

Advection Equation

- If we have a **flow** of the background fluid, mobile things can be transported

J: amount of some stuff (heat, chemicals, biology, etc.)
chain rule!

$$\frac{dJ}{dt} = \frac{\partial J}{\partial t} + \frac{\partial J}{\partial x} \left[\frac{\partial x}{\partial t} \right]$$

→ velocity of frame of reference (i.e. fluid flow)

$$\frac{\partial J}{\partial t} + u \frac{\partial J}{\partial x} = S(x, t)$$

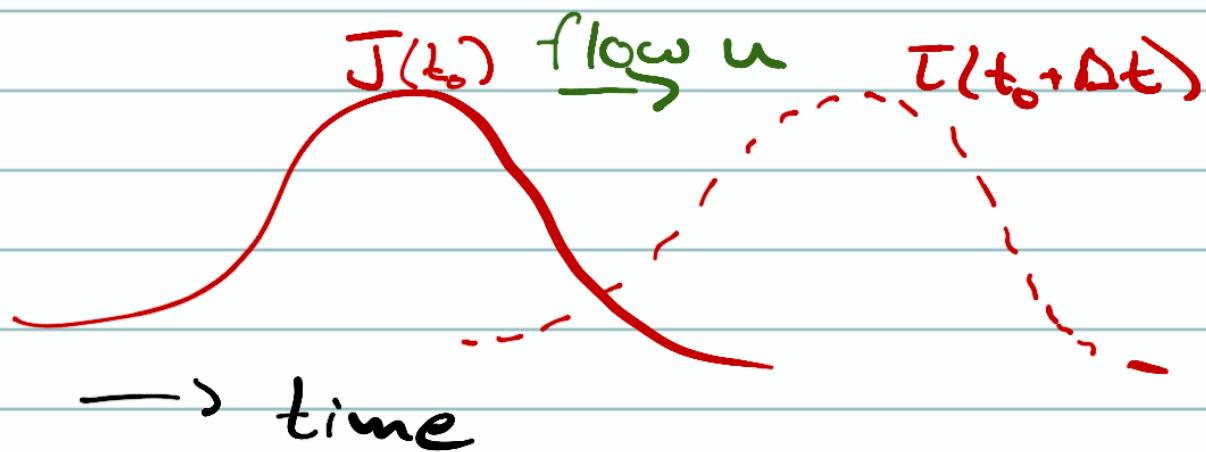
fluid velocity

source of stuff

Advection Equation!

This form assumes u is constant, but it can also be written in more general flux form

$$\frac{\partial J}{\partial t} + \frac{\partial}{\partial x}(uJ) = S(x, t)$$



Advection equation describes how things are transported by advection of background flow

Discretizing the advection equation

Consider $J_{i,k}^k \leftarrow \begin{matrix} \text{time index} \\ J(x_i, t_k) \\ \text{space index} \end{matrix}$

$i \in [1, \dots, n]$

$k \in [1, \dots, m]$

The strategy: discretize each derivative separately depending on type of variable and desired accuracy of method

Simplest; discretize time and space using forward Euler

Time

$$\frac{\partial J}{\partial t} = \frac{J_i^{k+1} - J_i^k}{\Delta t}$$

Space

$$\frac{\partial J}{\partial x} = \frac{J_i^k - J_{i-1}^k}{\Delta x}$$

Putting these together, the discretized advection eqn:

$$\frac{J_i^{k+1} - J_i^k}{\Delta t} = -u \frac{J_i^k - J_{i-1}^k}{\Delta x} + S(x_i, t_k)$$

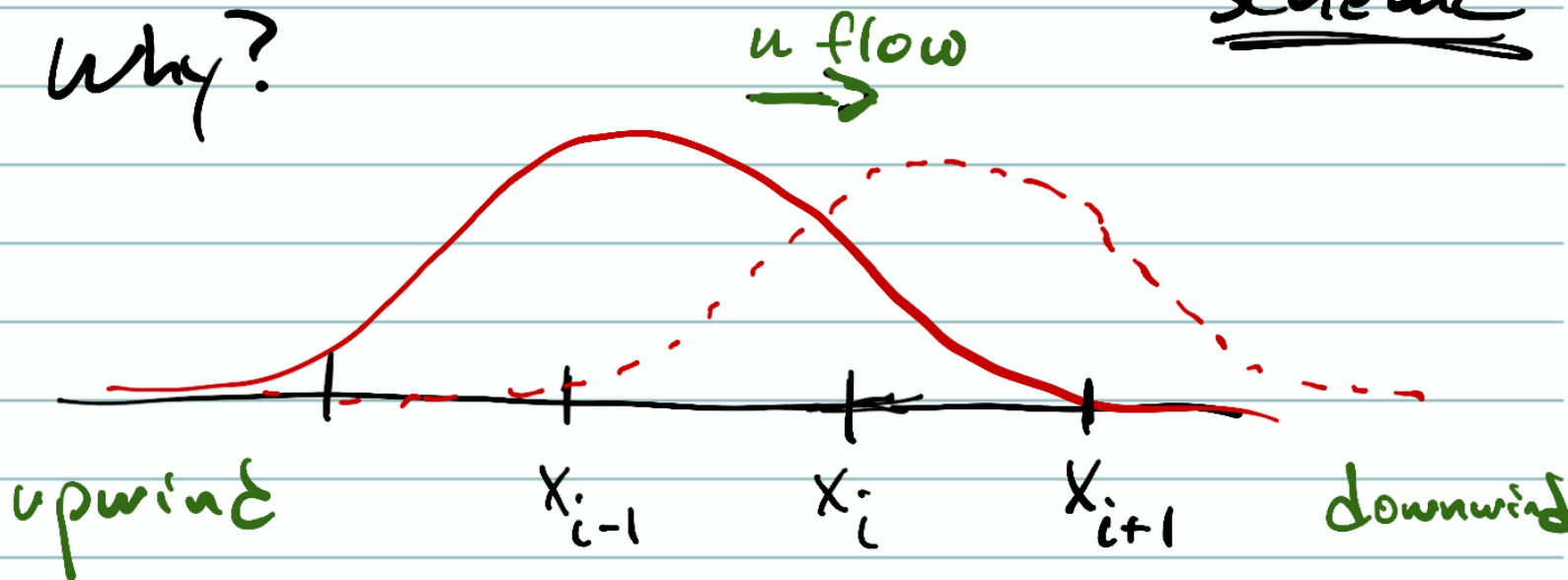
✓ In this scheme, this is the only unknown \rightarrow we can rewrite to solve for this

