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Abstract

The abstract of the report goes here. The abstract should state the topic(s) under investigation and the main results or conclusions. Methods or approaches should be stated if this is appropriate for the topic. The abstract should be self-contained, concise and clear. The typical length is one paragraph.

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1 Introduction

2 Background

2.1 Finance

A common problem in finance is to price financial derivatives, often referred just as derivatives. In essence, derivatives are contracts set between parties whose value in time derives from the price of their underlying assets. A notorious family of derivatives in financial markets are "options". Options are contracts set between two parties in which the holder has the right to sell or buy, commonly referred as exercise, an underlying stock at a preestablished price, also known as "strike price", in the future. Options are referred as "call options" or as "put options" if the exercise position is to buy or to sell, respectively. Similarly, options are classified depending on their exercise style. In that regard, the simplest of options are European options. European options give the right to exercise at the expiration date of the contract. Another well known type options are American options. American options give the right to exercise at any point in time between the beginning and expiration date of the contract. Let us define the payoff as

$$H(S, t) = \begin{cases} \max(S - K, 0) & \text{(call)} \\ \max(K - S, 0) & \text{(put)} \end{cases} \quad (2.1)$$

where K is the strike price, $S \in [0, \infty]$ is the stock price, and $t \in [0, T]$ is the current time. Note that t is measure in years, $t = 0$ and $t = T$ denotes the beginning of the contract and the expiration date, and the region $[0, T]$ is the life span of the option. While an American option's payoff is defined for all (S, t) , European options' payoff is only defined at $t = T$.

Obviously, options give greater flexibility to holders by removing their exposure of a negative payoff which is why writers of the option charge a premium to the buyers at the time they enter the contract. The premium is often referred as the price or value of the option and the problem of finding this value is called option pricing. When pricing options, it is important to find the just price because otherwise the writer or buyer of the option could set some scheme in which option will always be profitable to them. In other words, options pricing must follow the

principle of no-arbitrage. Therefore, we assume that the writer of the option uses the premium to construct a portfolio consisting of ϕ_0 units of the stock and invests ψ_0 units of cash into a risk-free asset such as US treasury bill, certificate of deposit, or bank account. Then, the writer rebalanced the portfolio (ϕ_0, ψ_0) to hedge any possible claims from the buyer of the option at any future time $0 < t \leq T$. Therefore, at any time t , the writer holds a portfolio $(\phi(t), \psi(t))$ with value

$$\Pi(t) = \phi(t)S(t) + \psi(t)B(t) \quad (2.2)$$

Moreover, the portfolio is self-financing. In other words, the changes in the value of the portfolio $V(t)$ depend on the changes in $S(t)$ and $B(t)$, and the current portfolio $(\phi(t), \psi(t))$

$$d\Pi(t) = \phi(t)dS(t) + \psi(t)dB(t) \quad (2.3)$$

Finally, the value of an option must satisfy the following

$$\Pi(t) = V(t) \quad (2.4)$$

at any time $0 \leq t \leq T$.

The black schole model is built upon the self-financing portfolio hedging strategy and expresses a mathematical model for dynamics of option's price. The black schole model makes with some assumptions about the market. For complete list of these assumptions look at (reference). We enumerates the one we believe are important for our task. First, the stock price $S(t)$ is a log-normal random variable

$$dS = r(t)Sdt + \sigma(t)SdW \quad (2.5)$$

where the risk-free interest $r(t)$ and the price volatility $\sigma(t)$ are deterministic functions of time during the life of the option. Secondly, the bank account $B(t)$ is a deterministic function of time

$$dB = r(t)B(t)dt \quad (2.6)$$

Finally, the stock does not pay dividends. From now on, we will assume that the risk-free interest rate and stock price volatility are constant during the life of the option. Later on, we will address the assumption about dividends.

By applying the black schole model to price European options, the famous black-schole PDE is obtained

$$\begin{cases} \frac{\partial V}{\partial t} + \frac{1}{2}\sigma^2 S^2 \frac{\partial^2 V}{\partial S^2} + rS \frac{\partial V}{\partial S} - rV = 0 & \text{for } t \in [0, T) \text{ and } S \in [0, \infty) \\ V(S, T) = H(S, T) & \text{for } S \in [0, \infty) \end{cases} \quad (2.7)$$

where $V(S, t)$ is a deterministic function. For the derivation of (2.7), we reference to REFERENCE.

