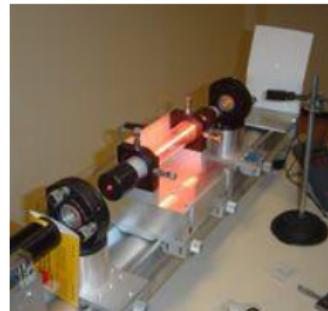


# The LASER

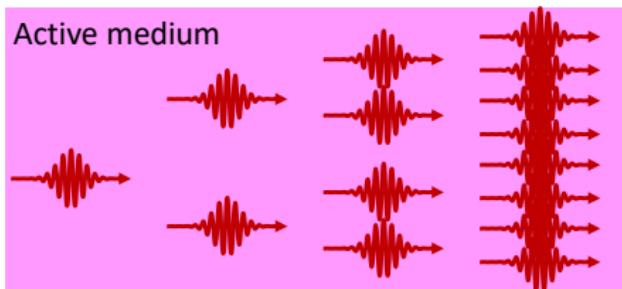
## Light Amplification by Stimulated Emission of Radiation

A bit of history....

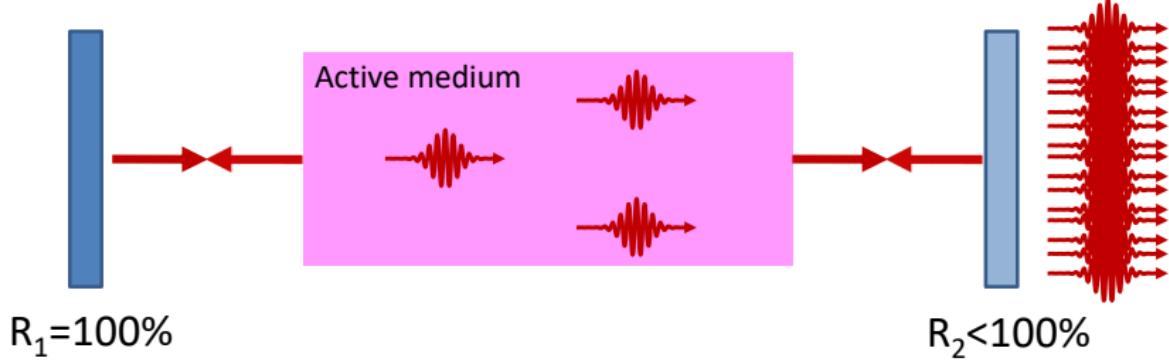
- 1917: A. Einstein studies the processes of spontaneous and stimulated emission;
- 1939 (and after the war in 1951): the Russian scientist V. A. Fabrikant suggests to use the stimulated emission process to amplify em radiation;
- 1952: the Russian scientists N.G. Basov e A.M. Prokhorov and the American physicist C.H. Townes suggest to use Fabrikant's idea in the microwaves range;
- 1954: the first MASER was built (by C. Townes and A. Schawlow);
- 1960: the American physicist Theodore Maiman put in operation the first LASER (ruby); A. Javan, W. Bennet and D. Harriott realize the first He-Ne laser at Bell's Labs;
- 1962-1963: the first semiconductors lasers were built.



