

## \*Electrons and holes in semiconductors

intrinsic and extrinsic carriers; variations in  $E_F$

---

July 1, 2021

- ✦ When an electron is excited from the valence band to the conduction band it leaves an hole in the valence band.

- ✦ When an electron is excited from the valence band to the conduction band it leaves a hole in the valence band.
- ✦ This excitation can occur by thermal energy, or either by phonon absorption, or both.

# Electrons and holes

- ✱ When an electron is excited from the valence band to the conduction band it leaves an hole in the valence band.
- ✱ This excitation can occur by thermal energy, or either by phonon absorption, or both.
- ✱ Hole is a missing electron in a valence band, and from the effective mass approximation it has positive effective mass and this is seen in the equation of motion. And is like considering a positive charged particle.

# Electrons and holes

- \* When an electron is excited from the valence band to the conduction band it leaves a hole in the valence band.
- \* This excitation can occur by thermal energy, or either by phonon absorption, or both.
- \* Hole is a missing electron in a valence band, and from the effective mass approximation it has negative effective mass and this is seen in the equation of motion. And is like considering a positive charged particle.
- \* Without the crystal lattice, the hole cannot exist. It is an artifact of the periodic potential ( $E_p$ ) created by the crystal

# Electrons and holes

- \* When an electron is excited from the valence band to the conduction band it leaves a hole in the valence band.
- \* This excitation can occur by thermal energy, or either by phonon absorption, or both.
- \* A hole is a missing electron in a valence band, and from the effective mass approximation it has negative effective mass and this is seen in the equation of motion. And is like considering a positive charged particle.
- \* Without the crystal lattice, the hole cannot exist. It is an artifact of the periodic potential ( $E_p$ ) created by the crystal
- \* A hole is positively charged and can move, and for this is different from ionised atoms don't.

## Intrinsic carrier concentration

- ✦ An important quantity to look at it is the concentration of carriers or holes as a function of temperature, in terms of the band gap.

# Intrinsic carrier concentration

- ✱ An important quantity to look at it is the concentration of carriers or holes as a function of temperature, in terms of the band gap.
- ✱ We can evaluate the carrier concentrations in the conduction band using the concept of DOS and the fermi-Dirac distribution:

$$\begin{aligned} n &= \int_{E_c}^{+\infty} \text{Dos}(\epsilon) f(\epsilon) d\epsilon \\ &= 2 \left( \frac{m_e^* k_b T}{2\pi \hbar^2} \right)^{\frac{3}{2}} e^{\frac{E_F - E_c}{k_b T}} = N_c e^{\frac{E_F - E_c}{k_b T}} \end{aligned} \quad (1)$$

The  $N_c$  term is temperature dependent called the effective density of states of the conduction band;



# Intrinsic carrier concentration

- ✱ An important quantity to look at it is the concentration of carriers or holes as a function of temperature, in terms of the band gap.
- ✱ We can evaluate the carrier concentrations in the conduction band using the concept of DOS and the fermi-Dirac distribution:

$$\begin{aligned} n &= \int_{E_c}^{+\infty} \text{Dos}(\epsilon) f(\epsilon) d\epsilon \\ &= 2 \left( \frac{m_e^* k_b T}{2\pi \hbar^2} \right)^{\frac{3}{2}} e^{\frac{E_{F_i} - E_c}{k_b T}} = N_c e^{\frac{E_{F_i} - E_c}{k_b T}} \end{aligned} \quad (1)$$

The  $N_c$  term is temperature dependent called the effective density of states of the conduction band;

- ✱ And proceeding as above we obtaining the concentration  $p$  of holes in the valence band:

$$\begin{aligned} p &= \int_{-\infty}^{E_v} \text{Dos}(\epsilon) (1 - f(\epsilon)) d\epsilon \\ &= 2 \left( \frac{m_h^* k_b T}{2\pi \hbar^2} \right)^{\frac{3}{2}} e^{\frac{E_v - E_{F_i}}{k_b T}} = N_v e^{\frac{E_v - E_{F_i}}{k_b T}} \end{aligned} \quad (2)$$

where  $N_v$  is a temperature dependent constant called the effective density of states of the valence band.