We previously mention that the one of the assumptions of the Black-Schole model is that the underlying stock does not pay dividends. In most cases, assets such as stock pay out dividends just a few times at year. Therefore, dividends are to be modelled discretely. However, there are certain assets that pay out a proportion of the current asset price during and interval of time. Thus, in such cases, it is useful to model dividends as a continuous yield. By arbitrage arguments [REFERENCES], it can be shown that the asset price with volatility $\sigma(t)$, rate of return $r(t)$, and continuous dividend yield $\delta(S, t)$ paid at instant of time dt is modeled as

$$dS = (r(t) - \delta(S, t))Sdt + \sigma(t)SdW \quad (2.8)$$

Similarly to as we did for the risk-free interest rate and the volatility of the asset price, we will assume that continuous dividends yield is as constant from now on. [REFERENCES] show that applying the Black-Schole model under the price model (2.8) to price European options, we obtain the slightly modified version of (2.7)

$$\begin{cases} \frac{\partial V}{\partial t} + \frac{1}{2}\sigma^2 S^2 \frac{\partial^2 V}{\partial S^2} + (r - \delta)S \frac{\partial V}{\partial S} - rV = 0 & \text{for } t \in [0, T) \text{ and } S \in [0, \infty) \\ V(S, T) = H(S, T) & \text{for } S \in [0, \infty) \end{cases} \quad (2.9)$$

Similarly, the Black-Schole model is applied to price American options. One important result is that the value of an American option $V_{\text{Ame}}(t)$ is bounded from below by the payoff function

$$V_{\text{Am}}(S, t) \geq H(S, t) \quad \text{for } t \in [0, T] \quad (2.10)$$

Moreover, the domain of $V(S, t)$ can be separated in the exercise region

$$\mathcal{S} := \{(S, t) : V(S, t) = H(S, t)\} \quad (2.11)$$

and the continuation region

$$\mathcal{C} := \{(S, t) : V(S, t) > H(S, t)\} \quad (2.12)$$

where the boundary of \mathcal{C} is defined as

$$\partial\mathcal{C} := \{(S, t) : S = \bar{S}(t)\} \quad (2.13)$$

Lastly, the price dynamics of American options behaves as European options within the continuation region. Since we know $V(S, t)$ at the stopping region, we only need to solve $V(S, t)$ at continuation region and determine its boundary ∂ at the same time. Therefore, this is known as the free boundary problem formulation of the American option pricing problem and is equivalent to solve.

$$\begin{cases} \frac{\partial V}{\partial t} + \frac{1}{2}\sigma^2 S^2 \frac{\partial^2 V}{\partial S^2} + (r - \delta)S \frac{\partial V}{\partial S} - rV = 0 & \text{for } (S, t) \in \mathcal{C} \\ V(S, t) = H(S, t) & \text{for } (S, t) \in \partial\mathcal{C} \end{cases} \quad (2.14)$$

Now, note that the bound of $V(S, t)$ in each region is

$$V(S, t) - H(S, t) > 0 \quad \text{for all } (S, t) \in \mathcal{C}$$

$$V(S, t) - H(S, t) = 0 \quad \text{for all } (S, t) \in \mathcal{S}$$

Similarly, the bound of $\mathcal{L}_{BS}(V)$ is

$$\begin{aligned} \frac{\partial V}{\partial t} + \frac{1}{2}\sigma^2 S^2 \frac{\partial^2 V}{\partial S^2} + (r - \delta)S \frac{\partial V}{\partial S} - rV &= 0 & \text{for } (S, t) \in \mathcal{C} \\ \frac{\partial V}{\partial t} + \frac{1}{2}\sigma^2 S^2 \frac{\partial^2 V}{\partial S^2} + (r - \delta)S \frac{\partial V}{\partial S} - rV &< 0 & \text{for } (S, t) \in \mathcal{S} \end{aligned}$$