## Laser as an AMPLIFIER

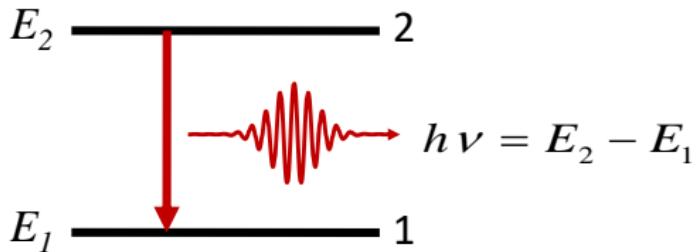


## Laser as an OSCILLATOR

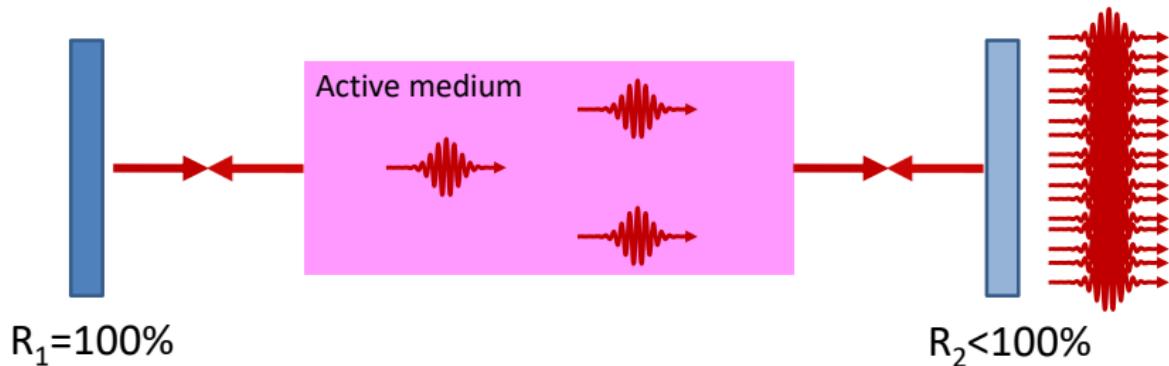


- Monochromaticity
- Directionality
- Brilliance
- Spatial and temporal coherence
- Ultra-short pulse generation

## ■ Monochromaticity

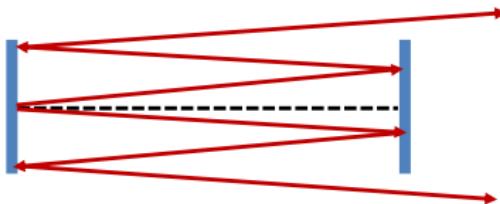


Only resonance frequencies of the laser feedback cavity are supported.

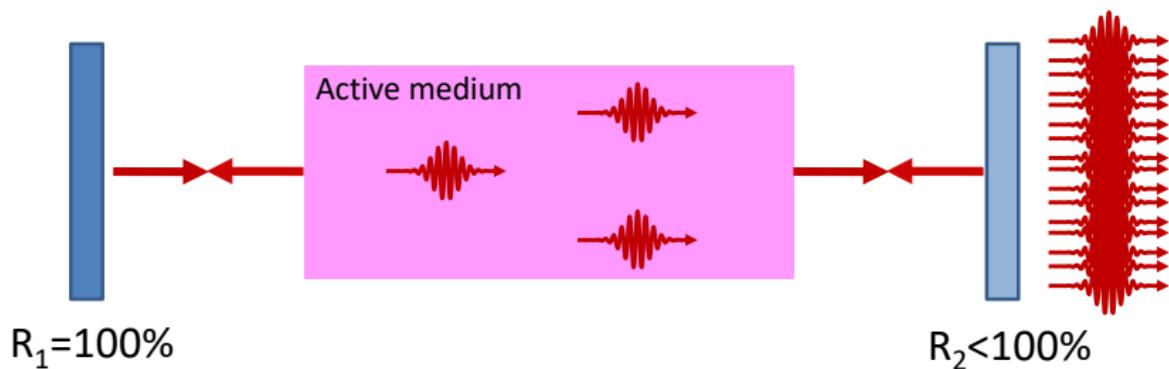


## ■ Directionality

Open feedback cavities: only modes that propagate along the cavity axis are sustained.



Only resonance frequencies of the laser feedback cavity are supported.



## ■ Brilliance

Emitted power per unit of surface ( $\Sigma$ ) and unit of solid angle ( $\Omega$ ):

$$B = \frac{P}{\Sigma \Omega} = \frac{I}{\Omega}$$

Owing to the high directionality, a laser beam of limited power (mW) has a brilliance of several orders of magnitude higher than the most powerful conventional sources.

# Thermal radiation

## SPECTRAL RADIANCE (Specific emissive power):

Emitted energy per unit of time, unit of surface and unit of solid angle, at frequency  $\nu$

$$R_\nu(T) = \frac{2\pi}{c^2} \frac{h\nu^3}{e^{h\nu/kT} - 1}$$

$$h = 6.63 \cdot 10^{-34} \text{ Js} \quad \text{Planck's constant}$$

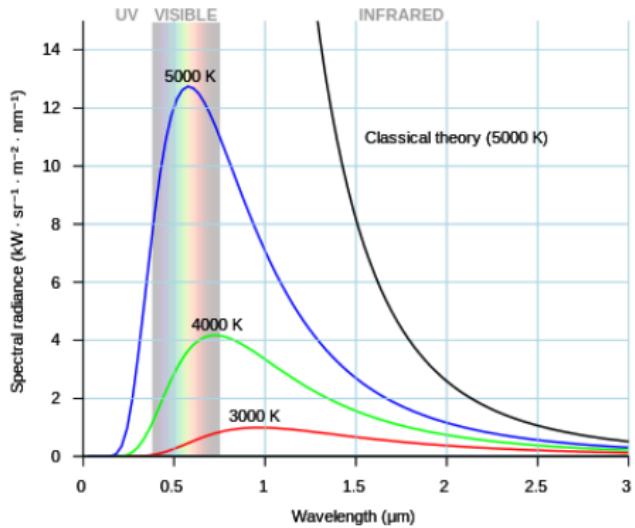
$$k = 1.38 \cdot 10^{-23} \text{ J/K} \quad \text{Boltzmann's constant}$$

