

Time-stepping methods (continued) and numerical derivatives

(Pencil Code School, Geneva/CERN, 21st of September 2025)

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Overview:

- * Summary on common time-stepping methods
- * “Creative” methods for time stepping
- * Numerical derivatives
- * Numerical curiosities due to precision
- * Numerical curiosities from the initial condition
- * Numerical curiosities from “switching on” the simulation

Time integration methods

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Euler method:

- * very simple to implement (single-step integration)
- * may have infinitely growing error

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Runge-Kutta scheme:

- * harder to implement (order of accuracy = number of integration substeps)
- * very accurate and computation efficient (larger dt is possible)

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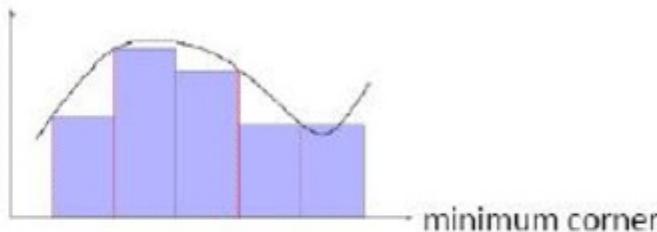
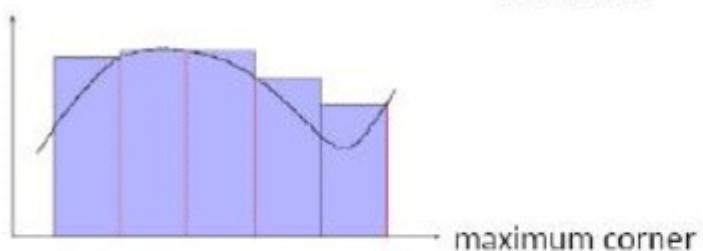
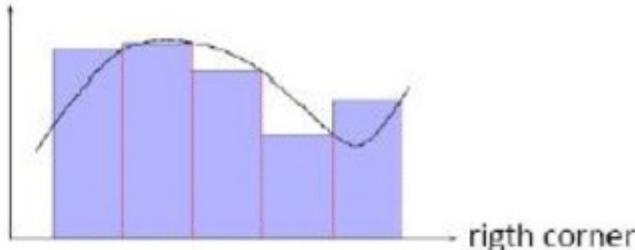
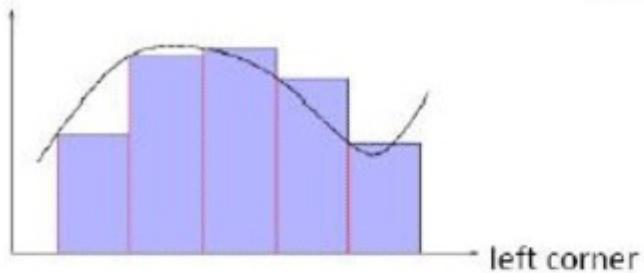
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Runge-Kutta scheme:

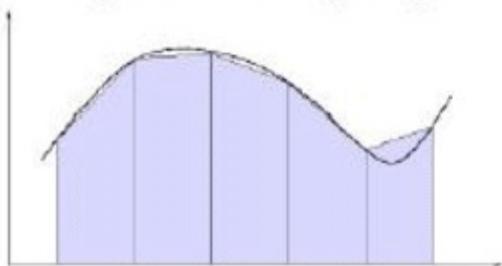
- * harder to implement (order of accuracy = number of integration substeps)
- * very accurate and computation efficient (larger dt is possible)
- * substeps require additional function evaluations and derivatives (\Rightarrow costs)

Time integration methods

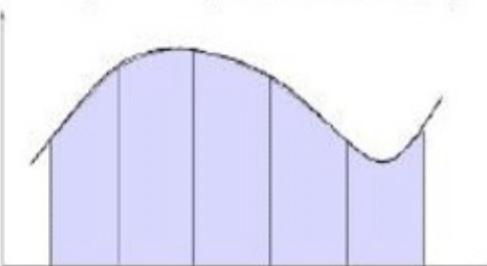
Rectangle
Step width * height



Trapezoid
Step width * average height

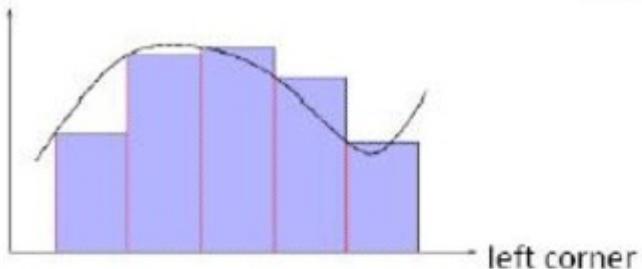


Simpson's
Step width * Simpson's height



Time integration methods

Rectangle
Step width * height



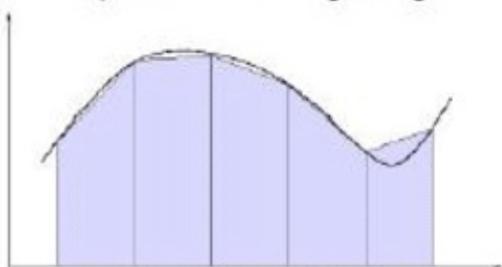
**Trapezoidal
Rule**



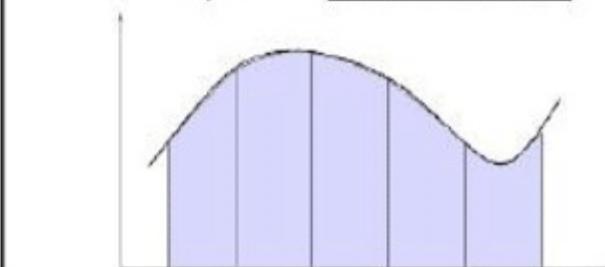
**Midpoint
Rule**



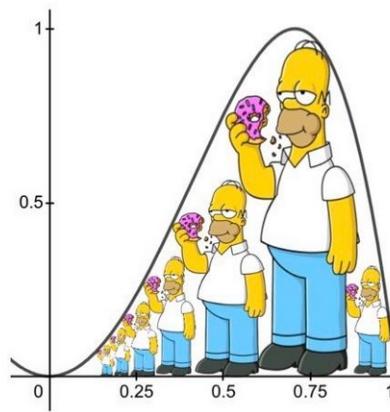
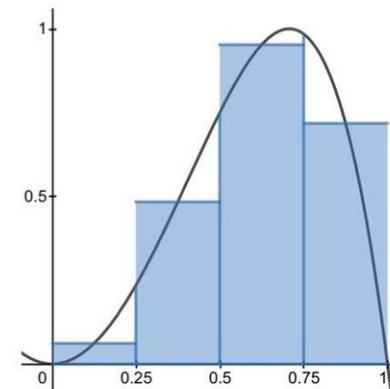
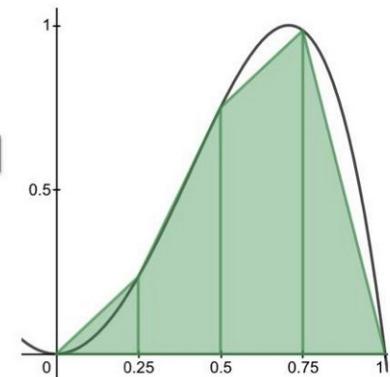
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**Simpson's
Rule**



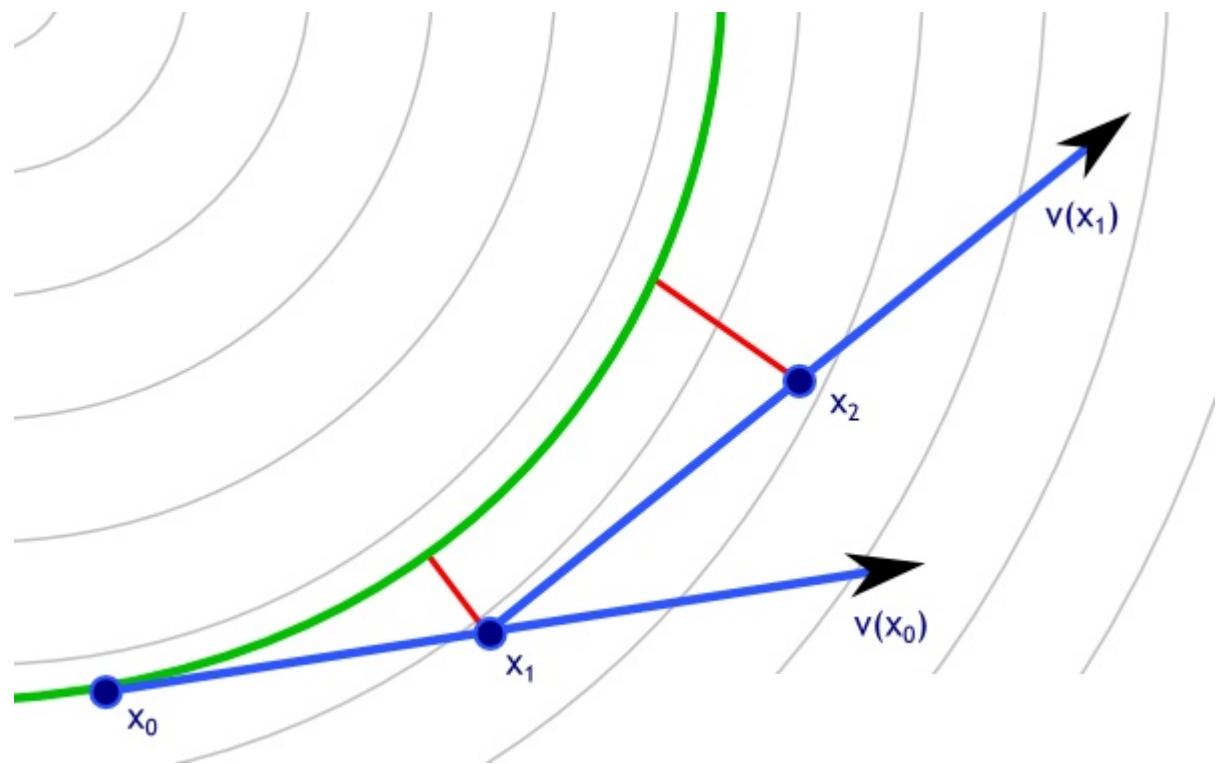
Euler method

Euler method

- 1) start from x_0
- 2) calculate derivatives
(tangential)

gray: true field

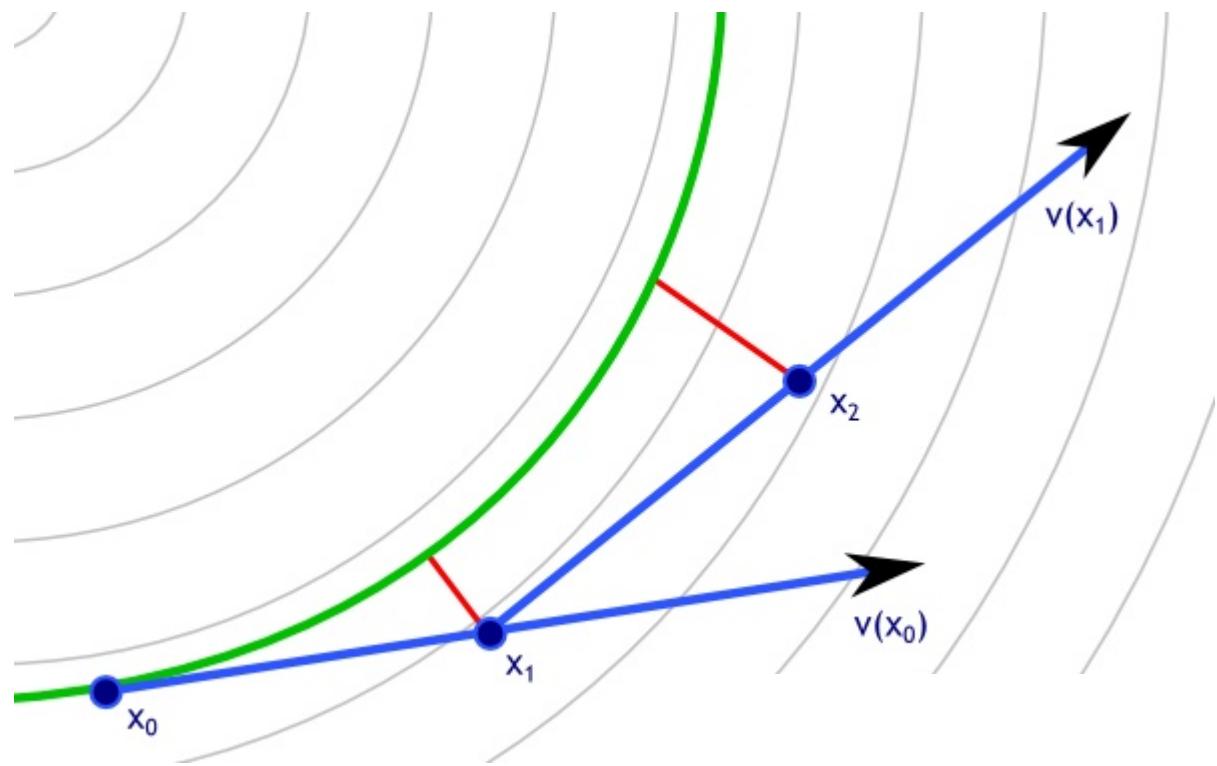
green: exact solution



Euler method

- 1) start from x_0
- 2) calculate derivatives
(tangential)
- 3) iterate to x_1
with $dt=0.5$

green: exact solution
blue: tangential derivative
red: deviation



$$x_{n+1} = x_n + dt \cdot v(x_n) + O(dt^2)$$

=> error grows infinitely...
=> make dt smaller?