Therefore, grouping the bounds above we form a linear complementary system of equations

$$\begin{cases} \mathcal{L}_{BS}(V) \cdot [V(S, t) - H(S, t)] = 0 & \text{for all } (S, t) \\ V(S, t) - H(S, t) \geq 0 & \text{for all } (S, t) \\ \mathcal{L}_{BS}(V) \leq 0 & \text{for all } (S, t) \\ V(S, T) = H(S, T) \end{cases}$$

By combining the expression above with the lower bound (2.10), we obtain a linearly complementary system of equations

The system of equation (2.14) could be reformulated as free a boundary problem, or as a linear complementary problem. To reframe the problem as free boundary problem, note that we know the value of $V(S, t)$ in the stopping region. Hence, we only need to solve the Black-Schole PDE at the continuation region but with unknown boundary condition given by $S = \bar{S}(t)$. Finally, at the terminal condition the unknown boundary condition is given by

$$\bar{S}(T) = K.$$

$$\begin{cases} \frac{\partial V}{\partial t} + \frac{1}{2}\sigma^2 S^2 \frac{\partial^2 V}{\partial S^2} + (r - \delta)S \frac{\partial V}{\partial S} - rV = 0 & \text{for } (S, t) \in \mathcal{C} \\ V(S, t) = H(S, t) & \text{for all } (t, \bar{S}(t)) \\ S = \bar{S}(T) \end{cases} \quad (2.15)$$

The system (2.15) is the free boundary formulation for the pricing problem for American options. A detailed derivation of (??) can be found at [REFERENCES].

The equation (??) is a parabolic PDE. Moreover, by applying the transformation,

$$S = Ke^x, \quad t = T - \frac{2\tau}{\sigma^2}, \quad q := \frac{2r}{\sigma^2}, \quad q_\delta := \frac{2(r - \delta)}{\sigma^2} \quad (2.16)$$

$$v(x, \tau) := V(S, t) = V\left(Ke^x, T - \frac{2\tau}{\sigma^2}\right) \quad (2.17)$$

$$\alpha := \frac{1}{2}(q_\delta - 1) \quad \beta := \frac{1}{4}(q_\delta - 1)^2 + q \quad (2.18)$$

$$v(x, \tau) := e^{-(\alpha x + \beta \tau)} y(x, \tau) \quad (2.19)$$

the equation (??) converts to the heat diffusion PDE.

$$h(x, \tau) := \frac{H(S, t)}{K} = \begin{cases} \max(e^x - 1, 0) \\ \max(1 - e^x, 0) \end{cases} \quad (2.20)$$

$$\bar{x}(\tau) := \log \bar{S}(t) - \log K \quad (2.21)$$

$$\begin{cases} \frac{\partial y}{\partial \tau} = \frac{\partial^2 y}{\partial x^2} & \text{for } \tau \in [0, \frac{\sigma^2}{2}T) \text{ and } x \in (\bar{x}(t), \infty) \\ y(x, \tau) = e^{(\alpha x + \beta \tau)} h(x, \tau) & \text{for } \tau \in [0, \frac{\sigma^2}{2}T] \text{ and } x \in (-\infty, \bar{x}(\tau)] \\ \bar{x}(0) = 0 \end{cases} \quad (2.22)$$

We can reformulate equation (2.22) as:

$$g := e^{\alpha x + \beta \tau} h(x, \tau) \quad (2.23)$$

$$\begin{cases} \left(\frac{\partial y}{\partial \tau} - \frac{\partial^2 y}{\partial x^2} \right) (y - g) = 0 \\ \frac{\partial y}{\partial \tau} - \frac{\partial^2 y}{\partial x^2} \geq 0 \quad y - g \geq 0 \\ y(x, 0) = g(x, 0) \end{cases} \quad (2.24)$$

By exploring the geometric properties of the value function $V(S, t)$, we can determine useful conditions that will later help on in solving the equation (??). Firstly, at any given time $0 \leq t \leq T$, American options match the linear segment of the payoff function within the stopping region. Therefore, we could say that

$$\frac{\partial V}{\partial S}(S, t) = \begin{cases} -1 & \text{(put)} \\ 1 & \text{(call)} \end{cases} \quad (2.25)$$