## STEFAN-BOLTZMANN's LAW:

(Total emitted power per unit of surface: intensity)

$$I = \int_0^\infty R_\nu(T) d\nu \int d\Omega = e\sigma T^4$$

$$\text{WIEN's LAW: } \lambda_{\max} T = \text{cost} = 2898 \text{ } \mu\text{m K}$$



$e$  = emissivity ( $e \leq 1$ )

$$\sigma = 5.67 \cdot 10^{-8} \text{ W/m}^2 \text{K}^4$$

Stefan's  
constant

Conventional sources: Intensity  $\sim 10^4$  K

LASER: Intensity  $\sim 10^6$ - $10^7$  K



# Spatial coherence

Consider two points  $P_1, P_2$  of a wavefront at  $t = 0$ .

By definition of wavefront, the phase difference between  $P_1, P_2$  is null at  $t = 0$ :

$$\Delta\phi_{P_1P_2}(t = 0) = 0$$

If  $\Delta\phi_{P_1P_2}(t) = 0 \quad \forall t > 0$

$\Rightarrow$  **Perfect coherence between  $P_1$  and  $P_2$**

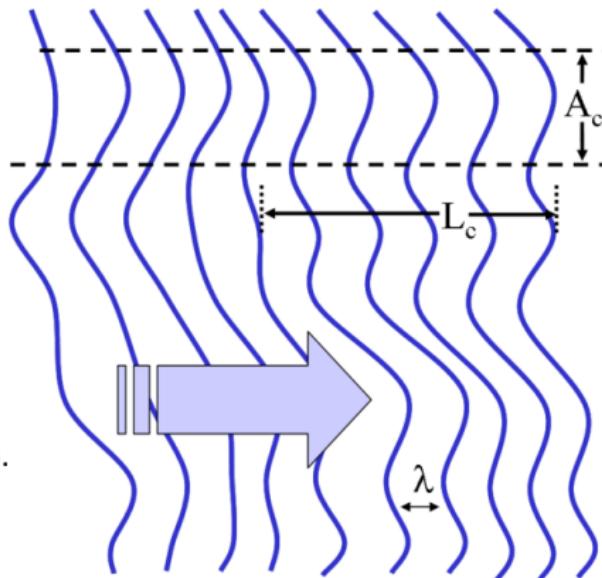
If this happens  $\forall (P_1, P_2)$ :

$\Rightarrow$  **Perfect spatial coherence**

Typically for a given  $P_1, P_2$  has to be within a given area  $A_C$  to have a good phase correlation.

$\Rightarrow$  **Partial spatial coherence (within  $A_C$ )**

$A_C$  = coherence area



$L_c$  = coherence length : propagation distance within which it is maintained a given degree of spatial coherence.

### Van-Cittert-Zernike theorem

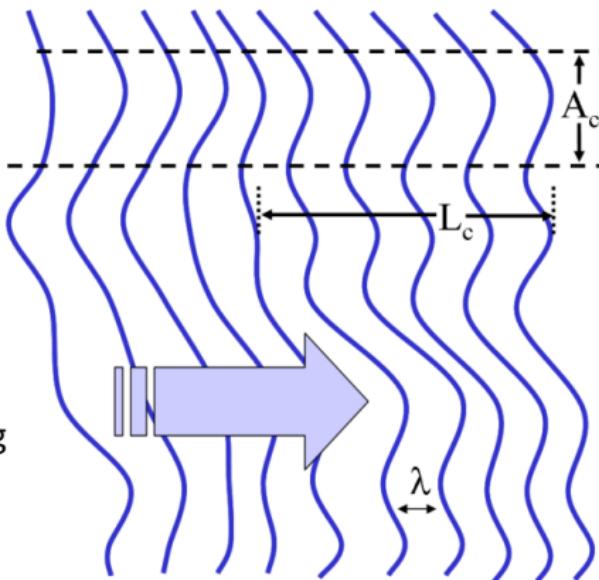
$$A_c = \frac{D^2 \lambda^2}{\pi d^2}$$

$D$  = distance from the source

$d$  = diameter of the source

⇒ Wavefronts tend to smooth-out going far away from the source

$A_c$  = coherence area



# Temporal coherence

Consider the electric field of an em wave in a given point  $P$  at time  $t$  and  $t + \tau$