# Intrinsic carrier concentration

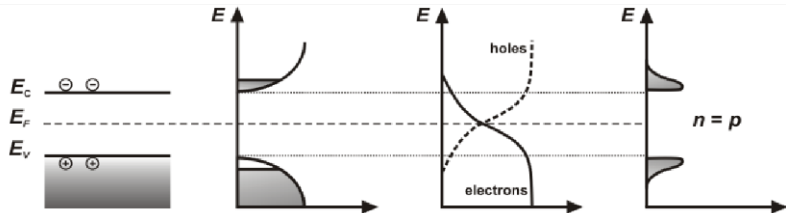


Figure 1: This figure shows the intrinsic behaviour of a semiconductor

- ✱ In physics we usually look for constant of the system, we obtain one of it if we multiply equation (1) and equation (2):

$$np = 4 \left( \frac{k_b T}{2\pi \hbar^2} \right)^3 (m_e^* m_h^*)^{\frac{3}{2}} e^{-\frac{E_g}{k_b T}} = n_i^2$$

This quantity does not depend on the Fermi energy  $E_F$ . In an intrinsic semiconductor  $n = p$ , because for each electron that has been excited from the valence band to the conduction band it leaves an hole in the valence band.

- ✱ In physics we usually look for constant of the system, we obtain one of it if we multiply equation (1) and equation (2):

$$np = 4 \left( \frac{k_b T}{2\pi \hbar^2} \right)^3 (m_e^* m_h^*)^{\frac{3}{2}} e^{-\frac{E_g}{k_b T}} = n_i^2$$

This quantity does not depend on the Fermi energy  $E_F$ . In an intrinsic semiconductor  $n = p$ , because for each electron that has been excited from the valence band to the conduction band it leaves a hole in the valence band.

- ✱ Then we have  $n = p = n_i$ . This is because the product of the electron and hole concentration is a constant, not depending on the impurity concentration at given temperature. The introduction of a small proportion of a suitable impurity to increase  $n$ , say, must decrease  $p$ .

- ✱ The property  $n = p$  in an intrinsic semiconductor gives us the possibility to estimate the so called intrinsic Fermi energy, equalling equation (1) and equation (2):

$$n = p$$
$$N_c e^{\frac{E_{F_i} - E_c}{k_b T}} = N_v e^{\frac{E_v - E_{F_i}}{k_b T}}$$
$$E_{F_i} = \frac{E_c + E_v}{2} + \frac{3}{4} k_b T \ln \left( \frac{m_h^*}{m_e^*} \right)$$

- ✱ The property  $n = p$  in an intrinsic semiconductor gives us the possibility to estimate the so called intrinsic Fermi energy, equalling equation (1) and equation (2):

$$n = p$$
$$N_c e^{\frac{E_{F_i} - E_c}{k_b T}} = N_v e^{\frac{E_v - E_{F_i}}{k_b T}}$$
$$E_{F_i} = \frac{E_c + E_v}{2} + \frac{3}{4} k_b T \ln \left( \frac{m_h^*}{m_e^*} \right)$$

- ✱ In an intrinsic semiconductor, the Fermi level is located close to the center of the band gap. Note that this quantity depends only on the band gap and on the electron and hole effective masses.

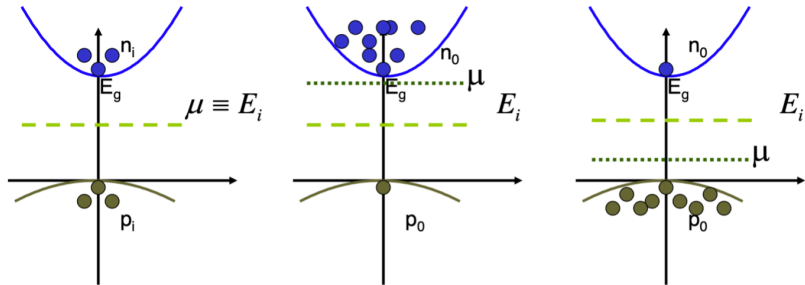
- ✱ In an extrinsic semiconductor, the general behaviour is slightly different, with the dopants fully ionized, there is an imbalance in the electron and hole concentration. This is reflected in the Fermi level position being shifted from the center of the band gap towards either the conduction band for an n-type or valence band for an p-type.