Moreover, as the price goes to infinity the value of the option tends to zero

$$\lim_{S \rightarrow \infty} V(S, t) = 0 \quad (2.26)$$

2.2 Numerical methods for the heat diffusion PDE

Pricing American options requires using numerical methods. The Black-Schole PDE in (XXX) can be converted to the heat diffusion equation

Therefore, we focus on analyzing the numerical solution of the heat diffusion equation. Suppose the equation (XXX) is defined within the rectangular region $[x_{\min}, x_{\max}] \times [\tau_{\min}, \tau_{\max}]$. By discretizing uniformly along the spatial direction x and temporal direction τ ,

$$M := \frac{x_{\max} - x_{\min}}{\Delta x} \quad (2.27)$$

$$N := \frac{\tau_{\max} - \tau_{\min}}{\Delta \tau} \quad (2.28)$$

$$x_i := x_{\min} + i\Delta x \quad \text{for } i = 0, \dots, M \quad (2.29)$$

$$\tau_i := \tau_{\min} + i\Delta \tau \quad \text{for } i = 0, \dots, N \quad (2.30)$$

where Δx and $\Delta \tau$ are the distance between two consecutive points, then, the grid is defined as the discrete region.

$$\mathcal{G} := \{(x_i, \tau_j) : (i, j) \in \{0, \dots, M\} \times \{0, \dots, N\}\} \quad (2.31)$$

Thus, solving numerically the equation (XXX) means finding an approximation for $y(x_i, \tau_n)$ at every (x_i, τ_n) within the grid \mathcal{G} ,

$$y_i^n \approx y(x_i, \tau_n) \quad (2.32)$$

To obtain such approximation, we rely on central difference approximations (REFERENCE).

$$f'(x_i) = \frac{f_{i+1} - f_i}{h} + O(h) \quad (2.33)$$

$$f'(x_i) = \frac{f_{i+1} - f_{i-1}}{2h} + O(h^2) \quad (2.34)$$

$$f''(x_i) = \frac{f_{i+1} - 2f_i + f_{i-1}}{h^2} + O(h^2) \quad (2.35)$$

2.2.1 Explicit scheme

An explicit scheme is one where we approximate the time partial derivative using a forward difference approximation. Hence, the PDE in (XXX) is approximated as

$$\frac{y_i^{n+1} - y_i^n}{\Delta\tau} = \frac{y_{i-1}^n - 2y_i^n + y_{i+1}^n}{(\Delta x)^2} \quad (2.36)$$

By rearranging the terms,

$$\lambda := \frac{\Delta\tau}{(\Delta x)^2} \quad (2.37)$$

$$y_i^{n+1} = \lambda y_{i-1}^n + (1 - 2\lambda)y_i^n + \lambda y_{i+1}^n \quad (2.38)$$

It is shown by reference [REFERENCE] that method (XXX) is stable and consistent under the following condition

$$0 < \Delta\tau \leq \frac{(\Delta x)^2}{2} \quad (2.39)$$

Moreover, the method has order of convergence $O(\Delta\tau, (\Delta x)^2)$.

2.2.2 Implicit scheme

The implicit scheme approximates the time derivative using a backward difference

$$\frac{y_i^{n+1} - y_i^n}{\Delta\tau} = \frac{y_{i-1}^{n+1} - 2y_i^{n+1} + y_{i+1}^{n+1}}{(\Delta x)^2} \quad (2.40)$$

$$y_i^{n+1} - \lambda(y_{i-1}^{n+1} - 2y_i^{n+1} + y_{i+1}^{n+1}) = y_i^n \quad (2.41)$$

$$K := \begin{bmatrix} 2 & -1 & & 0 \\ -1 & \ddots & \ddots & \\ & \ddots & \ddots & \ddots \\ 0 & & \ddots & \ddots \end{bmatrix} \quad (2.42)$$

