

Analytical Dynamics

Analysis of Lagrangian Dynamics

Admin: Mid Semester Survey over break

Last Time: Hamilton's Principle: $\int_{t_1}^{t_2} [ST + \delta W] dt = 0 \Rightarrow \int_{t_1}^{t_2} [\delta f + \sum_k Q_{k,nc} \delta q_k] dt = 0$ (δq sat. constr.)

Principle of Least Action: $\delta \int_{t_1}^{t_2} f dt = 0$ (δq vnc.)
(no non-conservative forces)

Lagrange's Equations: $\frac{d}{dt} \frac{\partial L}{\partial \dot{q}_k} - \frac{\partial L}{\partial q_k} = Q_{k,nc} \quad k=1, \dots, n$ (δq vnc.) [n second order Diff eqs.]

Today: • Analyze structure of Lagrange's equations
• ~~Vibration analysis~~ (next time)

Review: Consider a scalar function $f(q, \dot{q})$:

$$\nabla_q f = \begin{bmatrix} \frac{\partial f}{\partial q_1} \\ \vdots \\ \frac{\partial f}{\partial q_n} \end{bmatrix}$$

$$\nabla_q^2 f = \begin{bmatrix} \frac{\partial^2 f}{\partial q_1^2} & \cdots & \frac{\partial^2 f}{\partial q_1 \partial q_n} \\ \vdots & \ddots & \vdots \\ \frac{\partial^2 f}{\partial q_n \partial q_1} & \cdots & \frac{\partial^2 f}{\partial q_n^2} \end{bmatrix}$$

Fuzziest Point - 10-14-19

- (6) D'Alembert's Principle & the Generalized Inertial Force
- (5) Generalized forces and their physical meaning
- (4) Nonconservative forces vs. conservative forces
- (4) Constraints on generalized coordinates and on virtual displacements

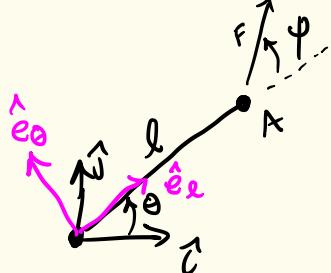
- (2) Virtual Work
- (2) Relationship between Euler-Lagrange equation and generalized forces
- (2) Analyzing virtual displacements with partial derivatives

- (1) Following notation
- (1) Virtual displacements
- (1) Euler Lagrange
- (1) Active vs. nonconservative force
- (1) Constrained Lagrangian
- (1) All the different derivative types (total, partial, variation)
- (1) MATLAB

Generalized Forces: A Case with a Clean Physical Interpretation

Polar Coordinates

$$q = [l, \theta]^T$$



Position & Velocity

$$\underline{r}_A = l c_\theta \hat{i} + l s_\theta \hat{j}$$

$$\begin{aligned}\dot{\underline{r}}_A &= \dot{l} [c_\theta \hat{i} + s_\theta \hat{j}] \\ &\quad + l \dot{\theta} [-s_\theta \hat{i} + c_\theta \hat{j}] \\ &= \dot{l} \hat{e}_r + l \dot{\theta} \hat{e}_\theta\end{aligned}$$

(Infinitesimal displacement over infinitesimal time)

Virtual Displacement

$$\delta \underline{r}_A = \delta l \hat{e}_r + l \delta \theta \hat{e}_\theta$$

(Infinitesimal Displacement with no change in time)

Virtual Work \underline{F}

$$\underline{F} = F(c_\theta \hat{e}_r + s_\theta \hat{e}_\theta)$$

Captures work done when l changes by δl and all other coords are fixed

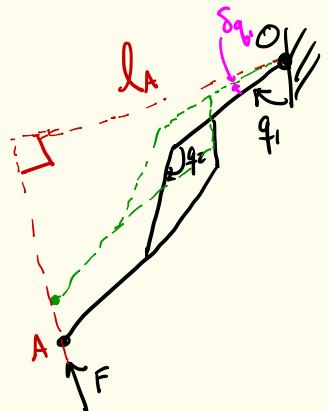
$$\delta W = \underline{F} \cdot \delta \underline{r}_A = Q_\theta \delta l + Q_\theta l \delta \theta$$

Generalized
Forces

$$Q_r = F c_\theta \quad (\text{N}) \quad \text{Force along } \hat{e}_r$$

$$Q_\theta = F l s_\theta \quad (\text{Nm}) \quad \text{Moment about the origin}$$