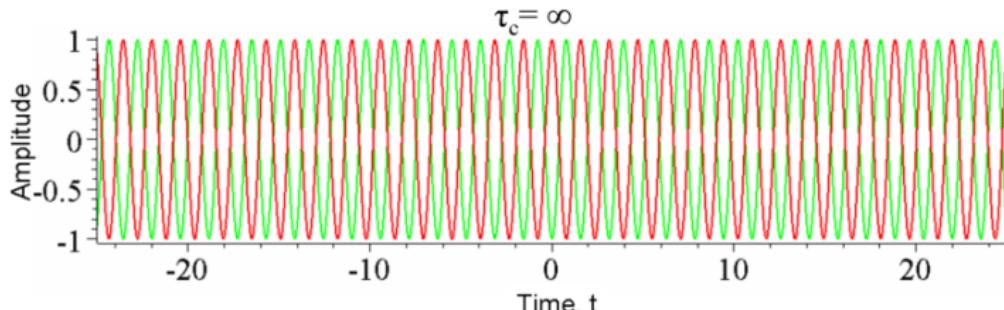
If for a given  $\tau$ , the phase difference between the two electric fields is the same  $\forall t$ :

⇒ **temporal coherence within time  $\tau$**

If this happens  $\forall \tau$  ⇒ **Perfect temporal coherence**

If it happens for  $0 < \tau < \tau_C$  ⇒ **Partial temporal coherence** (within  $\tau_C$ )

$$\tau_C = \text{coherence time} \quad L_C = c \tau_C$$



Consider the electric field of an em wave in a given point  $P$  at time  $t$  and  $t + \tau$

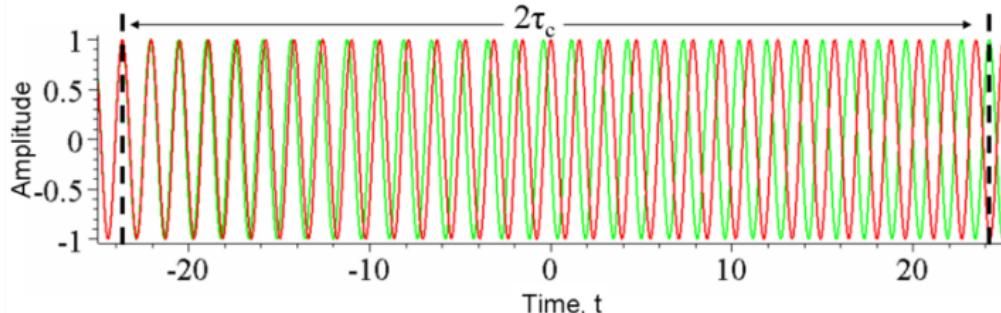
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⇒ **temporal coherence within time  $\tau$**

If this happens  $\forall \tau$  ⇒ **Perfect temporal coherence**

If it happens for  $0 < \tau < \tau_c$  ⇒ **Partial temporal coherence** (within  $\tau_c$ )

$$\tau_c = \text{coherence time} \quad L_c = c \tau_c$$



**Spatial and temporal coherence are independent concepts:** a wave can have perfect spatial coherence but partial temporal coherence, and viceversa.

**Spatial coherence** is associated with the wave's properties **transverse** to the direction of propagation.

Spatial coherence is associated with a **distribution of propagation vectors**  $\vec{k}$  associated with the wave, i.e., with a departure of the wave from the ideal plane wave.

**Temporal coherence** is associated with the wave's properties **along** the direction of propagation.

It is associated with the **frequency distribution** of the source.

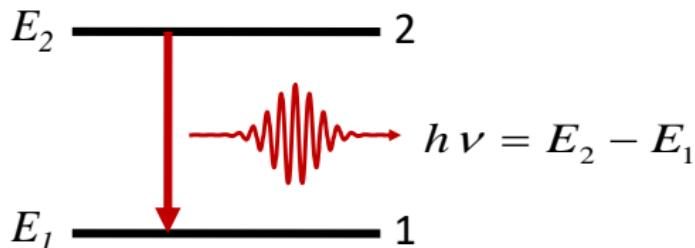
$$\tau_c \Delta\nu \approx 1 \quad \Rightarrow \quad \tau_c \approx \frac{1}{\Delta\nu} = \frac{\lambda^2}{c\Delta\lambda}$$

- Ultra-short pulse generation

It is not an intrinsic property of any laser, but with special techniques (as the mode-locking) some lasers can produce ultra-short pulses down to fs duration.

