

Compatible/equilibrium and incompatible/non-equilibrium initial states in OpenGeoSys

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Contents

1	Preliminary remarks	2
2	Small deformation	2
2.1	Case A: Total stress form—compatible/equilibrium initial states	2
2.2	Case B: Incremental stress form—incompatible/non-equilibrium initial states	3
2.3	Test case	3
3	[WIP] HM process	4
4	[WIP] RM process	4

1 Preliminary remarks

We distinguish between several concepts:

- **Constitutive reference states.** These are essentially material parameters and only active constitutively. Example: $T_{\text{ref}}, p_{\text{ref}}, \varrho_{\text{ref}}$ in order to calculate

$$\varrho = \varrho_{\text{ref}} + \beta_T(T - T_{\text{ref}}) + \beta_p(p - p_{\text{ref}}) \quad (1)$$

- **Compatible/equilibrium initial states.** These are initial values for primary and secondary variables that enter the weak form and are in equilibrium with the boundary conditions.
- **Incompatible/non-equilibrium initial states.** These are initial states which are not explicitly in equilibrium with any boundary conditions. They are not considered in the weak form but affect the constitutive behaviour of the materials.

2 Small deformation

A total stress is

Discrete linearized weak form (NR scheme):

$$\int_{\Omega} \mathbf{B}^T \mathbf{C} \mathbf{B} d\Omega \Delta \hat{\mathbf{u}} = \int_{\partial\Omega_N} \mathbf{N}^T \bar{\mathbf{t}} d\Gamma + \int_{\Omega} \mathbf{N}^T \varrho \mathbf{b} d\Omega - \int_{\Omega} \mathbf{B}^T (\boldsymbol{\sigma} - \boldsymbol{\sigma}_{0\text{neq}}) d\Omega \quad (2)$$

Where the last term has been included to account for incompatible initial states (non-equilibrium states). The constitutive model integration proceeds as follows:

$$\boldsymbol{\sigma} = \boldsymbol{\sigma}_{0\text{eq}} + \boldsymbol{\sigma}_{0\text{neq}} + \int_{t_0}^t \dot{\boldsymbol{\sigma}}(\dot{\boldsymbol{\epsilon}}, \boldsymbol{\sigma}, \boldsymbol{\kappa}, \dots) d\bar{t} \quad (3)$$

In other words, the constitutive integration always gets the total stress (including initial equilibrium and non-equilibrium contributions) as an argument.

2.1 Case A: Total stress form—compatible/equilibrium initial states

The user supplies an initial state which is in equilibrium with the initial tractions. All subsequent load changes are total loads in equilibrium with the total stress.

$$\int_{\Omega} \mathbf{B}^T \mathbf{C} \mathbf{B} d\Omega \Delta \hat{\mathbf{u}} = \int_{\partial\Omega_N} \mathbf{N}^T \bar{\mathbf{t}} d\Gamma + \int_{\Omega} \mathbf{N}^T \varrho \mathbf{b} d\Omega - \int_{\Omega} \mathbf{B}^T \boldsymbol{\sigma} d\Omega, \quad \boldsymbol{\sigma}_{0\text{neq}} = \mathbf{0} \quad (4)$$

such that a deformation-free state at $t = t_0$ requires

$$\int_{\partial\Omega_N} \mathbf{N}^T \bar{\mathbf{t}}(t = t_0) d\Gamma + \int_{\Omega} \mathbf{N}^T \varrho(t = t_0) \mathbf{b}(t = t_0) d\Omega - \int_{\Omega} \mathbf{B}^T \boldsymbol{\sigma}_{0\text{eq}} d\Omega = \mathbf{0} \quad (5)$$

