

Notes on curved-sky QE responses etc.

April 30, 2019

Contents

0.1	Semi-analytical QE $N_L^{(0)}$ calculations	1
0.2	QE response calculation	2
0.3	QE responses calculation	3
1	New derivation of generic response	3

JC: Document to be included with the pipeline release after submission of the revised L08.

Lensing and others quadratic estimators used in [?] are all built multiplying in position space spin transforms of spin-weighted fields. We may write all of these in the form

$$_{s_o+t_o}\hat{d}(\hat{n}) \equiv \left(\sum_{\ell m} w_{\ell}^{s_o s_i} \bar{X}_{\ell m} Y_{\ell m}(\hat{n}) \right) \left(\sum_{\ell m} w_{\ell}^{t_o t_i} \bar{X}_{\ell m} Y_{\ell m}(\hat{n}) \right) \quad (1)$$

where s_i, t_i are input spins, s_o, t_o outputs spins, and $w_{\ell}^{s_o s_i}, w_{\ell}^{t_o t_i}$ associated weights. The maps ${}_s\bar{X}_{\ell m}$ are the inverse variance filtered CMB maps,

$${}_0\bar{X}_{\ell m} = -\bar{T}_{\ell m}, \quad {}_{\pm 2}\bar{X}_{\ell m} = -(\bar{E}_{\ell m} \pm i\bar{B}_{\ell m}). \quad (2)$$

For purely analytical calculations, the filtering operation itself can be approximated as isotropic. For independently filtered temperature and polarization, the filtered $\bar{T}, \bar{E}, \bar{B}$ are directly proportional to T, E and B respectively. We keep the discussion focussed on generic fields \bar{X} of arbitrary spins in the following. The gradient (G) and curl (C) modes of definite parity are defined through

$$\begin{aligned} G_{LM}^s &= -\frac{1}{2} (|s| d_{LM} + (-1)^s {}_{-|s|}d_{LM}) \\ C_{LM}^s &= -\frac{1}{2i} (|s| d_{LM} - (-1)^s {}_{-|s|}d_{LM}). \end{aligned}$$

The formulae exposed here can be derived through simple application of this relation,

$$\sum_{m_1, m_2} \int d^2n \prod_{i=1}^3 {}_{s_i}Y_{\ell_i m_i}(\hat{n}) \int d^2n' \prod_{i=1}^3 {}_{t_i}Y_{\ell_i m_i}(\hat{n}') = \frac{2\ell_1+1}{4\pi} \frac{2\ell_2+1}{4\pi} 2\pi \int_{-1}^1 d\mu \prod_{i=1}^3 d_{s_i, t_i}^{\ell_i}(\mu) \quad (3)$$

0.1 Semi-analytical QE $N_L^{(0)}$ calculations

Q.E. noise (co)-variance can be evaluated very easily as was first demonstrated by Ref. []. For two generic estimators as defined in Eq. (1), we can jointly obtain their G and C co-variances with 4 one-dimensional integrals as we now describe.

Let $s = (s_i, s_o, w^{s_i s_o})$ collectively describes the in and out spins and weight function, and similarly for t, u and v . Let the covariance function $N_L^{st, uv}$ be defined through

$$\begin{aligned} \delta_{LL'} \delta_{MM'} N_L^{stuv} &\equiv \left\langle {}_{s_o+t_o}\hat{d}_{LM} {}_{u_o+v_o}\hat{d}_{L'M'}^* \right\rangle \Big|_{\text{Gauss}} \\ &= (-1)^{s_o+t_o+u_o+v_o} 2\pi \int_{-1}^1 d\mu d_{-s_o-t_o, -u_o-v_o}^L(\mu) [\xi^{su}(\mu) \xi^{tv}(\mu) + \xi^{sv}(\mu) \xi^{tu}(\mu)] \end{aligned} \quad (4)$$