The capability to produce ultra-short pulses, that implies the concentration of energy in time, can be considered as the counterpart of monochromaticity, that implies the concentration of energy in wavelength.

# Spontaneous emission



Planck's constant

$$h = 6.63 \cdot 10^{-34} \text{ Js}$$

$$\left( \frac{dN_2}{dt} \right)_{sp} = -A_{21} N_2$$

Einstein's coefficient for spontaneous emission

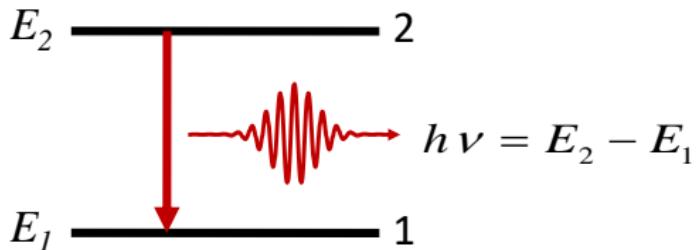
$$\tau_{sp} = \frac{1}{A_{21}}$$

$$\frac{dN_2}{dt} = -\frac{N_2}{\tau_{sp}} - \frac{N_2}{\tau'_{sp}} - \frac{N_2}{\tau_{nr}} = -\frac{N_2}{\tau}$$

$$N_2(t) = N_{20} e^{-t/\tau}$$

Lifetime of the level  $E_2$

## Spontaneous emission



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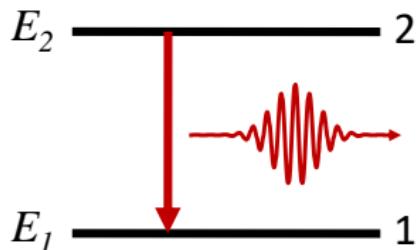
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$$N_2(t) = N_{20} e^{-t/\tau}$$

$$P(t) = \frac{N_2(t) h \nu}{\tau_{sp}} \propto e^{-t/\tau}$$

Lifetime of the level  $E_2$

## Spontaneous emission



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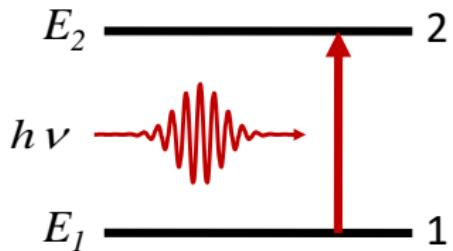
Spontaneous emission lifetime

$$\tau_{sp} = \frac{1}{A_{21}}$$

## Fluorescence Quantum Yield

$$\Phi_{21} = \frac{\text{number of photons with energy } h\nu = E_2 - E_1}{\text{number of atoms at level } E_2} = \frac{\tau}{\tau_{sp}}$$

# Absorption



Planck's constant

$$h = 6.63 \cdot 10^{-34} \text{ Js}$$

$$\frac{dN_1}{dt} = -W_{12} N_1$$

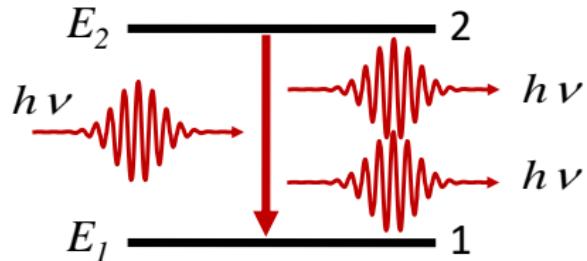
Photon flux  $F = \frac{I}{h\nu}$

$$W_{12} = B_{12} n(\nu) h \nu = B_{12} \rho(\nu) = \sigma_{12} F$$

Einstein's coefficient for  
absorption

Absorption  
cross-section

## Stimulated emission



Planck's constant

$$h = 6.63 \cdot 10^{-34} \text{ Js}$$

$$\frac{dN_2}{dt} = -W_{21}N_2$$

Photon flux  $F = \frac{I}{h\nu}$

$$W_{21} = B_{21}n(\nu)h\nu = B_{21}\rho(\nu) = \sigma_{21}F$$

Einstein's coefficient for  
stimulated emission

Stimulated emission  
cross-section