

ALPHA BETA AND GAMMA RADIATION LAB

INTRODUCTION

Unstable atomic nuclei may decay into less energetic nuclei by emission of alpha, beta, and gamma radiation. These forms of radiation are now known to consist of He^4 nuclei, electron (or positron), and high energy photons respectively. Which types of decay a nuclei may undergo depends on how it is unstable. Many materials may emit more than one form of radiation. This is especially true when the original parent radioactive nuclei produce daughter nuclei that are also radioactive. Such *decay chains* will be investigated further in one of the next labs, *Airborn Radiation: The Radon Decay Chain*. In such cases the various forms of radiation can be separated by a magnetic field as shown in Figure 1.

Nuclei with atomic numbers (number of protons, Z) greater than 83 may all theoretically undergo alpha decay (however, the rate of alpha decay may be so slow for some nuclei as to be effectively forever.) In alpha decay a helium nucleus consisting of two protons and two neutrons, is emitted reducing the atomic number of the original nuclei by two and its atomic mass by four.

Nuclei may also be unstable when they have too many or too few neutrons compared to protons. If there is an abundance of neutrons one may change into a proton while emitting an electron and antineutrino. The high energy electron is called a beta particle and this process is called β^- decay. When there are too many protons one may change into a neutron and emit a positron (negative electron) and a neutrino. The positron, which in this case is called a β^+

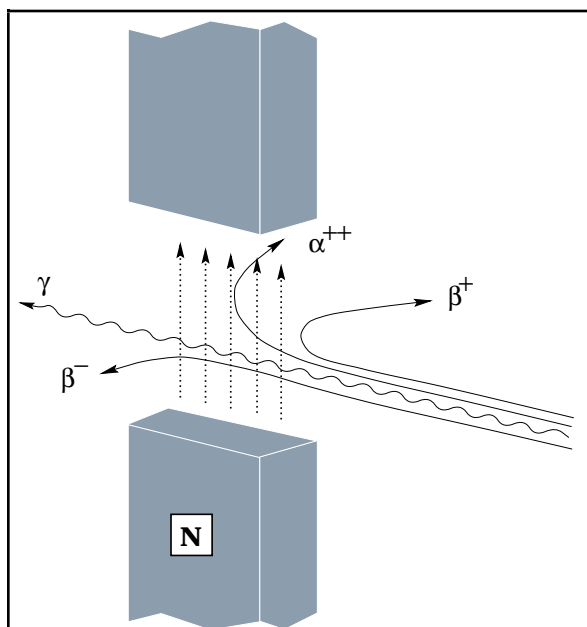


Figure 1

Deflection of Charged Particles by a Magnetic Field

particle is rarely detected directly. It usually quickly finds an electron and the two annihilate each other producing a pair of high energy photons called gamma rays.

Gamma radiation may also occur when beta or alpha decay leaves a nucleus in a high energy state. Just as an electron in a high energy atomic orbital may decay into a lower energy state by emitting a photon, a nucleus may also emit photons when changing from higher to lower energy states which in this case are called gamma rays.

The following next few sections of this lab describe the random nature of nuclear radiation emission and the natural decrease in intensity as radiation spreads out from a source. This is followed by specific sections for alpha, beta, and gamma radiation. These sections discuss the particular properties of each type of radiation and describe a set of experiments for each radiation type. Before performing these experiments the auxiliary lab handout, *Geiger-Müller Meters and Radiation Event Counting Software*, should also be read.

DECAY STATISTICS

Nuclear decays occur as random events. Each unstable nuclei in a radioactive sample has a constant slight probability of decaying during every infinitesimal instant in time. The number of decays occurring during a finite period of time in a sample follow *Poisson* statistics. For Poisson statistics if λ is the average number of decays per second then the probability that there will be n decays during a time interval of duration t is,

$$\text{Prob. of } n \text{ decays during } t \text{ seconds} = \frac{(\lambda t)^n e^{-\lambda t}}{n!} \quad (1)$$

where $n!$ means n factorial,

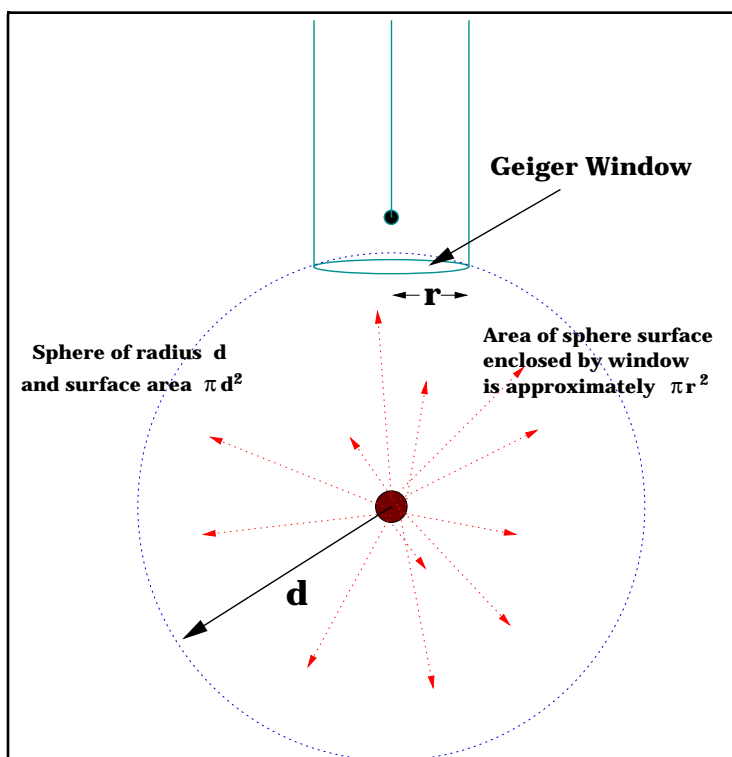


Figure 2 Inverse Square Law

The fraction of the total emissions emitted toward the detector window is inversely proportional to the square of the distance between the source and window.

i.e. $n! = (n) \cdot (n-1) \cdot (n-2) \cdot \dots \cdot (2) \cdot (1)$.

The standard deviation of the number of decays occurring during a time period is equal to the square-root of the average number of decays. Remember from statistics class that the standard deviation is defined as the square-root of the average of the square of the difference between the number of counts and the mean. So for some random quantity Q ,

$$\text{Standard Deviation } \sigma_Q = \sqrt{\langle [Q - \langle Q \rangle]^2 \rangle} \quad (2)$$

where $\langle \cdot \rangle$ denotes the average of some quantity. For Poisson statistics.

$$\sigma_n = \sqrt{\langle [n - \langle \lambda \rangle]^2 \rangle} = \sqrt{\lambda}. \quad (3)$$

The standard deviation gives a measure of how much the number of counts from one trial is likely to deviate from the mean. When the number of counts involved becomes large (λt large) Poisson statistics become very similar to Gaussian statistics. In this case about 68% of the time the number of counts measured will fall within plus or minus $\sqrt{\lambda}$ of the mean λ , i.e. within $\lambda \pm \sqrt{\lambda}$. About 95.5% of the time the number of counts will fall within $\lambda \pm 2\sqrt{\lambda}$. For example suppose the average number of counts is 100. Then 68% of the time the number of counts will fall between 90 and 100 counts and 95% of the time between 80 and 120 counts.