where ξ are position-space correlation functions

$$\xi^{st}(\mu) \equiv \sum_{\ell} \left(\frac{2\ell+1}{4\pi} \right) w_{\ell}^{s_o s_i} w_{\ell}^{t_o t_i} \bar{C}_{\ell}^{s_i t_i} d_{s_o, t_o}^{\ell}(\mu) \text{ with } \bar{C}_{\ell}^{s_i t_i} \equiv \langle s_i \bar{X}_{\ell m} t_i \bar{X}_{\ell m}^* \rangle \quad (5)$$

and $d_{mm'}^{\ell}$ are Wigner small d-matrices. Then

$$\boxed{\begin{aligned} \left\langle \hat{G}_{LM}^{s_o+t_o} \hat{G}_{L'M'}^{*, u_o+v_o} \right\rangle \Big|_{\text{Gauss.}} &= \delta_{LL'} \delta_{MM'} \frac{1}{2} [N_L^{stuv} + (-1)^{s_o+t_o} N_L^{-s-tuv}] \\ \left\langle \hat{C}_{LM}^{s_o+t_o} \hat{C}_{L'M'}^{*, u_o+v_o} \right\rangle \Big|_{\text{Gauss.}} &= \delta_{LL'} \delta_{MM'} \frac{1}{2} [N_L^{stuv} - (-1)^{s_o+t_o} N_L^{-s-tuv}] \\ \left\langle \hat{G}_{LM}^{s_o+t_o} \hat{C}_{L'M'}^{*, u_o+v_o} \right\rangle \Big|_{\text{Gauss.}} &= 0 \end{aligned}} \quad (6)$$

0.2 QE response calculation

Let the estimator respond to a source with spin-weight $r \geq 0$ as follows

$${}_s \delta X(\hat{n}) = {}_r \alpha(\hat{n}) \left(\sum_{\ell m} R_{\ell}^{rs} {}_s X_{\ell m} {}_{s-r} Y_{\ell m}(\hat{n}) \right) + {}_{-r} \alpha(\hat{n}) \left(\sum_{\ell m} R_{\ell}^{-rs} {}_s X_{\ell m} {}_{s+r} Y_{\ell m}(\hat{n}) \right) \quad (7)$$

for some harmonic space responses R_{ℓ}^{rs} . Let further the isotropic limit of the filtering procedure be the matrix F , defined through

$${}_s \bar{X}_{\ell m} = \sum_{s_2=0,2,-2} F_{\ell}^{ss_2} {}_{s_2} X_{\ell m} \quad (\text{isotropic approximation}). \quad (8)$$

Then the gradient and curl responses of estimator (1) are

$$\boxed{\begin{aligned} \mathcal{R}_L^{gg} &= R_L^{st,r} + (-1)^r R_L^{st,-r} \\ \mathcal{R}_L^{cc} &= R_L^{st,r} - (-1)^r R_L^{st,-r} \\ \mathcal{R}_L^{gc} &= 0 = \mathcal{R}_L^{cg}, \end{aligned}} \quad (9)$$

where $R_L^{st,r}$ is

$$R_L^{st,r} = (-1)^{s_o+t_o} 2\pi \int_{-1}^1 d\mu d_{-s_o-t_o, -r}^L(\mu) \sum_{\tilde{s}_i, \tilde{t}_i=0,2,-2} \left[\xi^{s_o s_i \tilde{s}_i}(\mu) \psi^{t_o t_i \tilde{t}_i \tilde{s}_i, r}(\mu) + \xi^{t_o t_i \tilde{t}_i}(\mu) \psi^{s_o s_i \tilde{s}_i \tilde{t}_i, r}(\mu) \right] \quad (10)$$

