

The wave function

Exercise 1: Dirac delta function

The Dirac delta function is a distribution defined by

$$\int_{-\infty}^{\infty} dx \delta(x - x_0) f(x) = \int_{x_0-\epsilon}^{x_0+\epsilon} dx \delta(x - x_0) f(x) = f(x_0)$$

for any smooth function in $x = x_0$. It also satisfies

$$\int_{-\infty}^{\infty} dx \delta(x - x_0) = 1 \quad \text{and} \quad \delta(x) = 0, \quad \forall x \neq 0.$$

i) Show that the Dirac delta function can be written as the limit when $L \rightarrow \infty$ or $\epsilon \rightarrow 0^+$ of the following functions:

$$(a) \delta_L(x) = \frac{1}{2\pi} \int_{-L}^L dk e^{ikx} = \frac{\sin Lx}{\pi x}$$

$$(b) \delta_\epsilon(x) = \frac{1}{(2\pi\epsilon^2)^{1/2}} e^{-x^2/2\epsilon^2}$$

$$(c) \delta_\epsilon(x) = \frac{\epsilon/\pi}{x^2 + \epsilon^2}$$

$$(d) \delta_\epsilon(x) = \frac{\theta(x + \frac{\epsilon}{2}) - \theta(x - \frac{\epsilon}{2})}{\epsilon} \quad \text{where } \theta(x) \text{ is the Heaviside step function.}$$

ii) Check the following properties of the Dirac delta function where the prime means derivative *in the sense of distributions*:

$$(a) \delta(x) = \delta(-x)$$

$$(b) \delta'(x) = -\delta'(-x)$$

$$(c) x\delta(x) = 0$$

$$(d) x\delta'(x) = -\delta(x)$$

$$(e) x^2\delta'(x) = 0$$

$$(f) \delta(ax) = \frac{1}{|a|} \delta(x), \quad a \neq 0$$

$$(g) \delta(f(x)) = \sum_{f(x_i)=0} \frac{1}{|f'(x_i)|} \delta(x - x_i)$$

$$(h) f(x)\delta(x - a) = f(a)\delta(x - a)$$

$$(i) \int_{-\infty}^{\infty} dx \delta(x - a) \delta(x - b) = \delta(a - b)$$

$$(j) \int_{-\infty}^{\infty} dx \delta'(x - a) \delta(x - b) = \delta'(b - a)$$

$$(k) \delta'(x) = \frac{i}{2\pi} \int_{-\infty}^{\infty} dk k e^{ikx}$$

$$(l) \delta(x) = \theta'(x) \quad (\text{compare with } i) \quad (d) \text{ of this exercise})$$

Exercise 2: Fourier transform

Let us consider the function $f(x)$ defined in the interval $-\infty < x < \infty$. When we can write $f(x)$ as

$$f(x) = \frac{1}{(2\pi)^{1/2}} \int_{-\infty}^{\infty} dk e^{ikx} g(k)$$

then $g(k)$ is the *Fourier transform* of $f(x)$.

- i) Using a representation of the Dirac delta function provided in the previous exercise, show that we can invert the relation:

$$g(k) = \frac{1}{(2\pi)^{1/2}} \int_{-\infty}^{\infty} dx e^{-ikx} f(x),$$

where now $f(x)$ is the *inverse Fourier transform* of $g(k)$.

- ii) Compute the Fourier transform of the following functions: [solutions]

i) $f(x) = \delta(x - x_0)$ $g(k) = (2\pi)^{-1/2} e^{-ikx_0}$

ii) $f(x) = c_0 e^{ik_0 x}$ $g(k) = (2\pi)^{1/2} c_0 \delta(k - k_0)$

iii) $f(x) = c_0 \cos k_0 x$ $g(k) = (2\pi)^{1/2} c_0 \frac{\delta(k - k_0) + \delta(k + k_0)}{2}$

iv) $f(x) = c_0 \sin k_0 x$ $g(k) = (2\pi)^{1/2} c_0 \frac{\delta(k - k_0) - \delta(k + k_0)}{2i}$

v) $f(x) = c_0 e^{-x^2/2\sigma^2} e^{ik_0 x}$ $g(k) = c_0 \sigma e^{-(k - k_0)^2 \sigma^2 / 2}$

vi) $f(x) = c_0 e^{-a|x|} e^{ik_0 x}$ $g(k) = (2/\pi)^{1/2} c_0 \frac{a}{a^2 + (k - k_0)^2}$

vii) $f(x) = \begin{cases} c_0 e^{ik_0 x}, & |x| < A \\ 0, & |x| > A \end{cases}, \quad A > 0$ $g(k) = (2/\pi)^{1/2} c_0 \frac{\sin(k - k_0)A}{k - k_0}$

viii) $f(x) = c_0 \frac{\sin ax}{x}, \quad a > 0$ $g(k) = \begin{cases} (\pi/2)^{1/2} c_0, & |k| < a \\ 0, & |k| > a \end{cases}$

Exercise 3: Wave functions and time evolution

The Fourier transform takes a wave function in the position representation into the momentum representation and *viceversa*:

$$\langle x|p\rangle = \frac{1}{(2\pi\hbar)^{1/2}} e^{ipx/\hbar} \Rightarrow \begin{cases} \langle x|\psi\rangle \equiv \psi(x) = \frac{1}{(2\pi\hbar)^{1/2}} \int_{-\infty}^{\infty} dp e^{ipx/\hbar} \hat{\psi}(p) \\ \langle p|\psi\rangle \equiv \hat{\psi}(p) = \frac{1}{(2\pi\hbar)^{1/2}} \int_{-\infty}^{\infty} dx e^{-ipx/\hbar} \psi(x). \end{cases}$$

