

Boundary conditions: continuity of tangential components at the interface

TE (s) polarization

$$\begin{cases} E_i + E_r = E_t \\ -B_i \cos\theta_i + B_r \cos\theta_r = -B_t \cos\theta_t \end{cases}$$

Fresnel's reflection and transmission coefficients

$$r_s = \left(\frac{E_r}{E_i} \right)_{TE} \quad t_s = \left(\frac{E_t}{E_i} \right)_{TE}$$

TM (p) polarization

$$\begin{cases} B_i - B_r = B_t \\ E_i \cos\theta_i + E_r \cos\theta_r = E_t \cos\theta_t \end{cases}$$

$$r_p = \left(\frac{E_r}{E_i} \right)_{TM} \quad t_p = \left(\frac{E_t}{E_i} \right)_{TM}$$

$$\theta_i = \theta_r \quad B = \frac{E}{v} = \frac{E}{c} n$$

Reflection coefficients

$$n = \frac{n_2}{n_1}$$

$$r_s = \left(\frac{E_r}{E_i} \right)_{TE} = \frac{n_1 \cos \theta_i - n_2 \cos \theta_t}{n_1 \cos \theta_i + n_2 \cos \theta_t} = \frac{\cos \theta_i - n \cos \theta_t}{\cos \theta_i + n \cos \theta_t}$$

$$r_p = \left(\frac{E_r}{E_i} \right)_{TM} = \frac{n_1 \cos \theta_t - n_2 \cos \theta_i}{n_1 \cos \theta_t + n_2 \cos \theta_i} = \frac{-n \cos \theta_i + \cos \theta_t}{n \cos \theta_i + \cos \theta_t}$$

Transmission coefficients

$$t_s = \left(\frac{E_t}{E_i} \right)_{TE} = \frac{2n_1 \cos \theta_i}{n_1 \cos \theta_i + n_2 \cos \theta_t} = \frac{2 \cos \theta_i}{\cos \theta_i + n \cos \theta_t}$$

$$t_p = \left(\frac{E_t}{E_i} \right)_{TM} = \frac{2n_1 \cos \theta_i}{n_1 \cos \theta_t + n_2 \cos \theta_i} = \frac{2 \cos \theta_i}{n \cos \theta_i + \cos \theta_t}$$

Fresnel's formulas

Snell's law

$$n = \frac{n_2}{n_1} = \frac{\sin \theta_i}{\sin \theta_t}$$

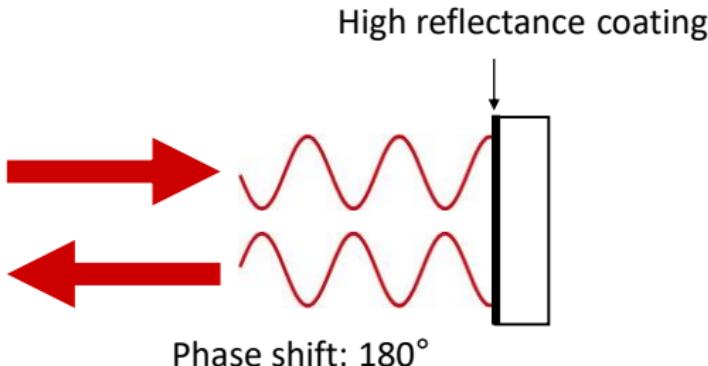
$$r_s = -\frac{\sin(\theta_i - \theta_t)}{\sin(\theta_i + \theta_t)} \quad t_s = \frac{2 \cos \theta_i \sin \theta_t}{\sin(\theta_i + \theta_t)}$$

$$r_p = -\frac{tg(\theta_i - \theta_t)}{tg(\theta_i + \theta_t)} \quad t_p = \frac{2 \cos \theta_i \sin \theta_t}{\sin(\theta_i + \theta_t) \cos(\theta_i - \theta_t)}$$

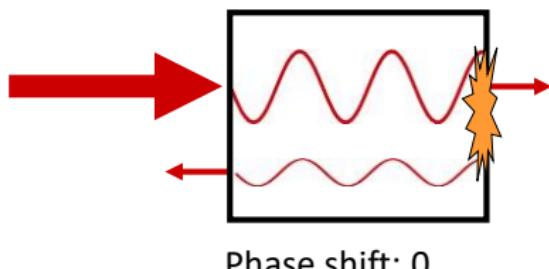
At normal incidence ($\theta_i = \theta_r = \theta_t = 0$):

$$r_s = r_p = \frac{E_r}{E_i} = \frac{1-n}{1+n} \quad t_s = t_p = \frac{E_t}{E_i} = \frac{2}{1+n}$$

$r < 0$ ($n > 1$, es. from air to glass) indicates that the phase of the reflected wave is varied of 180° with respect to the incident wave.