$$(I + \lambda K)\mathbf{y}^{n+1} = \mathbf{y}^n \quad (2.43)$$

2.2.3 Theta method

$$\frac{y_i^{n+1} - y_i^n}{\Delta\tau} = (1 - \theta) \frac{y_{i-1}^n - 2y_i^n + y_{i+1}^n}{(\Delta x)^2} + \theta \frac{y_{i-1}^{n+1} - 2y_i^{n+1} + y_{i+1}^{n+1}}{(\Delta x)^2} \quad (2.44)$$

$$y_i^{n+1} - \lambda\theta(y_{i-1}^{n+1} - 2y_i^{n+1} + y_{i+1}^{n+1}) = y_i^n + (1 - \theta)\lambda(y_{i-1}^n - 2y_i^n + y_{i+1}^n) \quad (2.45)$$

$$(1 + \lambda\theta K)\mathbf{y}^{n+1} = (1 - \lambda\theta K)\mathbf{y}^n \quad (2.46)$$

3 Front fixing method

3.1 Inverse transformation

In order to remove the free boundary in the system of equations, the following transformation is used:

$$x = \frac{S}{\bar{S}(t)} \quad (3.1)$$

Next,

$$v(x, t) := V(x\bar{S}(t), t) = V(S, t) \quad (3.2)$$

By computing the partial derivatives of V with respect to S and t

$$\frac{\partial V}{\partial t} = \frac{\partial v}{\partial t} + \frac{\partial v}{\partial x} \frac{\partial x}{\partial t} = \frac{\partial v}{\partial t} - x \frac{\bar{S}'(t)}{\bar{S}(t)} \frac{\partial v}{\partial x} \quad (3.3)$$

$$\frac{\partial V}{\partial S} = \frac{\partial v}{\partial x} \frac{\partial x}{\partial S} = \frac{1}{\bar{S}(t)} \frac{\partial v}{\partial x} \quad (3.4)$$

$$\frac{\partial^2 V}{\partial S^2} = \frac{1}{\bar{S}(t)^2} \frac{\partial^2 v}{\partial x^2} \quad (3.5)$$

an expression for (3.1) with respect to x is derived:

$$\frac{\partial v}{\partial t} + \frac{1}{2} \sigma^2 x^2 \frac{\partial^2 v}{\partial x^2} + \left[(r - \delta) - \frac{\bar{S}'(t)}{\bar{S}(t)} \right] x \frac{\partial v}{\partial x} - rv = 0 \quad \text{for } x > 1 \text{ and } 0 \leq t < T \quad (3.6)$$

Similarly, (3.2) is reformulated in terms of x to:

$$v(x, t) = K - x\bar{S}(t) \quad \text{for } 0 \leq x \leq 1 \text{ and } 0 \leq t < T \quad (3.7)$$

Next, the terminal condition (3.3) is re-written with respect to x :

$$v(x, T) = \max(K - x\bar{S}(T), 0) = K \max(1 - x, 0) = 0 \quad \text{for } x \geq 1 \quad (3.8)$$

Finally, the left and right boundary conditions are given with respect to x :

$$\frac{\partial v}{\partial x}(x, t) = -\bar{S}(t) \quad (3.9)$$

$$\lim_{x \rightarrow \infty} v(x, t) = 0 \quad (3.10)$$

In summary, a non linear system of PDEs is obtained:

$$\frac{\partial v}{\partial t} + \frac{1}{2}\sigma^2 x^2 \frac{\partial^2 v}{\partial x^2} + \left[(r - \delta) - \frac{\bar{S}'(t)}{\bar{S}(t)} \right] x \frac{\partial v}{\partial x} - rv = 0 \quad \text{for } x > 1 \text{ and } 0 \leq t < T \quad (3.11)$$

$$v(x, t) = K - x\bar{S}(t) \quad \text{for } 0 \leq x \leq 1 \text{ and } 0 \leq t < T \quad (3.12)$$

$$v(x, T) = 0 \quad \text{for } x \geq 1 \quad (3.13)$$

$$\frac{\partial v}{\partial x}(x, t) = -\bar{S}(t) \quad (3.14)$$

$$\lim_{x \rightarrow \infty} v(x, t) = 0 \quad (3.15)$$

$$\bar{S}(T) = K \quad (3.16)$$

3.1.1 Theta method

We discretize the system of equation containing (3.17) and (3.18) to solve it. Firstly, we define an uniform meshgrid within the region $[1, x_{\max}] \times [0, T]$ and with resolution Δx and Δt .

