CE394M: Critical State and Cam-Clay

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Overview

Critical State Soil Mechanics

Cam-Clay

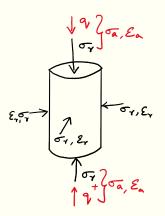
Critical State Soil Mechanics

Roscoe et al., (1958), Schofield & Worth (1968), Wood (1990):

- Provides a conceptual framework in which to interpret stress-strain-strength-volumetric strain response of soil.
- Started as a qualitative, rather than a mathematical model
- A unified framework of known or observed soil responses: drained / undrained / etc

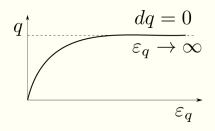
Critical state variables

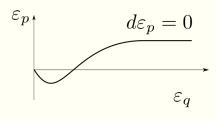
- Mean stress: $p' = \frac{\sigma_a' + 2\sigma_r'}{3} = p u$.
- Deviatoric stress: $q = \sigma'_a \sigma'_r = \sigma_a \sigma_r$
- Specific volume: $v = \frac{V_T}{V_s} = \frac{V_s + V_v}{V_s} = 1 + e$.



Roscoe, Schofield & Worth (1958): At shear-failure, soil exists at a unique state

- $d\varepsilon_q >> 0$ unlimited shear strain potential.
- $dp' = dq = d\varepsilon_p = 0$ no change in p', q, ε_p .
- Critical state stress ratio: $\eta = q/p' = const = M$ at failure q = Mp'.





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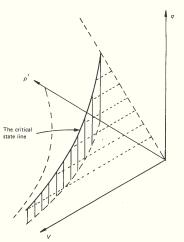
-Critical State Hypothesis: I

Rosson, Scholidd & Worth (1958). All shear-fathers, soil exists at a unique state, with q > 0 uniform dates retries parents. v + iq > 0 uniform dates retries parents. v + iq > 0 uniform dates retries parents. v + iq > 0 uniform dates are soil to v + iq > 0 uniform dates are soil to v + iq > 0 unique and in taken q = hdf. q = 0 eq > 0

Soil is sheared to a point where stresses are stationary (dq = dp' = 0) with no futher change in volume $(d\varepsilon_p = 0)$, unlimited shear strains $(d\varepsilon_q >> 0)$ and q/p' has a fixed value: **critical state**.

M can be related to phi': $M = \frac{6 \sin \phi'}{3 - \sin \phi'}$.

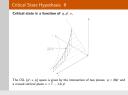
Critical state is a function of q, p', v.



The CSL (p', v, q) space is given by the intersection of two planes: q = Mp' and a cruved vertical plane $v = \Gamma - \lambda \ln p'$

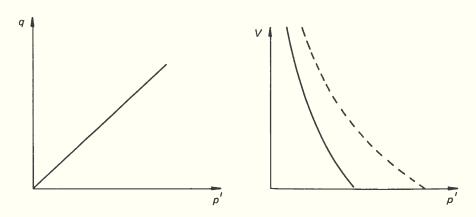
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Critical State Soil Mechanics

-Critical State Hypothesis: II

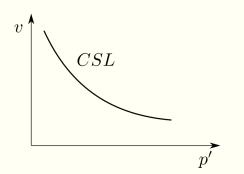


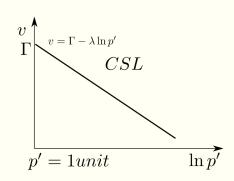
Critical state curve connecting critical state points:

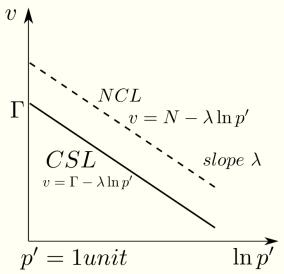
- Crticial state line
- Defined in 3D but we'll look at projections into $q-p^\prime$ and $v-p^\prime$ space



The CSL in (a) (p', q) plot and (b) (p', v) plot (isotropic normal compression line is shown in dashed)







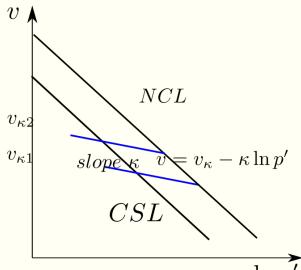
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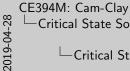
Critical State Soil Mechanics

—Critical State Hypothesis: II



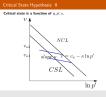
Isotropic virgin compression line (VCL) $\eta=0$. NCL is parallel to CSL. VCL is $\eta=0$, while CSL $\eta=M$. Oedometer falls between VCL and CSL at a constant η : $0<\eta< M$.





