

Abstract In this paper we summarize the procedure of estimating unpolarized pp cross section from beam current data in a transmission experiment. Also, the estimate.

I. TRIC PROGRAM

Tests of T-symmetry have been lacking in precision, due to the absence of null observables to observe. It was proved [1] that no vanishing observable exists for T-symmetry; however, the proof hinges on the bilinearity of observables with respect to probability amplitudes. There is, however, one relationship which, in one unique case, links a bilinear observable, and a linear one.

On the basis of the optical theorem, H.E. Conzett showed [2] that cross section asymmetry in polarized pd scattering does not violate P-symmetry, and does not depend on the reaction dynamics, and thus serves as a legitimate null observable of T-symmetry.

This is the physical foundation of TRIC (Time-Reversal Invariance at COSY).

II. THEORY

A. Physics

The intensity of a particle beam revolving inside an accelerator decreases according to the Beer-Lambert law:

$$I_{n+1} = I_n \cdot \exp \left(- \sum_{i=1}^N \sigma_i \cdot \int_0^L n_i(z) dz \right),$$

where L is the beam path length, N is the number of attenuating species, σ_i is the attenuation cross section, n is the number of passed revolutions.

Simplifying this expression in case of uniform attenuation, one obtains:

$$\begin{aligned} I_{n+1} &= I_n \cdot \exp \left(- \sum_{i=1}^N \sigma_i \cdot \Theta_i \right) \\ &= I_n \cdot \exp \left(- \sum_i \frac{1}{\tau_i} \right), \end{aligned}$$

where $\Theta_i = n_i \cdot \ell_i$ is the thickness of the corresponding attenuating species.

For average beam current, integration of the above yields

$$I_t = I_0 \cdot \exp(\beta \cdot t), \quad (1)$$

with $\beta = \sum_i \beta_i = -\nu \cdot \sum_i 1/\tau_i$, ν — the beam revolution frequency.

Within the confines of the experiment, an unpolarized proton beam interacts with an unpolarized deuterium

target with cross section σ_0 ; to that add all extra-target losses, to produce the following expression for beam loss:

$$\beta = \nu \sigma_0 \Theta + \nu \sigma_x \Theta_x. \quad (2)$$

Since $\sigma_x \Theta_x$ is independent from target state, an estimate of cross section is obtained from

$$\hat{\sigma}_0 = \frac{\hat{\beta}_{off} - \hat{\beta}_{on}}{\nu \Theta_{on}}. \quad (3)$$

B. Statistics

We estimate β by fitting a linear model to logarithmized beam current data: $\ln I_t = \ln I_0 + \beta \cdot t + \epsilon_t$; therefore, the data must satisfy the four Gauss-Markov assumptions in order for the slope estimator to be BLUE:

1. Linearity and additivity of the relationship;
2. Strict exogeneity;
3. No serial correlation of the errors;
4. Homoscedasticity of the errors.

Linearity is crucial for the validity of regression, while endogeneity causes bias in the estimate.

Heteroskedasticity and serial correlation do not cause estimator bias; however, they make it inefficient. This problem can be dealt with by generalized least squares.

III. OVERVIEW OF DATA

The analyzed data came from nine one-hour-long cycles, in four of which the target was turned on (FIG. 1). The experiments of this set were performed with an electron-cooled beam bunched using the barrier bucket.

The measurements were taken with a Beam Current Transformer (BCT). Based on inter-cycle data (baseline), it was concluded that the BCT offset systematically drifted toward zero by approximately 3.5 units per cycle during the cycles of interest.

IV. ESTIMATION OF SLOPE

To make correct slope estimates, we subtract offset Δ from the data. This is done because

$$\tilde{\beta} = \frac{d \ln \tilde{I}_t}{dt} = \frac{1}{\tilde{I}_t} \frac{d \tilde{I}_t}{dt},$$

where, if measured current

$$\tilde{I}_t = I_t + \Delta_t = I_0 \exp(\beta \cdot t) + \Delta_t,$$

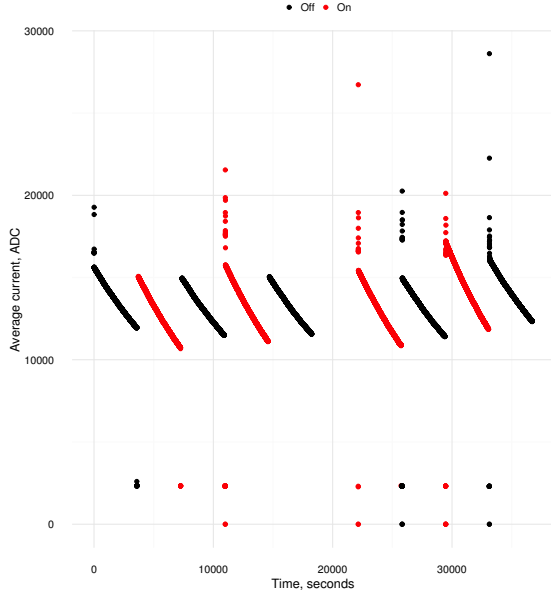


FIG. 1: The nine analyzed cycles.