- \* In an extrinsic semiconductor, the general behaviour is slightly different, with the dopants fully ionized, there is an imbalance in the electron and hole concentration. This is reflected in the Fermi level position being shifted from the center of the band gap towards either the conduction band for an n-type or valence band for an p-type.
- \* The doped semiconductor has to satisfy the charge neutrality:

$$p + N_d^+ = n + N_a^-$$



# Intrinsic and Extrinsic semiconductor



**Figure 2:** The first plot shows the intrinsic case. The second figure shows the case of n-type doped material. The last one shows the case of p-type doped material.

- ✱ The carrier concentration is calculated as in equation (1), but this time we have a different Fermi energy, depending on the donor concentration.

$$n = N_c e^{\frac{E_F - E_c}{k_b T}} = N_c e^{\frac{E_{F_i} - E_c}{k_b T}} e^{\frac{E_F - E_i}{k_b T}} = n_i e^{\frac{E_F - E_i}{k_b T}} \quad (3)$$

## Fermi energy extrinsic semiconductor

- \* The carrier concentration is calculated as in equation (1), but this time we have a different Fermi energy, depending on the donor concentration.

$$n = N_c e^{\frac{E_F - E_c}{k_b T}} = N_c e^{\frac{E_{F_i} - E_c}{k_b T}} e^{\frac{E_F - E_{F_i}}{k_b T}} = n_i e^{\frac{E_F - E_{F_i}}{k_b T}} \quad (3)$$

- \* Now we have to determine this new Fermi energy. To do that we start from the charge neutrality (the semiconductor must be neutral overall):

$$p + N_d^+ = n \quad (4)$$

## Fermi energy extrinsic semiconductor

- ✱ The carrier concentration is calculated as in equation (1), but this time we have a different Fermi energy, depending on the donor concentration.

$$n = N_c e^{\frac{E_F - E_c}{k_b T}} = N_c e^{\frac{E_{F_i} - E_c}{k_b T}} e^{\frac{E_F - E_{F_i}}{k_b T}} = n_i e^{\frac{E_F - E_{F_i}}{k_b T}} \quad (3)$$

- ✱ Now we have to determine this new Fermi energy. To do that we start from the charge neutrality (the semiconductor must be neutral overall):

$$p + N_d^+ = n \quad (4)$$

- ✱ Where  $N_d^+$  is the number of ionised donors:

$$N_d^+ = Nd(1 - f(E_D)) = N_d \left( 1 - \frac{1}{1 + g^{-1} e^{\frac{(E_D - E_F)}{k_b T}}} \right) = \frac{N_d}{1 + g e^{\frac{(E_F - E_D)}{k_b T}}}$$

Where  $g$  is a factor that take into account the degeneracy of the level. We use  $g = 2$  to account the spin degeneracy.

## Fermi energy extrinsic semiconductor

- \* The carrier concentration is calculated as in equation (1), but this time we have a different Fermi energy, depending on the donor concentration.

$$n = N_c e^{\frac{E_F - E_c}{k_b T}} = N_c e^{\frac{E_{F_i} - E_c}{k_b T}} e^{\frac{E_F - E_i}{k_b T}} = n_i e^{\frac{E_F - E_i}{k_b T}} \quad (3)$$

- \* Now we have to determine this new Fermi energy. To do that we start from the charge neutrality (the semiconductor must be neutral overall):

$$p + N_d^+ = n \quad (4)$$

- \* Where  $N_d^+$  is the number of ionised donors:

$$N_d^+ = N_d (1 - f(E_D)) = N_d \left( 1 - \frac{1}{1 + g^{-1} e^{\frac{(E_D - E_F)}{k_b T}}} \right) = \frac{N_d}{1 + g e^{\frac{(E_F - E_D)}{k_b T}}}$$

Where  $g$  is a factor that take into account the degeneracy of the level. We use  $g = 2$  to account the spin degeneracy.