Slightly increasing the intensity of a laser beam incident on a glass window, where will the window get damaged earlier (on the front or on the back)?

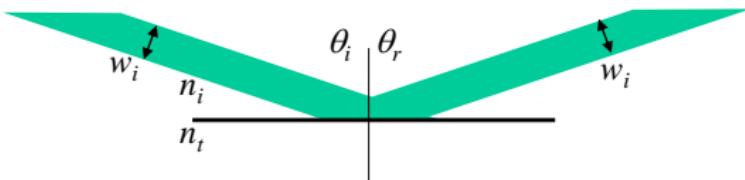


$$r = \frac{1-n}{1+n} = \frac{1-(1/1.5)}{1+(1/1.5)} = 0.2 \quad E_r = r E_i$$

$$I = |E_i + E_r|^2 = |(1+r)E_i|^2 = \\ = (1.2)^2 I_i = 1.44 I_i \quad 44\% \text{ higher!}$$

Reflectance (R)

$$I = \frac{1}{2} n \varepsilon_0 c |E|^2$$



$$\frac{A_r}{A_i} = \frac{w_r}{w_i} = 1$$

$$R = \frac{I_r A_r}{I_i A_i} = \left| \frac{E_r}{E_i} \right|^2 = |r|^2$$

$$R_s = \left| \frac{E_r}{E_i} \right|_{TE}^2 = |r_s|^2$$

$$R_p = \left| \frac{E_r}{E_i} \right|_{TM}^2 = |r_p|^2$$

At normal incidence ($\theta_i = \theta_r = \theta_t = 0^\circ$)

$$R_s = R_p = R = \left(\frac{1-n}{1+n} \right)^2$$

From air ($n_1 = 1$) to glass ($n_2 = 1.5$) :
 $R=0.04$ (4%)

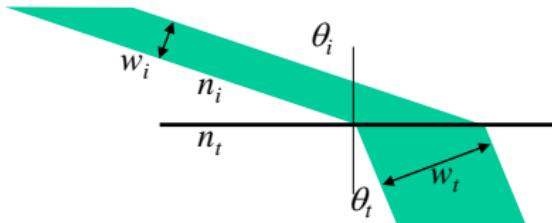
At grazing incidence ($\theta_i \sim 90^\circ$)

$$R_s = R_p = R = 1$$

R is independent of n.

Transmittance (T)

$$I = \frac{1}{2} n \varepsilon_0 c |E|^2$$



$$\frac{A_t}{A_i} = \frac{w_t}{w_i} = \frac{\cos \theta_t}{\cos \theta_i}$$

$$T = \frac{I_t A_t}{I_i A_i} = \frac{n_2}{n_1} \left| \frac{E_t}{E_i} \right|^2 \frac{w_t}{w_i} = n \left| t \right|^2 \frac{\cos \theta_t}{\cos \theta_i}$$

At normal incidence ($\theta_i = \theta_r = \theta_t = 0^\circ$)

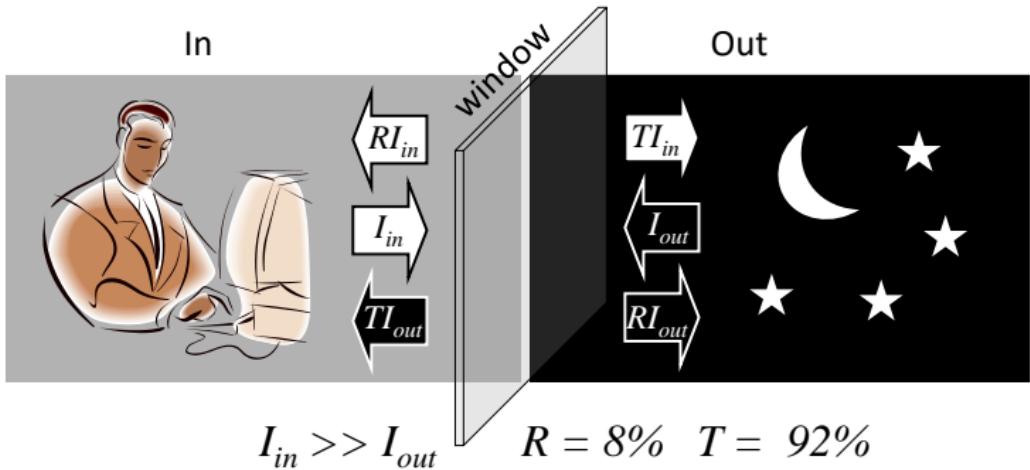
For any angle of incidence:

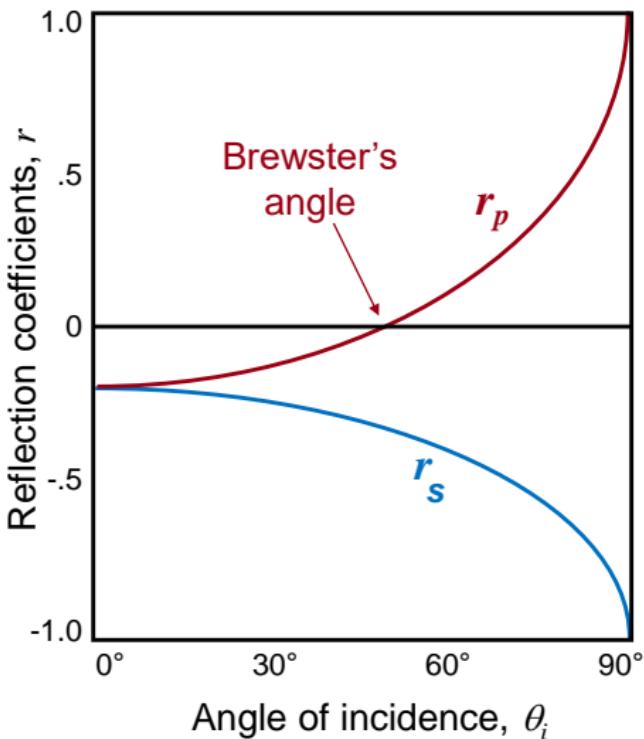
$$T_s = T_p = T = n \left(\frac{2}{1+n} \right)^2$$

$$R + T = 1$$

From air ($n_1 = 1$) to glass ($n_2 = 1.5$): $T=0.96$ (96%)

At night the windows seem mirrors when you are in an illuminated room!



(air-glass) $n > 1$ 

$$r_p = -\frac{\operatorname{tg}(\theta_i - \theta_t)}{\operatorname{tg}(\theta_i + \theta_t)} = 0$$

$$\theta_B + \theta_t = \frac{\pi}{2}$$

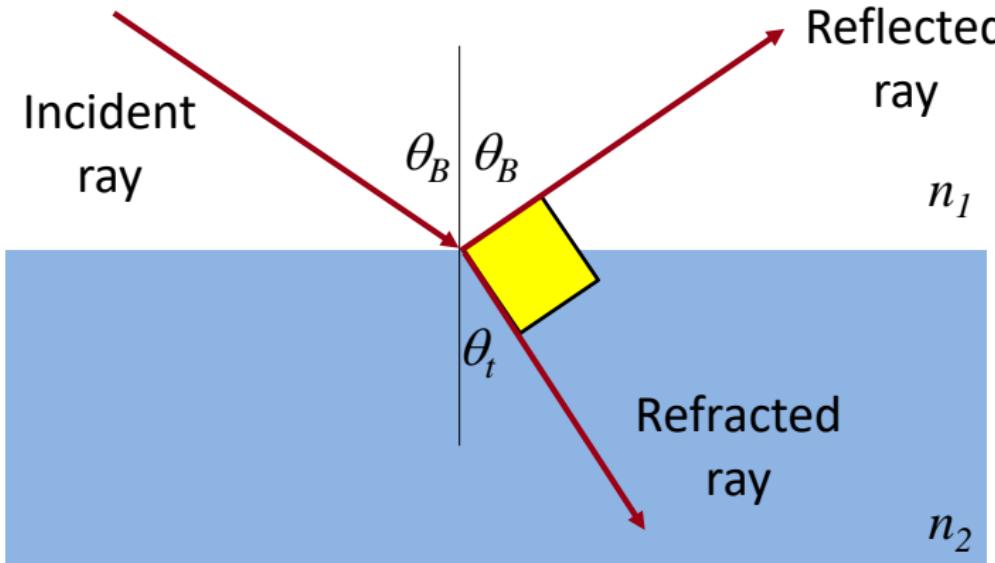
Brewster's angle

$$\theta_B = \operatorname{arctg}\left(\frac{n_2}{n_1}\right)$$

$$\text{air-glass} \quad \theta_B = \operatorname{arctg}\left(\frac{1.5}{1}\right) \cong 57^\circ$$

$$\text{glass-air} \quad \theta_B = \operatorname{arctg}\left(\frac{1}{1.5}\right) \cong 33^\circ$$