The relation between the standard deviation and the mean can be used to give an idea of how many counts are needed to obtain a certain degree of accuracy. Dividing the average difference from the mean is $\sqrt{\lambda}$ by the mean λ gives an rough estimate of the accuracy. Since $\sqrt{\lambda}/\lambda = 1/\sqrt{\lambda}$ the accuracy increases as the square root of the number of counts. If count rates of about 100 are measured about $\sqrt{100}/100 = 1/\sqrt{100} = 0.1$ or ten percent accuracy is obtained. In order to get one percent accuracy count

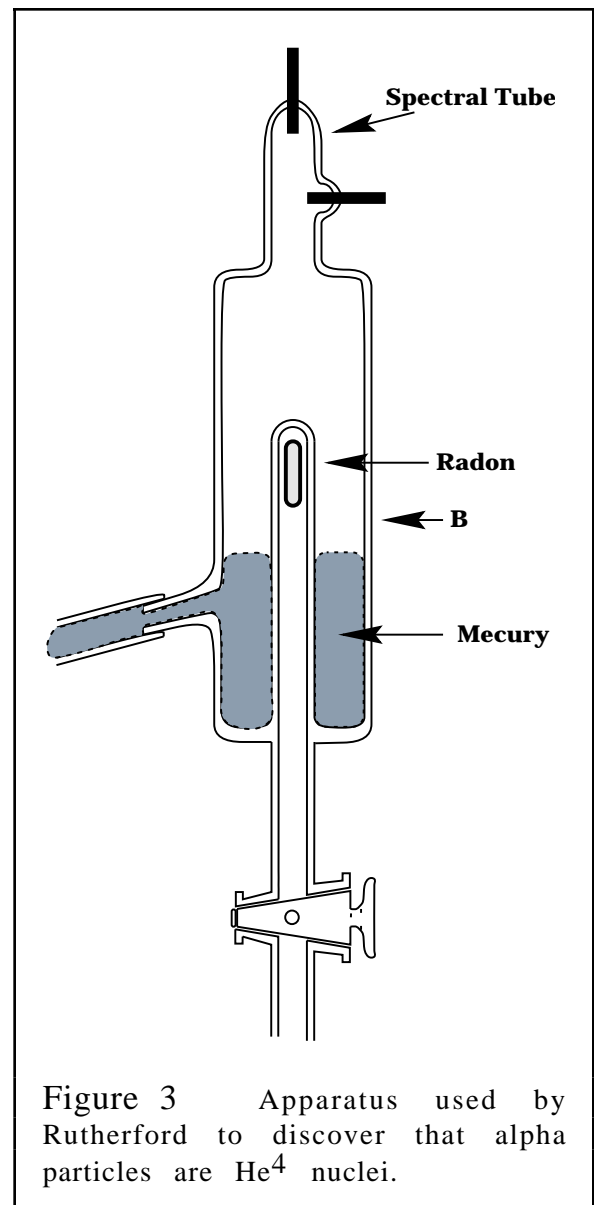
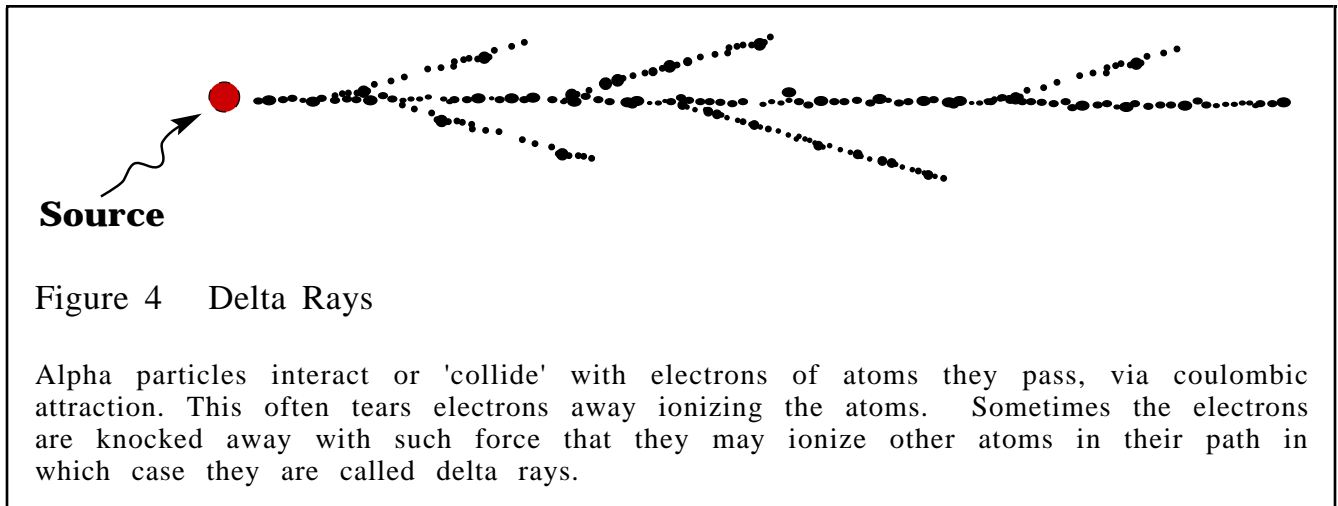


Figure 3 Apparatus used by Rutherford to discover that alpha particles are He^4 nuclei.



numbers of about 10,000 counts are needed since $1/\sqrt{10000} = 0.01$. The 'Natural Airborn Radiation: The Radon Decay Chain' lab provides a further description of Poisson statistics and nuclear decay rates.

THE INVERSE SQUARE LAW

When a nuclei decays the direction in which the particles are emitted is random. Therefore if absorption within a radioactive sample is neglected, the emissions are evenly distributed in all directions. Now consider a set of concentric spherical surfaces surrounding and centered on the source. If none of the particles are absorbed each sphere would have the same number of particles passing through it. However, the outer spheres have much more surface area than the inner ones and thus the density of radiation nearer to the source is much higher. Considering the surface area of a sphere,

$$S = 4\pi d^2 \quad (4)$$

it is seen that even without absorption radiation intensity levels decrease as the reciprocal of distance from the source squared, $1/d^2$. This is the inverse square law.

Now consider the window of a Geiger-Müller counter placed near a source as shown in Figure 2. The fraction of the total radiation emitted from the source in the direction of the window is given by the geometry or G factor. This factor is equal to the ratio of the area on a sphere enclosed by the counter window to the total area of the sphere.

$$G = \frac{S'}{4\pi d^2} \quad (5)$$

If, for example, the Geiger counter has a round window of radius r , the area of a sphere centered on the source enclosed by the window will be approximately πr^2 (as long as the window is not placed too close to the source.) This gives,

$$G \approx \frac{\pi r^2}{4\pi d^2} \approx \frac{r^2}{4d^2} \quad (6)$$

As the Geiger-Müller counter is moved away from the source the total area of the spherical surface the window is on increases by distance d squared but the area of the sphere enclosed by the window S' stays almost the same. Again we see that radiation intensity levels naturally fall off by an inverse square law as the particles spread away from the source. When measurements are taken at different distances from the source or different sized windows are

used the geometry or G factor must be used to correctly compare the measurements. The inverse square law may also be used to adjust the count rate. If the rate is too low, move the source nearer to the counter. If the rate is too high, move the source farther away.

BACKGROUND RADIATION

A percentage of any radiation events detected in an experiment will actually be due to outside sources. These may be natural sources such as cosmic rays and radioactive elements found in the surrounding air and building materials, or artificial sources such as unshielded radioactive chemicals stocked nearby or the luminous paint of a wristwatch. The background radiation count rate should always be measured as part of any radiation experiment. It should then be subtracted from the count rate data taken for the experimental source.

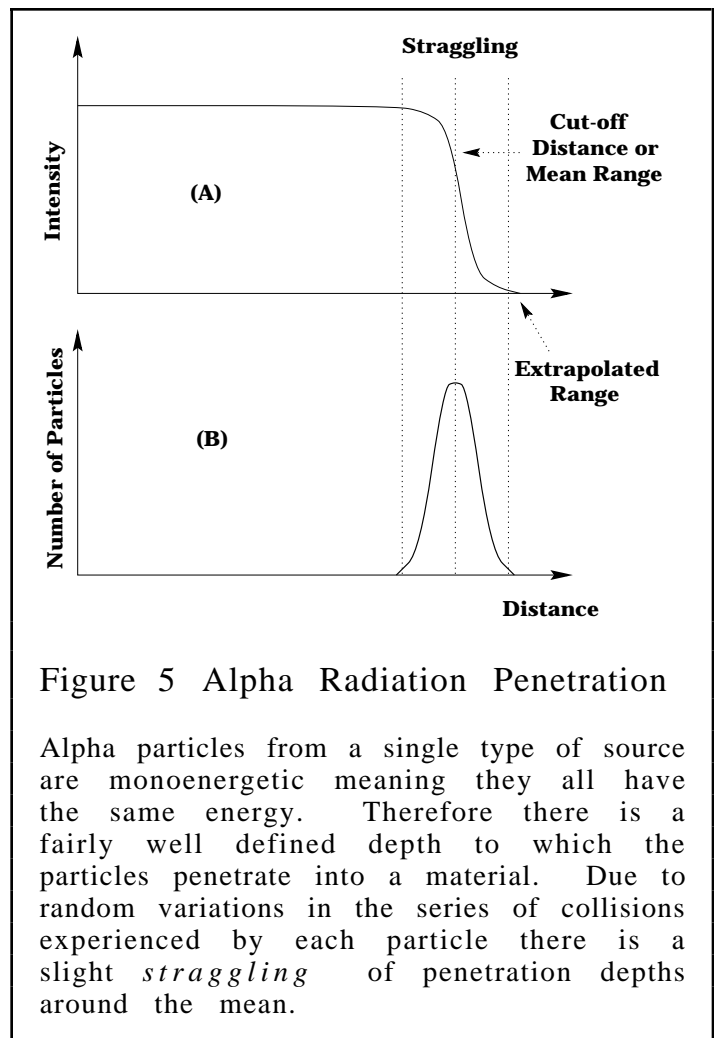


Figure 5 Alpha Radiation Penetration

Alpha particles from a single type of source are monoenergetic meaning they all have the same energy. Therefore there is a fairly well defined depth to which the particles penetrate into a material. Due to random variations in the series of collisions experienced by each particle there is a slight *straggling* of penetration depths around the mean.

