

Exact Solution of a Simple Adsorption Model of De-naturing DNA

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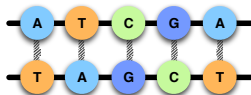
AUSTRALIAN RESEARCH COUNCIL
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ANZAMP meeting, 8th Aust-NZ Convention



DNA

- DNA is a polymer consisting of four repeating nucleic bases A,C,G,T.
- Two strands entwined with a helix structure
- **Denaturation**: At high T , strands pulled apart



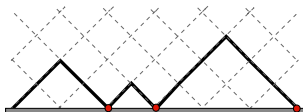
The other physical motivation is the **adsorption** phase transition where polymers in solution can **stick** to the surface of a container

- Second order phase transition
- Order parameter is coverage of the surface by the polymer

We use models in two dimensions and lattice models at that — Proven to be insightful and integrable

ADSORPTION: VERY SIMPLE ONE DIRECTED WALK MODEL

- Single Dyck path in a half space
- Energy $-\varepsilon_a$ for each time (number m_a) it visits the surface
- Boltzmann weight $a = e^{\varepsilon_a/k_B T}$



Consider the coverage, our **order parameter** (indicator for a phase):

$$\mathcal{A}(a) = \lim_{n \rightarrow \infty} \frac{\langle m_a \rangle}{n}$$

There exists two phases, desorbed and adsorbed, and a phase transition at a temperature T_a given by $a = 2$ between these:

- For $T > T_a$ (small a) the walk moves away entropically and $\mathcal{A} = 0$
- For $T < T_a$ (large a) the walk is adsorbed onto the surface and $\mathcal{A} > 0$

ADSORPTION AND UNZIPPING

Previously: *two friendly walks above a sticky wall with single and double interactions* in modelling ring polymers above an absorbing surface interacting with both sides of the ring.

Now we combine adsorption and unzipping:

(Our system — Adsorption and Unzipping)

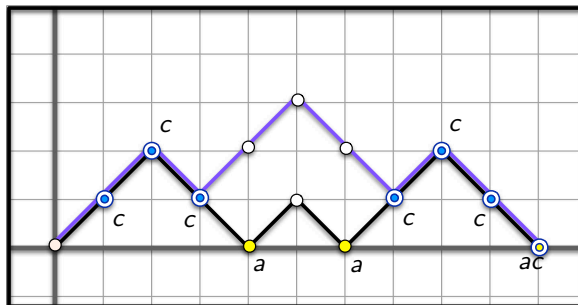
- *double DNA strand in a solvent*
- *near attractive surface*
- *assume aligned base sequence*
- *so expect both adsorption and denaturation (unzipping)*

ALLOWED WALKS

Consider **two directed walks** along the square lattice.

Let our model contain the class of allowed configs. with n steps as described:

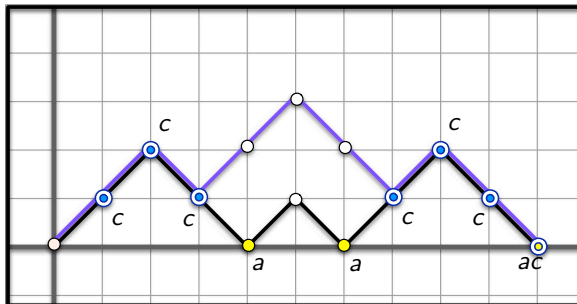
- both walks begin at $(0, 0)$, end at $(2n, 0)$.
- **directed**: can only take steps in the $(\pm 1, 0)$ directions.
- (∞) - **friendly**: walks can share sites, but cannot cross



UNZIPPING ADSORPTION MODEL

Let T be the system temperature, k_B the Boltzmann constant.

- **surface visit sites:** $a \equiv e^{\varepsilon_a/k_B T}$
- **shared site contacts:** $c \equiv e^{\varepsilon_c/k_B T}$
- trivial walk consisting of zero steps has weight 1.