- \* The neutrality equations becomes:

$$n_i e^{\frac{E_{F_i} - E_F}{k_b T}} + \frac{N_d}{1 + 2 e^{\frac{E_F - E_D}{k_b T}}} = n_i e^{\frac{E_F - E_{F_i}}{k_b T}} \quad (5)$$

# Fermi energy extrinsic semiconductor

- \* An analytic solutions can only be found in special limiting cases. But it is possible to solve it numerically. In the case of silicon we have that:

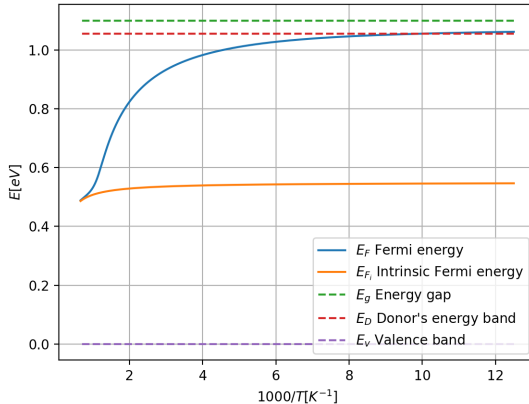


Figure 3: Caption

# Intrinsic and extrinsic carrier concentration

- ✳ once we obtained the Fermi energy we can compute the electron carrier concentration in a doped semiconductor using equation (3).

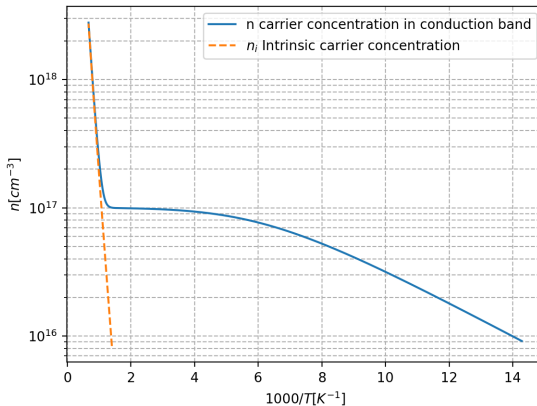


Figure 4: Caption

- ✦ From the previous figure we can define 3 regions, hence 3 different behaviour of the doped semiconductor:



- ✦ From the previous figure we can define 3 regions, hence 3 different behaviour of the doped semiconductor:
  - ✦ Low temperature (freeze-out region)

- ✦ From the previous figure we can define 3 regions, hence 3 different behaviour of the doped semiconductor:
  - ✦ Low temperature (freeze-out region)
  - ✦ Medium temperatures (Saturation region)

- ✦ From the previous figure we can define 3 regions, hence 3 different behaviour of the doped semiconductor:
  - ✦ Low temperature (freeze-out region)
  - ✦ Medium temperatures (Saturation region)
  - ✦ High temperature (Intrinsic range)

# Intrinsic and extrinsic carrier concentration

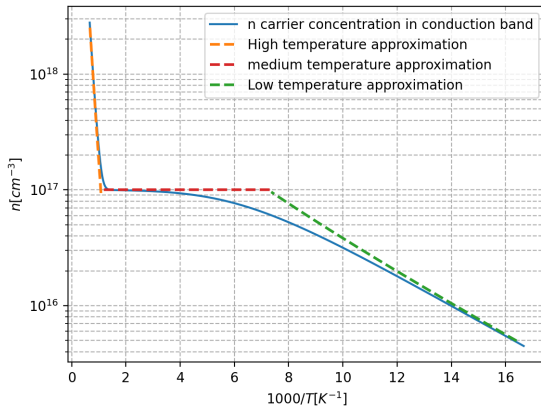


Figure 5: Caption

## Low temperature (freeze-out region)

- ✱ At absolute zero there are no ionized carriers. Valence band is full and the donor level is full and conduction band is empty.