with

$$\begin{aligned} \xi^{s_o s_i \tilde{s}_i}(\mu) &\equiv \sum_{\ell} \left(\frac{2\ell+1}{4\pi} \right) w_{\ell}^{s_o s_i} F_{\ell}^{s_i \tilde{s}_i} d_{s_o, \tilde{s}_i}^{\ell}(\mu) \\ \psi^{s_o s_i \tilde{s}_i \tilde{t}_i, r}(\mu) &\equiv \sum_{\ell} \left(\frac{2\ell+1}{4\pi} \right) w_{\ell}^{s_o s_i} F_{\ell}^{s_i \tilde{s}_i} R_{\ell}^{r t_i} C_{\ell}^{\tilde{s}_i - \tilde{t}_i} d_{s_o, -\tilde{t}_i+r}^{\ell}(\mu) \end{aligned} \quad (11)$$

0.3 QE responses calculation

JC: should define the response at the covariance matrix level... The covariance matrices is

$$\langle {}_s X(\hat{n}) {}_t X^*(\hat{n}') \rangle = \langle {}_s X(\hat{n}) {}_{-t} X(\hat{n}') \rangle = \sum_{\ell m} C_\ell^{st} {}_s Y_{\ell m}(\hat{n}) {}_t Y_{\ell m}^*(\hat{n}') \quad (12)$$

How does this responds to a source of anisotropy (with spin r), ${}_r \alpha(\hat{n})$? For all cases in this work, we can parametrize this as follows

$$\delta \langle {}_s X(\hat{n}) {}_t X(\hat{n}') \rangle = {}_r \alpha(\hat{n}') \sum_{\ell m} {}_r R_\ell^{st} {}_s Y_{\ell m}(\hat{n}) {}_{r-t} Y_{\ell m}^*(\hat{n}') + ((-t, \hat{n}') \leftrightarrow (s, \hat{n})) \quad (13)$$

for some set of isotropic response functions R_ℓ . What is the response to the estimator Eq. (1)?

Examples include:

- Lensing: The source of anisotropy is the spin-1 field ${}_1 \alpha(\hat{n})$, with response

$$\delta_s X(\hat{n}) = \frac{1}{2} \alpha_1 \bar{\partial}_s X(\hat{n}) + \frac{1}{2} \alpha_{-1} \bar{\partial}_s X(\hat{n}) \quad (14)$$

where $\bar{\partial}$ and $\bar{\partial}$ are the spin lowering and spin raising operator **JC: check notation** respectively. From their action on the spin-weighted harmonics, the harmonic space responses are **JC: ...**

$$s_r = s - 1, s + 1 \quad R^{s, s-1} = \text{JC} : \dots, R^{s, s+1} = \text{JC} : \dots \quad (15)$$

- Modulation estimator: The source is spin 0, with response

$$\delta_s X(\hat{n}) = {}_0 \alpha(\hat{n}) {}_s X(\hat{n}) \quad (16)$$

Hence,

$$s_r = s, R_\ell^{ss} = 1 \quad (17)$$

- Point source:

$$\delta_s X(\hat{n}) = {}_0 \alpha(\hat{n}) \delta^D(\hat{n}) \quad (18)$$

- Noise anisotropies:

intermediate steps for response calc. Then $(R_\ell^{-r, -s} = (-1)^r R_\ell^{r, s})$

$${}_{s_o+t_o} \hat{d}_{LM} = \mathcal{R}_L^{st, r} {}_{-r} \alpha_{LM} + \mathcal{R}_L^{ts, -r} {}_r \alpha_{LM} \text{ with } \mathcal{R}_L^{st, r} = 2\pi \int_{-1}^1 d\mu \Xi^{st, r}(\mu) \Xi^t(\mu) d_{r, s_o+t_o}^L(\mu) \quad (19)$$

where

$$\Xi^{st, r}(\mu) = \sum_l \left(\frac{2\ell+1}{4\pi} \right) C_\ell^{s_i, -t_i} R_\ell^{r-t_i} w_\ell^{s_i, s_o} d_{t_i+r, s_o}^\ell(\mu) \text{ and } \Xi^t(\mu) = \sum_l \left(\frac{2\ell+1}{4\pi} \right) w_\ell^{t_i, t_o} d_{-t_i, t_o}^\ell(\mu) \quad (20)$$

