当 y(0) = C 时, 解得 C = C。

因此通解:

$$y(x) = Ce^{Ax} + \int_0^x f(x)e^{A(x-x)}dx$$
 (2.35)

求解上半平面无源静电场内电势的边值问题 2.4

例题 2.4. 求解上半平面无源静电场内电势的边值问题 (拉普拉斯方程):

$$\begin{cases} \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} = 0, x \in \mathbb{R}, y > 0 \\ u|_{y=0} = f(x) \end{cases}$$
 (2.36a)

(2.36b)

解 2.4. 记 $V(w,y) = \mathcal{F}_1 u(x,y)(w)$ 。

将原方程化为:

$$\begin{cases}
-w^{2}V(w,y) + \frac{\partial^{2}}{\partial y^{2}}V(w,y) = 0 & (2.37a) \\
V(w,y)|_{y=0} = \hat{f}(w) & (2.37b) \\
\lim_{y \to +\infty} V(w,y) = 0 & (2.37c)
\end{cases}$$

$$V(w,y)|_{y=0} = \hat{f}(w)$$
 (2.37b)

$$\lim_{u \to +\infty} V(w, y) = 0 \tag{2.37c}$$

视 w 为常数, 这是一个二阶常微分方程。特征方程为:

$$\lambda^2 - w^2 = 0 \implies \lambda = \pm |w| \tag{2.38}$$

通解为:

$$V(w,y) = C_1 e^{|w|y} + C_2 e^{-|w|y}$$
(2.39)

由边界条件:

$$V(w,y)|_{y=0} = \hat{f}(w) \implies C_1 + C_2 = \hat{f}(w)$$
 (2.40)

以及

$$\lim_{y \to +\infty} V(w, y) = 0 \implies C_1 = 0 \tag{2.41}$$

因此:

$$V(w,y) = \hat{f}(w)e^{-|w|y}$$
 (2.42)

利用傅里叶逆变换:

$$\mathcal{F}^{-1}\left(e^{-|w|y}\right)(x) = \sqrt{\frac{\pi}{y}} \frac{y}{x^2 + y^2}$$
 (2.43)

上式子可由下式得出

$$\mathcal{F}\left(e^{-\frac{1}{r}}\right) \implies \sqrt{\frac{\pi}{2}} \frac{r}{1 + r^2 \xi^2} \tag{2.44}$$

最终得解

$$u(x,y) = \frac{1}{\sqrt{2\pi}} \left(f * \sqrt{\frac{2}{\pi}} \frac{y}{x^2 + y^2} \right) (x)$$

$$= \frac{1}{\pi} \int_{-\infty}^{\infty} \frac{y}{x^2 + y^2} f(x - t) dt$$
(2.45)

上半平面 Poisson 积分核

$$P_y(x) = \frac{1}{\pi} \cdot \frac{y}{y^2 + x^2} \tag{2.46}$$