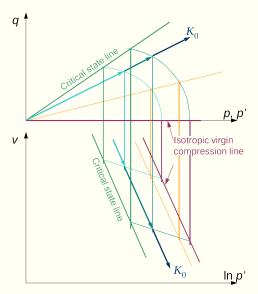
Critical State Soil Mechanics

-Critical State Hypothesis: II

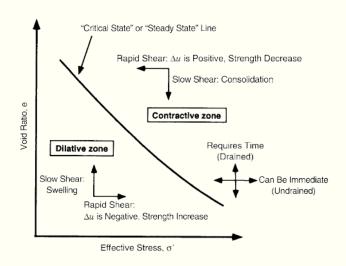


 v_{κ} depends on which κ line you are on. $\kappa \neq c_r$ and $\lambda \neq C_c$

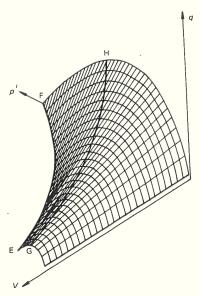
Stress paths $\sigma_3'/\sigma_1' = K_c = const$



Clay behavior



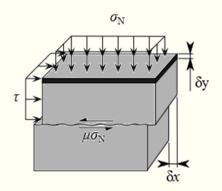
Critical state boundary surface



Summary of critical state behavior

- Can only traverse NCL in one direction
- Can traverse RCL (κ -line) in both directions
- To move from one κ -line to another must move along NCL. Hence, plastic volumetric strains must occur.
- Critical state line is **NOT** a yield surface. It's where it's going but a lot of plastic straining is needed to get there. (if CSL = F = 0) then with associative flow rule $d\varepsilon_p^p \neq 0$ at critical state. Real F is horizontal at critical state.

Stress - dilatancy theory (Taylor, 1948)



Work in friction and dilation:

$$\tau dx - \sigma_n' dy = \mu \sigma_n' dx$$

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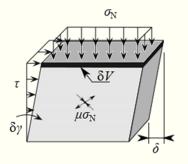
 $^{igspace-}$ Stress - dilatancy theory (Taylor, 1948)



Taylor (1948) proposed a stress-dilatancy theory based on the work balance equation: The external work corresponds to the product of the measured displacements and forces (assuming that the elastic deformation is negligible). The internal work corresponds to the frictional force.

Formulation of elasto-plastic Cam-Clay (OCC): Yield function

Derived from work consideration:



External work: $\delta w_{ext}^p = p' d\varepsilon_p^p + q d\varepsilon_q^p$ Assume that the internal work is dissipated by internal friction only: $\delta w_{int}^p = Mp' d\varepsilon_q^p$

$$\delta w_{\rm ext}^p = p' d\varepsilon_p^p + q d\varepsilon_q^p = Mp' d\varepsilon_q^p = \delta w_{\rm int}^p$$

Formulation of elasto-plastic Cam-Clay (OCC):
Yield function

This dissipation function can be regarded simply as generalisation of Taylor's equation. It should be noted that both Taylor's equation and CamClay dissipation function equation assume that when there is some combination of volumechange (dy or $\partial \varepsilon_v$) and of shear distortion (dx or $\partial \varepsilon_s$) it is the shear strain that determines the dissipation rate. The dilation or volume change is a geometrical consequence of interlocking, and does not appear explicitly in the dissipation function.