1 New derivation of generic response

Let

$${}_s X(\hat{n}') = {}_a \alpha(\hat{n}') \left(\sum_{\ell'' m''} R_{\ell''}^{a, s} X_{\ell'' m''} {}_{s-a} Y_{\ell'' m''}(\hat{n}') \right) \quad (21)$$

where a sum over spin a is implicit. Let further the spin-weight spectra $C_\ell^{s_1 s_2}$ be defined as

$$\left\langle {}_{s_1} X_{\ell m} {}_{s_2} X_{\ell' m'}^\dagger \right\rangle \equiv C_\ell^{s_1 s_2} \delta_{\ell, \ell'} \delta_{m, m'} \quad (22)$$

Let further be the filtering $\mathcal{B}^\dagger \text{Cov}^{-1}$ defined by a matrix

$${}_s \bar{X}_{\ell m} = F_\ell^{ss'} {}_{s'} X_{\ell m} \quad (23)$$

and the QE be (**JC: NB: not exactly as above**)

$${}_{s_o+t_o} \hat{d}(\hat{n}) \equiv \left(\sum_{\ell m} w_\ell^{s_o s_i} {}_{s_i} \bar{X}_{\ell m} {}_{s_o} Y_{\ell m}(\hat{n}) \right) \left(\sum_{\ell' m'} w_{\ell'}^{t_o t_i} {}_{t_i} \bar{X}_{\ell' m'} {}_{t_o} Y_{\ell' m'}(\hat{n}) \right). \quad (24)$$