Although, subtracting the background count rate data from the measured count rate gives a better idea of the amount of a source's emission rate, since both the source and background emissions are random events one can never know exactly how many of the counts were due to each source. This introduces some extra uncertainty or *noise* into the data. Since the background radiation also follows Poisson statistics the uncertainty will be about the square-root of

the average number of background counts. For example suppose intervals of five minutes are used to sample the emission rate of a source. A background radiation check produces one thousand one hundred and four background counts measured over one hundred minutes. This produces an estimate of fifty five and one fifth counts per five minutes with a standard deviation or uncertainty of about seven and a half counts ($\sqrt{55.2}=7.43$). If during one of the five minute intervals 324 counts are measured, the total uncertainty is $\approx\sqrt{324} = 18$ with the number of counts due to the source being around $324 - 55.2 \pm 7 \approx 269 \pm 7$. Note that uncertainties from different sources add as follows,

$$\sigma_{\text{Total}}^2 = \sigma_{\text{Source 1}}^2 + \sigma_{\text{Source 2}}^2 + \sigma_{\text{Source 3}}^2 + \dots \quad (7)$$

or as in this case if the average source rate is 269 counts per minute,

$$\sigma_{\text{Total}}^2 = \sigma_{\text{Source}}^2 + \sigma_{\text{Background}}^2 = (\sqrt{269}=16.4)^2 + (7.43)^2 = 269 + 55 \quad (8)$$

giving, $\sigma_{\text{Total}}^2 = 18 \quad (= \sqrt{324})$.

The effects of noise added by background radiation on count rate accuracy can be decreased by increasing the ratio of experimental source rates to the background

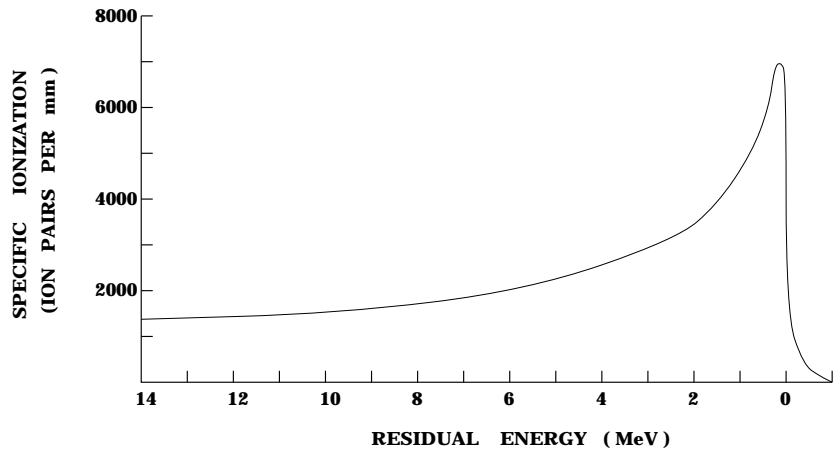


Figure 6 Bragg Curve for Alpha Particles

This curve shows the specific ionization as a function of alpha particle energy. Specific ionization is defined as the number of pairs produced per unit length of path and characterizes the interaction strength of the alpha particle with atoms it passes. At slower speeds and thus lower energies the alpha particle spends more time in the vicinity of each atom it passes giving it a greater opportunity to attract an electron.

level. Shielding the experiment from outside sources, moving the counter closer to the source and using larger or more radioactive sources will all improve this ratio.

ALPHA RADIATION

PROPERTIES OF ALPHA RADIATION

Alpha particles consist of two protons and two neutrons the same as the nucleus of a helium atom. This was confirmed with an experiment performed Rutherford and Royds in 1909. They placed a bit of radon inside a thin walled glass tube. This tube was surrounded by a second tube in which a vacuum had been created using a mercury column (as shown in Figure 3.) Alpha particles emitted by the radon passed through the first glass wall and became trapped in the second tube. By raising the level of the mercury Rutherford and Royds concentrated the alpha

particles in the top portion of the tube. The concentrated alpha particles were found to have the same spectroscopic signature as helium.

Alpha particles emitted from a single type of source are monoenergetic. That is they are all emitted with the same kinetic energy. Depending on the source this energy ranges from 3-8MeV corresponding to velocities of $1.4\text{-}2\cdot 10^7$ m/s. Due to their double charge alpha particles interact strongly with the electrons of the molecules of the material they travel through. These 'collisions' often tear the electrons away leaving an ion pair trail along the alpha particles path. Sometimes the electron is knocked away with enough speed that it may ionize some more

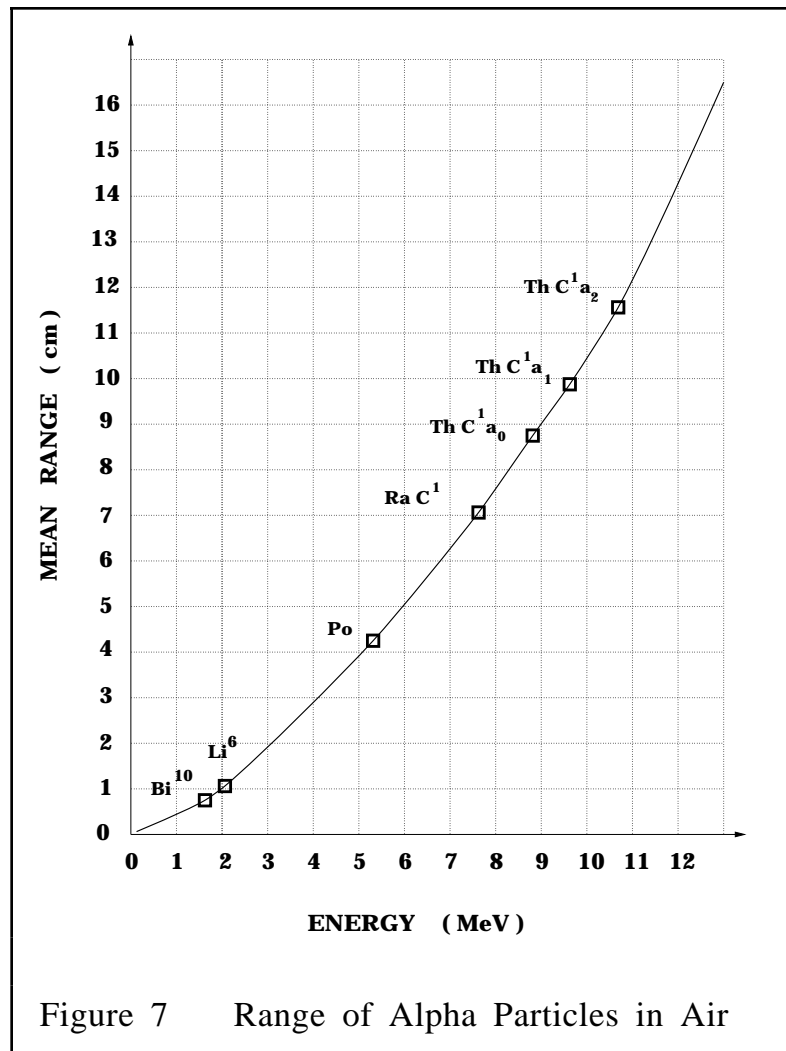


Figure 7 Range of Alpha Particles in Air

d	Counts	Counting Duration	Count-Rate ($\frac{\text{Counts}}{\text{Duration}} - \text{Background}$)	$\frac{\text{Count-Rate}}{d^2}$	Standard Deviation

Table 1 Alpha Radiation Absorption Data

molecules and is then called a delta ray (see Figure 4.) Because alpha particles have such a high mass collisions with electrons have little effect on their direction of travel. Therefore they travel in a fairly straight line. It takes between 25-40 electron Volts to strip away an electron (a value of 32.5eV is usually used for air) so each collision saps a little of the alpha particles kinetic energy. Dividing the energy of a normal alpha particle by the ionization energy it is seen that around 100,000 ion pairs are created. Quite a few! Although alpha particles are not very penetrating they may cause a lot of damage to a material at the depth they reach.