$$M := \frac{x_{\max} - 1}{\Delta x} \quad N := \frac{T}{\Delta t}$$

$$\begin{aligned}
x_i &:= 1 + i\Delta x && \text{for } i = 0, \dots, M \\
t_n &:= n\Delta t && \text{for } n = 0, \dots, N
\end{aligned}$$

Now we define the approximations

$$v_i^n \approx v(x_i, t_n) \quad \text{for } (x_i, t_n) \in \{x_k\}_0^M \times \{t_k\}_0^N \quad (3.17)$$

$$\bar{S}^n \approx \bar{S}(t_n) \quad \text{for } t_n \in \{t_k\}_0^N \quad (3.18)$$

By the boundary conditions, we can derive an expression for

$$v_0^n = K - \bar{S}^n \quad \text{for } n = 0, \dots, N-1, N \quad (3.19)$$

$$v_{M+1}^n = 0 \quad \text{for } n = 0, \dots, N-1, N \quad (3.20)$$

Additionally by using the smoothness condition, we get:

$$\frac{v_1^n - v_0^n}{\Delta x} = -\bar{S}^n \quad (3.21)$$

or

$$v_1^n = K - (1 + \Delta x)\bar{S}^n \quad (3.22)$$

Next, we equation (3.XX) discretize using centered finite difference. The discretization method, we use is the theta method which a interpolation between an implicit and explicit scheme.

$$\begin{aligned}
& v_i^{t+1} - v_i^t + \theta \left\{ \frac{1}{2} \sigma^2 x_i^2 \frac{\Delta t}{(\Delta x)^2} (v_{i-1}^t - 2v_i^t + v_{i+1}^t) + \left[(r - \delta) - \frac{1}{\bar{S}^t} \frac{\bar{S}^{t+1} - \bar{S}^t}{\Delta t} \right] \frac{\Delta t}{2\Delta x} (v_{i+1}^t - v_{i-1}^t) - r v_i^t \Delta t \right\} \\
& + (1 - \theta) \left\{ \frac{1}{2} \sigma^2 x_i^2 \frac{\Delta t}{(\Delta x)^2} (v_{i-1}^{t+1} - 2v_i^{t+1} + v_{i+1}^{t+1}) \right. \\
& \left. + \left[(r - \delta) - \frac{1}{\bar{S}^{t+1}} \frac{\bar{S}^{t+1} - \bar{S}^t}{\Delta t} \right] \frac{\Delta t}{2\Delta x} (v_{i+1}^{t+1} - v_{i-1}^{t+1}) - r v_i^{t+1} \Delta t \right\} = 0
\end{aligned} \tag{3.23}$$

To simplify the expression above, we introduce the following terms

$$\lambda := \frac{\Delta t}{(\Delta x)^2} \tag{3.24}$$

$$\alpha_i := 1 + \theta(\lambda \sigma^2 x_i^2 + r \Delta t) \tag{3.25}$$

$$\beta_i := -\frac{1}{2} \lambda \theta \left[\sigma^2 x_i^2 - x_i \Delta x (r - \delta) \right] - \frac{1}{2} \lambda \theta x_i \Delta x \frac{\bar{S}^{n+1} - \bar{S}^n}{\Delta t \bar{S}^n} \tag{3.26}$$

$$\gamma_i := -\frac{1}{2} \lambda \theta \left[\sigma^2 x_i^2 + x_i \Delta x (r - \delta) \right] + \frac{1}{2} \lambda \theta x_i \Delta x \frac{\bar{S}^{n+1} - \bar{S}^n}{\Delta t \bar{S}^n} \tag{3.27}$$

$$a_i := 1 - (1 - \theta)(\lambda \sigma^2 x_i^2 + r \Delta t) \tag{3.28}$$

$$b_i := \frac{1}{2} (1 - \theta) \lambda \left[\sigma^2 x_i^2 - x_i \Delta x \left((r - \delta) - \frac{1}{\Delta t} \right) \right] \tag{3.29}$$

$$c_i := \frac{1}{2} (1 - \theta) \lambda \left[\sigma^2 x_i^2 + x_i \Delta x \left((r - \delta) - \frac{1}{\Delta t} \right) \right] \tag{3.30}$$

$$d_i^{n+1} := (1 - \theta) \frac{x_i}{2\Delta x} \frac{v_{i+1}^{n+1} - v_{i-1}^{n+1}}{\bar{S}^{n+1}} \tag{3.31}$$