Then, with $A \equiv wF$,

$$\begin{aligned} {}_{s_o+t_o} \hat{d}_{LM} &= {}_a \alpha_{L'M'} R_{\ell'}^{a, s_i} A_\ell^{s_o s_i} A_{\ell'}^{t_o t_i} \langle {}_{s_i} X_{\ell'' m''} {}_{t_i} X_{\ell' m'} \rangle \int d^2 n {}_{s_o} Y_{\ell m} {}_{t_o} Y_{\ell' m'} {}_{s_o+t_o} Y_{LM}^* \int d^2 n' {}_{s_i} Y_{\ell m}^* {}_{s_i-a} Y_{\ell'' m''} {}_a Y_{L'M'} \\ &+ {}_a \alpha_{L'M'} R_{\ell'}^{a, t_i} A_\ell^{s_o s_i} A_{\ell'}^{t_o t_i} \langle {}_{t_i} X_{\ell'' m''} {}_{s_i} X_{\ell m} \rangle \int d^2 n {}_{s_o} Y_{\ell m} {}_{t_o} Y_{\ell' m'} {}_{s_o+t_o} Y_{LM}^* \int d^2 n' {}_{t_i-a} Y_{\ell'' m''} {}_{t_i} Y_{\ell' m'}^* {}_a Y_{L'M'} \\ &= {}_a \alpha_{L'M'} R_{\ell'}^{a, s_i} A_\ell^{s_o s_i} A_{\ell'}^{t_o t_i} (-1)^{s_i-m'} C_{\ell'}^{t_i, -s_i} \int d^2 n {}_{s_o} Y_{\ell m} {}_{t_o} Y_{\ell' m'} {}_{s_o+t_o} Y_{LM}^* \int d^2 n' {}_{s_i} Y_{\ell m}^* {}_{s_i-a} Y_{\ell', -m'} {}_a Y_{L'M'} \\ &+ {}_a \alpha_{L'M'} R_{\ell'}^{a, t_i} A_\ell^{s_o s_i} A_{\ell'}^{t_o t_i} (-1)^{t_i-m} C_\ell^{s_i, -t_i} \int d^2 n {}_{s_o} Y_{\ell m} {}_{t_o} Y_{\ell' m'} {}_{s_o+t_o} Y_{LM}^* \int d^2 n' {}_{t_i-a} Y_{\ell, -m} {}_{t_i} Y_{\ell' m'}^* {}_a Y_{L'M'} \\ &= {}_a \alpha_{L'M'} R_{\ell'}^{a, s_i} A_\ell^{s_o s_i} A_{\ell'}^{t_o t_i} (-1)^a C_{\ell'}^{t_i, -s_i} \int d^2 n {}_{s_o} Y_{\ell m} {}_{t_o} Y_{\ell' m'} {}_{s_o+t_o} Y_{LM}^* \int d^2 n' {}_{s_i} Y_{\ell m}^* {}_{-s_i+a} Y_{\ell' m'}^* {}_a Y_{L'M'} \\ &+ {}_a \alpha_{L'M'} R_\ell^{a, t_i} A_\ell^{s_o s_i} A_{\ell'}^{t_o t_i} (-1)^a C_\ell^{s_i, -t_i} \int d^2 n {}_{s_o} Y_{\ell m} {}_{t_o} Y_{\ell' m'} {}_{s_o+t_o} Y_{LM}^* \int d^2 n' {}_{-t_i+a} Y_{\ell m}^* {}_{t_i} Y_{\ell' m'}^* {}_a Y_{L'M'} \\ &= (-1)^{M-M'} {}_a \alpha_{L'M'} R_{\ell'}^{a, s_i} A_\ell^{s_o s_i} A_{\ell'}^{t_o t_i} (-1)^{s_o+t_o} C_{\ell'}^{t_i, -s_i} \int d^2 n {}_{s_o} Y_{\ell m} {}_{t_o} Y_{\ell' m'} {}_{-s_o-t_o} Y_{L-M} \int d^2 n' {}_{s_i} Y_{\ell m}^* {}_{-s_i+a} Y_{\ell' m'}^* {}_{-a} Y_{L'-M} \\ &+ (-1)^{M-M'} {}_a \alpha_{L'M'} R_\ell^{a, t_i} A_\ell^{s_o s_i} A_{\ell'}^{t_o t_i} (-1)^{s_o+t_o} C_\ell^{s_i, -t_i} \int d^2 n {}_{s_o} Y_{\ell m} {}_{t_o} Y_{\ell' m'} {}_{-s_o-t_o} Y_{L-M} \int d^2 n' {}_{-t_i+a} Y_{\ell m}^* {}_{t_i} Y_{\ell' m'}^* {}_{-a} Y_{L'-M} \\ &= \delta_{LL, MM'} {}_a \alpha_{L'M'} R_{\ell'}^{a, s_i} A_\ell^{s_o s_i} A_{\ell'}^{t_o t_i} (-1)^{s_o+t_o} C_{\ell'}^{t_i, -s_i} \left(\frac{2\ell+1}{4\pi} \right) \left(\frac{2\ell'+1}{4\pi} \right) 2\pi \int_{-1}^1 d\mu d_{s_o s_i}^\ell(\mu) d_{t_o, -s_i+a}^{\ell'}(\mu) d_{-s_o-t_o, -a}^L(\mu) \\ &+ \delta_{LL, MM'} {}_a \alpha_{L'M'} R_\ell^{a, t_i} A_\ell^{s_o s_i} A_{\ell'}^{t_o t_i} (-1)^{s_o+t_o} C_\ell^{s_i, -t_i} \left(\frac{2\ell+1}{4\pi} \right) \left(\frac{2\ell'+1}{4\pi} \right) 2\pi \int_{-1}^1 d\mu d_{s_o, -t_i+a}^\ell(\mu) d_{t_o, t_i}^{\ell'}(\mu) d_{-s_o-t_o, -a}^L(\mu) \end{aligned} \quad (25)$$