Since all alpha particle from a single source type are emitted with the same energy one might expect them to penetrate to the same depth in a given material. This is close to true. However, slight variations in the random collisions experienced along the path cause a slight 'straggling' around the mean penetration range with a Gaussian distribution. The radiation intensity vs. depth and corresponding penetration depth histogram are graphed in Figure 5. As an alpha particle penetrates more deeply it is continually slowed down by each collision. As its speed slows the alpha particle stays in the vicinity of each electron passed a little longer. Extra interaction time increases the chance that the electron will be dislodged. Figure 6 shows the dependency of the interaction strength on the kinetic energy ($\frac{1}{2} m v^2$) of an alpha particle.

Alpha particles have a range of 2-8cm in air depending on their initial velocity as shown in Figure 7. This may be expressed by the formula,

$$v_0^3 = a R \quad (9)$$

where v_0 is velocity, $a \approx 10^{27}$, and R is the range. If R is expressed in centimeters and energy E in MeV then $R \approx E$.

The penetration power of alpha particles into a material is often characterized by the materials *linear relative stopping power* which relates the average penetration depth to that in air.

$$S_{\text{material}} = \frac{R_{\text{air}}}{R_{\text{mat}}} \quad (10)$$

Penetration depth is nearly inversely proportional to material density.

$$R_{\text{air}} \rho_{\text{air}} = R_{\text{mat}} \rho_{\text{mat}} \quad (11)$$

Therefore the linear relative stopping power may be expressed by relative density to air.

$$S_{\text{material}} = \frac{\rho_{\text{mat}}}{\rho_{\text{air}}} \quad (12)$$

For aluminum,

$$S_{\text{Al}} = \frac{\rho_{\text{Al}}}{\rho_{\text{air}}} = \frac{2.74 \cdot 10^3}{1.29} = 2.1 \cdot 10^3 \quad (13)$$

ALPHA RADIATION EXPERIMENT

Polonium²¹⁰ (Po²¹⁰) will be used as your alpha radiation source. It emits alpha particles with an energy of 5.4MeV. You will be measuring the range of alpha particles in air by detecting the variation in count rate as a function of distance between the source and Geiger-Müller detector window. As discussed in an earlier section of this handout (The Inverse Square Law), since the source emits radiation evenly in all directions, neglecting absorption, the count rate should decrease as the reciprocal of distance d from the source squared. This factor will need to be accounted for (by multiplying count rate by d^2) in order to see the true absorption by air. As

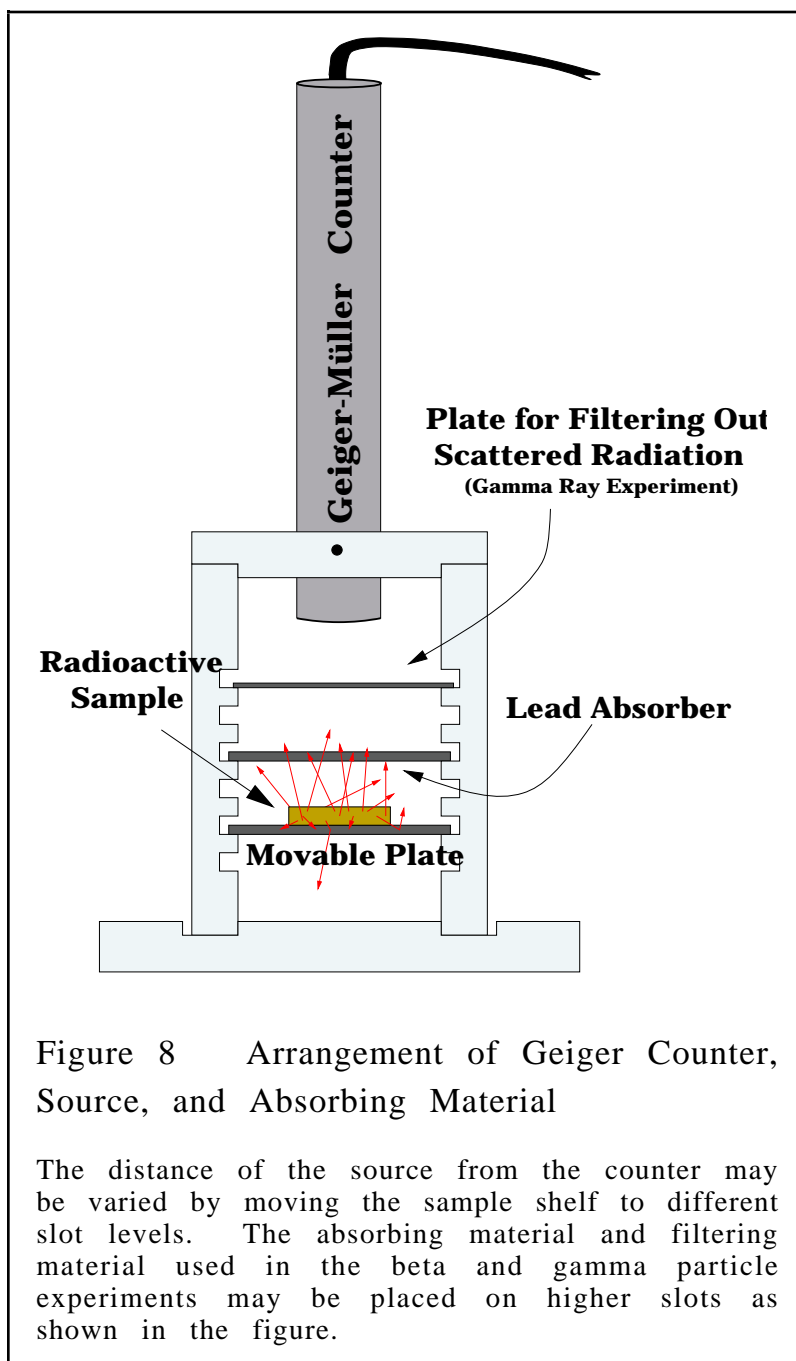
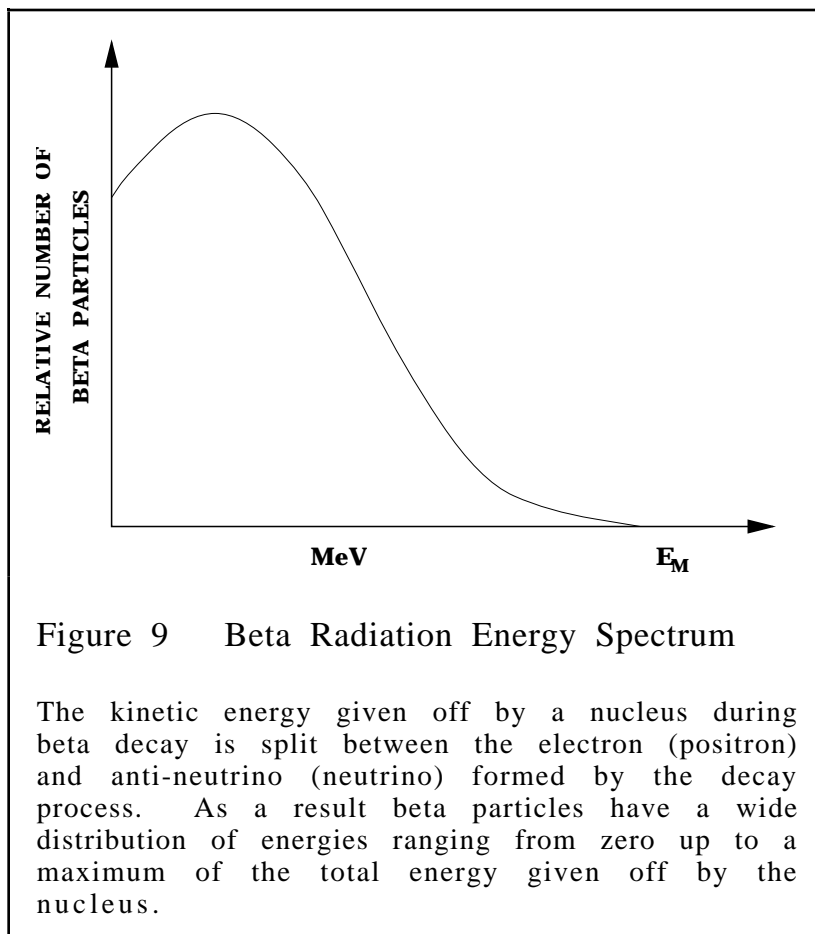


Figure 8 Arrangement of Geiger Counter, Source, and Absorbing Material

The distance of the source from the counter may be varied by moving the sample shelf to different slot levels. The absorbing material and filtering material used in the beta and gamma particle experiments may be placed on higher slots as shown in the figure.

stated before the range in air in centimeters is approximately the same as the energy of the alpha particles in mega-electron Volts. Therefore a sharp drop in count rate should be expected as the source is moved through a distance of about five centimeters away from the Geiger counter. Remember to measure and subtract the background as discussed in the earlier section of this handout, Background Radiation.