Cam-Clay (OCC): Stress dilatancy relation

$$p'd\varepsilon_p^p + qd\varepsilon_q^p = Mp'd\varepsilon_q^p$$

Rearranging the terms (divide by $p'd\varepsilon_a^p$):

$$\frac{d\varepsilon_p^p}{d\varepsilon_q^p} = M - \frac{q}{p'} = M - \eta$$

Where $\eta = q/p'$ is defined as the stress-ratio. This equation is known as the dilatancy expression and expresses the ratio in plastic volumetric and deviatoric components.

$$q/p < M: \quad \frac{darepsilon_p^p}{darepsilon_q^p} > 0 o \quad darepsilon_p^p > 0 \quad ext{Contractive response}$$
 $q/p > M: \quad \frac{darepsilon_p^p}{darepsilon_q^p} > 0 o \quad darepsilon_p^p < 0 \quad ext{Dilative response}$ $q/p = M: \quad darepsilon_p^p = 0 \quad ext{No volume change}$

$$q/p>M: \quad rac{darepsilon_p^p}{darepsilon_q^p}>0 o \quad darepsilon_p^p<0 \quad ext{Dilative response}$$

$$q/p=M: \quad darepsilon_p^p=0 \quad ext{No volume change}$$

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-Cam-Clay (OCC): Stress dilatancy relation

Rearranging the term ($\phi^{2} = \phi^{2} + \phi^{2}$

The critical state is defined by an absence of volume change or, in other words, a nil dilatancy conditions. Therefore, at critical state, the stress-dilatancy rule yields to the critical state condition $\eta = M$.

Cam-Clay (OCC): flow-rule

The original idea was very simple. The yield locus must be such that each associated flow rule $(\delta \varepsilon_p, \delta \varepsilon_q)$ would be orthogonal to the tangent to the yield locus.

$$\frac{\delta\varepsilon_{\textit{p}}}{\delta\varepsilon_{\textit{q}}} = -\frac{\delta\varepsilon_{\textit{q}}}{\delta\varepsilon_{\textit{p}}'}$$

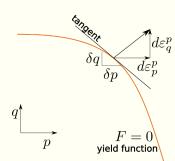
From stress dilation condition:

$$\frac{dq}{dp'} = -(M - \eta) = -M + \eta$$

Integrating we obtain:

$$q = Mp' \ln \left(rac{p_c'}{p'}
ight)$$

Where p'_{c} is the value of p' at q=0.





Original Cam-Clay integration

$$\eta = q/p \quad o d\eta = rac{\partial \eta}{\partial q} dq + rac{\partial \eta}{\partial p'} dp'$$

Which gives:

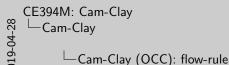
$$d\eta = \frac{dq}{p} - \frac{q}{p^2}dp \rightarrow dq = pd\eta + \eta dp$$

We know from flow rule and orthogonality: $dq = dp(-M + \eta)$ Equating the above 2 equations:

$$dp = pd\eta + \eta dp = dp(-M + \eta)$$
 $pd\eta = -Mdp o deta = -M rac{dp}{p}$

Integrating this expression we obtain:

$$\eta = -M \ln p + C \tag{1}$$



The original idea was very simple. The yield locus must be such that acc associated from rule, $b_1 a_2^2$ vanish we be orthogonal to the tangent to the yield locus. $\frac{\delta x_2}{\delta x_1} = \frac{\delta x_2}{\delta x_2}$ From stress Galation constraints, $\frac{\delta x_2}{\delta x_2} = \frac{\delta x_2}{\delta x_2}$ From stress Galation constraints, $\frac{\delta x_2}{\delta x_2} = -(M-\eta) = M+\eta$ $\frac{\delta x_2}{\delta x_2} = -(M-\eta) = M+\eta$ locations, where $\frac{\delta x_2}{\delta x_2} = -(M-\eta) = \frac{\delta x_2}{\delta x_2}$ between $\frac{\delta x_2}{\delta x_2} = -(M-\eta) = \frac{\delta x_2}{\delta x_2}$. Where x_1^2 is the value of x_1^2 at x_2^2 is $\frac{\delta x_2}{\delta x_2} = \frac{\delta x_2}{\delta x_2} = \frac{\delta x_2}{\delta x_2}$. Where x_1^2 is the value of x_1^2 at x_2^2 is the value of x_1^2 at x_2^2 in the value of x_1^2 at x_2^2 is the value of x_1^2 at x_2^2 in the value of x_2^2 and x_2^2 is the value of x_1^2 at x_2^2 in the value of x_2^2 and x_2^2 is the value of x_2^2 and x_2^2 in the value of x_2^2 and x_2^2 is the value of x_2^2 and x_2^2 a

Original Cam-Clay integration

$$\eta = -M \ln p + C$$

To find the constants, for $\eta = 0$, we get $p = p_c$:

$$0 = -M \ln p_c + C \quad C = M \ln p_c$$

Which gives:

$$\eta = M \ln p_c - M \ln p$$
$$q/p = M \ln (p_c/p)$$

Yield function:

$$F = q - Mp' \ln(p'_c/p') = 0$$

(2)

Cam-Clay (OCC): Elastic properties

Swelling: $\delta v_{\kappa} = \kappa \ln(p_1'/p_2')$

Elastic bulk modulus: $K = \frac{dp'}{d\varepsilon_p}$.

