

[Brief Review\*]

# Towards a Reconstruction of General Bulk Metrics[1]

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A new approach for the bulk reconstruction, the method using the light-cone cuts, is introduced. The past (future) light-cone cut of the bulk point  $p$  is defined as

$$C^-(p) = \partial J^-(p) \cap \partial M \quad (C^+(p) = \partial J^+(p) \cap \partial M), \quad (1)$$

where  $M$  is the spacetime assumed to be asymptotically  $\text{AdS}_n$  with some natural conditions for mathematical reasons. From now on, we focus on past cuts, but the same discussion can be applied to future cuts.

First, using the “bulk-point singularity” [2], we can get the past cuts of the points in  $J^+(\partial M) \cap J^-(\partial M)$ . The divergence of a boundary to boundary correlator, whose fields are massless, is caused by that the interaction vertex is a bulk point which is null-related to all points in the correlator, under certain conditions. The divergence of this class is called bulk-point singularity. Let us consider  $n$  boundary points  $x_1, \dots, x_n$  which are spacelike separated each other, and two boundary points  $z_1$  and  $z_2$  in the past of  $x_i$ . Then moving  $z_1$  and  $z_2$  while keeping

$$\langle \mathcal{O}(z_1) \mathcal{O}(z_2) \mathcal{O}(x_1) \cdots \mathcal{O}(x_n) \rangle \rightarrow \infty \quad (2)$$

for fixed  $x_i$  traces out the past cut of the unique bulk point which is null-separated from  $x_i$ . Since we do not know about the bulk geometry, we cannot know whose past cut it is. In principle, however, the past cut can be labeled by some  $n$  parameters related to  $x_i$ . Let us set the parameters  $\lambda = (\lambda^1, \dots, \lambda^n)$  and denote the cut by  $C_\lambda^-$ . By repeating this step, all past cuts related one-to-one to points in  $J^+(\partial M) \cap J^-(\partial M)$  can be collected, then we define the vector space of  $\lambda$ 's,  $\mathcal{M}^- := \{\lambda\}$ .

Let  $\lambda_1, \lambda_2, \dots$  be vectors on  $\mathcal{M}^-$ , meaning that  $C_{\lambda_1}^-, C_{\lambda_2}^- \cdots$  are the past cuts. In the appendix, it is shown that, if  $C_{\lambda_1}^- = C^-(p)$  and  $C_{\lambda_2}^- = C^-(q)$  are tangent each other precisely at one point, then  $p$  and  $q$  are null related. Therefore, taking a boundary point  $x \in C_{\lambda_1}^-$  and collecting the past cuts tangent to  $C_{\lambda_1}^-$  at  $x$ , the unique null geodesic from  $p$  to  $x$  via  $q$  is obtained. Then we can identify this null geodesic to the one connecting  $\lambda_1$  and other  $\lambda$ 's corresponding to other cuts tangent to  $C_{\lambda_1}^-$  at  $x$ , because of the existence of the one-to-one map from  $J^+(\partial M) \cap J^-(\partial M)$  to  $\mathcal{M}^-$ . Repeating this process for points on  $C_{\lambda_1}^-$ , we get null geodesics through  $\lambda_1$ , and they form light-cone of  $\lambda_1$  in  $\mathcal{M}^-$ .

Once the causal structure is introduced in  $\mathcal{M}^-$ , the conformal metric (the metric up to conformal factors) can be determined. The null generators at  $\lambda \in \mathcal{M}^-$  are obtained by looking for other past cuts tangent to  $C_\lambda^-$ . Let  $\sigma = (\sigma^1, \dots, \sigma^{n-2})$  be the parameter describing the points on  $C_\lambda^-$  as  $C_\lambda^-(\sigma) \in \partial M$ . The condition which  $C_{\lambda_2}^-$  is tangent to  $C_{\lambda_1}^-$  is as follows:

$$\exists \sigma_1, \exists \sigma_2, \quad C_{\lambda_1}^-(\sigma_1) = C_{\lambda_2}^-(\sigma_2) \quad \text{and} \quad \nabla_\sigma C_{\lambda_1}^-(\sigma)|_{\sigma=\sigma_1} = \nabla_\sigma C_{\lambda_2}^-(\sigma)|_{\sigma=\sigma_2}. \quad (3)$$

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A null generator  $N$  at  $\lambda$  satisfies the above condition for  $\lambda_1 = \lambda$  and  $\lambda_2 = \lambda + \varepsilon N$ , where  $\varepsilon$  is any infinitesimal real variable. By solving the conditions, we can find out  $n(n+1)/2$  generators at  $\lambda$ . Imposing the condition that the  $n(n+1)/2$  null generators are null at  $\lambda$  determines the metric at  $\lambda$  up to conformal factors. Since  $\mathcal{M}^-$  can be regarded as the copy of  $J^+(\partial M) \cap J^-(\partial M)$ , the reconstruction of the bulk metric has been completed now.

## References

- [1] N. Engelhardt and G.T. Horowitz, *Towards a Reconstruction of General Bulk Metrics*, *Class. Quant. Grav.* **34** (2017) 015004 [1605.01070].
- [2] J. Maldacena, D. Simmons-Duffin and A. Zhiboedov, *Looking for a bulk point*, *JHEP* **01** (2017) 013 [1509.03612].