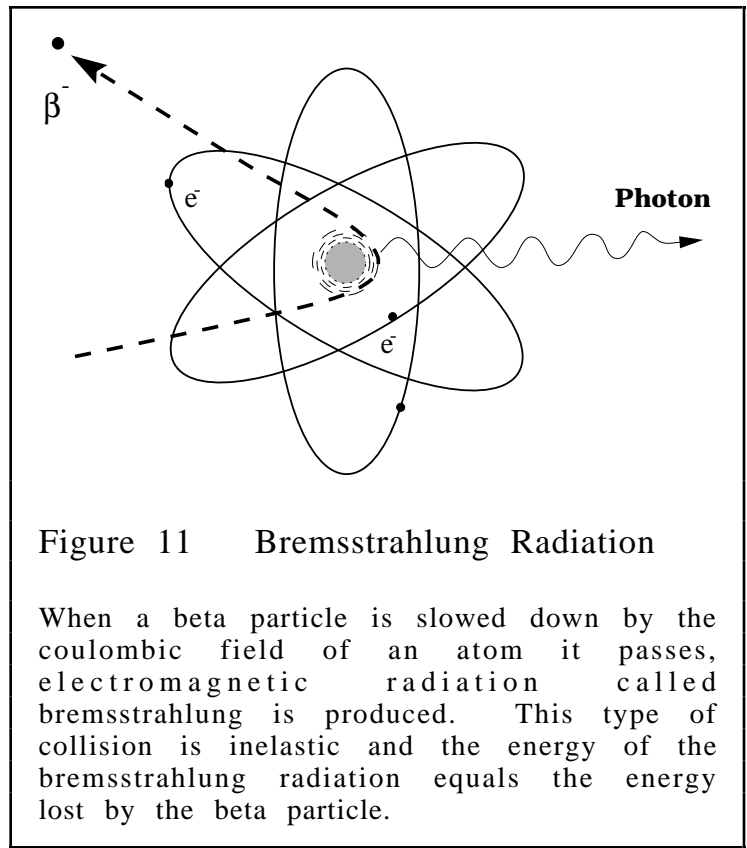
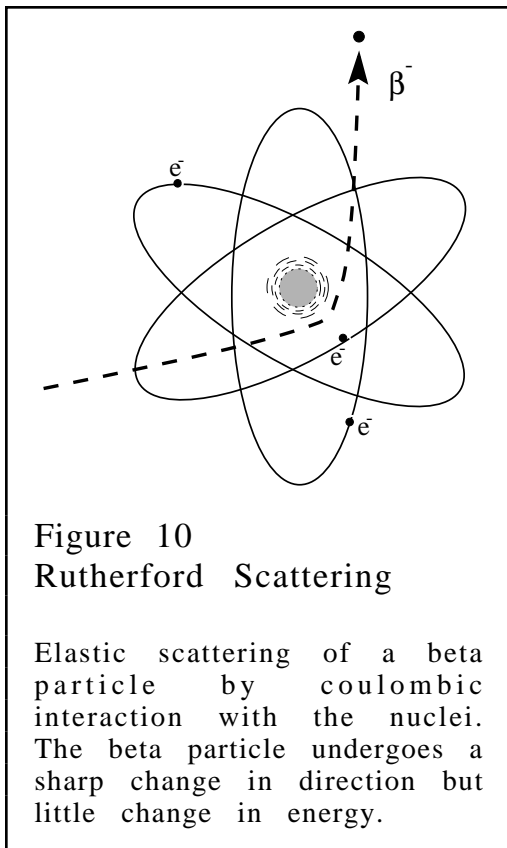
The distance between the source and counter may be varied by placing the source on a plate and inserting it into the slots at different levels in the Geiger meter holder (see Figure 8.) Make sure to count for a long enough period of time in order to get good data. Remember count accuracy is equal to about plus or minus the square-root of the number of counts. The computer counting programs you may use are described further in the *Radiation Experiments Auxiliary Handout: Geiger Counters and Radiation Event Counting Software*. An example of how the data from this experiment might be tabulated is given in Table 1. Discuss the results.



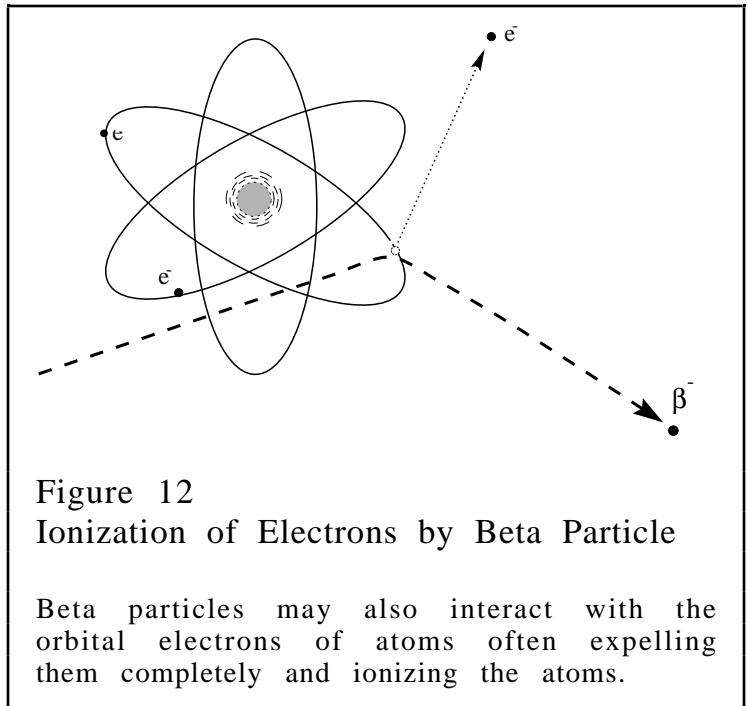
BETA RADIATION

PROPERTIES OF BETA RADIATION

Beta particles have been experimentally shown to be electrons emitted from a nucleus. The velocity of an electron is related to the energy with which it was emitted. Emission of an beta particle is always accompanied by a neutrino emission. The energy given off by the decaying nucleus is split between the electron and



neutrino according to momentum conservation. As a result the energy of beta radiation is continuously distributed from zero up to the total decay energy given off by the nucleus (in which case the neutrino receives zero kinetic energy.) The energy spectrum of beta particles is explored further in the *C35 Beta Energy Spectrum Lab*. As shown in Figure 9 the average energy received by the beta particle is between one forth and one third the maximum energy (the maximum being the total energy given off by the nucleus.)



Because of their small mass and the energies involved in nuclear decay, beta particles often are emitted at relativistic speeds. The total energy of a particle is given by Einstein's formula.

$$E = m c^2 \quad (14)$$

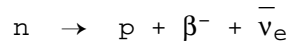
The effective mass m is related to the rest mass m_0 by,

$$m = \frac{m_0}{\sqrt{1-\beta^2}} \quad (15)$$

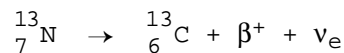
where $\beta = \frac{v}{c}$, v is the velocity of the particle, and c is the velocity of light. Subtracting the rest energy $U = m_0 c^2$ from the total energy gives the kinetic energy.

$$\text{K.E.} = (m - m_0) c^2 = m_0 c^2 \left(\frac{1}{\sqrt{1-\beta^2}} - 1 \right) \quad (16)$$