x_n : current position

x_{n+1} : next position

dt : time step

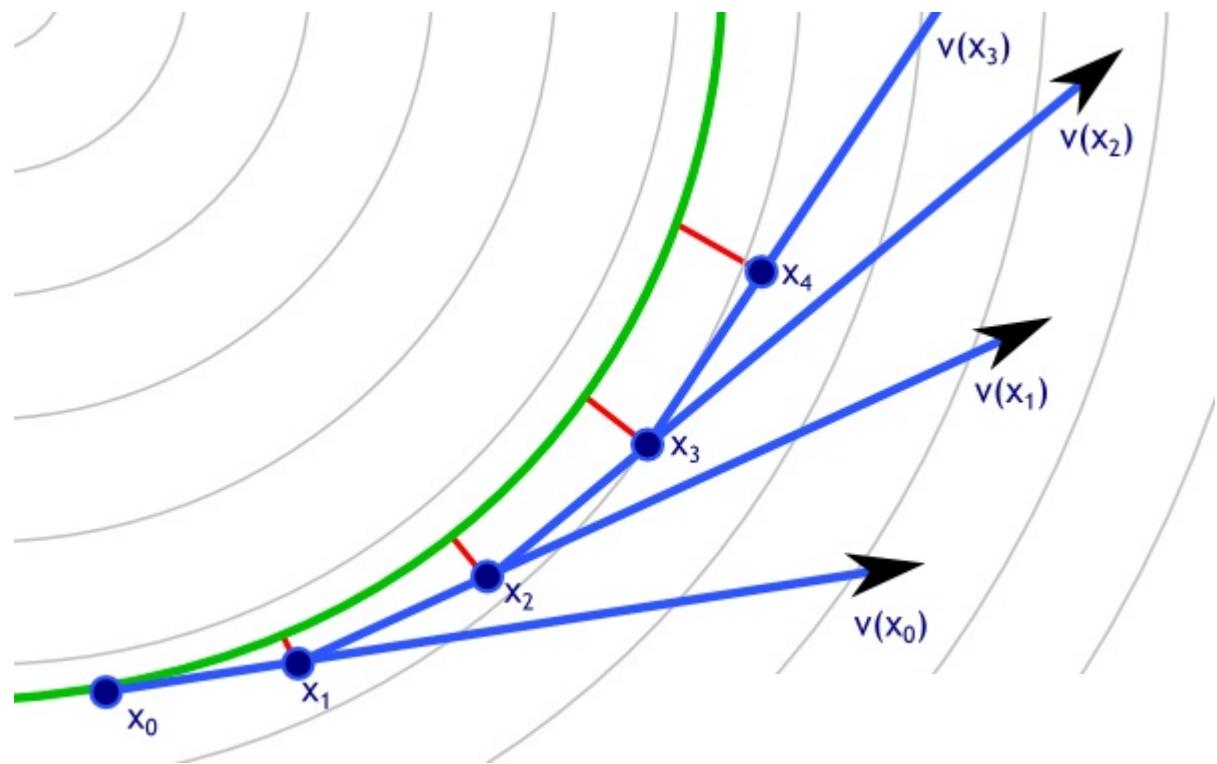
$v(x_n)$: vector field at position x_n

$O(dt^2)$: error

Euler method

- 1) start from x_0
- 2) calculate derivatives
(tangential)
- 3) iterate to x_1
with $dt = \mathbf{0.25}$

green: exact solution
blue: tangential derivative
red: deviation



$$x_{n+1} = x_n + dt \cdot v(x_n) + O(dt^2)$$

=> accuracy can be improved
by making dt smaller
=> but error still grows...

x_n : current position

x_{n+1} : next position

dt : time step

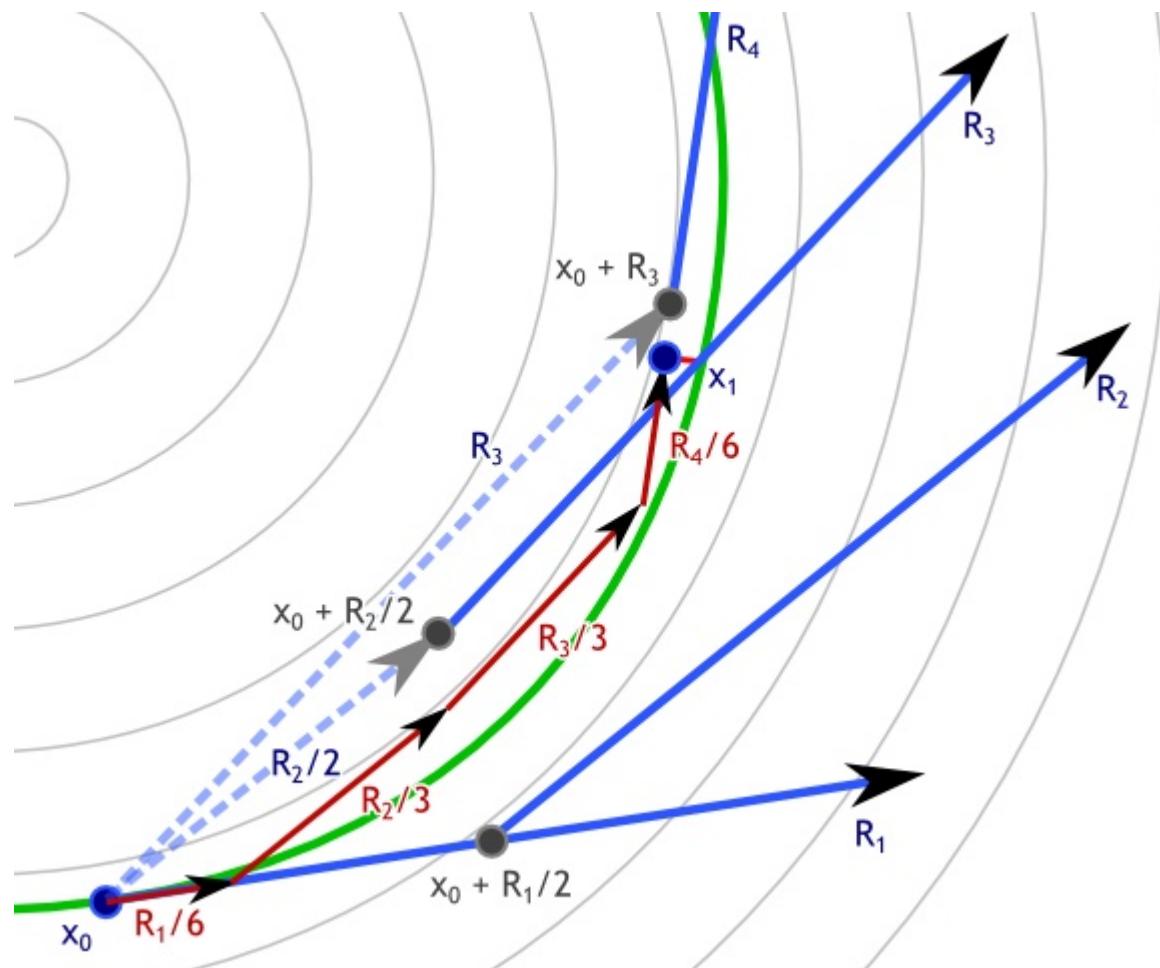
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Runge-Kutta method (4th order)

- 1) start from x_0
- 2) calculate derivatives
(tangential)
- 3) iterate substeps
and recalculate
derivatives (tangential)
- 4) reach x_1 with substeps

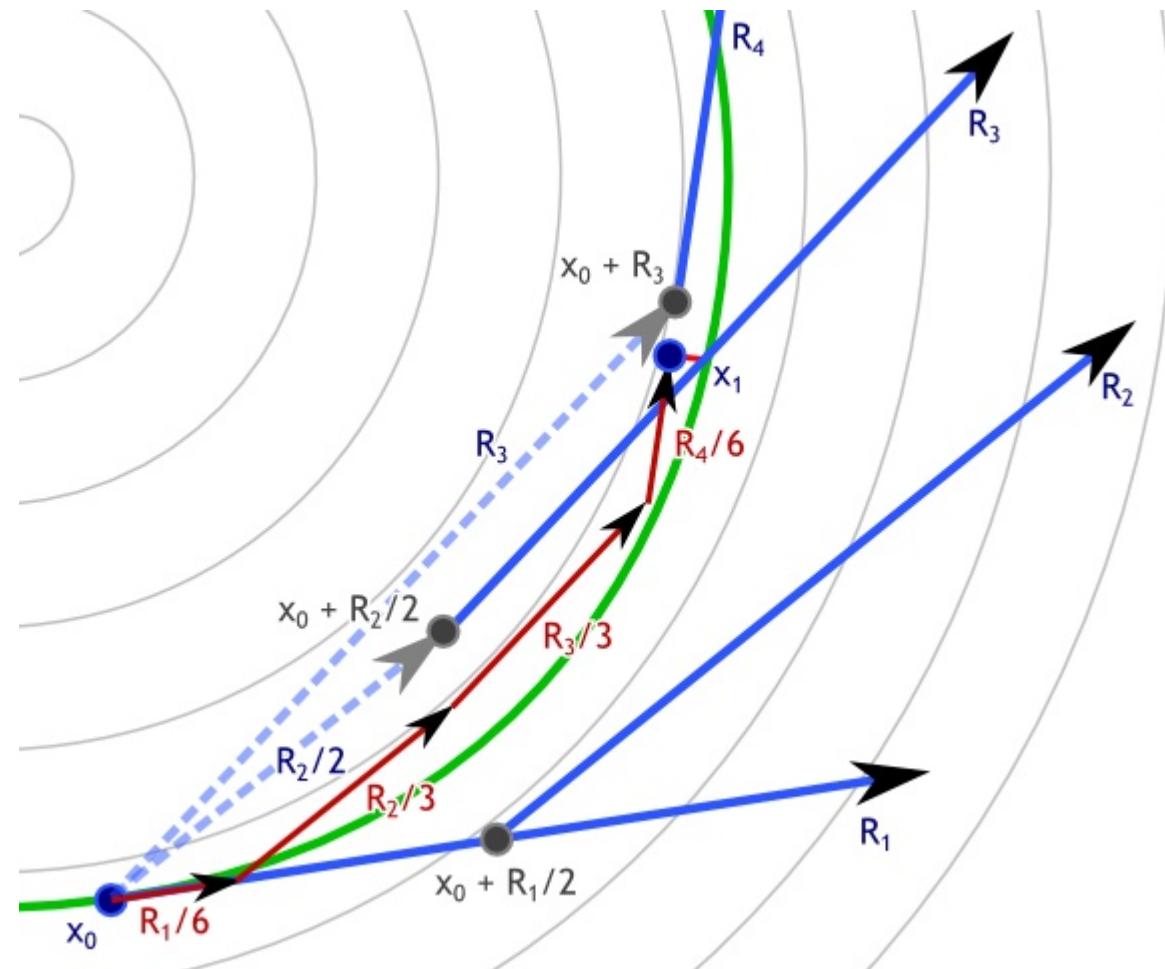
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blue: tangential derivative
red: deviation



=> much better accuracy!

=> larger computational cost

Runge-Kutta method (4th order)

RK4 scheme:

$$k_1 = dt \cdot v(x_n)$$

$$k_2 = dt \cdot v(x_n + \frac{k_1}{2})$$

$$k_3 = dt \cdot v(x_n + \frac{k_2}{2})$$

$$k_4 = dt \cdot v(x_n + k_3)$$

$$x_{n+1} = x_n + \frac{k_1}{6} + \frac{k_2}{3} + \frac{k_3}{3} + \frac{k_4}{6} + O(dt^5)$$

x_n : current position

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dt : time step

$v(x)$: vector field at position x

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=> effect: back-reaction to system is delayed => slower overall dynamics!
more timesteps needed & finally not such a big saving of computing time.

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=> 3D: either huge diffusion or it will need the same amount of time steps.

=> Speed up and accuracy granted only in 1D.

“Creative” time-stepping methods

→ Split Runge-Kutta time step into two halves (**timestep_strang**).

Idea: ...?

=> effects?

“Creative” time-stepping methods

→ Split Runge-Kutta time step into two halves (**timestep_strang**).

Idea: ...?

=> effects?

→ Algorithms for stiff PDEs (**timestep_stiff** from Numerical Recipies).

Idea: ...?

=> effects?

Numerical derivatives (6th order)

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Pencil Code manual:

H Numerical methods

H.1 Sixth-order spatial derivatives

Spectral methods are commonly used in almost all studies of ordinary (usually incompressible) turbulence. The use of this method is justified mainly by the high numerical accuracy of spectral schemes. Alternatively, one may use high order finite differences that are faster to compute and that can possess almost spectral accuracy. Nordlund & Stein [32] and Brandenburg et al. [16] use high order finite difference methods, for example fourth and sixth order compact schemes [28]¹⁹

The sixth order first and second derivative schemes are given by

$$f'_i = (-f_{i-3} + 9f_{i-2} - 45f_{i-1} + 45f_{i+1} - 9f_{i+2} + f_{i+3})/(60\delta x), \quad (235)$$

$$f''_i = (2f_{i-3} - 27f_{i-2} + 270f_{i-1} - 490f_i + 270f_{i+1} - 27f_{i+2} + 2f_{i+3})/(180\delta x^2), \quad (236)$$

Numerical derivatives (6th order)

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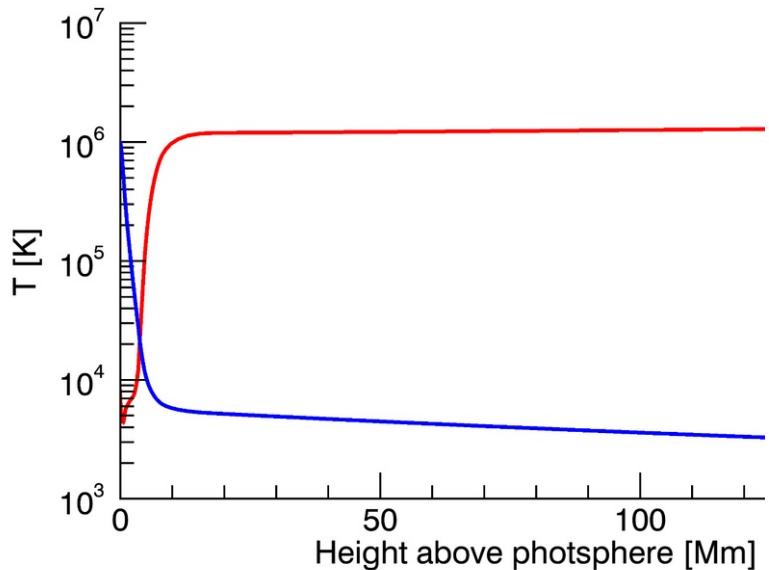
=> Coefficients can be obtained analytically.