Now, the expression above becomes

$$\beta_i^n v_{i-1}^n + \alpha_i^n v_i^n + \gamma_i^n v_{i+1}^n = b_i v_{i-1}^{n+1} + a_i v_i^{n+1} + c_i v_{i+1}^{n+1} + d_i^{n+1} \bar{S}^n \tag{3.32}$$

$$\gamma_1^n v_2^n = b_1 v_0^{n+1} + a_1 v_1^{n+1} + c_1 v_2^{n+1} + d_1^{n+1} \bar{S}^n - \beta_1^n (K - \bar{S}^n) - \alpha_1^n (K - (1 + \Delta x) \bar{S}^n) \quad (3.33)$$

$$\alpha_2^n v_2^n + \gamma_2^n v_3^n = b_2 v_1^{n+1} + a_2 v_2^{n+1} + c_2 v_3^{n+1} + d_2^{n+1} \bar{S}^n - \beta_2^n (K - (1 + \Delta x) \bar{S}^n) \quad (3.34)$$

$$\beta_M^n v_{M-1}^n + \alpha_M^n v_M^n = b_i v_{i-1}^{n+1} + a_i v_i^{n+1} + c_i v_{i+1}^{n+1} + d_i^{n+1} \bar{S}^n \quad (3.35)$$

3.2 Log transformation

4 Linear complementary problem

5 Conclusion

6 Another section

6.1 A subsection

Subsections may be used. Use a clear structure in your report.

We denote the set of real numbers by \mathbb{R} , the set of integers by \mathbb{Z} and the set of complex numbers by \mathbb{C} . Our analysis is based on the equation $e^{\pi i} = -1$ and the relation

$$\frac{2}{4} = \frac{1}{2} \tag{6.1}$$

which we verify in the appendix B. Useful consequences are

$$\frac{4}{8} = \frac{1}{2} \tag{6.2}$$

$$\frac{4}{12} + \frac{1}{\Gamma(s)} \int_0^\infty \frac{t^{s-1}}{e^t - 1} dt = \frac{1}{3} + \sum_{n=1}^\infty \frac{1}{n^s} \tag{6.3}$$

$$\frac{2}{10} = \frac{1}{5} \tag{6.4}$$

For any $0 \neq a \in \mathbb{Z}$, the equality

$$\frac{2a}{4a} = \frac{1}{2}$$

follows from equation (6.1).

6.2 Another subsection

6.2.1 A subsubsection

Sometimes subsubsections may be appropriate.

6.2.2 Another subsubsection

This could contain a table of interesting numbers

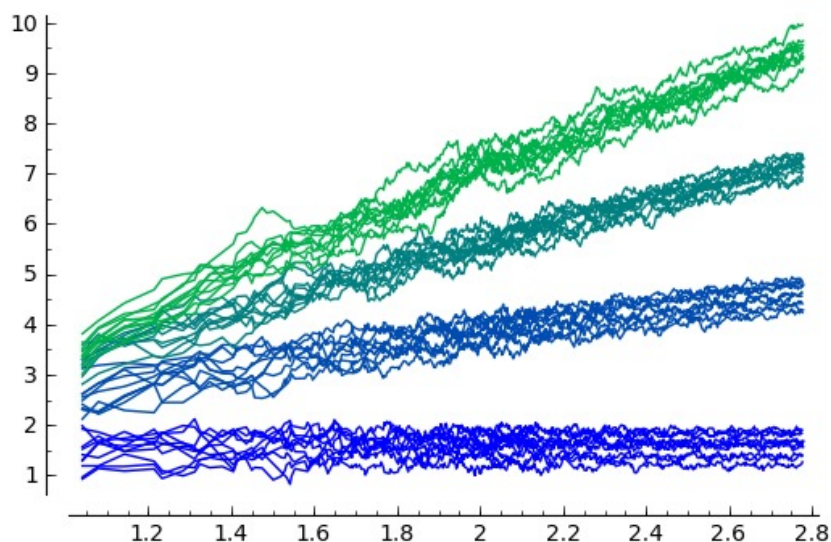


Figure 1: Oh look, something happens here !