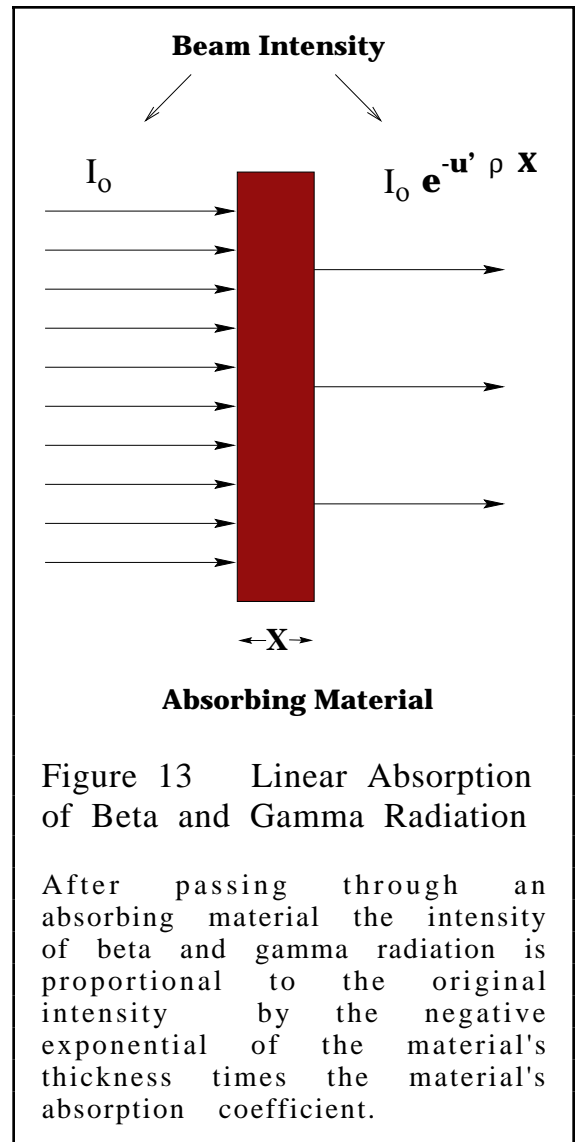
Beta decay may occur in two ways. In β^- decay a neutron decays into a proton, an electron (beta particle) and an anti-neutrino. This is given by the equation,



A neutron has a slightly higher rest mass than a proton. As a result a proton cannot decay into a neutron without absorbing some energy from an outside source. The extra energy of an unstable nuclei can provide the energy needed to turn a proton into a neutron if this change puts the nucleus into a lower energy state. Such a process is called β^+ decay. As the nucleus changes into a new nucleus with one less proton and one more neutron it also emits a positron and a neutrino.



Positrons usually quickly annihilate with a nearby electron producing a pair of gamma rays in the process.



BETA INTERACTIONS

Beta particles may interact with both the nuclei and electrons of a material they pass through. A beta particle is said to have collided with a nuclei when it is strongly deflected by the nuclei's coulombic field (see Figure 10.) If the collision is elastic the process is called Rutherford scattering and the energy of the electron undergoes little change. The same type of result is seen if one throws a ping-pong ball at a bowling ball. As a result of momentum conservation the ping-pong ball changes direction (a large change in momentum) but imparts little of its energy to the bowling ball. Most back scattering of electrons is due to Rutherford scattering.

As electrons are accelerated or decelerated by the coulombic field of a nuclei, electromagnetic radiation called *bremsstrahlung* is produced (see Figure 11.) This type of collision is inelastic. Only about one percent of the energy of a beta particle beam incident on a target is converted into bremsstrahlung radiation. The percentage increases with atomic number of the target atoms.

Beta particles may also interact with orbital electrons via their mutually repulsive charge. This often ionizes atoms by knocking their electrons out of orbit producing characteristic X-rays when higher orbital electrons decay down to the empty state (see Figure 12.)

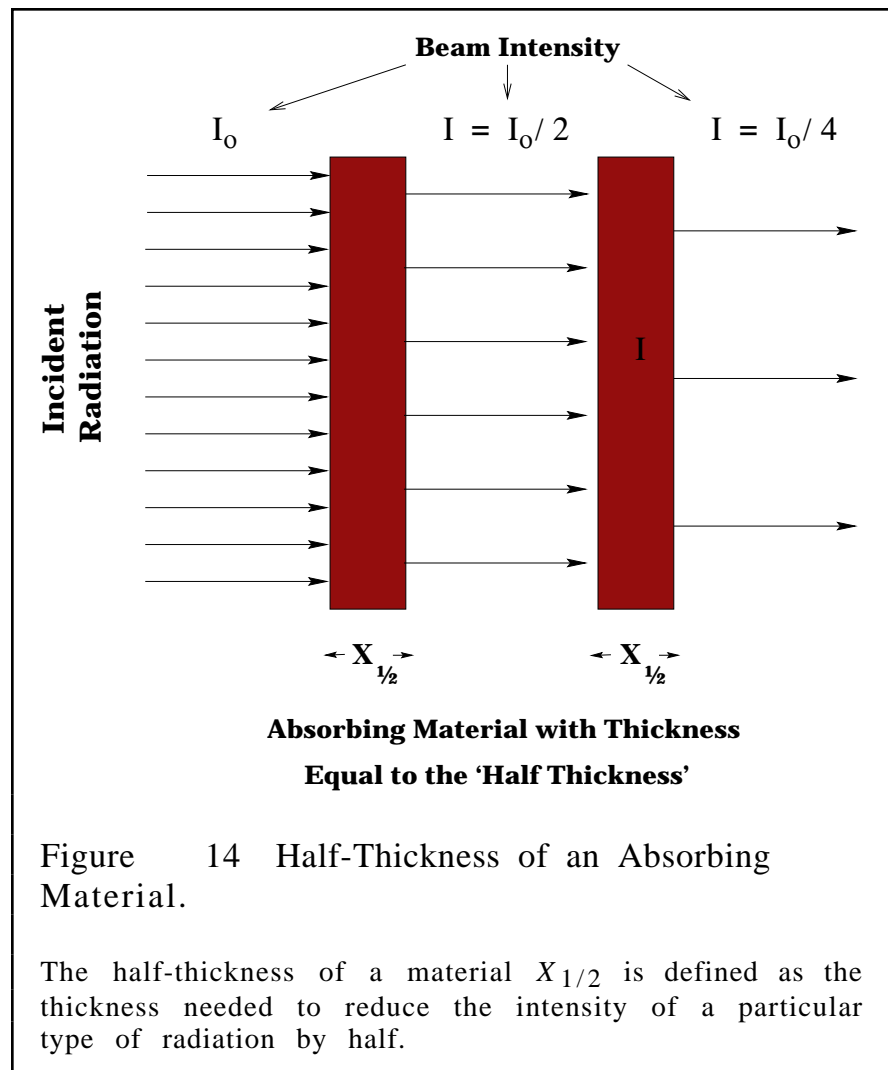


Figure 14 Half-Thickness of an Absorbing Material.

The half-thickness of a material $X_{1/2}$ is defined as the thickness needed to reduce the intensity of a particular type of radiation by half.

Although absorption of beta particles by a material is a random process, on a macroscopic scale (for which random fluctuations have averaged out) the intensity of an electron beam is found to exponentially decay with distance traveled according to the formula.