Resolving steep gradients in the solar corona

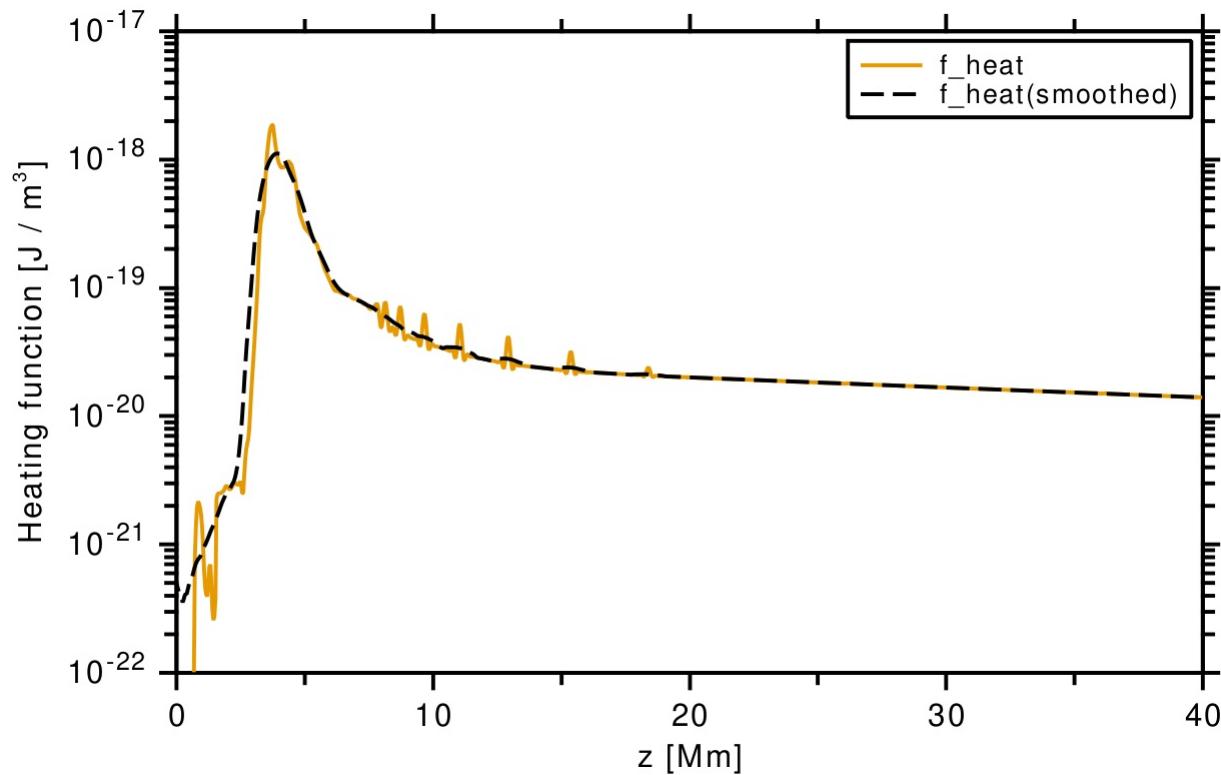
(Vartika Pandey)

Coronal heating in 1D MHD models:

- Initial condition:

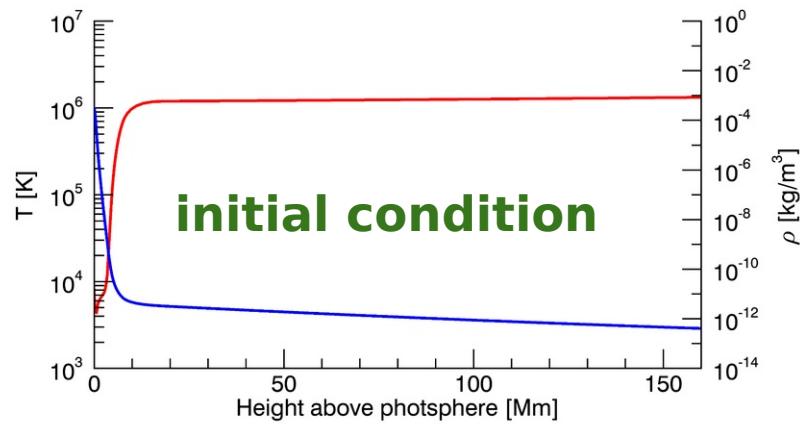


=> After one iteration:

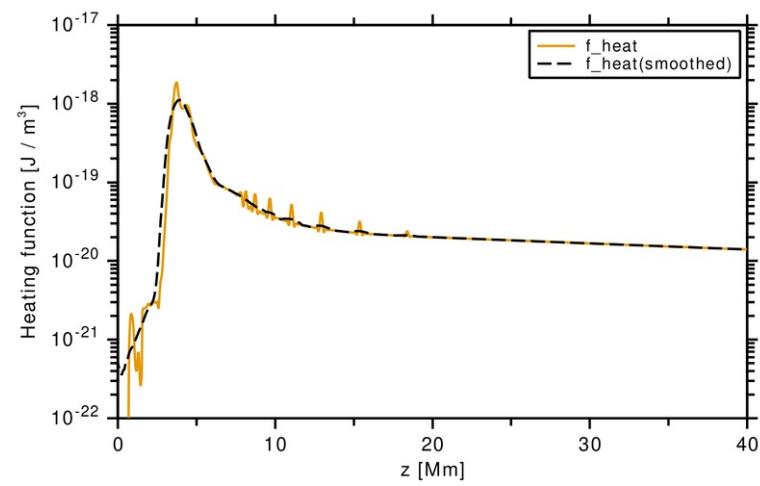


(Pandey & Bourdin, 2025)

Numerical stability analysis:

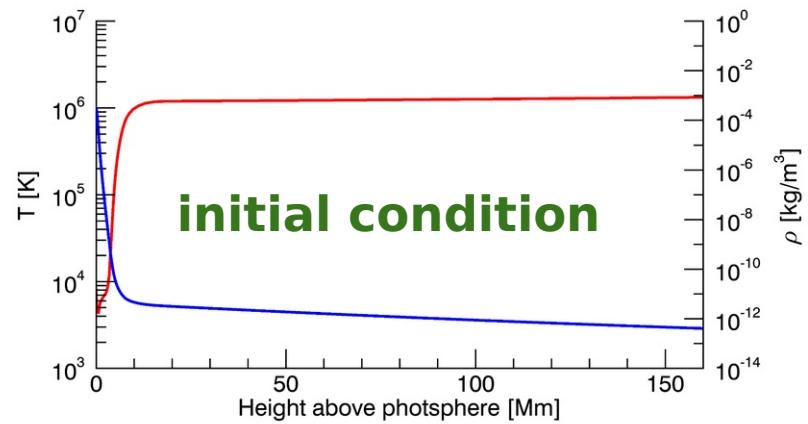


+ heating

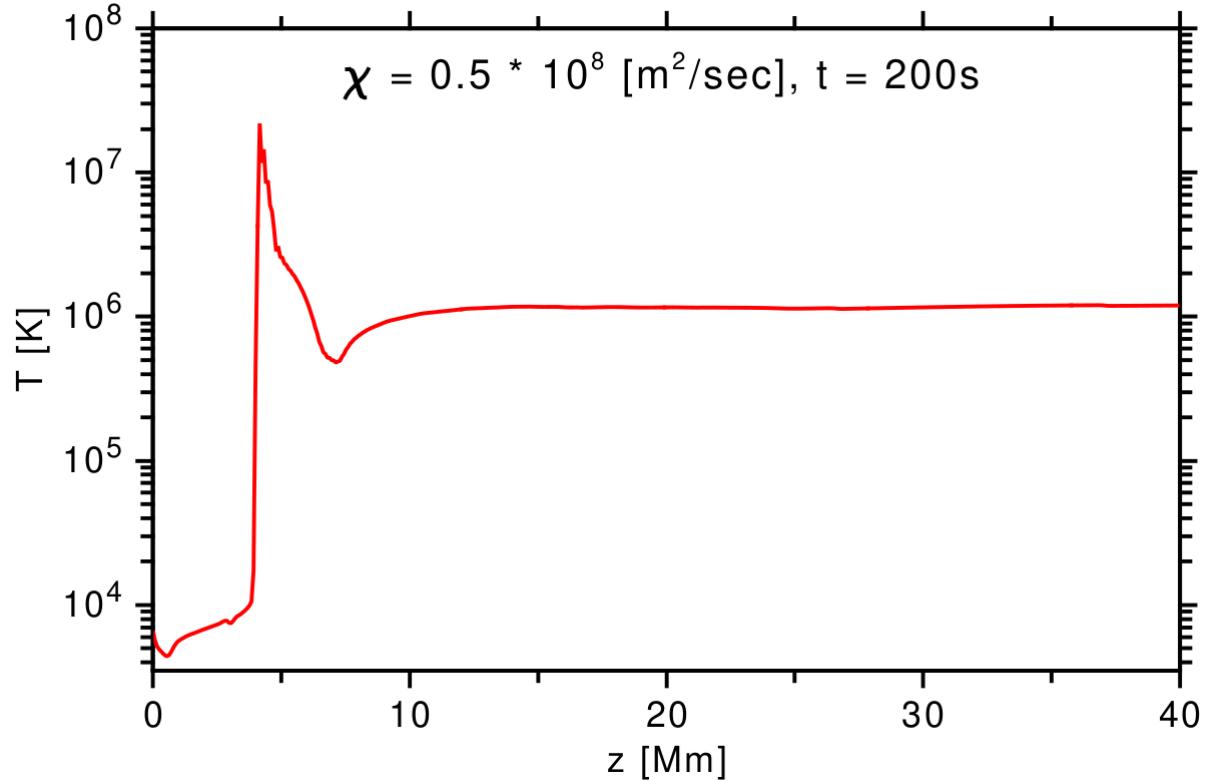
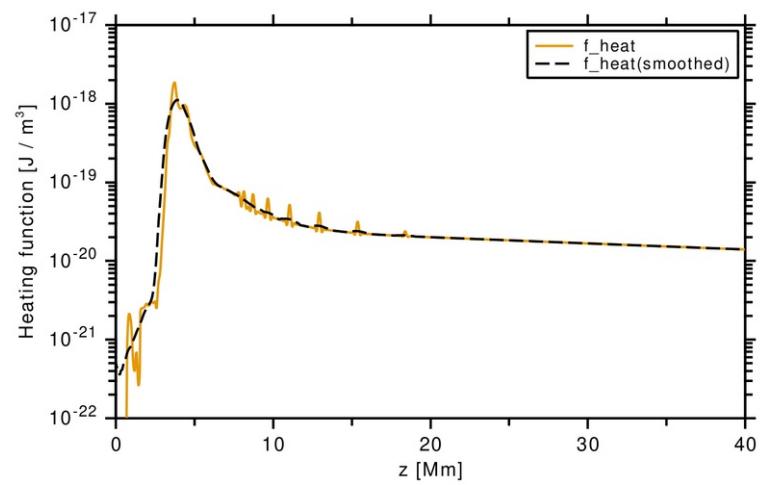


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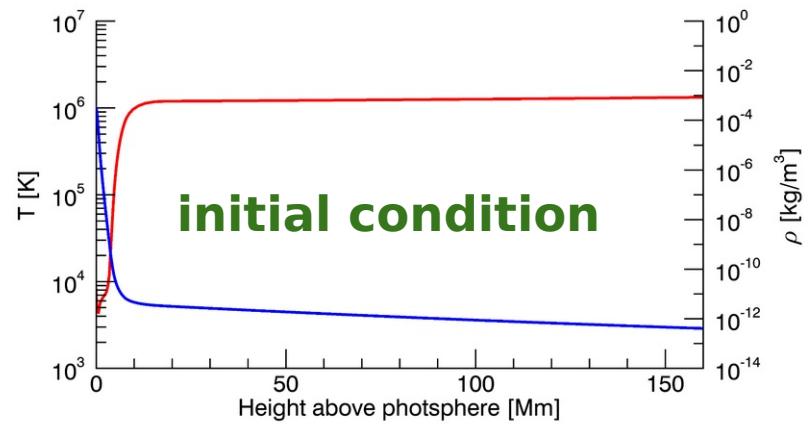


+ heating \Rightarrow

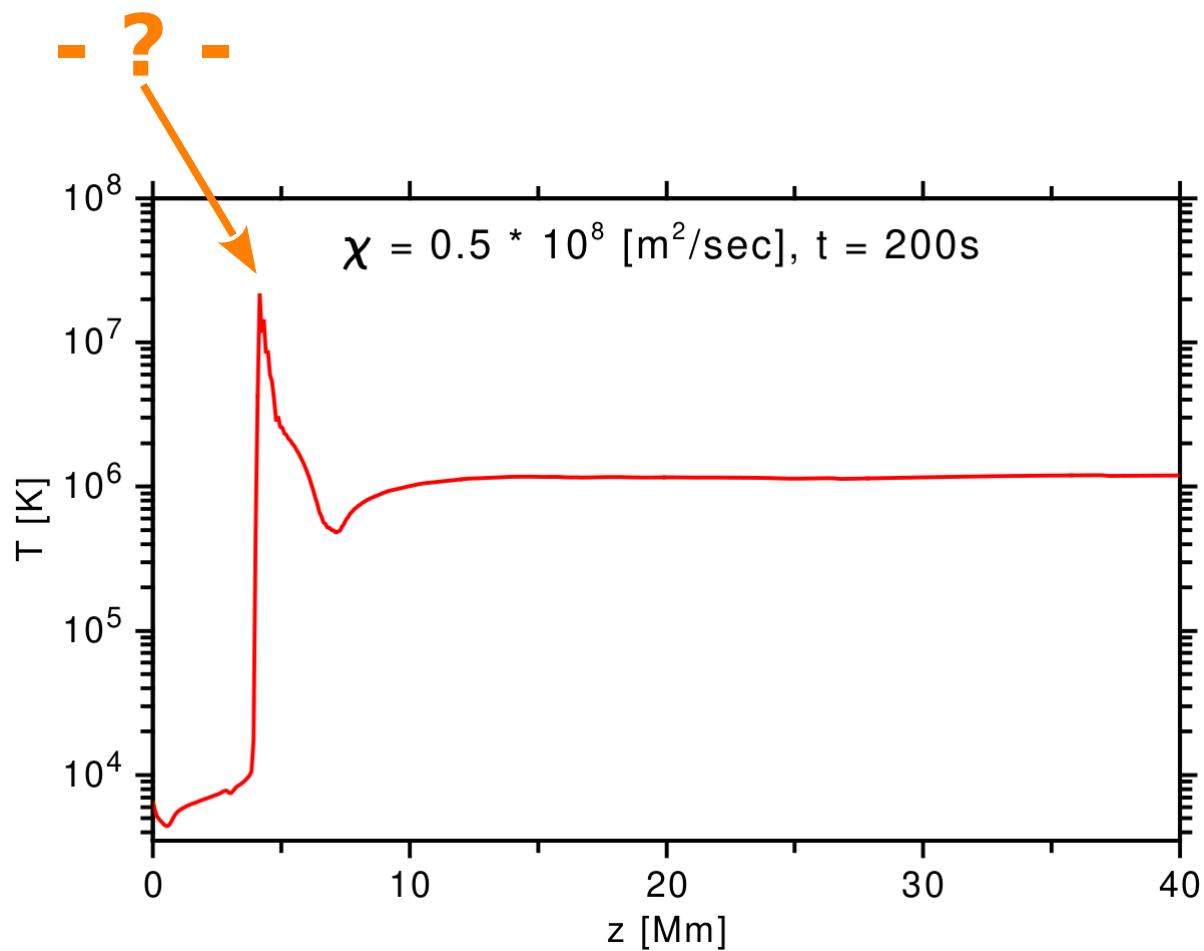
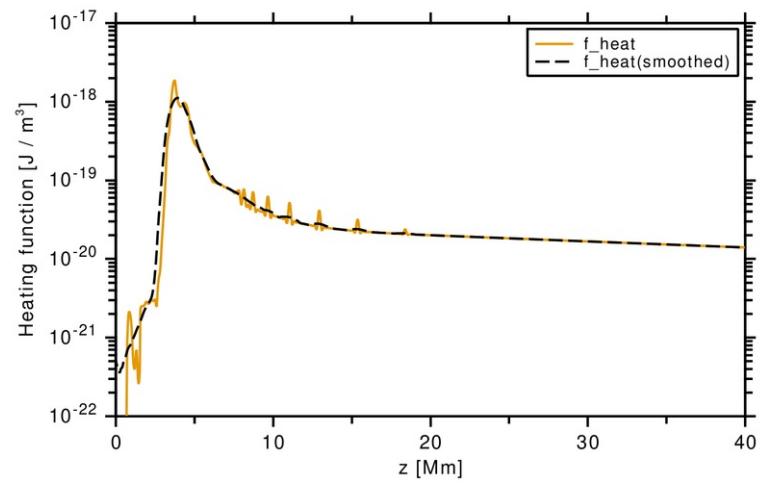