An allowed configuration of length 10. The overall weight is $a^3 c^7$

GENERATING FUNCTION

- **Partition function:** $Z_n(a, c) = \sum_{\hat{\varphi} \ni |\hat{\varphi}|=n} a^{m_a(\hat{\varphi})} c^{m_c(\hat{\varphi})}$
- **Generating function:** $G(a, c) \equiv G(a, c; z) = \sum_{n \geq 1} Z_n(a, c) z^n$
- **Reduced free energy:**

$$\kappa(a, c) = \lim_{n \rightarrow \infty} n^{-1} \log Z_n(a, c) = \log z_s(a, c)$$

where $z_s(a, c)$ is dominant singularity of G w.r.t. z

Two order parameters:

$$\mathcal{A}(a, c) = \lim_{n \rightarrow \infty} \frac{\langle m_a \rangle}{n} \quad \text{and} \quad \mathcal{C}(a, c) = \lim_{n \rightarrow \infty} \frac{\langle m_c \rangle}{n},$$

CONTEXT

No wall/interaction:

- **Vicious dir. walks:** *Lindström-Gessel-Viennot thm. ('85), Essam & Guttmann ('95), Guttmann, Owczarek & Viennot ('98)*
- **Friendly walks & Osculating walks:** *Brak ('97), Guttmann & Vöge ('02), Bousquet-Mélou ('06)*

With wall (but no interaction)

- **Vicious:** *Krattenhaler, Guttmann & Viennot ('00)*

With wall (interaction)

- **Vicious:** *Brak, Essam & Owczarek ('99, '01)*
- **Friendly (two walks):** *Owczarek, Rechnitzer & Wong ('12) - last year's talk!*

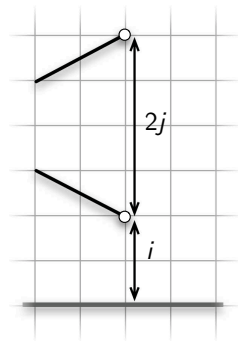
GENERALISED GENERATING FUNCTION

We consider walks φ in the larger set, where each walk can end at any possible height.

- To find $G(a, c)$, consider larger class of configs.
- **Generalised generating function:**

$$F(r, s) \equiv F(r, s, a, c; z) \\ = \sum_{\varphi \in \Omega} a^{m_a(\varphi)} c^{m_c(\varphi)} r^i s^j z^n$$

- $G(a, c) = F(0, 0)$



ESTABLISHING A FUNCTIONAL EQUATION

- By considering the addition of a single column onto a configuration, and the types of walks obtained, we can find a decomposition of all configurations
- Translating back to generating functions we end up with

$$\begin{aligned} K(r, s)F(r, s) &= \frac{1}{ac} + \left(C(c) - \frac{zr}{s}\right) F(r, 0) \\ &\quad + \left[A(a) - \frac{z}{r}(s+1)\right] F(0, s) - A(a)C(c)F(0, 0) \end{aligned}$$

where

$$A(a) = \frac{a-1}{a}, \quad C(c) = \frac{c-1}{c}$$

and the kernel $K(r, s)$ is

$$K(r, s) \equiv K(r, s; z) = \left(1 - z \left[r + \frac{s}{r} + \frac{r}{s} + \frac{1}{r}\right]\right).$$

OBSTINATE KERNEL METHOD

- Equation is written as “bulk = boundary terms” where bulk term is product of kernel and bulk generating function
- Answer needed is one of the boundary generating functions so try to remove bulk by setting the value of a catalytic variable to a value that makes the kernel vanish
- Standard kernel method due to *Knuth* (1968): use values of “catalytic variable” to “kill” kernel
- From \approx early '00's applied to a number of dir. walk problems
- Sometimes need multiple values of catalytic variable(s): **obstinate kernel method**
- More than one catalytic variable requires this
- Earliest combinatorial application of the **obstinate kernel method** due to *Bousquet-Mélou* ('02).
- See Bousquet-Mélou *Math. and Comp. Sci* 2 (2002)), Bousquet-Mélou, Mishna *Contemp. Math.* **520** (2010)

