Transparent Boundary Conditions for the Equation of Rod Transverse Vibrations and Compact Approximation Scheme

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Abstract

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1. Introduction

$$\rho \frac{\partial^2 u}{\partial t^2} - R^2 \rho \frac{\partial^4 u}{\partial x^2 \partial t^2} + E R^2 \frac{\partial^4 u}{\partial x^4} = 0.$$
 (1)

2. Implicit Compact Scheme

$$d(u_{m-2}^n + u_{m+2}^n) + a(u_{m-1}^{n-1} + u_{m+1}^{n-1} + u_{m+1}^{n+1} + u_{m+1}^{n+1}) + c(u_{m-1}^n + u_{m+1}^n) + bu_m^n + u_m^{n-1} + u_m^{n+1} = 0,$$
 (2)

where $a = \frac{3}{12\mu+4} - \frac{1}{2}$, $b = \frac{9\nu}{3\mu+1} - 2$, $c = 1 - \frac{12\nu+3}{6\mu+2}$, $d = \frac{3\nu}{6\mu+2}$, and the dimensionless parameters $v = \frac{ER^2}{\rho} \cdot \tau^2 h^{-4}$, $\mu = R^2 \cdot h^{-2}$.

3. DTBCs for Compact Approximation Scheme

After Z-transform of (2) the homogeneous characteristic equation is

$$dz\lambda^4 + (a(z^2+1) + cz)\lambda^3 + (z^2+1 + bz)\lambda^2 + (a(z^2+1) + cz)\lambda + dz = 0.$$
 (3)

We rearrange terms in (3) and divide both parts by $\lambda^2 \neq 0$ and $z^2 \neq 0$ to get

$$dz^{-1}(\lambda + \lambda^{-1})^2 + (a(1+z^{-2}) + cz^{-1})(\lambda + \lambda^{-1}) + 1 + z^{-2} + bz^{-1} - 2dz^{-1} = 0.$$
(4)

Finally, we introduce change of variables $n = \lambda + \lambda^{-1}$ and $\omega = z^{-1}$ to get

$$d\omega\eta^2 + (a(1+\omega^2) + c\omega)\eta + \omega^2 + (b-2d)\omega + 1 = 0.$$
 (5)

The roots of (5) are

$$\eta_{1,2}(\omega) = -\frac{a(1+\omega^2) + c\omega}{2d\omega} \pm \frac{1}{2} \sqrt{\left(\frac{a(1+\omega^2) + c\omega}{d\omega}\right)^2 - 4 \cdot \frac{\omega^2 + (b-2d)\omega + 1}{d\omega}}.$$
 (6)

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3.1. Eta series

$$\eta_{1,2}(\omega) = \frac{(6\mu - 1)\omega^2 - 2(-12\nu + 6\mu - 1)\omega + 6\mu - 1}{12\nu\omega} \pm \frac{(6\mu - 1)(1 - \omega)\sqrt{\omega^2 - 2\left(1 + \frac{72\nu}{(6\mu - 1)^2}\right)\omega + 1}}{12\nu\omega}$$
(7)

$$\eta_{1,2}(\omega) = \frac{1}{12\nu\omega} \left[(6\mu - 1)\omega^2 - 2(-12\nu + 6\mu - 1)\omega + 6\mu - 1 \pm (6\mu - 1)(1 - \omega)\sqrt{\omega^2 - 2\left(1 + \frac{72\nu}{(6\mu - 1)^2}\right)\omega + 1} \right]$$

Now we use the following formula for the Legendre polynomials generating function:

$$(\omega^2 - 2\varepsilon\omega + 1)^{-1/2} = \sum_{n=0}^{\infty} P_n(\varepsilon)\omega^n$$

and obtain for $\mu > 1/6$

$$\eta_{1,2}(\omega) = \frac{6\mu - 1}{12\nu\omega} \left[\omega^2 - 2\frac{-12\nu + 6\mu - 1}{6\mu - 1}\omega + 1 \pm (1 - \omega) \left(\omega^2 - 2\left(1 + \frac{72\nu}{(6\mu - 1)^2}\right)\omega + 1\right) \sum_{n=0}^{\infty} P_n \left(1 + \frac{72\nu}{(6\mu - 1)^2}\right)\omega^n \right]$$
(8)

3.2. Eta branching

Consider the points ω were functions $\eta_{1,2}(\omega)$ have branching. From Eq. (7), this happens when

$$(6\mu - 1)(1 - \omega)\sqrt{\omega^2 - 2\left(1 + \frac{72\nu}{(6\mu - 1)^2}\right)\omega + 1} = 0,$$
(9)

i.e.

$$\omega_1 = 1$$
, $\omega_{2,3} = 1 + \frac{72\nu}{(6\mu - 1)^2} \pm \sqrt{\left(1 + \frac{72\nu}{(6\mu - 1)^2}\right)^2 - 1}$.