Generalized Forces: Case with less clear physical interpretation

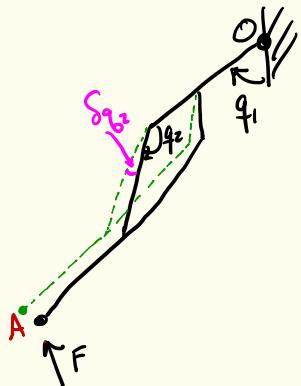


$$\delta W = Q_1 \delta q_1 + Q_2 \delta q_2$$

- Find Q_1 by looking at work done when q_1 changes by δq_1 and all other coords are fixed

$$\delta W = Q_1 \delta q_1 = \underbrace{l_A F}_{\substack{\text{Nm} \\ \text{moment about } O \text{ (Nm)}}} \delta q_1$$

✓ Clear meaning for Q_1



- Find Q_2 by looking at work done when q_2 changes by δq_2 and all other coords are fixed

- $Q_2 \delta q_2$ gives δW when only q_2 is varied

- Mathematically

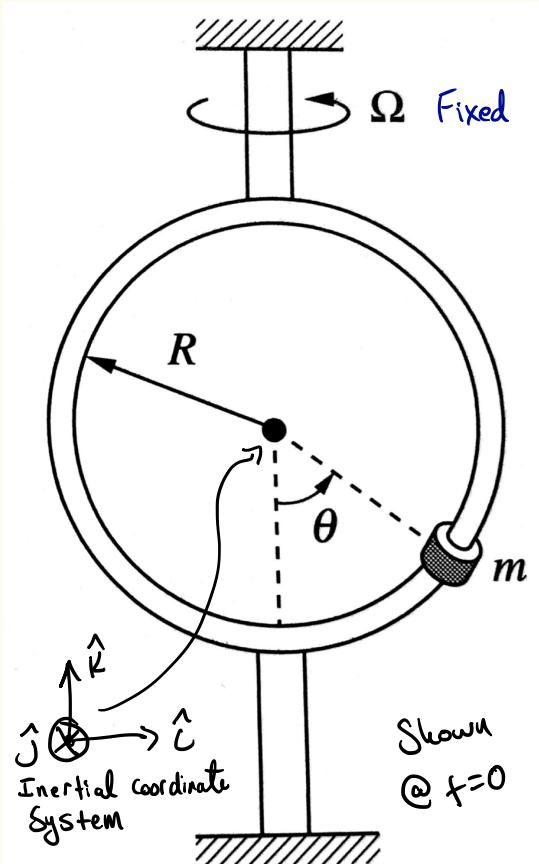
$$Q_2 \delta q_2 = \delta l_A \cdot F = \left(\frac{\partial l_A}{\partial q_2} \delta q_2 \right) \cdot F$$

$$\Rightarrow Q_2 = \frac{\partial l_A}{\partial q_2} \cdot F \quad (\text{units, Nm})$$

No clear form
since kinematics of the four
bar are complex.

X No clear meaning
for Q_2 . It is
still valid mathematically
but harder to think about
physically.

Natural vs. Non-Natural Systems



$$\left\{ \Gamma_{\text{Read}}(\theta, t) \right\} = \begin{bmatrix} (\cos \Omega t) & -\sin \Omega t & 0 \\ \sin \Omega t & \cos \Omega t & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} R \dot{\theta} \\ 0 \\ -R \ddot{\theta} \end{bmatrix}$$

$$= R_3(\Omega t)$$

$$= \begin{bmatrix} R \dot{\theta} \cos(\Omega t) \\ R \dot{\theta} \sin(\Omega t) \\ -R \ddot{\theta} \end{bmatrix}$$

- When some motion is prescribed we say the system is non-natural
- Equivalent to $\Sigma_i(q, t)$ depends on time for some particle in the system.

Structure of the Kinetic Energy

$$q = [q_1, \dots, q_n]^T$$

- System of particles $\Sigma_i(q, t)$

- Velocity of each particle $\sum_i \dot{v}_i = \sum_{k=1}^n \left\{ \frac{\partial r_i}{\partial q_k} \right\} \dot{q}_k + \left\{ \frac{\partial r_i}{\partial t} \right\}$

$$= \sum_i J_i \dot{q} + \sum_i \frac{\partial r_i}{\partial t}$$

$$J_i = \begin{bmatrix} \left\{ \frac{\partial r_i}{\partial q_1} \right\} & \dots & \left\{ \frac{\partial r_i}{\partial q_n} \right\} \end{bmatrix}$$