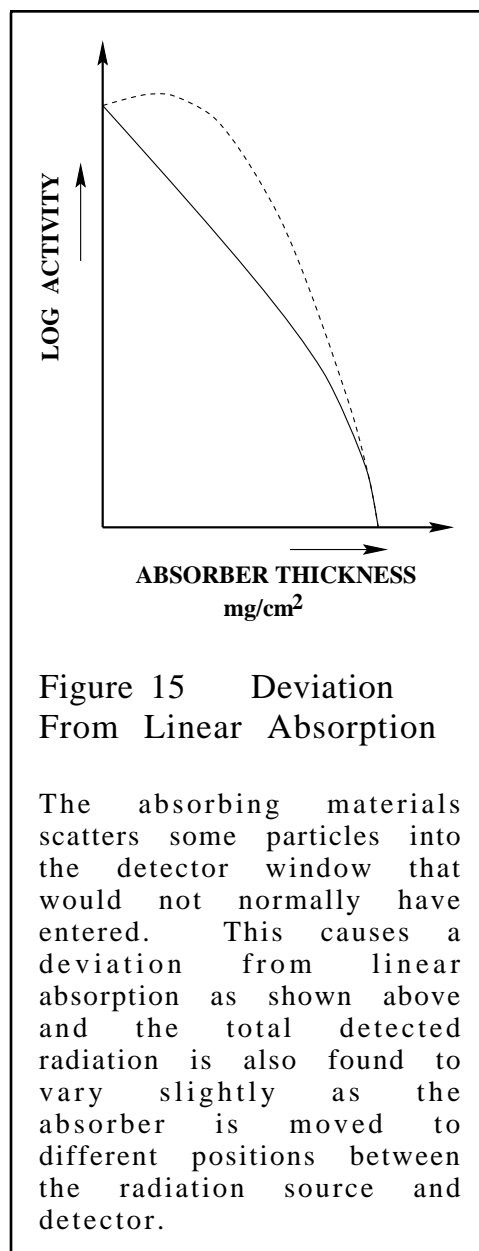
$$I = I_0 e^{-u'\rho X} \quad (17)$$

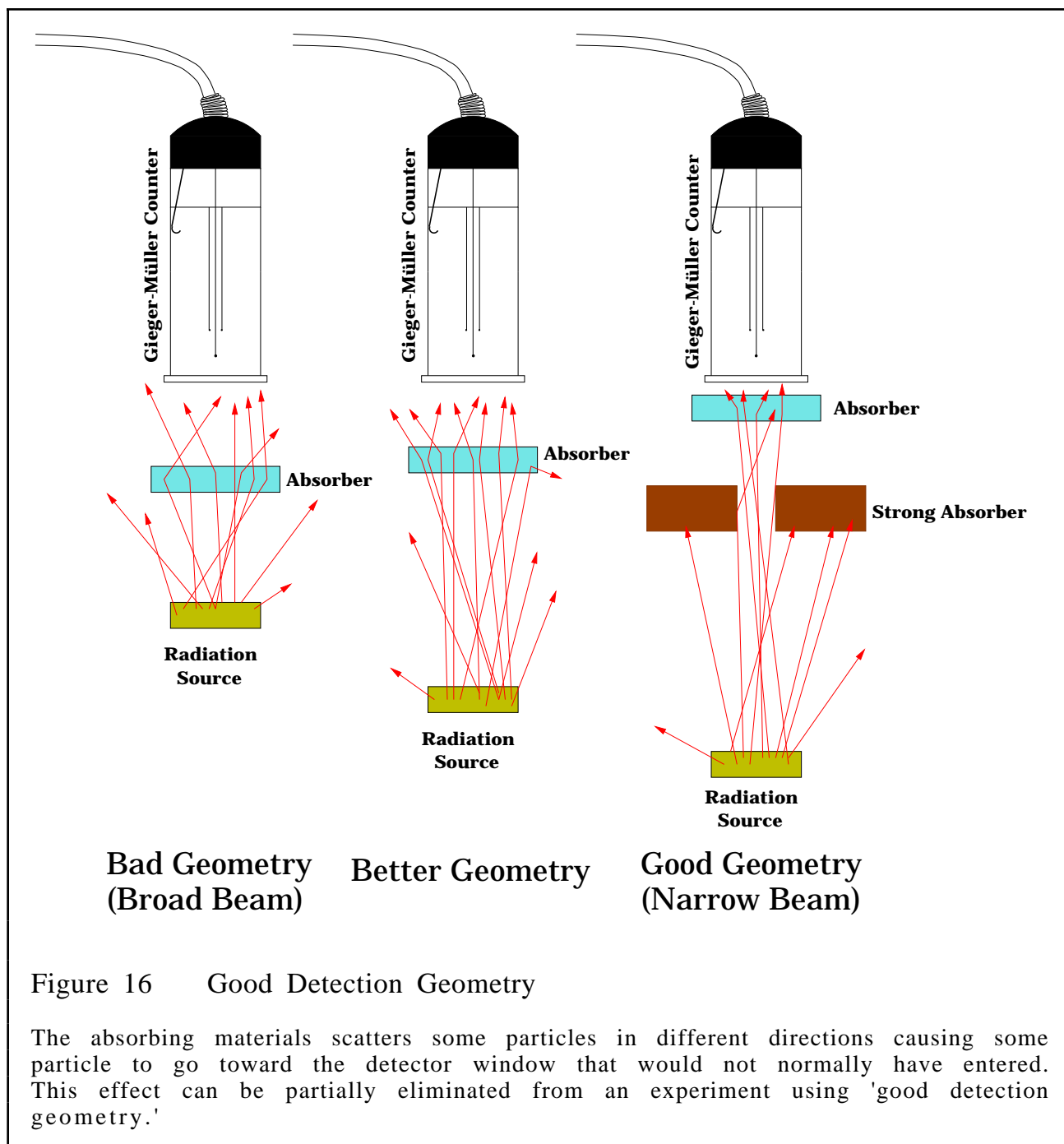
In equation 17, X is distance traveled, ρ is the density of the absorber and u' is the *mass absorption coefficient* (see Figure 13.) This is closely similar to the absorption of light and gamma rays. Often a material may be characterized by its half-thickness which is the thickness of material needed to reduce the intensity of a beam by one half. If a beam is passed through two half-thicknesses of material the emission on the other side would be one fourth of the original intensity. With three half-thicknesses intensity is reduced to one eighth as shown in Figure 14.

In actuality the resulting intensity measured by a Geiger counter is slightly dependent on the position of the absorbing material between the beta source and the detector. This is because some beta particles not initially emitted toward the detector window may be scattered back into it as they pass through the absorber. As a result the total detected radiation increases slightly as the position of the absorber with respect to the detector and source. This effect also causes deviation from equation 17 as shown in Figure 15. This effect may be ameliorated if the experiment is set up with 'good geometry' as discussed further in the gamma radiation section and shown in Figure 16.

BETA RADIATION EXPERIMENT

The beta source is either Thallium²⁰⁴ (Tl²⁰⁴) or Strontium⁹⁰ (S⁹⁰). In this experiment we will be calculating the half thickness of aluminum. You will have





many thin films of aluminum of varying thickness (1,2 5, 10, 20 microns) which you can stack together to get a large range of thicknesses. As in the alpha radiation experiment the source emits in all directions so you will need to account for decreases in count rate due to $\frac{1}{(\text{distance})^2}$ loss if the distance between source and detector is changed. You will also need to subtract the background rate.

It is probably easiest to place the source on the lowest rung of the Geiger meter holder and leave it there. Then the aluminum films may be stacked together on the intermediate rungs. Again make sure that enough counts are measured at each thickness in order to get accurate count rates. About five minutes is a good time.

Tabulate the count rate versus absorber thickness using a chart like the one suggested in Table 2. Does it follow an decaying exponential curve? Make a linear-least-squares fit to the data and calculate the half thickness. This may be easily performed using MatLab or with some calculators. You can check out Dr. Erlenmeyer's Least-Squares fit web page (<http://www.astro.nwu.edu/PAGES/erlenmeyer.html>) and the corresponding Theory page (<http://www.astro.nwu.edu/PAGES/lsgerr.html>) for more details.

It is possible for the count rate to get so low that it is below the background. At this point some of the count rate samples will give estimated rates below zero due to noise fluctuations. This data obviously is unphysical and indicates that the number of counts is so low that the data is completely inaccurate and basically useless. Either longer counting times are needed to reduce the relative noise level and obtain better measurements or data past this point should be neglected.

Thickness	Counts Measured	Counting Time	Counts/Time - Bckgrd (Rate - BG)	$\frac{\text{Rate}-\text{BG}}{d^2}$	$\ln\left(\frac{\text{Rate}-\text{BG}}{d^2}\right)$

Table 2 Beta Radiation Absorption Data

GAMMA RADIATION

PROPERTIES OF GAMMA RADIATION

Gamma radiation consists of extremely high energy photons with wavelengths usually much smaller than 10^{-10} meters (compare to $\approx 10^{-6}$ meters for visible light.) Gamma

emission is usually preceded by an alpha or beta emission which left the nucleus in an excited state. An excited nucleus may change to a lower energy state by emitting a photon which in this case is called a gamma ray. This is similar to emission of a photon when an electron changes from a higher to lower orbital. The spacing between nuclear energy states is on the order of millions of electron Volts (compared to electron volts for electron orbitals.) Substituting the energy of a photon $E=hf$ into the wavelength-frequency relation $c=f \lambda$ gives,

$$\lambda = \frac{c}{f} = \frac{hc}{E} = \frac{1240 \text{ MeV} \cdot \text{femtometers}}{1 \text{ MeV}} = 1240 \text{ fm} \quad (18)$$

As for visible light, absorption of gamma rays by a material is closely linear. That is for every unit of distance the radiation travels through the material an equal fraction of the remaining light is absorbed. For example suppose the radiation is passed through a series of absorbing plates each of the same material and thickness. Also suppose each plate is of a thickness that it absorbs half the gamma rays passing through it, i.e. the *half-thickness*. After the first plate the intensity is half of the original. Passing through the second plate half the remaining rays are removed leaving one fourth the original. After the third plate only one eighth is left and so on (see Figure 15) This is an exponential type decay which may be expressed by Lambert's law.

$$I(X) = I_0 e^{-\mu X} \quad (19)$$

or

$$\ln(I/I_0) = -\mu X \quad (20)$$

where μ is the linear absorption coefficient. The half-thickness of a material is thus given by,

$$I(X_{1/2}) = 0.5 I_0 = I_0 e^{-\mu X_{1/2}} \quad (21)$$

and

$$X_{1/2} = \frac{-\ln(1/2)}{\mu} = \frac{0.693}{\mu}$$

Absorption of gamma rays occurs by several means, the most prominent being the photoelectric effect at low energies, the Compton effect at intermediate energies and pair production at high energies. The linear absorption coefficient is a sum of the coefficients for each of these processes.

$$\mu = \tau (\text{photoelectric}) + \sigma (\text{Compton}) + K (\text{pair production}) \quad (22)$$

Beam energy is also lost in negligible amounts to Rayleigh scattering, Bragg scattering, photodisintegration, and nuclear resonance scattering.