## Low temperature (freeze-out region)

- ✱ At absolute zero there are no ionized carriers. Valence band is full and the donor level is full and conduction band is empty.
- ✱ As temperature is increased, electrons are excited from the valence band and the donor level to the conduction band. But since the valence band ionization energy is of the order of eV, at low temperature the number of electrons excited from it are negligible compared to the electrons from the donor level. Hence:

$$N_d^+ \approx n$$

This equation is simpler and it can be solved analytically. Find the fermi energy and then substitute it in the carrier concentration formula.

$$n = N_c e^{\frac{E_F - E_c}{k_b T}} = \sqrt{\frac{N_d N_c}{2}} e^{-\frac{E_c - E_D}{2k_b T}}$$

## Low temperature (freeze-out region)

- ✱ At absolute zero there are no ionized carriers. Valence band is full and the donor level is full and conduction band is empty.
- ✱ As temperature is increased, electrons are excited from the valence band and the donor level to the conduction band. But since the valence band ionization energy is of the order of eV, at low temperature the number of electrons excited from it are negligible compared to the electrons from the donor level. Hence:

$$N_d^+ \approx n$$

This equation is simpler and it can be solved analytically. Find the fermi energy and then substitute it in the carrier concentration formula.

$$n = N_c e^{\frac{E_F - E_c}{k_b T}} = \sqrt{\frac{N_d N_c}{2}} e^{-\frac{E_c - E_D}{2 k_b T}}$$

- ✱ For low temperatures we expect that the concentration of holes is much lower than that of electrons (given mainly from the donor), and becomes very low when  $N_d$  becomes very large, in fact  $p = \frac{n_i^2}{N_d^+} \ll N_d^+$ .

## Medium temperatures (Saturation region)

- ✦ At medium temperatures the energy  $k_b T$  is still not enough to promote an electron from the valence band to the conduction band, but all the electrons given by the donor have been promoted to the conduction band.



## Medium temperatures (Saturation region)

- ✦ At medium temperatures the energy  $k_b T$  is still not enough to promote an electron from the valence band to the conduction band, but all the electrons given by the donor have been promoted to the conduction band.
- ✦ Hence, we expect that the carrier concentration is flattened out at  $N_d$ , as it is shown in figure 5 .

## Medium temperatures (Saturation region)

- ✦ At medium temperatures the energy  $k_b T$  is still not enough to promote an electron from the valence band to the conduction band, but all the electrons given by the donor have been promoted to the conduction band.
- ✦ Hence, we expect that the carrier concentration is flattened out at  $N_d$ , as it is shown in figure 5 .
- ✦ There is a saturation range where the concentration is determined by doping and all the impurities are ionised: this is where usually lies the working range of the device.

- ✦ In the end, when we pass to high temperature regime we cannot ignore anymore the contributions from the intrinsic concentration  $n_i$ .

## High temperature (Intrinsic range)

- ✦ In the end, when we pass to high temperature regime we cannot ignore anymore the contributions from the intrinsic concentration  $n_i$ .
- ✦ If this is the case, in this temperature domain the energy is enough to promote an electron from the valence band to the conduction band.

## High temperature (Intrinsic range)

- ✱ In the end, when we pass to high temperature regime we cannot ignore anymore the contributions from the intrinsic concentration  $n_i$ .
- ✱ If this is the case, in this temperature domain the energy is enough to promote an electron from the valence band to the conduction band.
- ✱ Now the main contribution on the carrier concentration is given by  $n_i$  ( $n \approx N_d^+ + p \approx n_i$ ). In fact at too high temperature, the concentration is essentially given by the intrinsic one ( $n_i$ ), making the doping useless.