$3 \times n$
matrix

- $T = \sum_{i=1}^n \frac{1}{2} m_i \sum_i \dot{v}_i^T \sum_i \dot{v}_i = \sum_{i=1}^n \frac{1}{2} m_i \left[\sum_i J_i \dot{q} + \sum_i \frac{\partial r_i}{\partial t} \right]^T \left[\sum_i J_i \dot{q} + \sum_i \frac{\partial r_i}{\partial t} \right]$

$$= \frac{1}{2} \dot{q}^T \left[\sum_{i=1}^n m_i \left\{ \sum_i J_i \dot{q} \right\}^T \sum_i J_i \dot{q} \right] \dot{q}$$

(T_2) quadratic in \dot{q}

$$+ \left[\sum_{i=1}^n m_i \left\{ \frac{\partial r_i}{\partial t} \right\}^T \sum_i J_i \dot{q} \right] \dot{q}$$

(T_1) Linear in \dot{q}

$$+ \frac{1}{2} \sum_{i=1}^n m_i \left\{ \frac{\partial r_i}{\partial t} \right\}^T \sum_i \frac{\partial r_i}{\partial t}$$

(T_0) No gen. vels

0 if a natural system

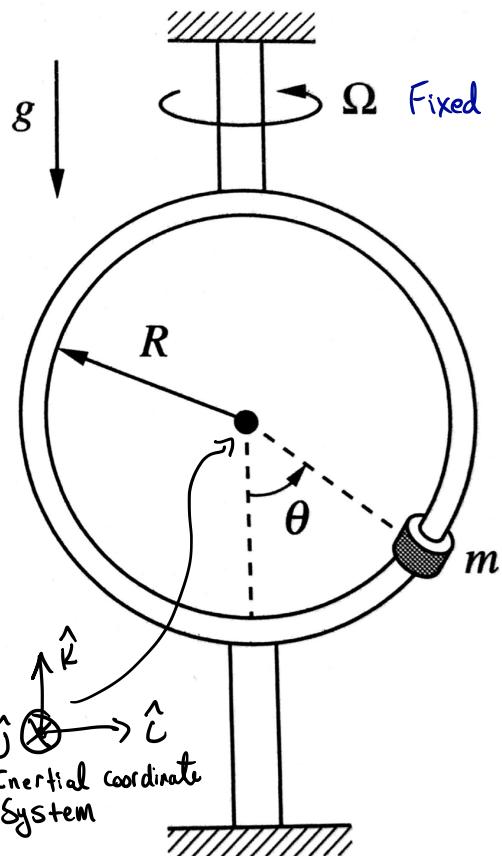
$$T = \frac{1}{2} \dot{q}^T M(q, t) \dot{q} + B(q, t)^T \dot{q} + T_0(q, t)$$

Mass matrix
 $n \times n$ matrix

called the gyroscopic vector
($n \times 1$)

Exercise
 $M = \nabla^2 T$

Kinetic Energy Example: Non-Natural System



$$\begin{aligned} \left\{ \Gamma_{BEAD}(\vec{q}, t) \right\} &= \begin{bmatrix} R\dot{\theta} c(\Omega t) \\ R\dot{\theta} s(\Omega t) \\ -R\ddot{\theta} \end{bmatrix} \\ \left\{ \vec{J}^2 = \frac{\partial \Gamma}{\partial \dot{\theta}} \right\} &= \begin{bmatrix} R\dot{\theta} c(\Omega t) \\ R\dot{\theta} s(\Omega t) \\ R\ddot{\theta} \end{bmatrix} \quad \left\{ \frac{\partial \Gamma}{\partial \ddot{\theta}} \right\} = \begin{bmatrix} -R\dot{\theta} s(\Omega t) \Omega \\ R\dot{\theta} c(\Omega t) \Omega \\ 0 \end{bmatrix} \end{aligned}$$

$$\begin{aligned} T_2 &= \frac{1}{2} \dot{\theta}_m^2 \left\{ \vec{J}^2 \right\}^T \left\{ \vec{J} \right\} = \frac{1}{2} \dot{\theta}_m^2 \left[R\dot{\theta} c(\Omega t)^2 + R\dot{\theta} s(\Omega t)^2 \right] \\ &= \frac{1}{2} \dot{\theta}_m^2 [mR^2] \quad + R^2 S^2 \dot{\theta} \end{aligned}$$

$$T_1 = \textcircled{0}$$

$$\begin{aligned} T_3 &= \frac{1}{2} m \left\{ \frac{\partial \Gamma}{\partial \ddot{\theta}} \right\}^T \left\{ \frac{\partial \Gamma}{\partial \ddot{\theta}} \right\} = \frac{1}{2} m R^2 S_\theta^2 \Omega^2 \\ &= \frac{1}{2} m (R S_\theta \Omega)^2 \end{aligned}$$

$$T = \frac{1}{2} \dot{\theta}_m^2 (mR^2) + \frac{1}{2} m (R S_\theta \Omega)^2$$