2.2 Case B: Incremental stress form—incompatible/non-equilibrium initial states

The user supplies an initial state which is *not* in equilibrium with the initial tractions. Since these stresses are not considered in the equilibrium formulation, all tractions supplied are considered as *additional* loads

$$\int_{\Omega} \mathbf{B}^T \mathbf{C} \mathbf{B} d\Omega \Delta \hat{\mathbf{u}} = \int_{\partial\Omega_N} \mathbf{N}^T \bar{\mathbf{t}} d\Gamma + \int_{\Omega} \mathbf{N}^T \varrho \mathbf{b} d\Omega - \int_{\Omega} \mathbf{B}^T (\boldsymbol{\sigma} - \boldsymbol{\sigma}_{0\text{neq}}) d\Omega \quad (6)$$

such that a deformation-free state at $t = t_0$ requires

$$\int_{\partial\Omega_N} \mathbf{N}^T \bar{\mathbf{t}}(t = t_0) d\Gamma + \int_{\Omega} \mathbf{N}^T \varrho \mathbf{b}(t = t_0) d\Omega = \mathbf{0} \quad (7)$$

i.e. that no (additional) external loads are acting.

Remark on the body-force term: Typically, the specific body force will be constant, e.g. gravity. Therefore, additional *specific* body forces will not occur. However, density changes can cause changes in the body forces:

$$\varrho \mathbf{b} = (\varrho_0 + \varrho_{\text{incr}})(\mathbf{b}_0 + \mathbf{b}_{\text{incr}}) \quad \text{with} \quad \varrho_{\text{incr}} = \int_{t=t_0}^{t_1} \dot{\varrho} d\bar{t} \quad (8)$$

Thus, even if a constant specific source term (e.g. gravity) is used, i.e. $\mathbf{b}_{\text{incr}} = \mathbf{0}$, one would not have a vanishing body force term in Eq. (6). Instead, the acting force causing deformation beyond the initial state would consist of a specific body-force increment acting on the total density and a body force stemming from the density alteration:

$$\varrho \mathbf{b}_{\text{incr}} + \varrho_{\text{incr}} \mathbf{b}_0 = \varrho \mathbf{b} - \varrho_0 \mathbf{b}_0 \quad (9)$$

This is essential in unsaturated settings where during de-/resaturation strong density changes occur. In the small deformation process we assume this to be negligible. Thus, one would arrive at a typical setting that looks like this:

$$\int_{\Omega} \mathbf{B}^T \mathbf{C} \mathbf{B} d\Omega \Delta \hat{\mathbf{u}} = \int_{\partial\Omega_N} \mathbf{N}^T \bar{\mathbf{t}} d\Gamma - \int_{\Omega} \mathbf{B}^T (\boldsymbol{\sigma} - \boldsymbol{\sigma}_{0\text{neq}}) d\Omega \quad (10)$$

such that a deformation-free state at $t = t_0$ requires

$$\int_{\partial\Omega_N} \mathbf{N}^T \bar{\mathbf{t}}(t = t_0) d\Gamma = \mathbf{0} \quad (11)$$

i.e. that no (additional) external tractions are acting and body forces are zero.

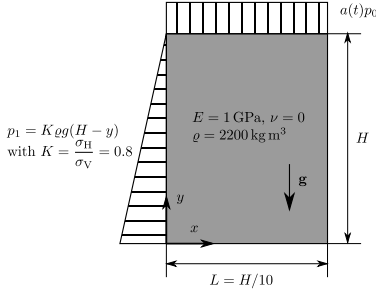
2.3 Test case

The test cases illustrated in Fig. 1 using the above formulations have been added to Tests/Data/Mechanics/Linear/Initial_States.

Case 1: Calculate into initial state

$$t_1 = [0, 2] \text{ s} \quad \sigma_{0\text{eq}} = \sigma_{0\text{neq}} = \mathbf{0}$$

$$a(t) = 0.5t$$

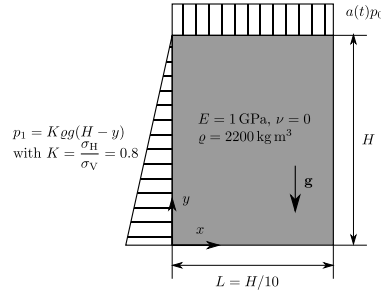


Case 2: Calculate from equilibrium initial state

$$t_2 = [1, 2] \text{ s} \quad \sigma_{0\text{eq}} = -p_1 \mathbf{e}_x \otimes \mathbf{e}_x - (0.5p_0 + \rho g(H-y)) \mathbf{e}_y \otimes \mathbf{e}_y$$