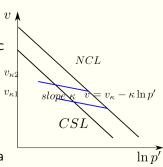
We know the volumetric compression on elastic reloading line:

$$dv = -\kappa \frac{dp'}{p'}$$

$$d\varepsilon_p = \frac{-de}{1 + e_0} = \frac{-dv}{v_0} = \frac{\kappa}{v_0} \frac{dp'}{p'}$$

K' is not constant: K' = K'(p'). Assuming a constant poisson ratio: ν , so G, K vary.

$$K = \frac{dp}{d\varepsilon_p} = \frac{v_o p'}{\kappa} = \frac{(1+e_0)p'}{\kappa}$$





Observation:

- Stiffness *K* increases with *p*': correct.
- Stiffness increases with void ratio (not right)!

Note: The original derivation assumed that there were no recoverable (elastic) shear strains so $G=\infty$. We can find the stress-strain relationships for a single element in this case, but for a finite element forumulation we need to have a finite G^e . So there are two options:

- Define G = f(e, p').
- Use a constant "elastic" Poisson ratio. Ratio between the shear and bulk modulus is constant. 2G/K = const.

The first alternative has the shortcoming that depending on the choice of G we may have unreasonable values of the "elastic" Poisson's ratio. I prefer the second choice.

Cam-Clay (OCC): Hardening law

We need to define how the yield surface hardens as plastic work is being performed. Only "memory" parameter in our yield surface is the size: p'_c . From the isotropic NCL:

$$d\varepsilon_p = \frac{-dv}{v} = \frac{-de}{1+e} = \frac{+\lambda}{v} \frac{dp_c'}{p_c'}$$

But the increment in elastic volumetric strain is:

$$d\varepsilon_{p}^{e} = \left(\frac{-dv}{v}\right)^{elastic} = +\frac{\kappa}{v}\left(\frac{dp_{c}'}{p_{c}'}\right)$$

Therefore the increment of p_c can be related to the increment of plastic volumetric strain:

$$d\varepsilon_p^p = d\varepsilon_p - d\varepsilon_p^e = (\lambda - \kappa) \left(\frac{dp_c'}{p_c'}\right) \rightarrow dp_c' = \left(\frac{v \cdot p_c'}{(\lambda - \kappa)}\right) \cdot d\varepsilon_p^p$$

Cam-Clay (OCC): Hardening law

We have seen that the hardening law:

$$H = -\left(\frac{\partial F}{\partial Wp}\right) \left(\frac{\partial Wp}{\partial \varepsilon^p}\right)^T \cdot \frac{\partial G}{\partial \sigma}$$

 W_p is the vector of memory parameters. In our case, the CC model has only one parameter: p'_c and it's variation is only a function of the plastic volumetric strain. So:

$$H = -\left(\frac{\partial F}{\partial p_c'}\right) \left(\frac{\partial p_c'}{\partial \varepsilon^p}\right)^T \cdot \frac{\partial G}{\partial \sigma}$$

Cam-Clay (OCC): Hardening law

We know:

$$\begin{split} \frac{\partial F}{\partial p'_c} &= -Mp'/p'_c \\ \frac{\partial p'_c}{\partial \varepsilon^p} &= \frac{v}{(\lambda - \kappa)} p'_c \\ \frac{\partial G}{\partial \sigma} &= P_p = Q_p = M - \eta \end{split}$$

$$H = -\left(\frac{\partial F}{\partial p_c'}\right) \left(\frac{\partial p_c'}{\partial \varepsilon^p}\right)^T \cdot \frac{\partial G}{\partial \sigma} = M \frac{(M - \eta)}{(\lambda - \kappa)} \cdot (1 + e) \cdot p'$$