## Tisserand Plot

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## 1 Introduction

Tisserand plot can help understand the effects and benefits of (one or multiple) flybys.

This analysis assumes that the secondary body (pedix "s") flies in a circular orbit of radius  $a_s$  about the primary body (pedix "p").

Also, we denotes distance from the primary body with upper-case R, wherease distances in the Sphere of Influence (SOI) of the secondary body with lower-case r.

In this document we will only consider the  $(R_P, R_P)$  Tisserand plot, which is the simplest one to deal with elliptic orbits about the primary body. Note that hyperbolic orbits are not allowed in this plot, since the apocenter radius  $R_A$  does not exist. In that case, other kinds of Tisserand plots could be used such as the  $R_P - \mathcal{E}$  (that is pericenter radius vs period) plot.

First, we show that we can cover the plot by iso- $v_{\infty}$  contour lines. Indeed, each pair  $(R_P, R_A)$  corresponds to a unique pair  $(v_{\infty}, \alpha)$ .

Let  $V_{sc}$  be the spacecraft velocity and V

$$V_{sc}^{2} = V_{s}^{2} + v_{\infty}^{2} + 2V_{s}v_{\infty}\cos\alpha\tag{1}$$

with  $V_s = \sqrt{\mu_p/r_s}$ 

$$a = (R_P + R_A)/2 \tag{2}$$

$$e = 1 - rp/a \tag{3}$$

$$\cos \nu = (p - r_s)/(er_s) \tag{4}$$

$$V_{sc,r} = \sqrt{\frac{\mu_p}{p}} e \sin \nu \tag{5}$$

$$V_{sc,t} = \sqrt{\frac{\mu_p}{p}} (1 + e \cos \nu) \tag{6}$$

$$V_{sc} = \sqrt{V_{sc,r}^2 + V_{sc,t}^2} \tag{7}$$

$$v_{\infty,r} = V_r \tag{8}$$

$$v_{\infty,t} = V_t - V_s \tag{9}$$

$$v_{\infty} = \sqrt{v_{\infty,r}^2 + v_{\infty,t}^2} \tag{10}$$

$$\alpha = \arccos \frac{V_{sc}^2 - V_s^2 - v_\infty^2}{2V_{sc}v_\infty} \tag{11}$$

The inverse transform follows the same reasoning

$$V_{sc}^2 = V_s^2 + v_\infty^2 + 2V_s v_\infty \cos \alpha \tag{12}$$

$$V_{sc,r} = v_{\infty} \sin \alpha \tag{13}$$

$$V_{sc,t} = V_s + v_\infty \cos \alpha \tag{14}$$

$$a = \frac{\mu_p/2}{\frac{V_{sc}^2}{2} - \frac{\mu_p}{R_s}} \tag{15}$$

$$h = R_s V_{sc,t} \tag{16}$$

$$e = \sqrt{1 - \frac{h^2}{\mu a}} \tag{17}$$

$$R_P = a(1 - e) \tag{18}$$

$$R_A = a(1+e) \tag{19}$$

Using the inverse transform it is possible to draw lines (iso-contours) of constant  $v_{\infty}$ , fixing  $v_{\infty} = \bar{v}_{\infty}$  and varying  $\alpha \in [0, \pi]$ .

Iso-contours of alpha can also be plotted as dashed lines. These lines are obtained fixing the values of  $\alpha$  (e.g.,  $\alpha=0,\ 15^\circ,\ 30^\circ,\ ...,\ 180^\circ$ ). and varying  $v_\infty\in[0,v_{\infty_{\max}}]$ .

## FIGURA 1

As an alternative, one can draw lines of constant period (of the spacecraft orbit around the primary body), for a given n:m resonance. These iso-contours appears a straight lines in the  $(R_A,R_P)$ . In fact,  $R_P+R_A=2a$ , with a constant if the period is constat, being  $a=\sqrt[3]{\left(\frac{T}{2\pi}\right)^2\mu_p}$ .

For a resonant orbit, one hase  $T_{sc}: T_s = n: m$ , hence

$$a_{sc} = a_s (\frac{n}{m})^{2/3} \tag{20}$$

Thus, the iso-resonant lines are found for a specific choice of n:m calculating the pairs  $(R_A,\,R_P)$  varying the eccentricity from 0 to 1 (or  $e_{\rm max}$ .

$$R_P = a_s(\frac{n}{m})^{2/3} (1 - e) \tag{21}$$

$$R_A = a_s(\frac{n}{m})^{2/3} (1+e)$$
 (22)

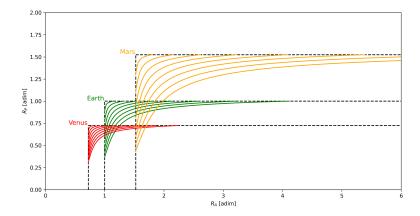


Figure 1: Tisserand plot for the Solar system.

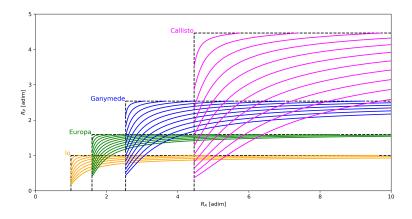


Figure 2: Tisserand plot for the Jovian system.

The effect of a flyby corrsponds to an instantaneous change of velocity. This means a jump on the Tisserand plot. Note that the  $v_{\infty}$  is not altered by a flyby, thus the points before and after the flyby must be located on the same iso- $v_{\infty}$  contour line.

The amount of displacement along the  $v_{\infty}$  contour line depends on the change of the pump angle between the conditions after  $(\alpha_{+})$  and before  $(\alpha_{-})$  the flyby. Their difference is equal to the deflection angle  $\delta$ 

$$\boldsymbol{v}_{\infty-}^T \cdot \boldsymbol{v}_{\infty+} = v_{\infty}^2 \cos \delta \tag{23}$$

figura del triangolo di velocita'

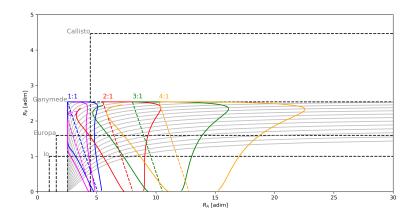


Figure 3: Tisserand plot for the Solar system.

The deflection angle  $\delta$ , can be found as

$$\sin\frac{\delta}{2} = \frac{\mu_s/r_{flyby}}{v_\infty^2 + \mu_s/r_{flyby}} \tag{24}$$

where  $r_{flyby}$  is the pericenter radius of the planetocentric hyperbola, that is, the flyby radius.