

PHY 831: Statistical Mechanics

Homework 3

Due September 28th, 2020

1. Show that in the canonical ensemble we have

$$C_V = T^{-2} \langle (E - \langle E \rangle)^2 \rangle \quad (1)$$

which shows that the heat capacity is related to the magnitude of fluctuations of the energy of a system about its mean value.

2. Consider a one-dimensional classical gas of distinguishable particles moving in a single-particle potential $U_i = \kappa x_i^2$, so that the energy of the system is $E = \sum_i^N (p_i^2/2m + \frac{m}{2}\omega^2 x_i^2)$ with $\omega = \sqrt{2\kappa/m}$. Assume that the system can move anywhere in the $2N$ -dimensional phase space.

- (a) Calculate the volume of phase space with energy below some energy E , and use this to calculate the number of states with energy at or below E , $\Sigma(E)$, assuming the fiducial phase space volume is h for this one-dimensional system. You may need the result

$$\int_{-\infty}^{\infty} dx_1 \dots \int_{-\infty}^{\infty} dx_M \Theta(R^2 - \sum_{i=1}^M x_i^2) = \frac{\pi^{M/2}}{\Gamma(M/2 + 1)} R^M, \quad (2)$$

which is just the M -dimensional volume of an M -sphere of radius R . Here

$$\Theta(x) = \begin{cases} 1 & \text{for } x \geq 0 \\ 0 & \text{for } x < 0 \end{cases} \quad (3)$$

is the Heaviside step function.

- (b) Calculate the entropy of the gas in the large N limit working in the micro-canonical ensemble.
- (c) Write down the energy of the gas in terms of the temperature using your result from part (b).
- (d) Calculate the Helmholtz free energy, entropy and energy of the gas using the canonical ensemble. These should be the same as the results you found in the microcanonical ensemble.
- (e) Calculate the Helmholtz free energy from the canonical partition function for a system of N one-dimensional quantum harmonic oscillators with single-particle energies $\epsilon_i = \frac{\hbar}{2}\omega(2n_i + \frac{1}{2})$ (with $n_i = \{0, 1, 2, \dots, \infty\}$) and verify that you arrive at the same result as for the classical expression in the low density limit aside from a zero-point energy offset.

3. If the “free volume” \bar{V} of a classical gas is defined by the equation

$$\bar{V}^N = \int d^3r_1 \dots d^3r_N \exp[\beta(\langle U \rangle - U(\vec{q}))]$$

where $\langle U \rangle$ is the average potential energy of the system and $U(\vec{r}_1, \dots, \vec{r}_N) = \sum_{i < j} u(\vec{r}_i - \vec{r}_j)$ is the total potential energy for a particular position of the particle in configuration space, then show that

$$S = k_b N \ln \left[\frac{\bar{V}}{N} \left(\frac{mk_b T}{2\pi\hbar^2} \right)^{3/2} \right] + \frac{5}{2} k_b N$$

In what sense is it justified to refer to the quantity \bar{V} as the free volume? Substantiate your answer by considering a gas of hard spheres (i.e. particles with a two-body potential given by

$$u(r_{ij}) = \begin{cases} 0 & \text{for } r_{ij} > a \\ \infty & \text{for } r_{ij} < a \end{cases}$$

where a is the radius of the spheres and r_{ij} is the distance between particles i and j .)

4. Consider a low-density, relativistic gas of particles moving in one-dimension (i.e. $\epsilon = \sqrt{p^2 + m^2}$). Show that

$$\left\langle \frac{p^2}{\epsilon} \right\rangle = T$$

both by the equipartition theorem and by integration over the phase-space density in the canonical ensemble.