Photoelectric absorption occurs when a gamma ray ionizes an atom. The photoelectric absorption rate τ is proportional to the density of the material and also increases strongly with atomic number.

$$\tau = 0.0089 \rho \frac{Z^{4.1}}{A} \lambda^n \quad (23)$$

In equation (23) ρ is the density of the material, Z is atomic number, A is atomic mass, λ is wavelength and the number n is

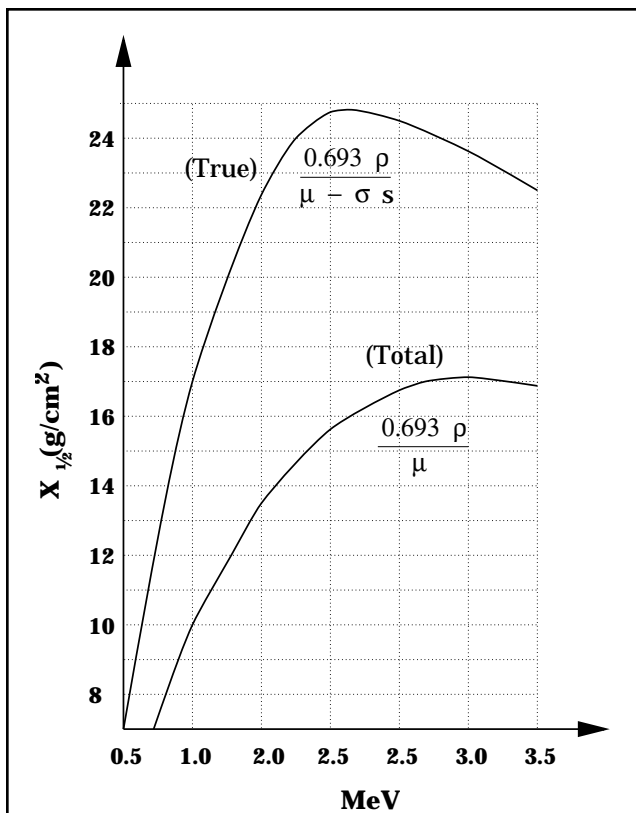


Figure 18 Half-Thickness Widths for Lead Absorbers

"True" represents the lower limit for broad-beam geometry. "Total" represents the upper limit for narrow-beam geometry. All values should fall between these limits.

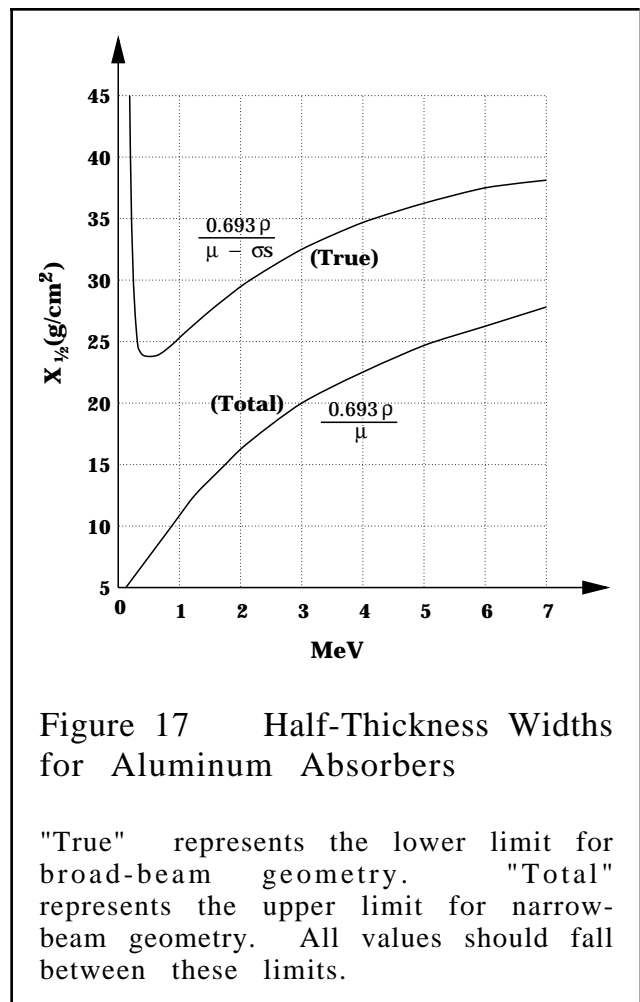


Figure 17 Half-Thickness Widths for Aluminum Absorbers

"True" represents the lower limit for broad-beam geometry. "Total" represents the upper limit for narrow-beam geometry. All values should fall between these limits.

equal to three for Nitrogen, Oxygen, and Carbon and 2.85 for all other elements with atomic numbers less than iron. Photoelectric absorption is strongest at longer wavelengths (i.e. less energetic gamma rays) and increases very rapidly with atomic number.

When a gamma ray disintegrates into an electron positron pair it is called pair production. For this to occur the energy of the gamma ray must be greater than the resting energy of an electron and a

positron, $hf \geq 2 \cdot 0.51 \text{ MeV}$. The pair formation coefficient is given by,

$$K = a N Z^2 (E - 1.02 \text{ MeV}) \quad (24)$$

where a is a constant, N is Avogadro's number, and Z is the atomic number.

Loss by Compton scattering consists of two components,

$$\sigma = \sigma_a + \sigma_b \quad (25)$$

where σ_a is the scattering component due to energy loss via collisions with electrons and σ_b is energy loss due to scattering of photons scattered away from the beam direction. The σ_b loss is strongly dependent on the geometry of the source, absorber and detector. This is because in some geometries many gamma rays that were not initially headed toward the detector window may be redirected into it by Compton scattering. For 'narrow beam' geometries this is unlikely, but for broad beams or wide radiation sources and when the absorber is not directly against the detector window count rates will be found to be slightly higher due to the redirected rays. (see Figure 16) Because of this effect the total absorption coefficient μ is often divided into a 'true' absorption component and a scattered component.

$$\mu = (\tau + K + \sigma_a) + \sigma_b \quad (26)$$

$$\text{Total} = \text{True} + \text{Scatter}$$

'Good geometry' must be established in order to measure the total absorption coefficient accurately.

GAMMA RADIATION EXPERIMENT

In this experiment the half-thickness of aluminum and lead will be measured for the gamma rays emitted by a Cobalt⁶⁰ (Co⁶⁰) source which radiates at 1.1MeV and 1.3MeV. It is suggested that a thin piece of aluminum is placed as close to the detector window as possible as shown in Figure 8. This will help to absorb and filter out some of the unwanted scattered radiation. Otherwise this experiment is identical to the beta radiation experiment except that thicker plates will be used since gamma radiation is more penetrating.

- 1) Measure the count rate at many different thicknesses for both lead and aluminum and compile a data table similar to Table 2. As in the beta decay experiment use the method of linear least-squares fitting to calculate the

half-thicknesses. Which has the higher half-thickness, aluminum or lead? Why?

- 2) Is there any absorption due to photon pair production for this source?
- 3) At lower energies the most loss is due to the photoelectric effect which is proportional to the density of the material. Divide the half-thicknesses you found for aluminum and lead by their respective densities (2.7 g/cm^3 for aluminum and 11.4 g/cm^3 for lead.) This gives the half-thickness in (g/cm^2) which is often used as a more true parameter of the absorption strength of a material type since it accounts for variations in density from sample to sample. Does the lead or the aluminum sample have the higher adjusted (g/cm^2) half-thickness? Why?
- 4) Compare the half-thicknesses calculated in Question 3 to Figures 17 and 18. What type of geometry matches this experiment?