Expanding the argument of the square root as difference of a squares, obtain

$$\omega_{2,3} = 1 + \frac{72\nu}{(6\mu - 1)^2} \pm \sqrt{\frac{72\nu}{(6\mu - 1)^2} \left(2 + \frac{72\nu}{(6\mu - 1)^2}\right)}.$$

Therefore, we obtain three branching points:

$$\omega_1 = 1$$
, $\omega_{2,3} = 1 + \frac{72\nu}{(6\mu - 1)^2} \pm \frac{1}{(6\mu - 1)^2} \sqrt{72\nu(72\mu^2 - 24\mu + 2 + 72\nu)}$, (10)

Note that from Eq. (9) according to Vieta's theorem the product $\omega_2 \cdot \omega_3 = 1$, $\omega_2 > 1$ (ω_2 is taken from Eq. (10) with positive sign), and both points are real valued. Therefore, point $0 < \omega_3 < 1$ belongs to the unit circle on the complex plain.

3.3. Lambda series

$$\lambda_{1,3} = \frac{\eta_1(\omega)}{2} \mp \sqrt{\frac{\eta_1^2(\omega)}{4} - 1}$$
 (11)

$$\lambda_{2,4} = \frac{\eta_2(\omega)}{2} \mp \sqrt{\frac{\eta_2^2(\omega)}{4} - 1} \tag{12}$$

$$\lambda_{1,3} = \frac{\eta_1(\omega)}{2} \mp \sqrt{\frac{\eta_1(\omega)}{2} - 1} \cdot \sqrt{\frac{\eta_1(\omega)}{2} + 1}$$

As mentioned before, we have $\eta_1(\omega) \to -\frac{1}{a} + r(\omega)$, where $r(\omega) \to 0$ as $\omega \to 0$. Therefore,

$$\lambda_{1,3} = \frac{\eta_1(\omega)}{2} \mp \sqrt{-\frac{1}{2a} + \frac{r(\omega)}{2} - 1} \cdot \sqrt{-\frac{1}{2a} + \frac{r(\omega)}{2} + 1}$$

$$\lambda_{1,3} = \frac{\eta_1(\omega)}{2} \mp \sqrt{-\frac{1}{2a} - 1} \cdot \sqrt{1 + \frac{r(\omega)}{-\frac{1}{a} - 2}} \cdot \sqrt{-\frac{1}{2a} + 1} \cdot \sqrt{1 + \frac{r(\omega)}{-\frac{1}{a} + 2}}$$

$$\lambda_{1,3} = \frac{\eta_1(\omega)}{2} \mp \sqrt{\frac{1}{4a^2} - 1} \cdot \sqrt{1 - \frac{r(\omega)}{\frac{1}{a} + 2}} \cdot \sqrt{1 + \frac{r(\omega)}{-\frac{1}{a} + 2}}$$

Then we use the Maclaurin series expansion of as square root

$$\sqrt{1+x} = \sum_{n=0}^{\infty} \frac{(-1)^n (2n)!}{(1-2n) \, n! \, 4^n} x^n, \quad |x| < 1$$

to obtain

$$\lambda_{1,3} = \frac{\eta_1(\omega)}{2} \mp \sqrt{\frac{1}{4a^2} - 1} \cdot \sum_{n=0}^{\infty} \frac{(2n)!}{(1 - 2n) \, n! \, 4^n \left(\frac{1}{a} + 2\right)^n} r^n(\omega) \cdot \sum_{n=0}^{\infty} \frac{(-1)^n \, (2n)!}{(1 - 2n) \, n! \, 4^n \left(2 - \frac{1}{a}\right)^n} r^n(\omega).$$

- 4. Initial Data Construction
- 5. Results of Numerical Experiments
- 6. Conclusion

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References