$$a(t) = 0.5t$$

$$\sigma_{0\text{neq}} = \mathbf{0}$$



Case 3: Calculate from non-equilibrium initial state

$$t_2 = [1, 2] \text{ s} \quad \sigma_{0\text{neq}} = -p_1 \mathbf{e}_x \otimes \mathbf{e}_x - (0.5p_0 + \rho g(H-y)) \mathbf{e}_y \otimes \mathbf{e}_y$$

$$a(t) = 0.5(t-1)$$

$$\sigma_{0\text{eq}} = \mathbf{0}$$

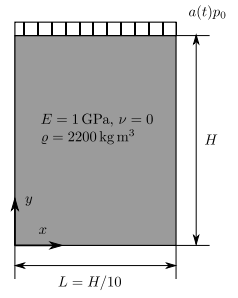


Figure 1: Test cases for the small deformation process.

3 [WIP] HM process

The linearised discretized weak forms of the two governing equations read:

$$\begin{aligned} \int_{\Omega} \mathbf{B}_u^T \mathbf{C} \mathbf{B}_u \, d\Omega \, \Delta \hat{\mathbf{u}} - \int_{\Omega} \mathbf{B}_u^T \alpha_B \mathbf{I} \mathbf{N}_p \, d\Omega \, \Delta \hat{\mathbf{p}} = \\ = \int_{\partial\Omega_t} \mathbf{N}_u^T \bar{\mathbf{t}} \, d\Gamma + \int_{\Omega} \mathbf{N}_u^T \rho \mathbf{b} \, d\Omega - \int_{\Omega} \mathbf{B}_u^T (\boldsymbol{\sigma}' - \boldsymbol{\sigma}_{0\text{neq}}) \, d\Omega + \int_{\Omega} \mathbf{B}_u^T \alpha_B (p - p_{0\text{neq}}) \mathbf{I} \, d\Omega \end{aligned} \quad (12)$$

$$\begin{aligned} \int_{\Omega} \mathbf{N}_p^T \frac{\alpha_B}{\Delta t} \mathbf{I}^T \mathbf{B}_u \, d\Omega \, \Delta \hat{\mathbf{u}} + \left[\int_{\Omega} \mathbf{N}_p^T \frac{S}{\Delta t} \mathbf{N}_p \, d\Omega + \int_{\Omega} \nabla \mathbf{N}_p^T \frac{\mathbf{K}}{\mu_{\text{LR}}} \nabla \mathbf{N}_p \, d\Omega \right] \Delta \hat{\mathbf{p}} = \\ = \int_{\partial\Omega_w} \mathbf{N}_p^T \dot{m}_n \, d\Gamma - \int_{\Omega} \mathbf{N}_p^T S \mathbf{N}_p \, d\Omega (\hat{\mathbf{p}})'_S - \int_{\Omega} \mathbf{N}_p^T \alpha_B \mathbf{I}^T \mathbf{B}_u \, d\Omega (\hat{\mathbf{u}})'_S - \int_{\Omega} \nabla \mathbf{N}_p^T \frac{\mathbf{K}}{\mu_{\text{LR}}} \nabla \mathbf{N}_p \, d\Omega (\hat{\mathbf{p}} - \hat{\mathbf{p}}_{0\text{neq}}) + \\ + \int_{\Omega} \nabla \mathbf{N}_p^T \rho_{\text{LR}} \frac{\mathbf{K}}{\mu_{\text{LR}}} \mathbf{b} \, d\Omega \end{aligned} \quad (13)$$

Remark on the body-force term: Similar rationale as described under small deformation process in current implementation.

4 [WIP] RM process

to be done ...

Remark on the body-force term: Here, this term is crucial (uplift upon saturation increase).