then

$$\tilde{\beta} = \frac{1}{1 + \lambda_t} \left(\beta + \frac{1}{I_t} \frac{d\Delta_t}{dt} \right),$$

$$\lambda_t = I_0^{-1} \cdot \Delta_t \cdot \exp(-\beta \cdot t).$$

In the favorable case of a constant offset, it still must be removed in order to have a constant slope to estimate. A time-dependent offset violates the exogeneity assumption, and hence biases the slope estimates. At this stage, estimation was done assuming offset was constant within a cycle, and only changed from cycle to cycle.

Since we did not have baseline data for two cycles, their offsets were estimated assuming a linear drift and fitting the cycles' baseline medians (FIG. 2). Each cycle's data were then subtracted its respective offset, and fitted a linear model via ordinary least squares.

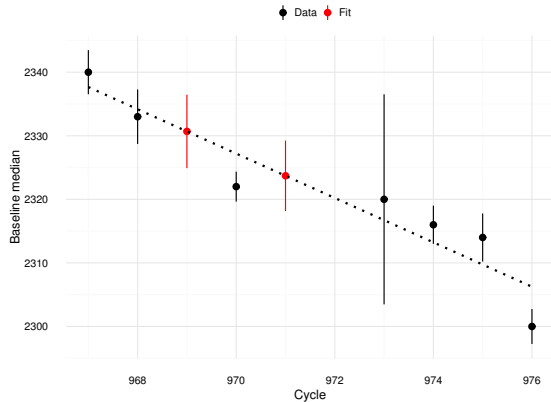


FIG. 2: Offset model.

The models' reduced chi-squares deviate from one in the fourth decimal place; however, one should note that

the data do not pass linearity tests, and are likely to have structural slope changes as well (see TABLE I). Since the model residuals exhibit serial correlation (FIG. 3), the slope estimates' standard errors are estimated with robust estimators.

TABLE I: Characteristics of a typical cycle.

Characteristic	Test	P-value
Linearity	Harvey-Collier	0%
-	Rainbow	0%
Constant slope	Chow ^a	100%
-	Moving estimates	1%
Homoskedasticity	Breusch-Pagan	82%
Autocorrelation	Durbin-Watson	0%

^a The Chow test was performed at every point in the fitting range. The average of F-statistics is used as the test statistic.

V. ESTIMATION OF CROSS SECTION

In estimating cross section, only the estimates from adjacent cycles are used. This is done to minimize the effect of drifts of environmental variables such as target thickness, which is estimated to increase by .5 %/hour. (The thickness by which the slope differences are divided, assumed constant, was provided by a Schottky measurement.)

This reduces the number of estimates from 20 to seven, of which three are deemed unsound according to Tukey's range test (with the range parameter corresponding to ten standard deviations).

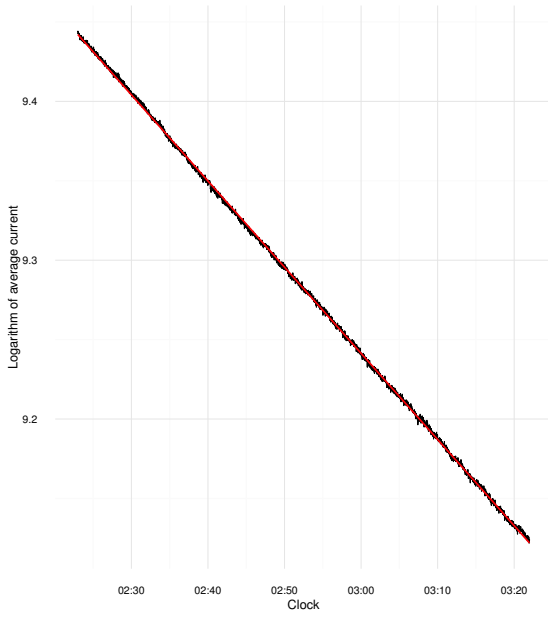
An estimate of a cross section estimate's standard error (SE) is made by adding the squared standard errors of the paired slopes, not taking account of the covariance term:

$$\hat{\sigma}[\hat{\sigma}_0] = \sqrt{\hat{\sigma}[\hat{\beta}_{off}]^2 + \hat{\sigma}[\hat{\beta}_{on}]^2}. \quad (4)$$