SYMMETRIES OF THE KERNEL

The kernel is symmetric under the following two transformations, which are involutions:

$$(r, s) \mapsto \left(r, \frac{r^2}{s}\right), \quad (r, s) \mapsto \left(\frac{s}{r}, s\right)$$

Transformations generate a family of 8 symmetries ('group of the walk')

$$(r, s), \left(r, \frac{r^2}{s}\right), \left(\frac{s}{r}, \frac{s}{r^2}\right), \left(\frac{r}{s}, \frac{1}{s}\right), \left(\frac{1}{r}, \frac{1}{s}\right), \left(\frac{1}{r}, \frac{s}{r^2}\right), \left(\frac{r}{s}, \frac{r^2}{s}\right), \text{ and } \left(\frac{s}{r}, s\right)$$

- *We make use of four of these which only involve positive powers of r .*
- *This gives us four equations.*
- *One can eliminate many of the unknown generating functions by a clever choice of adding these equations*

ROOTS OF THE KERNEL

- The kernel has two roots as function of either r or s
- choose the one which gives a positive term power series expansion in z
- with Laurent polynomial coefficients in s (r):

$$\hat{r}(s; z) \equiv \hat{r} = \frac{s \left(1 - \sqrt{1 - 4 \frac{(1+s)^2 z^2}{s}} \right)}{2(1+s)z} = \sum_{n \geq 0} C_n \frac{(1+s)^{2n+1} z^{2n+1}}{s^n},$$

where $C_n = \frac{1}{n+1} \binom{2n}{n}$ is a Catalan number.

- *Make the substitution $r \mapsto \hat{r}$*

FINDING THE SOLUTION

Key idea

- Treat K as fn. of r or s to get roots \hat{r} and \hat{s}
- Then use subset of \mathcal{F} to get system of eqns. E.g. Using \hat{r} :

(\hat{r}, s)	$F(\hat{r}, 0)$	$F(0, s)$	$F(0, 0)$
$(\hat{r}, \hat{r}^2/s)$	$F(\hat{r}, 0)$	$F(0, \hat{r}^2/s)$	$F(0, 0)$
$(\hat{r}/s, \hat{r}^2/s)$	$F(\hat{r}/s, 0)$	$F(0, \hat{r}^2/s)$	$F(0, 0)$
$(\hat{r}/s, 1/s)$	$F(\hat{r}/s, 0)$	$F(0, 1/s)$	$F(0, 0)$

- Combine these eqns. to get new fn. eqn

$$N_1^*(s; z)F(0, 1/s) + N_2^*(s; z)F(0, s) = \left[M^*(s) - c^2 H^*(s; z) \right] \left(\frac{1}{ac} - ACF(0, 0) \right),$$

- Can do the same using \hat{s} !
- Nice things happen when $a = 1$ or $c = 1$ to $N_1^*(s; z)$ etc

SOLUTION FOR $G(a, 1)$

Exact solution for $G(a, 1)$ is known and can be found using the method described

- Brak, Essam & Owczarek (1998, 2001): Partition fn. using Lindström-Gessel-Viennot Thm.
- Owczarek, Rechnitzer & Wong (2012): Gen. fn calculated by employing same kernel method techniques.

Specifically:

$$G(a, 1) = \sum_{n \geq 0} z^{2n} \sum_{k=0}^n a^k \frac{k(k+1)(k+2)}{(2n-k)(n+1)^2(n+2)} \binom{2n-k}{n} \binom{2n}{n}.$$

SOLUTION FOR $G(1, c)$

- No known previous solution for $G(1, c)$

We can write functional equation as

$$G(1, c) = F(0, 0, 1, c; z) = [r^1] \frac{\hat{s} (r^2 - 1) [r - cr + cz (1 + r^2 - \hat{s})]}{(c - 1) (\hat{s} - c\hat{s} + crz)},$$

expanding RHS as power series in c and so obtain, after some work:

$$G(1, c; z) = 1 + c^2 z^2 + c^3 (1 + 2z) z^4 + \sum_{i=3}^{\infty} z^{2i} \sum_{m=3}^{2i} c^m \sum_{k=3}^m (-1)^{k+1} \frac{k(k-1)(k-2)(2i-k+1)(i-k+2)}{i^2(i-1)^2(i+1)(i-2)} \binom{m}{k} \binom{2i-k}{i-2} \binom{2i-k-1}{i-3}.$$