n	1	2	3	4	5	6
F_n	1	1	2	3	5	8
B_n	$\frac{1}{2}$	$\frac{1}{6}$	0	$-\frac{1}{30}$	0	$\frac{1}{42}$
p_n	2	3	5	7	11	13

7 Yet another section

Graphics can be included. Figure 1 shows an example. Learn about floats and pictures in the \LaTeX wikibook to place the figures at the right place.

8 Conclusions

Further help on \LaTeX can be found easily on the internet. The \LaTeX wikibook¹ contains a lot. For instance you would find there how to type theorems and proofs nicely. Or how to include source code written in some programming language like python. There are long lists available with all sorts of common mathematical symbols like ξ , ∇ , ∞ , \log , \iff , etc.

¹<http://en.wikibooks.org/wiki/LaTeX>

A Raw data

Material that needs to be included but would distract from the main line of presentation can be put in appendices. Examples of such material are raw data, computing codes and details of calculations.

But note tha the maximal number of pages includes the appendix and the references.

B Calculations for section 6

In this appendix we could verify equation (6.1) or present the code that was used.

```
def gcd(a, b):  
    """  
    Return the greatest common divisor  
    of a and b  
    """  
    while b > 0:  
        (a, b) = (b, a % b)  
    return a
```


References

- [1] N. L. Alling and N. Greenleaf, *Foundations of the Theory of Klein Surfaces*, Lecture Notes in Mathematics Vol. 219 (Springer, Berlin, 1971).
- [2] J. W. Barrett and R. A. Dawe Martins, “Non-commutative geometry and the standard model vacuum”, *J. Math. Phys.* **47**, 052305 (2006). (arXiv:hep-th/0601192)
- [3] R. Bott and L. W. Tu, *Differential Forms in Algebraic Topology* (Springer, New York, 1982).
- [4] B. S. DeWitt, “Quantum theory of gravity. I. The canonical theory”, *Phys. Rev.* **160**, 1113–1148 (1967).
- [5] D. Giulini, “3-manifolds in canonical quantum gravity”, PhD Thesis, University of Cambridge (1990).
- [6] A. Hatcher, *Algebraic Topology* (Cambridge University Press, Cambridge, 2002), Proposition 1.40 and Exercise 1.3.24.
- [7] S. W. Hawking and G. F. R. Ellis, *The Large Scale Structure of Space-Time* (Cambridge University Press, Cambridge, 1973).
- [8] K. Krasnov and J. Louko, “ $SO_0(1, d + 1)$ Racah coefficients: Type I representations”, *J. Math. Phys.* **47**, 033513 (2006). (arXiv:math-ph/0502017)
- [9] P. Langlois, “Imprints of spacetime topology in the Hawking-Unruh effect”, PhD Thesis, University of Nottingham (2005). (arXiv: gr-qc/0510127)
- [10] E. Poisson, “The motion of point particles in curved spacetime”, *Living Rev. Relativity* **7** 6 (2004), URL : <http://www.livingreviews.org/lrr-2004-6> (cited on 31 August 2006). (arXiv: gr-qc/0306052)
- [11] E. Poisson, “The gravitational self-force”, in *Proceedings of the 17th International Conference on General Relativity and Gravitation* (Dublin, Ireland, July 18–23, 2004), edited

by P. Florides, B. Nolan and A. Ottewill (World Scientific, Singapore, 2005) 119–141.
(arXiv:gr-qc/0410127)

[12] J. A. Wheeler, “Geons”, *Phys. Rev.* **97**, 511–536 (1955).

[13] J. A. Wolf, *Spaces of Constant Curvature*, 5th edition (Publish or Perish, Wilmington, 1984).

[14] Website <http://www.ligo.caltech.edu/>, visited 14 August 2007.