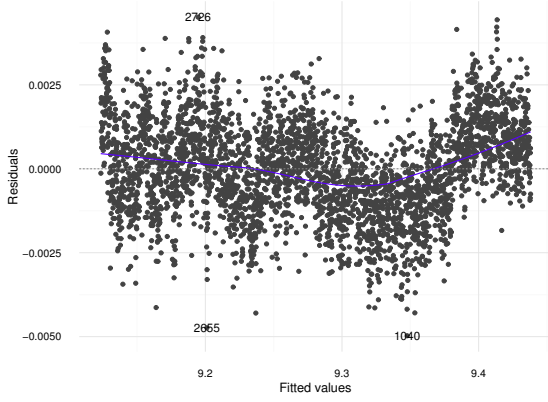
This is done so because depending on whether an on-slope is paired with the preceding or succeeding off-slope, the covariance term changes sign. Since there's no criterion favoring either of the two mappings, the covariance term was omitted.

VI. RESULTS

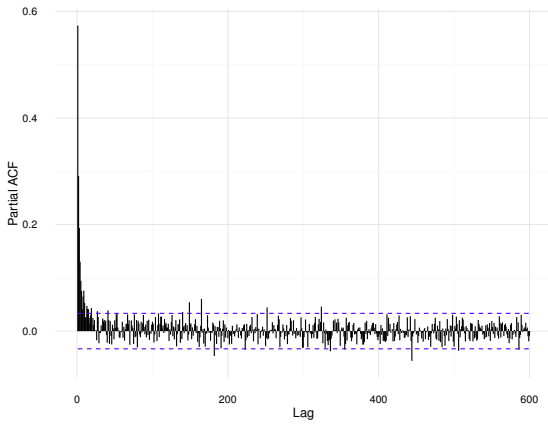
The summary statistics of cross section estimates, grouped by soundness and closeness of the slope estimates they are based on, are presented in TABLE II and FIG. 4a. The slopes themselves are shown in FIG. 4b.



(a) Data and fit.



(b) Residuals vs fitted values.

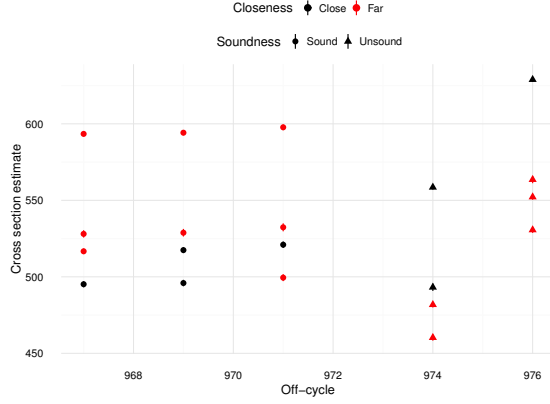


(c) The partial Auto-Correlation Function of the residuals.

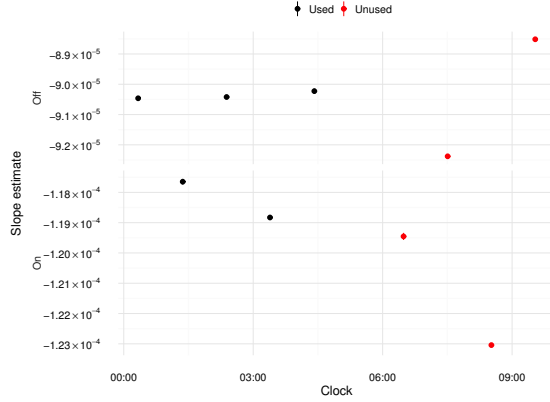
FIG. 3: A typical cycle.

TABLE II: Cross section estimates' statistics

Soundness	Closeness	#	Mean (a.u.)	SE (a.u.)
Sound	Close	4	507	7
Sound	Far	8	549	14
Unsound	Close	3	560	39
Unsound	Far	5	518	20



(a) Cross section estimates.



(b) Slope estimates.

FIG. 4: Results.

[1] Arash, Firooz, Michael J. Moravcsik, and Gary R. Goldstein. Dynamics-Independent Null Experiment for Testing Time-Reversal Invariance. Physical Review Letters 54, no. 25 (1985): 2649. <https://journals.aps.org/prl/pdf/10.1103/PhysRevLett.54.2649>

[2] Homer E. Conzett. On Null Tests of Time-Reversal Invariance, 6. Paris, France, 1990. <https://publications.lbl.gov/islandora/object/ir%3A93728/datastream/PDF/download/citation.pdf>

[3] <http://people.duke.edu/~rnau/testing.htm>