If the wave function has a well defined momentum p_0 , then its transform gives a plane wave:

$$\hat{\psi}(p) = \delta(p - p_0), \quad \psi(x) = \frac{1}{(2\pi\hbar)^{1/2}} e^{ip_0x/\hbar}.$$

Suppose that at time $t = 0$ the state of a 1-dimensional free massive particle is given by the Gaussian wave function

$$\xi(x) = \frac{1}{\pi^{1/4} \sqrt{d}} e^{ik_0x - \frac{x^2}{2d}},$$

where d and k_0 are positive real numbers.

- i) Plot the wave function ξ .
- ii) Normalize ξ to one.
- iii) Write the wave function in the momentum representation.
- iv) Find the expectation values and the uncertainties in \hat{x} and \hat{p} , and show that the product of uncertainties is minimal.
- v) Find the time evolution of the wave function.
- vi) Comment on the limit $d \rightarrow \infty$.

Exercise 4: Wave function, probability and time evolution

A particle of mass m is inside a 1-dimensional box of length a (i.e., an infinite quantum well at $0 \leq x \leq a$). Let us suppose that at $t = 0$ its wave function is

$$\psi(x, t = 0) = \sqrt{\frac{8}{5a}} \left[1 + \cos\left(\frac{\pi x}{a}\right) \right] \sin\left(\frac{\pi x}{a}\right).$$

- i) Solve Schrodinger equation and find the energy eigestates and eigenvalues.
- ii) Find the wave function of the particle at an arbitrary time t .
- iii) Determine the probability to find the particle in the left half of the box at time t .

Exercise 5: Probability density and density current

- i) Prove that a stationary state in a non-degenerate energy level implies a vanishing density current $\vec{j}(x)$.
- ii) Consider the following eigenstate in a 3-dimensional harmonic oscillator of frequency ω and mass M :

$$\psi(x) = \pi^{-3/4} \alpha^{5/2} (x + iy) \exp(-\alpha^2 r^2 / 2), \quad \alpha = (M\omega/\hbar)^{1/2}.$$

Show that its probability density current is not zero. Is its energy level degenerate? Is this state stationary?

Exercise 6: Retarded propagator

- i) Calculate the retarded propagator for evolution of a free particle in one dimension.
- ii) Calculate the evolution of a wave packet with minimum dispersion at $t = 0$, i.e. a plane wave modulated with a Gaussian function,

$$\psi(x, t = 0) = \frac{1}{(\pi a^2)^{1/4}} e^{-x^2/2a^2} e^{ip_0 x/\hbar},$$

where a and p_0 are positive real numbers. What is the dispersion of the wave packet at $t > 0$?

- iii) A particle of mass M moves freely in a straight line, with wave function at $t = 0$ given by the expression above, and an initial uncertainty in position of 1 \AA . At which time is the uncertainty in position equal to 1 mm ? Calculate this time for an electron with mass $m_e = 9 \times 10^{-28} \text{ g}$ and a particle with mass $M = 1 \text{ g}$.

Exercise 7: Test exercise, September 2012 [2.5P]

A free particle of mass m evolving in one dimension is in a state described at $t = 0$ by the wave function $\Psi(x, 0) = A \sin(k_0 x)$.

- i) Find the wave function at a time t .
- ii) What are the possible values for its momentum and what is the probability they will be measured?
- iii) We measure the momentum at $t = t_0$ to be $\hbar k_0$. What is the wave function at $t > t_0$?

Exercise 8: Ehrenfest theorem

Consider a 1-dimensional particle in a region with potential energy $V(x)$.

- i) Find the time evolution of the expectation values of its position and its momentum [Ehrenfest theorem]. Integrate these equations and compare them with the classical equations of motion for a massive particle in a potential $V(x) = -Fx$, where F is a constant.
- ii) Does the center of the wave function follow the classical laws of motion for a generic potential $V(x) = -\lambda x^n$?

Exercise 9: Virial theorem

The *virial theorem* establishes that $2\langle K \rangle_\psi = \langle \vec{x} \cdot \nabla V(\vec{x}) \rangle_\psi$, where K is the kinetic energy of a stationary state ψ with potential energy V . Show that

- i) The theorem applied to a potential $V(x) = kx^n$ implies $2\langle K \rangle = n\langle V \rangle$.
- ii) If $V(\vec{x})$ decreases radially $\forall x \neq 0$, then H has no eigenvectors.
- iii) If $\vec{x} \cdot \nabla V(\vec{x}) \leq -\gamma V(\vec{x})$, $\forall \vec{x} \neq 0$, $0 < \gamma < 2$, then H has no eigenvalues $E > 0$.
- iv) If $V(r) = -ge^{-kr}/r$ (where g and k are positive real numbers), then $\vec{x} \cdot \nabla V(r) \leq -V(r) + gk$ and all the energy eigenvalues verify $E \leq gk$.
- v) Use that the observable $\vec{X} \cdot \vec{P}$ is a constant of motion of an arbitrary stationary state (why?) to prove the virial theorem.