SOLUTION FOR $G(1, c)$

- While we have an explicit solution for $G(1, c)$ it is advantageous for analysis to directly read off the singularities
- Alternative — find differential equation satisfied by generating function
- Use Zeilberger-Gosper algorithm: **Maple**: DETools package, Zeilberger hyperexp. implementation
- Result: DE for $G(1, c)$ is order 6 with poly. coeff of $\deg_z = 12$

FORTUNATE DECOMPOSITION OF $G(a, c)$

Using various combinatorial relationships between the generating functions we can re-write $G(a, c)$ in terms of $G(a, 1)$ and $G(1, c)$:

$$G(a, c) = \frac{1}{(a-1)(c-1)} + \frac{p_1(a, c, z)}{p_2(a, c, z) + p_3(a, c, z)G(a, 1) + p_4(a, c, z)G(1, c)}$$

where p_i are polynomials in a, c and z : quadratics in z^2 .

Key point: With solutions to $G(a, 1)$ and $G(1, c)$ we additionally have solved for $G(a, c)$.

SINGULARITIES OF $G(a, 1)$ & $G(1, c)$

- Recall, free energy $\kappa(a, c) = \log z_s(a, c)$
- For $G(a, 1)$, prev. known:

$$z_s(a, 1) = \begin{cases} z_b \equiv 1/4, & a \leq 2 \\ z_a \equiv \frac{\sqrt{a-1}}{2a}, & a > 2 \end{cases}$$

- For $G(1, c)$, we use the DE (roots of leading poly. coeff.):

$$z_s(1, c) = \begin{cases} z_b \equiv 1/4, & c \leq 4/3 \\ z_c \equiv \frac{1-c+\sqrt{c^2-c}}{c}, & c > 4/3 \end{cases}$$

TRANSITIONS OF $G(a, 1)$ & $G(1, c)$

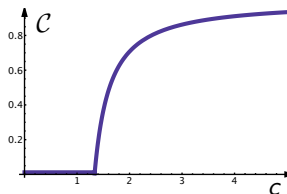
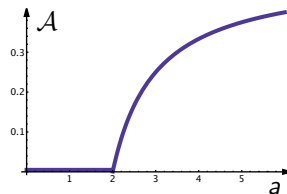
- For $G(a, 1)$: the order parameter associated with the phase transition is the **surface coverage**

$$\mathcal{A}(a, 1) = \begin{cases} 0, & a \leq 2 \\ \frac{a-2}{2(a-1)}, & a > 2 \end{cases}$$

- For $G(1, c)$: the order parameter associated with the phase transition is the **shared site density**

$$\mathcal{C}(1, c) = \begin{cases} 0, & c \leq 4/3 \\ \frac{c-2+\sqrt{c(c-1)}}{2(c-1)}, & c > 4/3 \end{cases}$$

- Second-order** adsorption and zipping phase trans. resp.



SINGULARITIES AND PHASES

This leads us to associate the singularities of $G(a, 1)$ and $G(1, c)$ with the phases as

- $z_b = 1/4$ with a **desorbed** phase where $\mathcal{A} = 0$ and $\mathcal{C} = 0$
- $z_a = \frac{\sqrt{a-1}}{2a}$ with an **adsorbed** phase where $\mathcal{A} > 0$
- $z_c = \frac{1-c+\sqrt{c^2-c}}{c}$ with a **zipped** phase where $\mathcal{C} > 0$

ORDER PARAMETERS FOR THE FULL MODEL

Four possible phases:

- **Desorbed:** $\mathcal{A} = \mathcal{C} = 0$
- **Adsorbed:** (a-rich) $\mathcal{A} > 0, \mathcal{C} = 0$
- **Zippered:** (c-rich) $\mathcal{A} = 0, \mathcal{C} > 0$
- **Zippered & Adsorbed:** (ac-rich) $\mathcal{A} > 0, \mathcal{C} > 0$