(Pandey & Bourdin, 2025)

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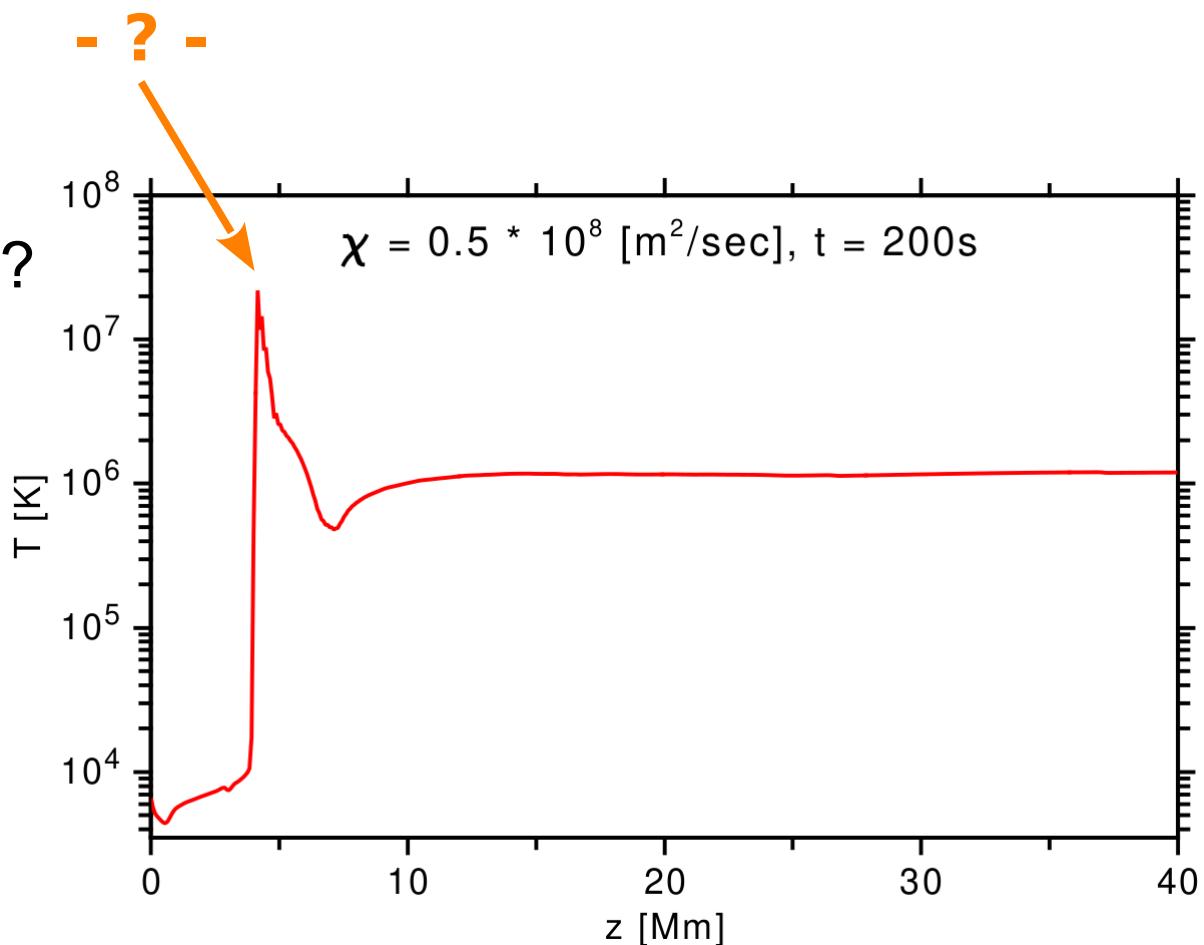


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Numerical stability analysis:

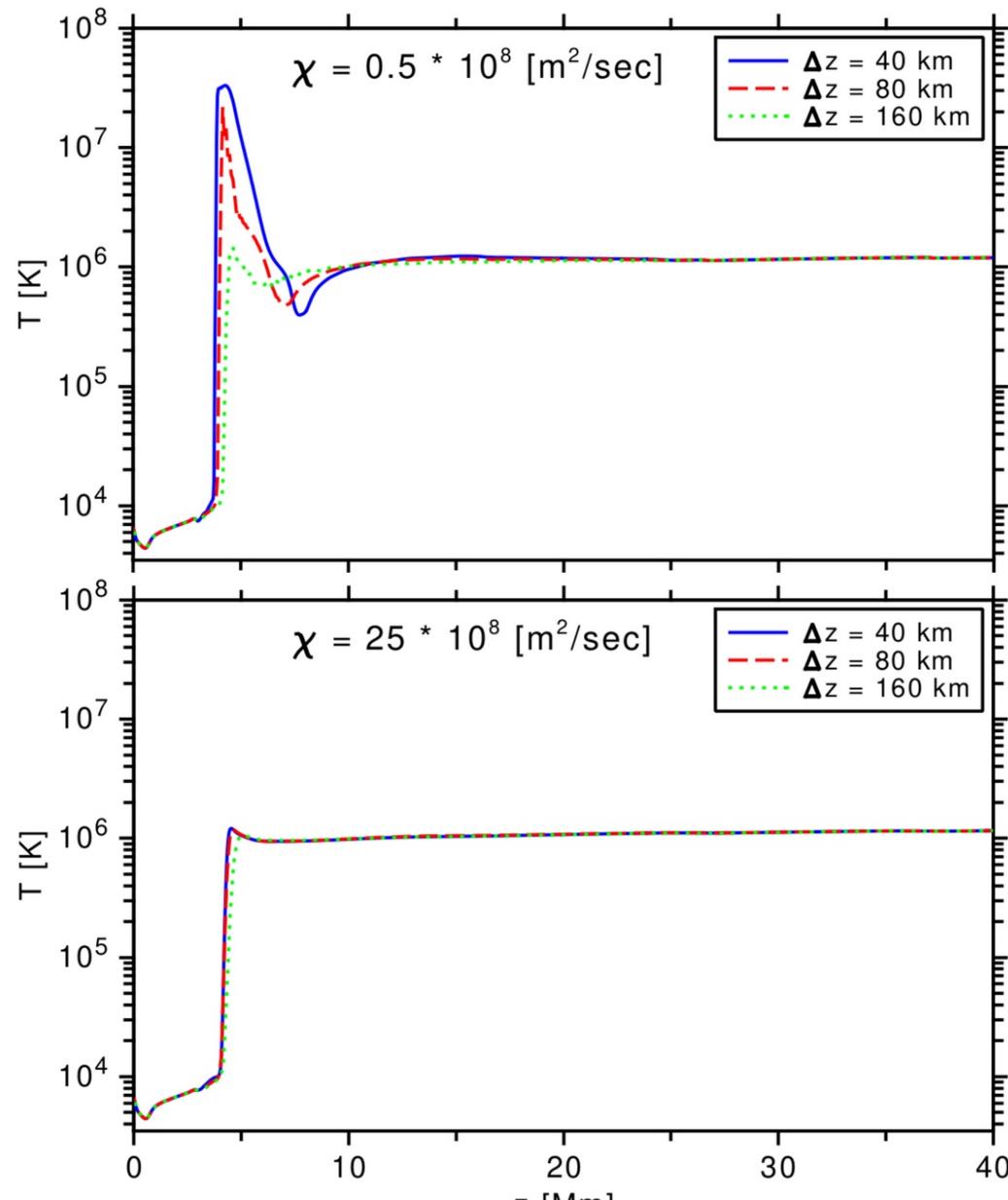
Possible problems:

- 1) too coarse grid?
- 2) insufficient diffusion?
- 3) absence of magnetic field?



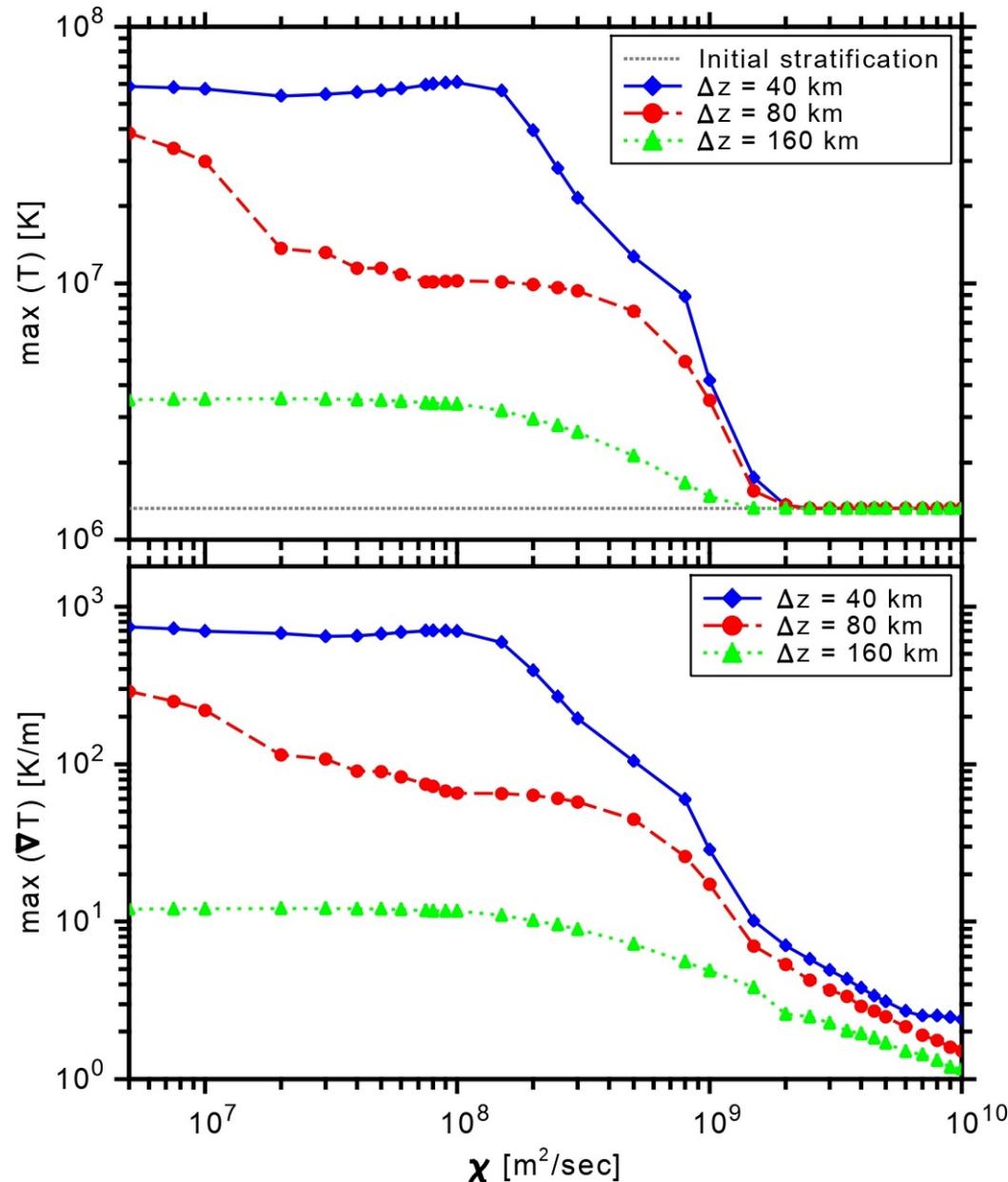
(Pandey & Bourdin, 2025)

Effects due to the grid resolution and heat conduction:



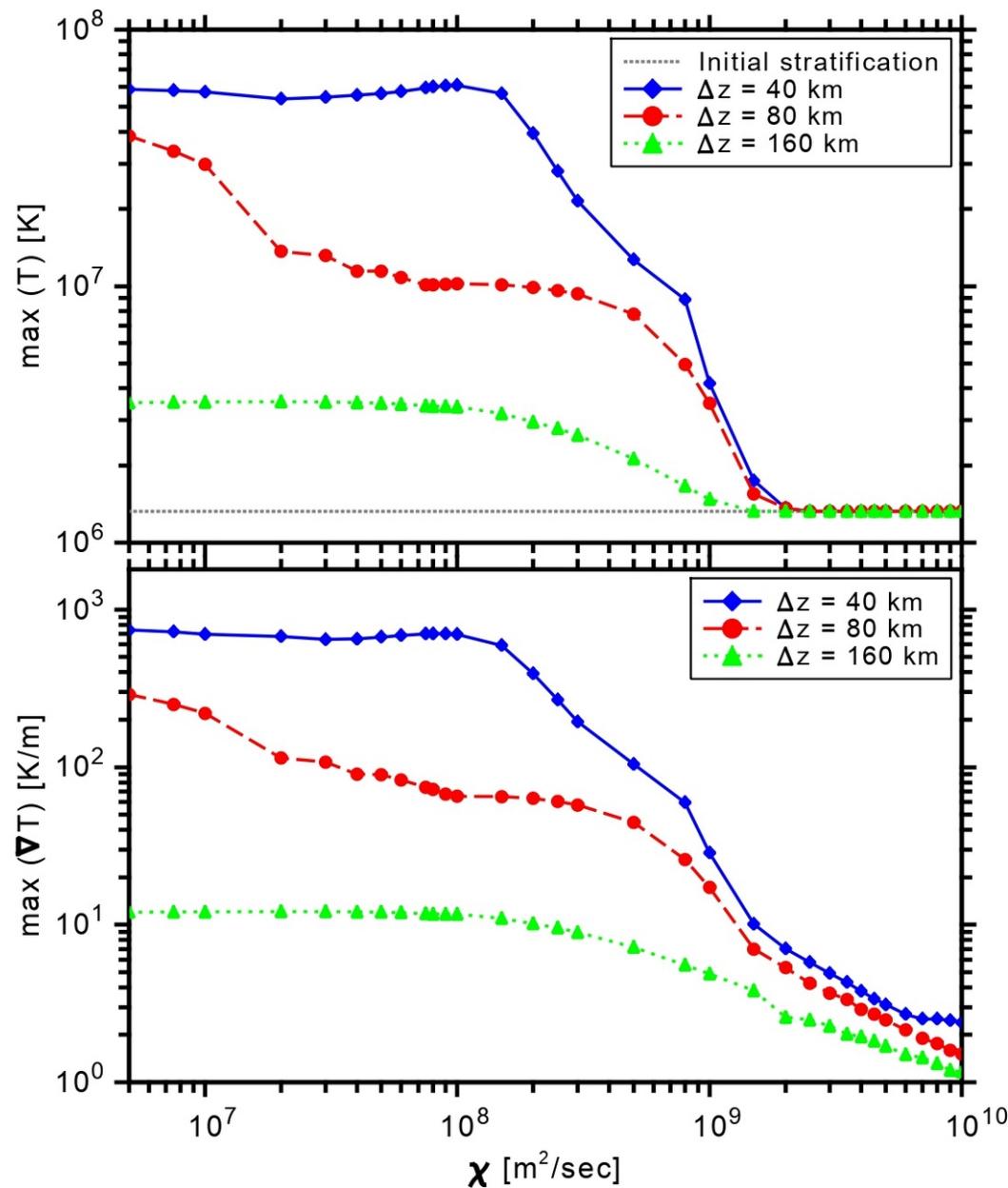
(Pandey & Bourdin, 2025)