$$\bar{J}_i^{k+1} = \bar{J}_i^k - u \frac{\Delta t}{\Delta x} [\bar{J}_i^k - \bar{J}_{i-1}^k]$$

Marching Equation

$C = \frac{u \Delta t}{\Delta x} \rightarrow$ A dimensionless parameter known as the Courant number (must be small) to ensure stability

This method is known as the upwind scheme
Why?



When $u > 0$ and $\bar{J}_i - \bar{J}_{i-1} < 0$ then $\frac{\partial \bar{J}}{\partial t} > 0$

Or in other words, information is carried from upwind to downwind (the dir of u)

→ To ensure that information is always upwind → we must specify that if $u < 0$, then

$$\frac{\partial J}{\partial x} \approx \frac{\bar{J}_{i+1} - \bar{J}_i}{\Delta x}$$

Pros and cons of upwind scheme

$$\begin{aligned}
 J_i^{k+1} - J_i^k &= -u \frac{\Delta t}{\Delta x} (J_i^k - J_{i-1}^k) \\
 &\quad - \frac{u \Delta t}{2 \Delta x} (2J_i^k - 2J_{i-1}^k) \\
 &= \frac{u \Delta t}{2 \Delta x} \left(\underbrace{J_{i+1}^k}_{\text{Term 1}} - \underbrace{J_{i-1}^k}_{\text{Term 2}} \right) + \frac{u \Delta t}{2 \Delta x} \left(\underbrace{J_{i+1}^k}_{\text{Term 1}} - \underbrace{2J_i^k}_{\text{Term 2}} + \underbrace{J_{i+1}^k}_{\text{Term 2}} \right)
 \end{aligned}$$

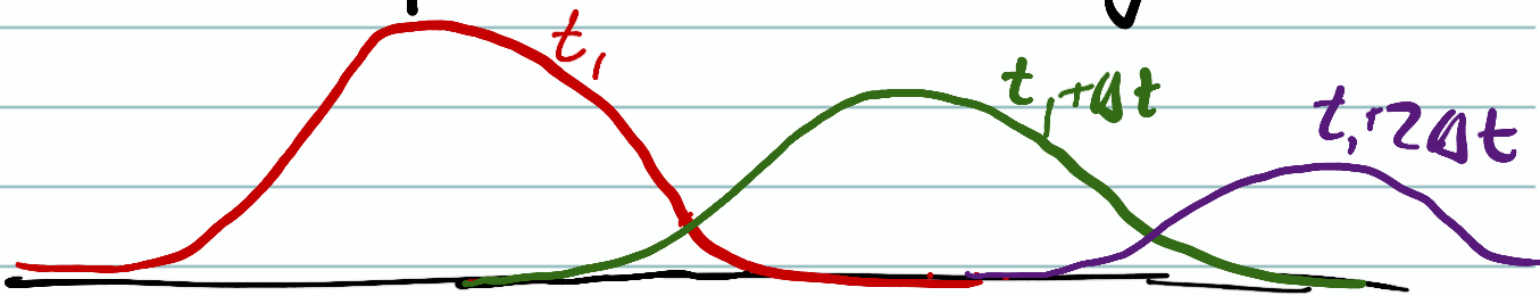
Term 1

Centered difference
approx for $\frac{\partial \bar{U}}{\partial x}$

Term

Diffusion
 $\frac{\partial^2 \bar{U}}{\partial x^2} \Delta x^2$

The prominence of this diffusion term in the upwind scheme can be a problem - leading to:



Dissipation of concentrations of \bar{U} \rightarrow time, x

Solution: be careful with picking $\Delta x, \Delta t$

$$C = u \frac{\Delta t}{\Delta x} \rightarrow$$

when $C > 1 \rightarrow$ method unstable
when $C < 1 \rightarrow$ method diffusive
when $C = 1 \rightarrow$ method stable
+ non-diffusive