ANALYSING $G(a, c)$

Recall

$$G(a, c) \sim \frac{p_1(a, c, z)}{p_2(a, c, z) + p_3(a, c, z)G(a, 1) + p_4(a, c, z)G(1, c)}$$

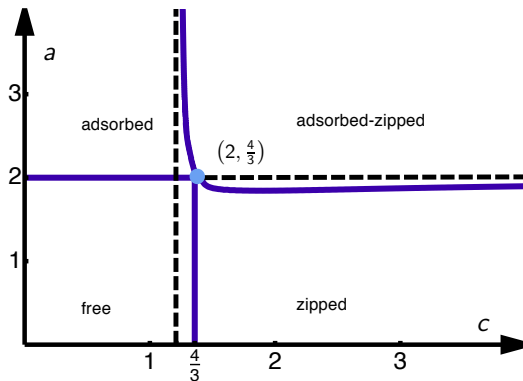
- \Rightarrow **Singularities**: Look at $G(a, 1)$, $G(1, c)$ and root of above denom.
- root of denominator is associated with the **zipped-adsorbed** phase

The dominant singularity $z_s(a, c)$ of the generating function $G(a, c; z)$ is one of four types associated with the four phases

$$z_s(a, c) = \begin{cases} z_b \equiv 1/4, & a \leq 2, c \leq 4/3 \\ z_a(a) \equiv \frac{\sqrt{a-1}}{2a}, & a > 2, c \leq \alpha(a) \\ z_c(c) \equiv \frac{1-c+\sqrt{c^2-c}}{c}, & a \leq \gamma(c), c > 4/3 \\ z_{ac}(a, c), & a > \gamma(c), c > \alpha(a) \end{cases}$$

- $\alpha(a)$ is boundary between **adsorbed** and **zipped-adsorbed** phases
- $\gamma(c)$ is the boundary between **zipped** and **zipped-adsorbed** phases

PHASE DIAGRAM



All transitions found to be second order

Low-temp argument gives

- $c \rightarrow \infty, \gamma(c) \rightarrow 2$
- $a \rightarrow \infty, \alpha(a) \rightarrow \sqrt{5} - 1$

ASYMPTOTICS

Table : The growth rates of the coefficients $Z_n(a, c)$ modulo the amplitudes of the full generating function $G(a, c; z)$ over the entire phase space.

phase region	$Z_n(a, c) \sim$
free	$4^n n^{-5}$
free to adsorbed boundary	$4^n n^{-3}$
free to zipped boundary	$4^n n^{-3}$
$a = 2, c = 4/3$	$4^n n^{-3}$
adsorbed	$z_a(a)^{-n} n^{-3/2}$
zipped	$z_c(c)^{-n} n^{-3/2}$
adsorbed to adsorbed-zipped boundary ($\alpha(a)$)	$z_a(c)^{-n} n^{-1/2}$
zipped to adsorbed-zipped boundary ($\gamma(c)$)	$z_c(c)^{-n} n^{-1/2}$
adsorbed-zipped	$z_{ac}(a, c)^{-n} n^{-1}$

CONCLUSION

- Simple model of DNA as two friendly walks near a boundary
- Used combinatorial decomposition to obtain linear functional equation
- Used **obstinate kernel method** to solve functional equations (using symmetries to provide sufficient information)
- Explicit series solutions for $G(a, 1)$ and $G(1, c)$
- Combined these equations to relate $G(a, c)$ to both $G(a, 1)$ and $G(1, c)$
- Also used **Zeilberger-Gosper** algorithm to find linear DE for $G(1, c)$
- Full analysis of asymptotics and phase diagram
- [R. Tabbara, A. L. Owczarek and A. Rechnitzer, *J. Phys. A.: Math. Theor.*, **47**, 015202 \(34pp\), 2014](#)

FUTURE WORKS

- Consider three walks with multiple unzipping interactions: work in progress – maybe next year's talk!
- Combine single, double surface and unzipping interactions: work in progress
- Also in progress is work in a slit