Scaling the heat conduction:



(Pandey & Bourdin, 2025)

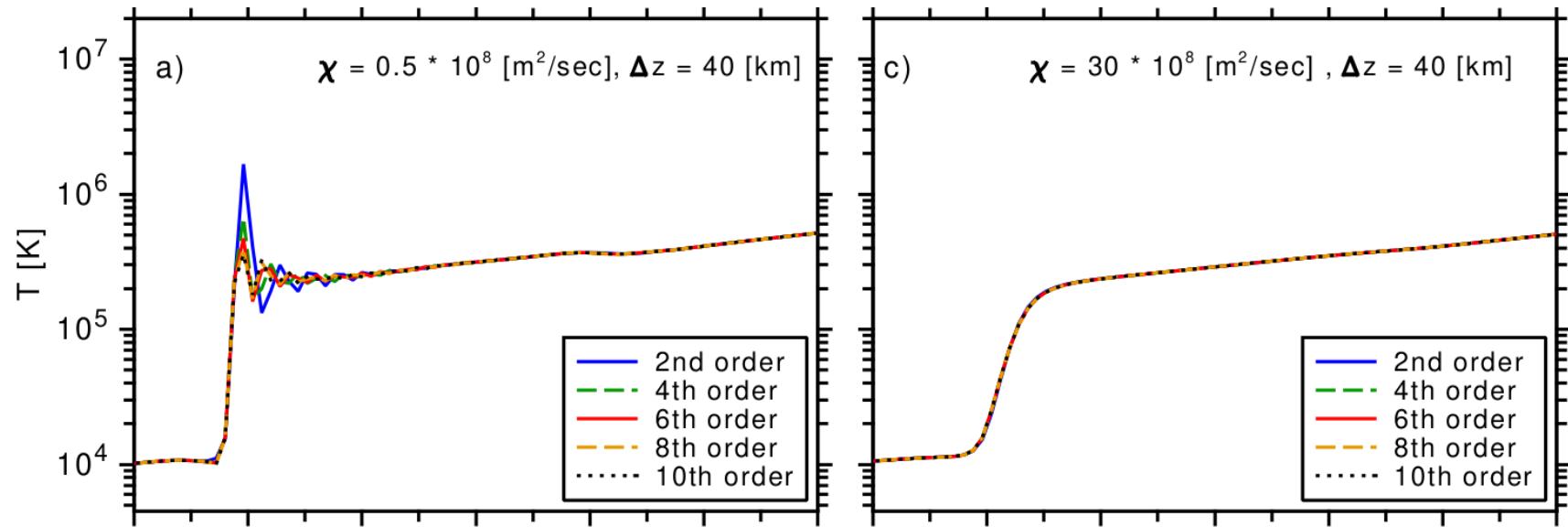
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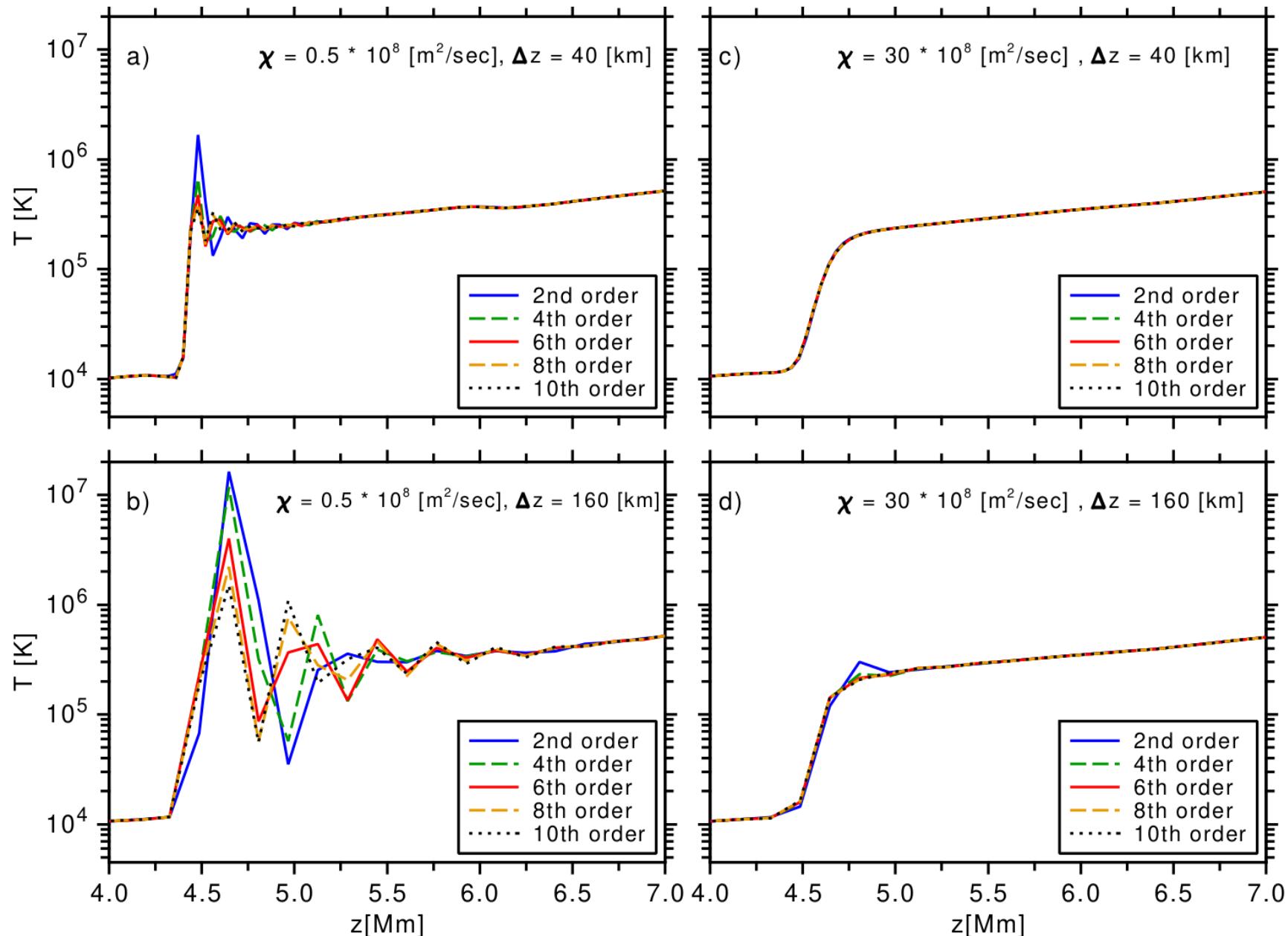
=> higher resolutions
do not always imply
models are “better”

=> diffusion constants
need to fit to the
grid resolution

Effects due to the order of numerical derivatives:



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Numerical curiosities (due to precision & truncation errors)

Precision and machine epsilon

Numerical curiosities cabinet:

$$(x + y) + z \neq x + (y + z)$$

Precision and machine epsilon

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Machine epsilon:

$$1.0 + \epsilon = 1.0$$

Precision and machine epsilon

Numerical curiosities cabinet:

$$(x+y)+z \neq x+(y+z)$$

Machine epsilon:

$$1.0 + \epsilon = 1.0 \quad (*)$$

- Estimation of the granularity:

$$(b = 2; p_o = 0.5)$$

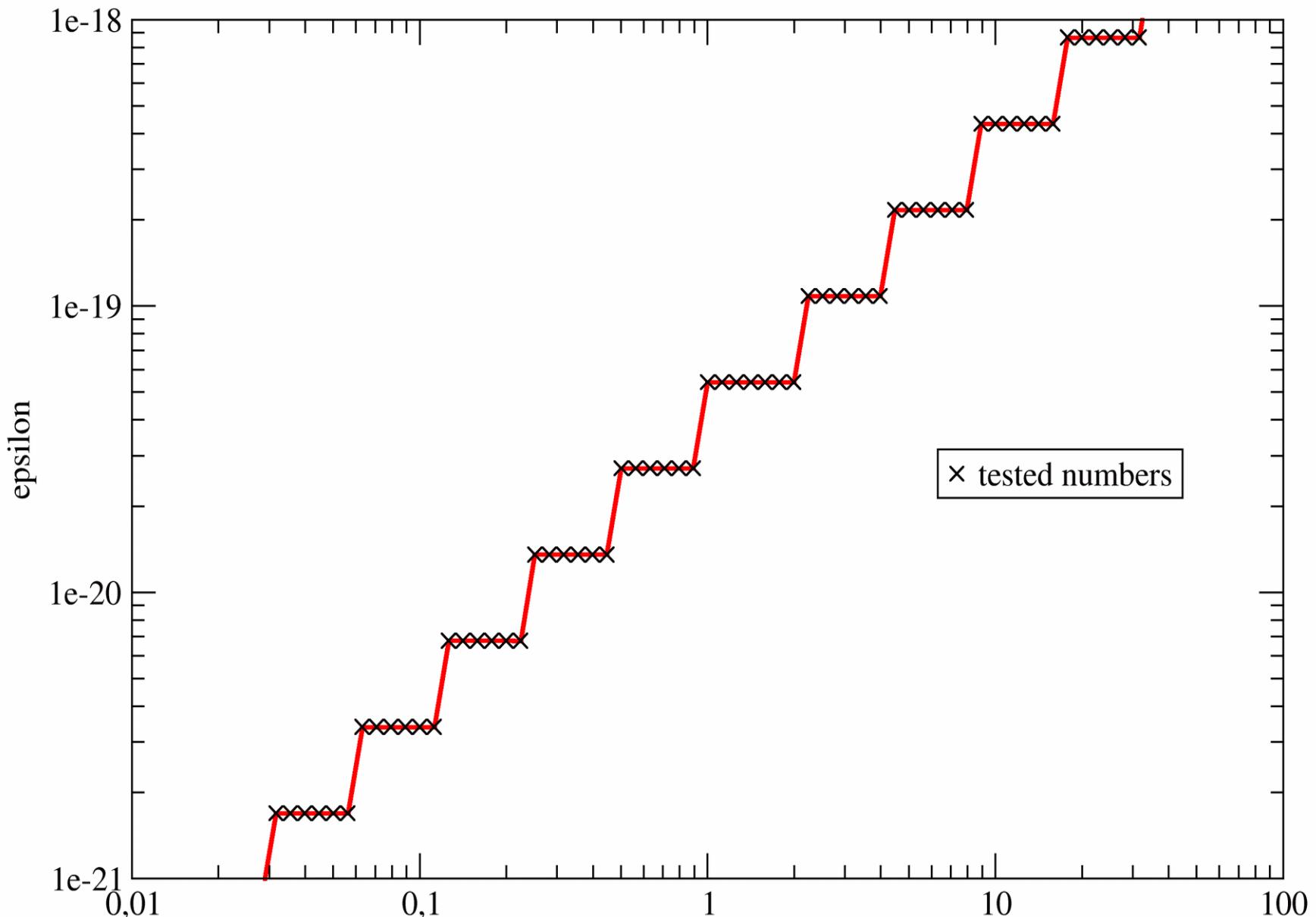
$$\epsilon_N = \frac{b}{2} \cdot b^{-p_N}$$

$$p_{N+1} = p_N \cdot b$$

Stop, if (*) is true

Precision and machine epsilon

Machine epsilon in Fortran 90 (double precision).



Granularity of numbers

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Numerical curiosities cabinet: (continued...)

Example: simple increment of simulation time:

$$T_N + dT = T_{N+1}$$

Granularity of numbers

Numerical curiosities cabinet: (continued...)

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$$4096 + 4 \cdot 10^{-4} = 4096$$

Granularity of numbers

Numerical curiosities cabinet: (continued...)

Example: simple increment of simulation time:

$$T_N + dT = T_{N+1}$$

$$4096 + 4 \cdot 10^{-4} = 4096$$

$$\Rightarrow \frac{T_N}{dT} \approx 10^{-7}$$

Granularity of numbers

Numerical curiosities cabinet: (continued...)

Example: simple increment of simulation time:

$$T_N + dT = T_{N+1}$$

$$4096 + 4 \cdot 10^{-4} = 4096$$

$$\Rightarrow \frac{T_N}{dT} \approx 10^{-7}$$

But:

$$4096 + 8 \cdot 10^{-4} \neq 4096$$

Granularity effects in physical observables

Granularity effects in physical observables

Density in a coronal 3D MHD simulation:

131

 absolute scaling subtract averages

RESET

55

 show crosshairs show horizontal averages

LOAD

103

data set LOGARITHMIC_DENSITY

SAVE

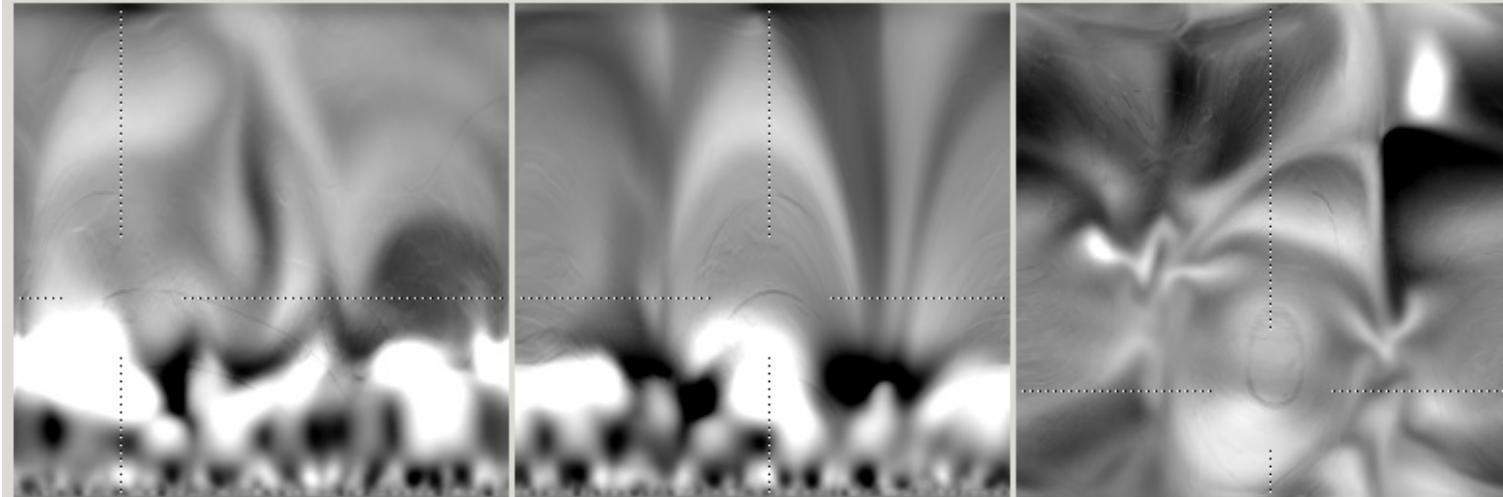
time step VRPC

PLAY

QUIT

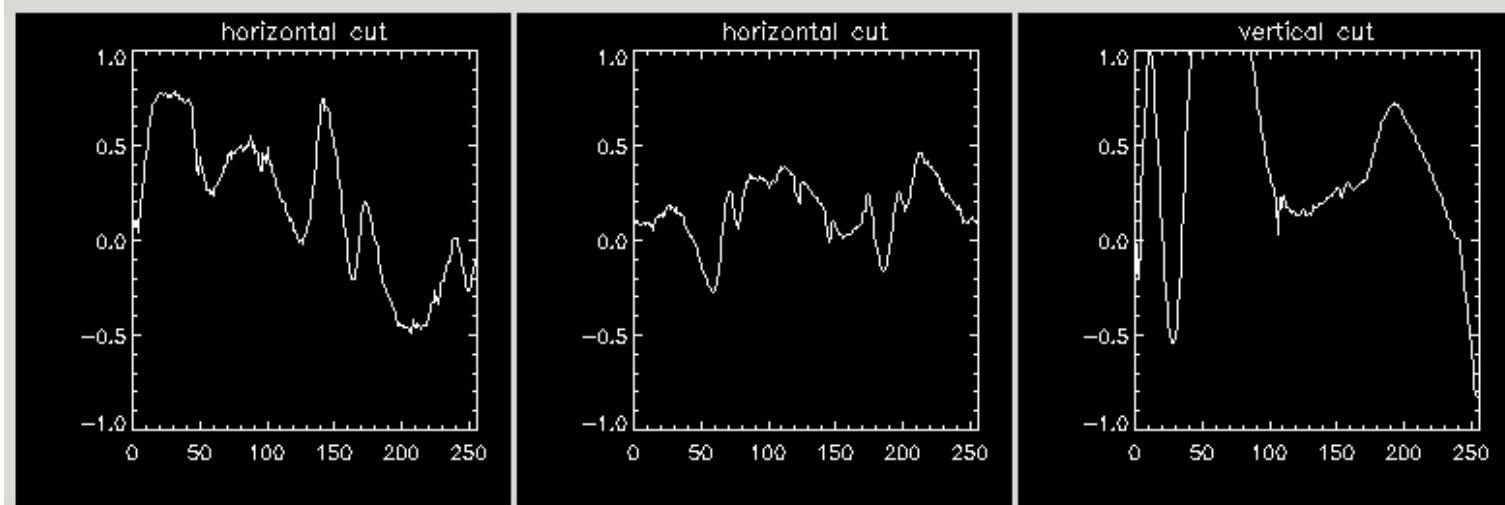
In(ρ):

SP



-1.00000

1.00000



131

 absolute scaling subtract averages

RESET

55

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LOAD

103

data set LOGARITHMIC_DENSITY =

SAVE

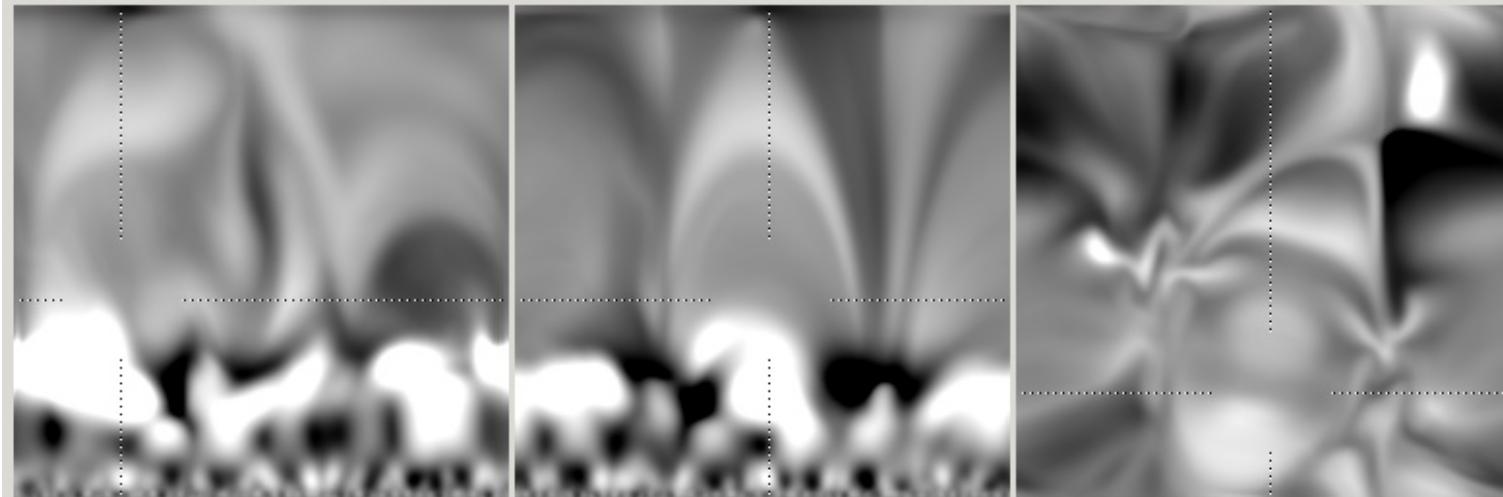
time step VRPC =

PLAY

In(ρ):

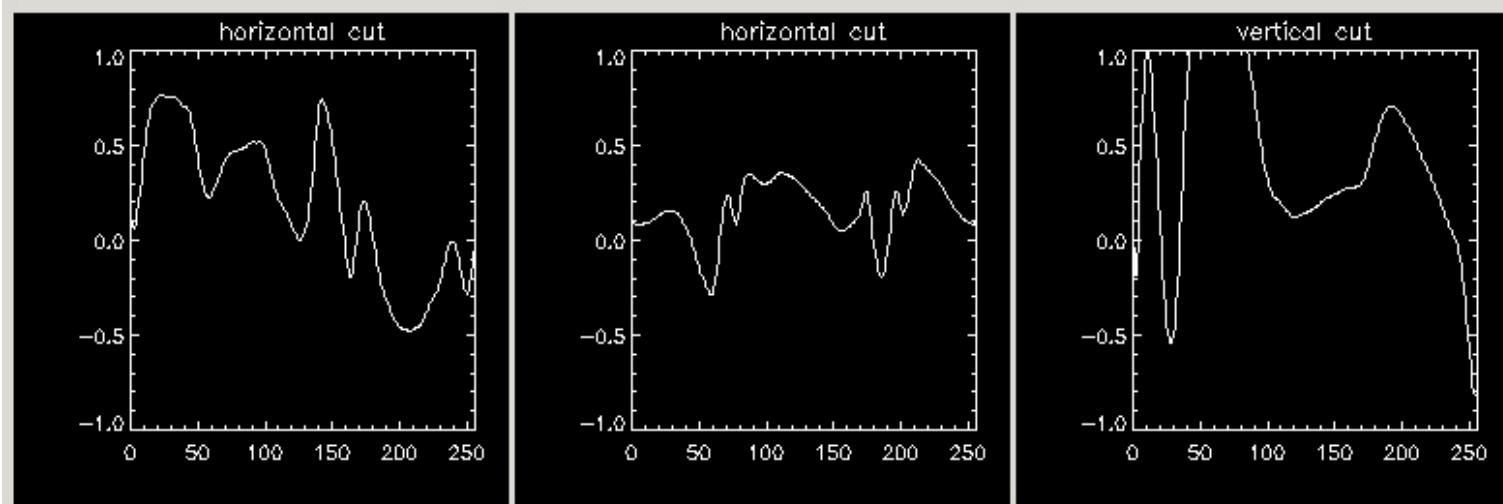
DP

QUIT



-1.00000

1.00000



Granularity effects in physical observables

Vertical velocity:

136

222

132

data set

VELOCITY_Z

time step

VAPC

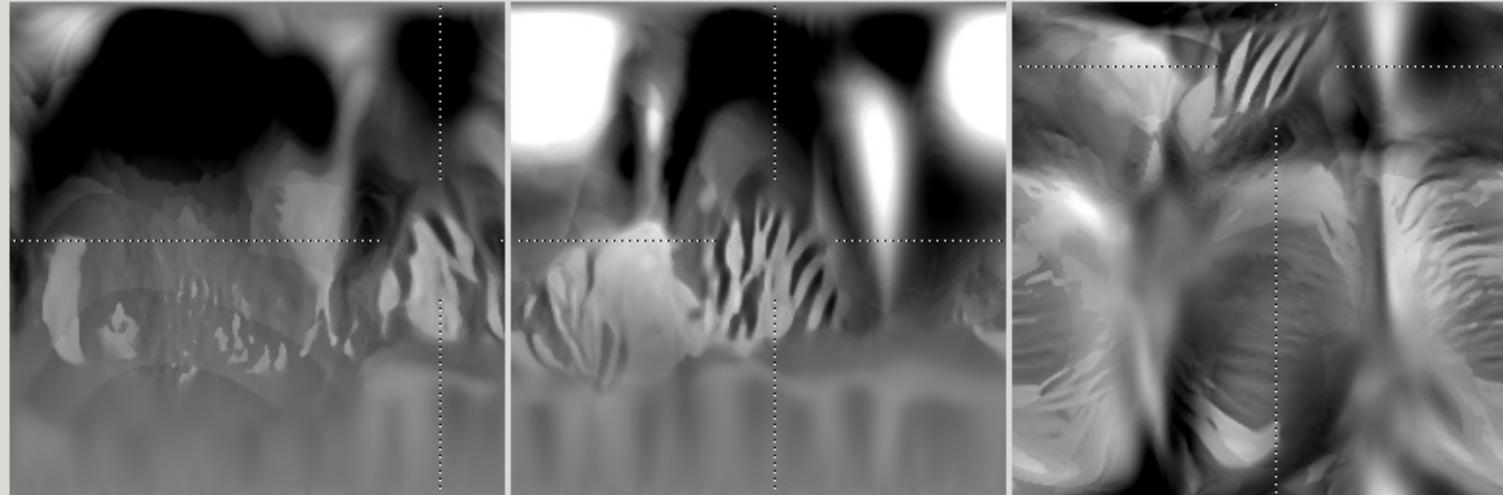
RESET

LOAD

SAVE

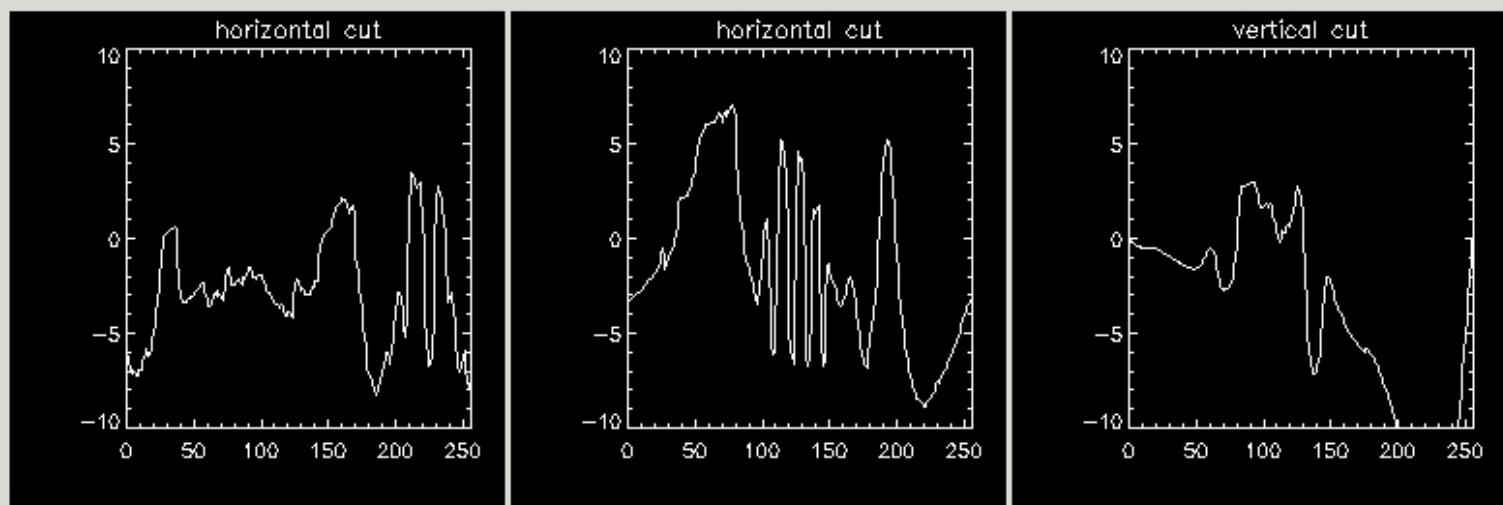
PLAY

QUIT

Vz:**SP**

-10.0000

10.0000



136

222

132

data set

VELOCITY_Z

time step

VAPC

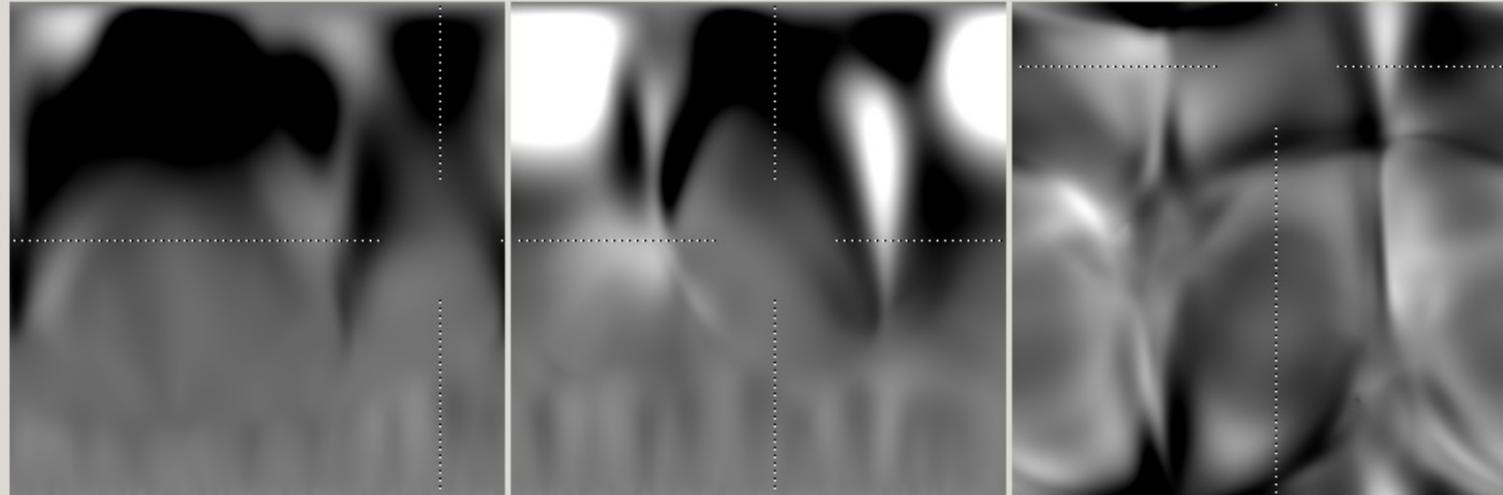
RESET

LOAD

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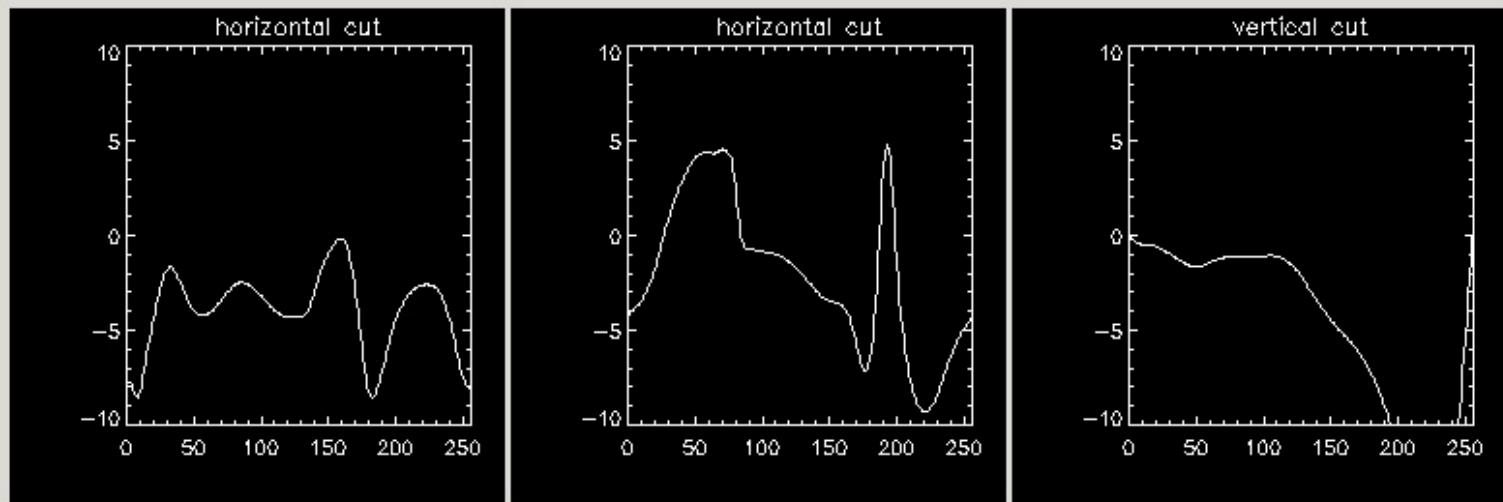
PLAY

QUIT

Vz:**DP**

I -10.0000

10.0000



Granularity effects in physical observables

Current density / Heating rate: $H = \mu_0 \eta j^2 \sim \nabla \times (\nabla \times A)$

|j|:
SP

79

 79

105

 105

170

 170 absolute scaling subtract averages

RESET

 show crosshairs show horizontal averages

LOAD

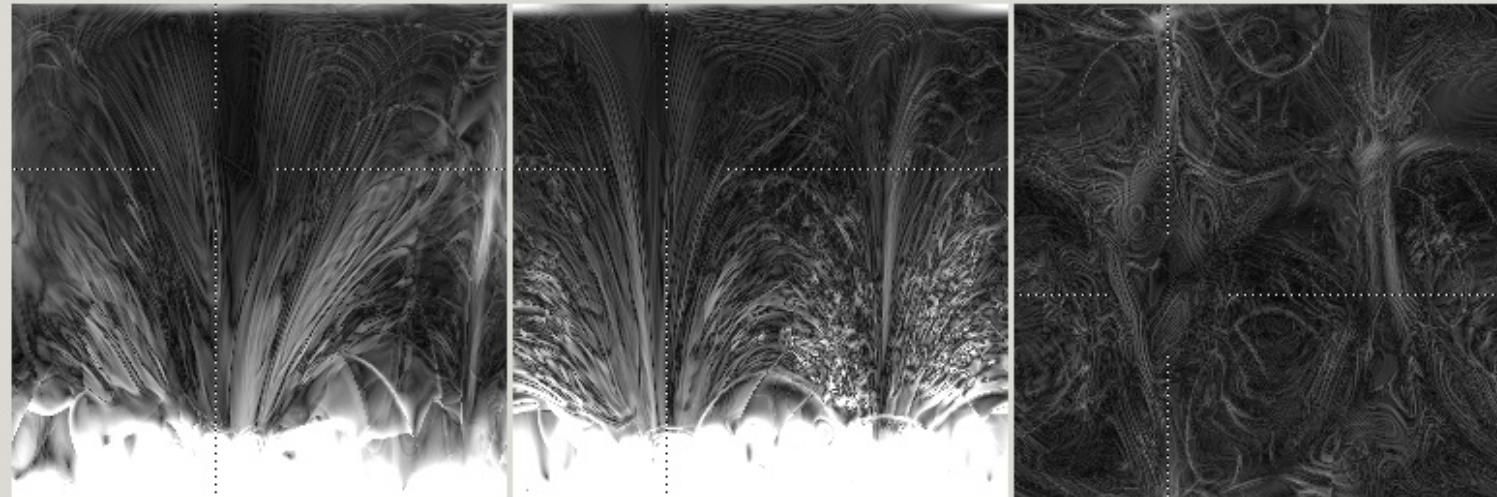
data set CURRENTDENSITY

SAVE

time step VARIOUS

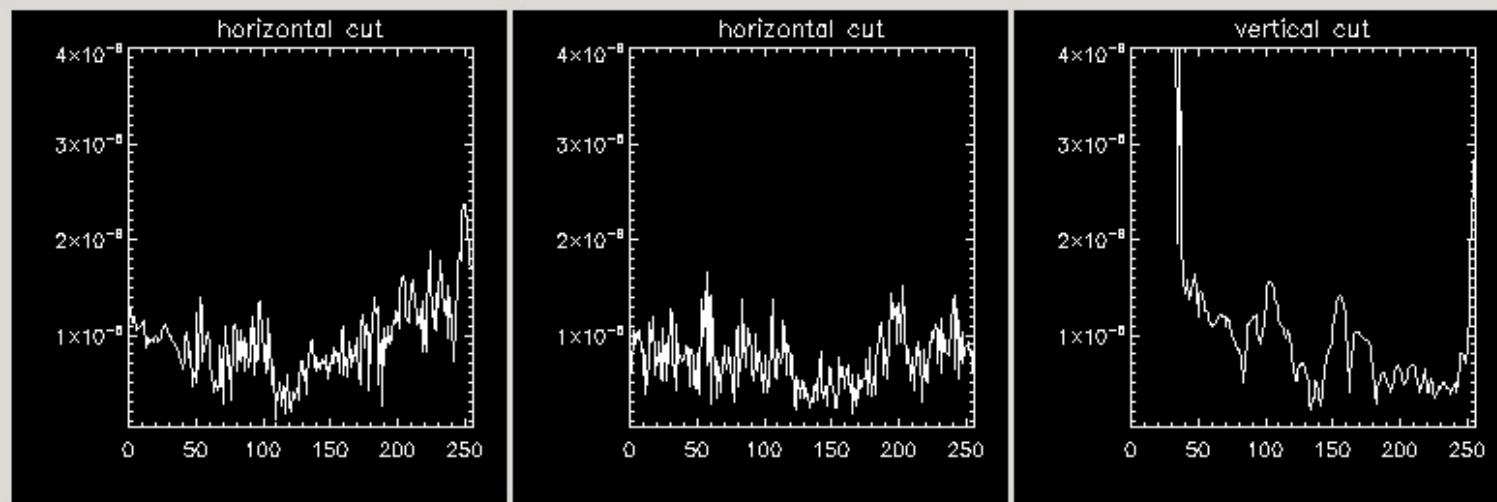
PLAY

QUIT



4.10000E-08

4.00000E-06



|j|:

DP

79

 79

105

 105

170

 170

absolute scaling subtract averages

RESET

show crosshairs show horizontal averages

LOAD

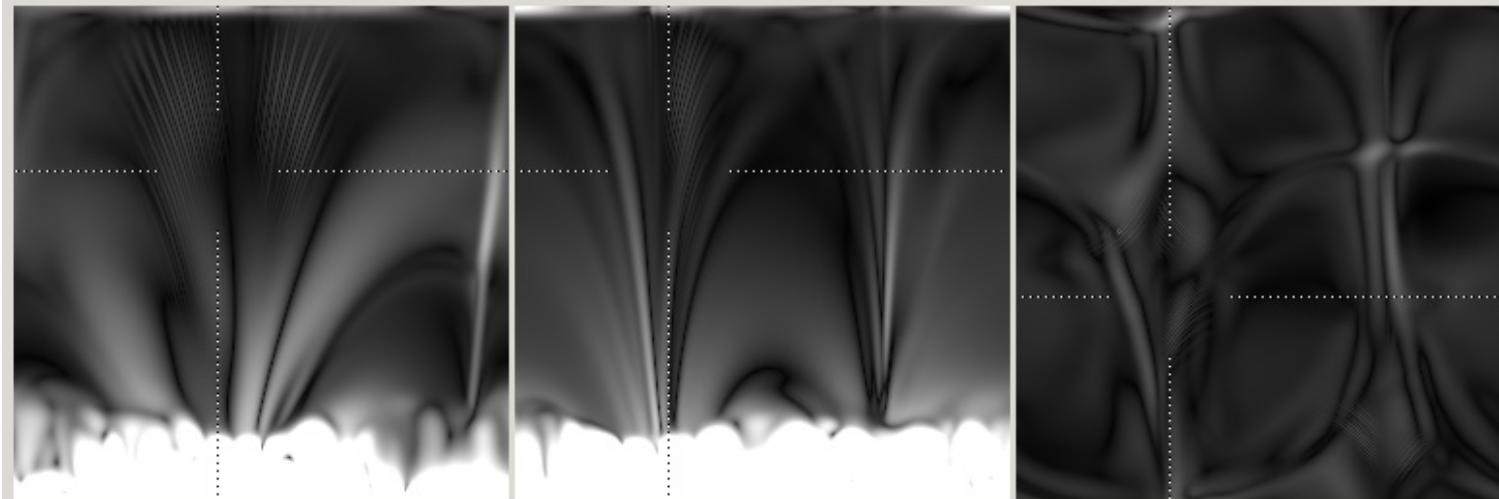
data set CURRENTDENSITY

SAVE

time step VRSC

PLAY

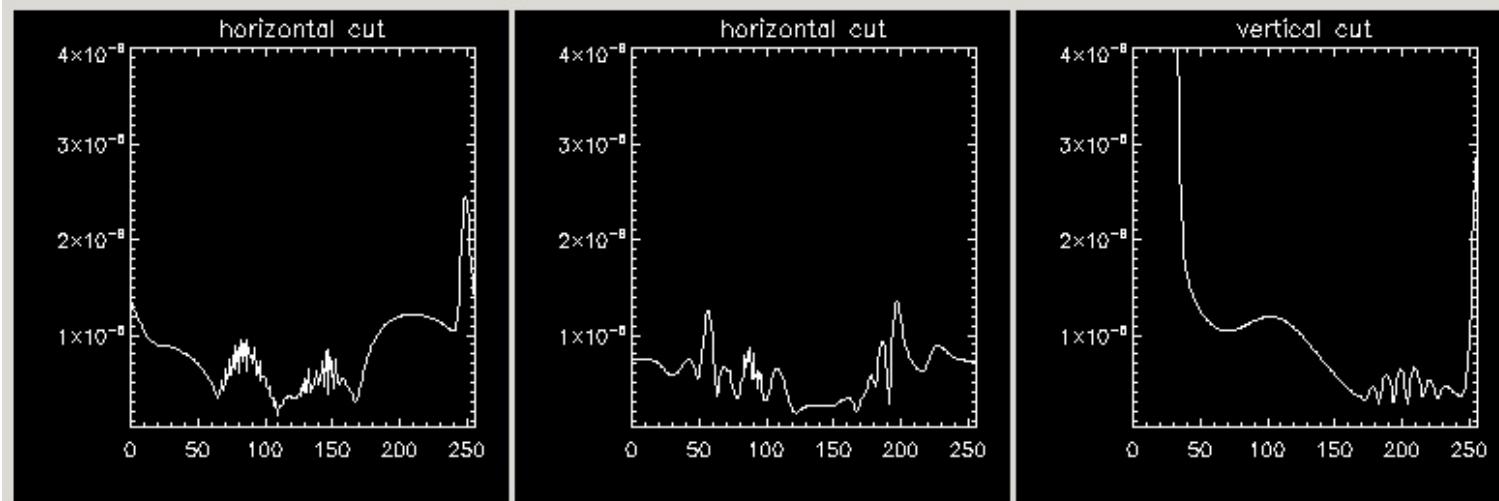
QUIT



I 4.10000E-08

 I

x 4.00000E-06

 x

Granularity effects in physical observables

Catastrophic cancellation:

“What Every Computer Scientist Should Know About Floating-Point Arithmetic”

(David Goldberg, 1991)

$$\nabla \times A \sim B \quad ; \quad \nabla \times B \sim j$$

Granularity effects in physical observables

Catastrophic cancellation:

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$$\nabla \times A \approx B \quad ; \quad \nabla \times B \approx j$$

$$j_z \approx \frac{\partial B_y}{\partial x} - \frac{\partial B_x}{\partial y}$$

$$B_x \approx B_y$$

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$$B_x \approx B_y \quad \Rightarrow \quad \partial B_x \approx \partial B_y$$

$$\Rightarrow \quad j_z \approx 2 \cdot \epsilon$$

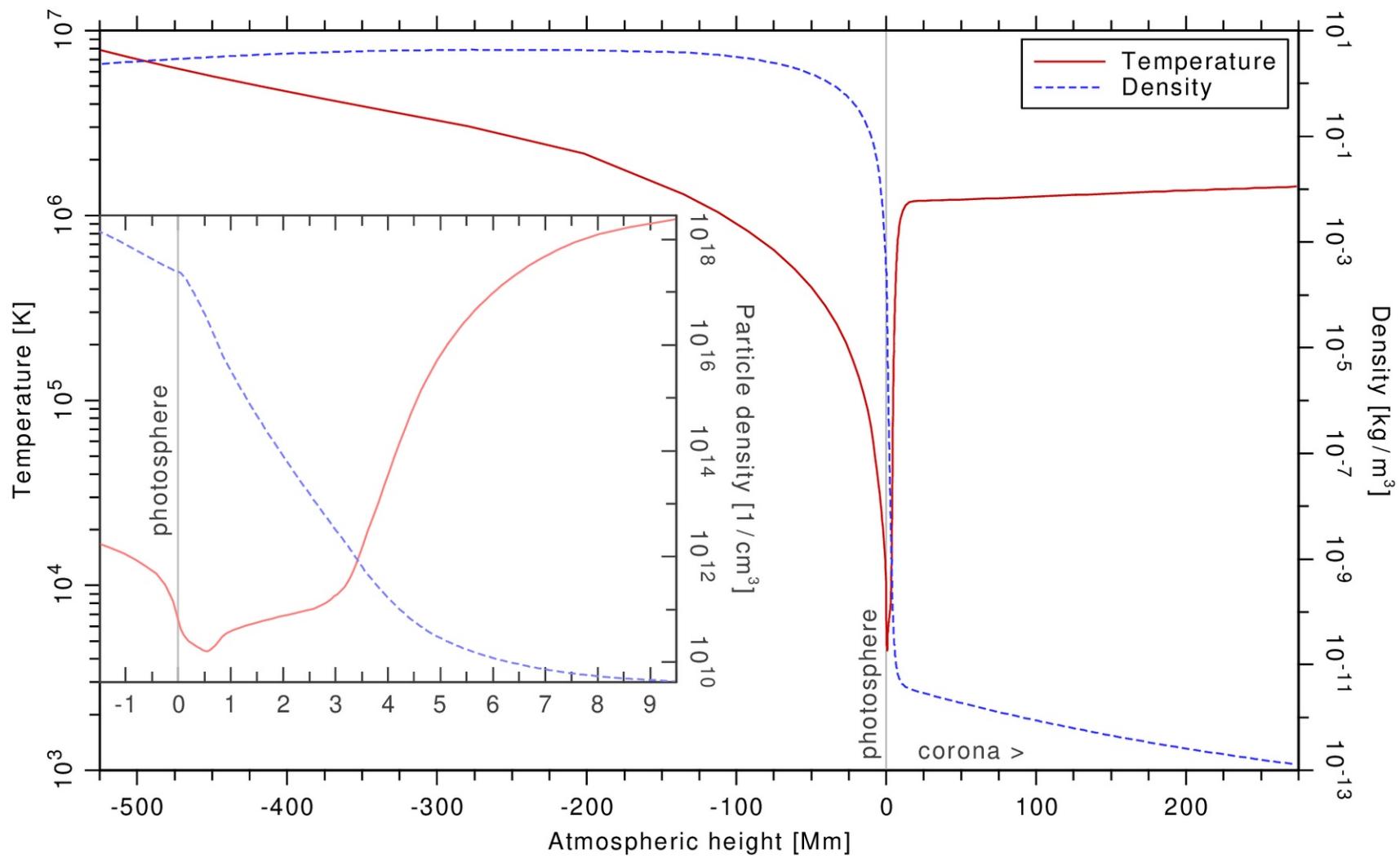

Initial condition

Initial condition

Problem: initial condition might not be in „numerical equilibrium“

Initial condition

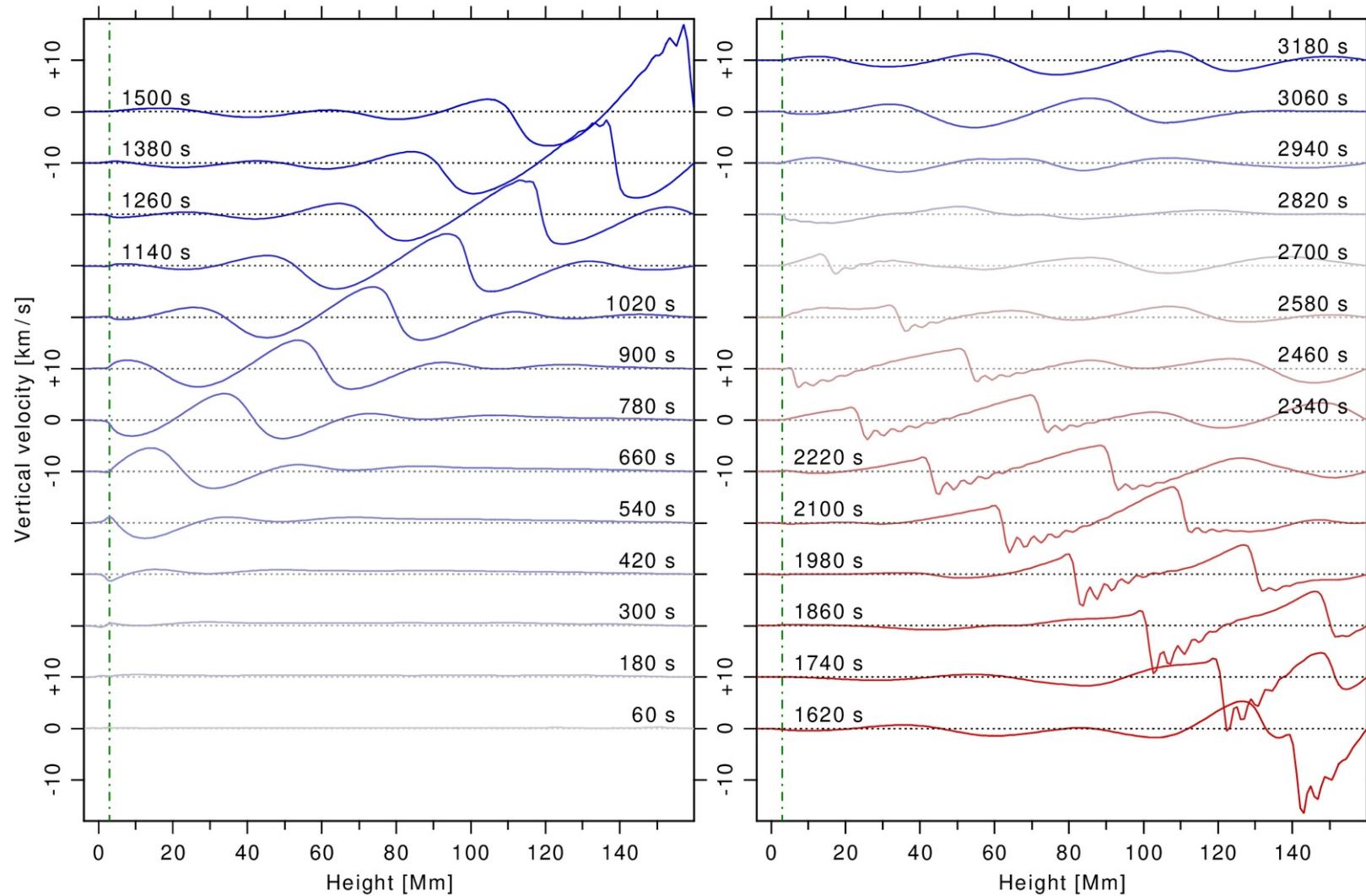
Problem: initial condition might not be in „numerical equilibrium“



=> Solar atmosphere in hydrostatic equilibrium (Bourdin, 2014, CEAB)

Initial condition

Problem: initial condition might not be in „numerical equilibrium“



=> Analytic solution needs time to equilibrate

(Bourdin, 2014, CEAB)

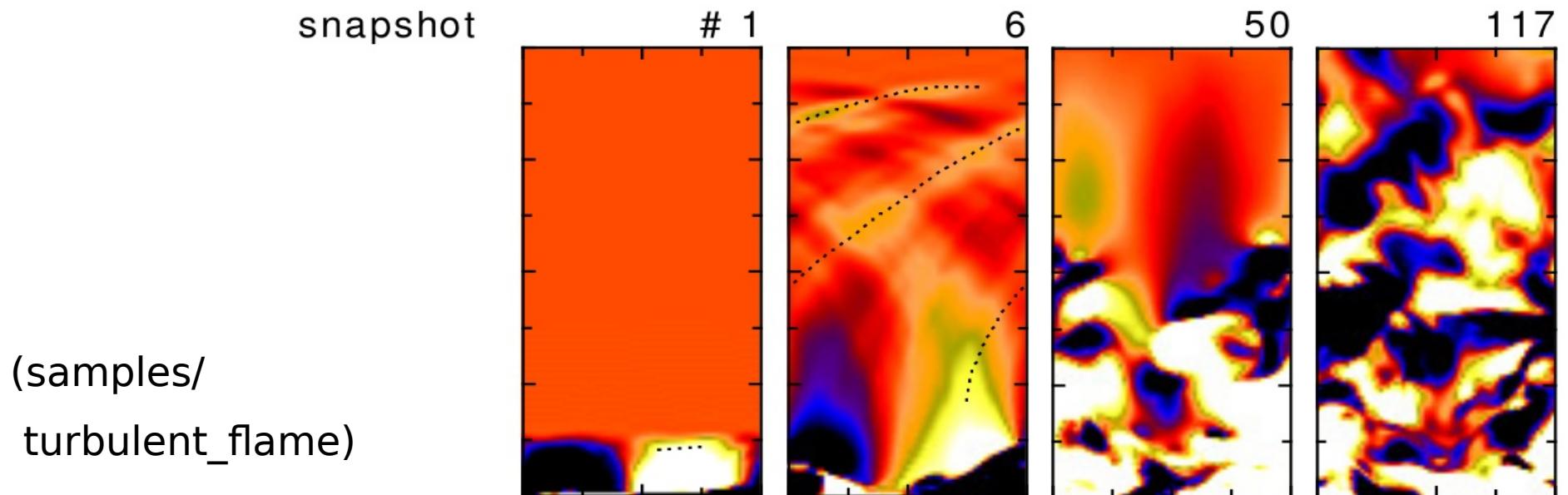
Driving a simulation with external forces

Driving a simulation with external forces

Problem: switching on over small dt => infinite momentum

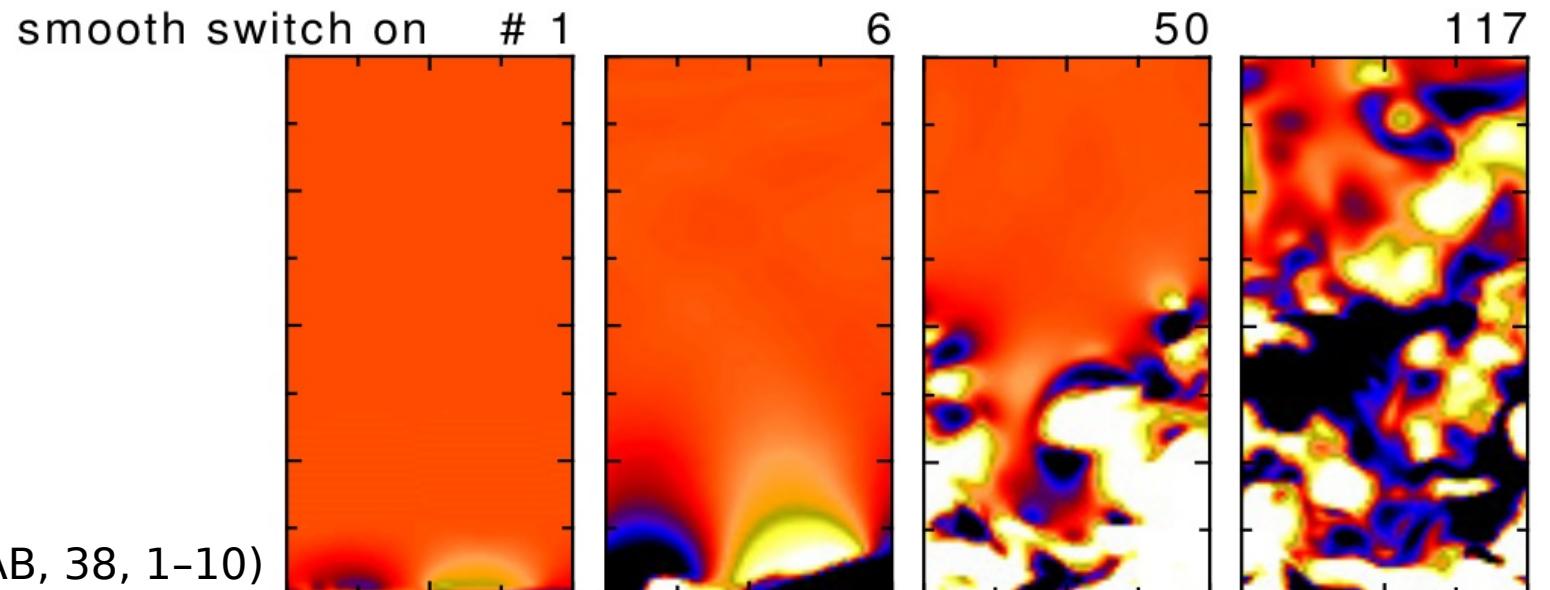
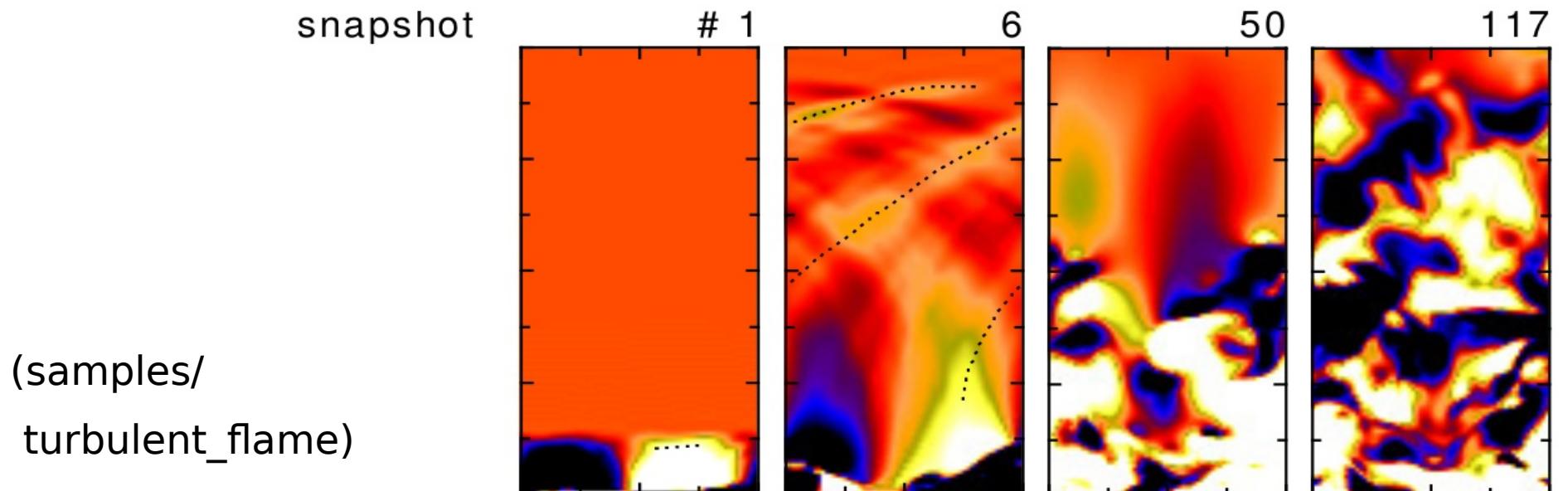
Driving a simulation with external forces

Problem: switching on over small dt => infinite momentum => shock!



Driving a simulation with external forces

Problem: switching on over small dt => infinite momentum => shock!



(Bourdin, 2014, CEAB, 38, 1-10)