Brewster's angle

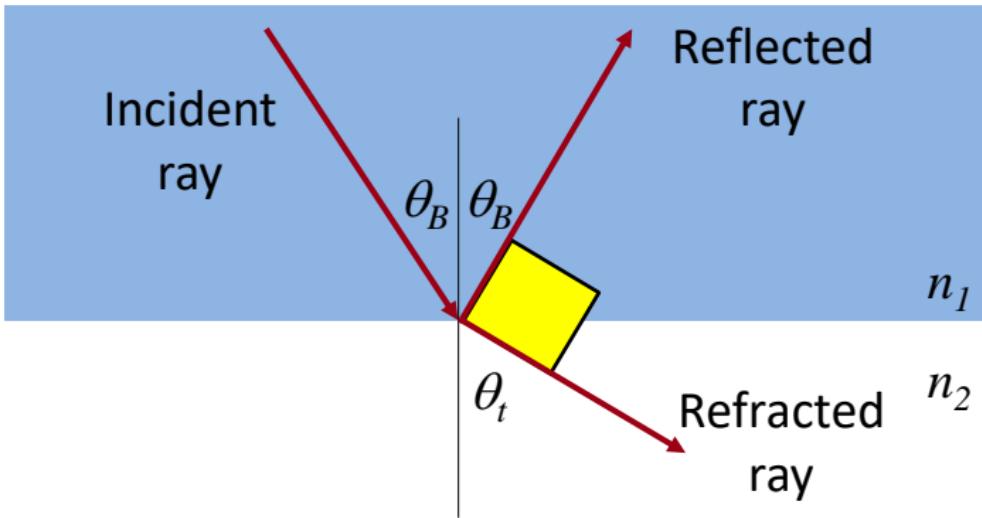


Brewster's
angle

$$\theta_B = \arctg\left(\frac{n_2}{n_1}\right) = \arctg\left(\frac{1.5}{1}\right) = 57^\circ$$

The reflected and refracted rays are perpendicular!

Brewster's angle

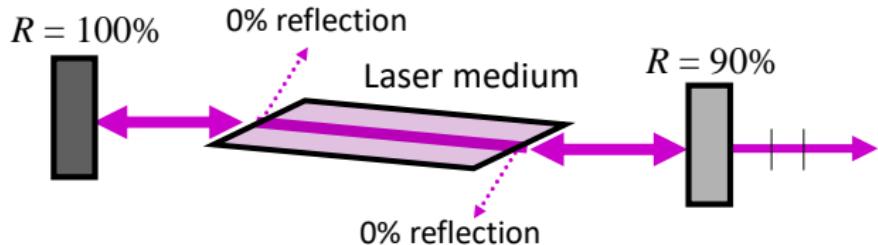
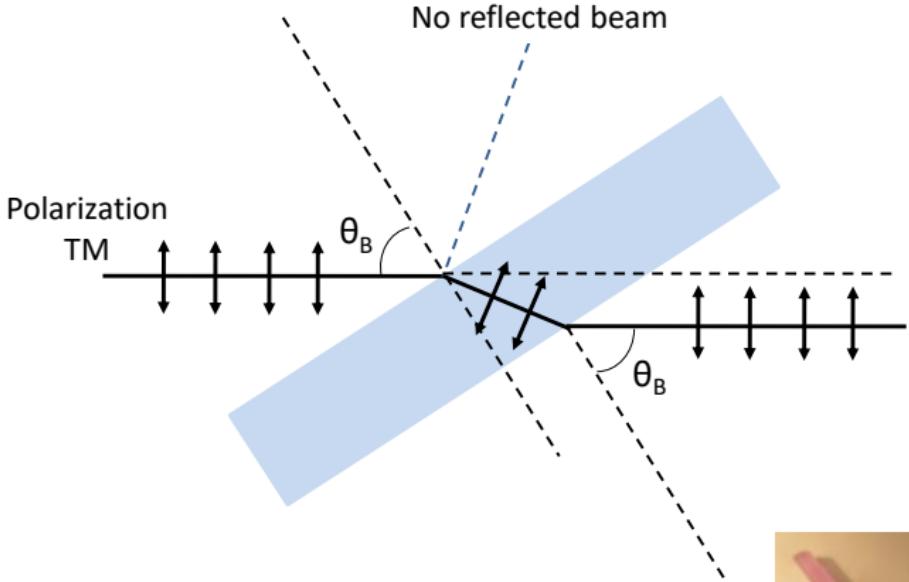


Brewster's
angle

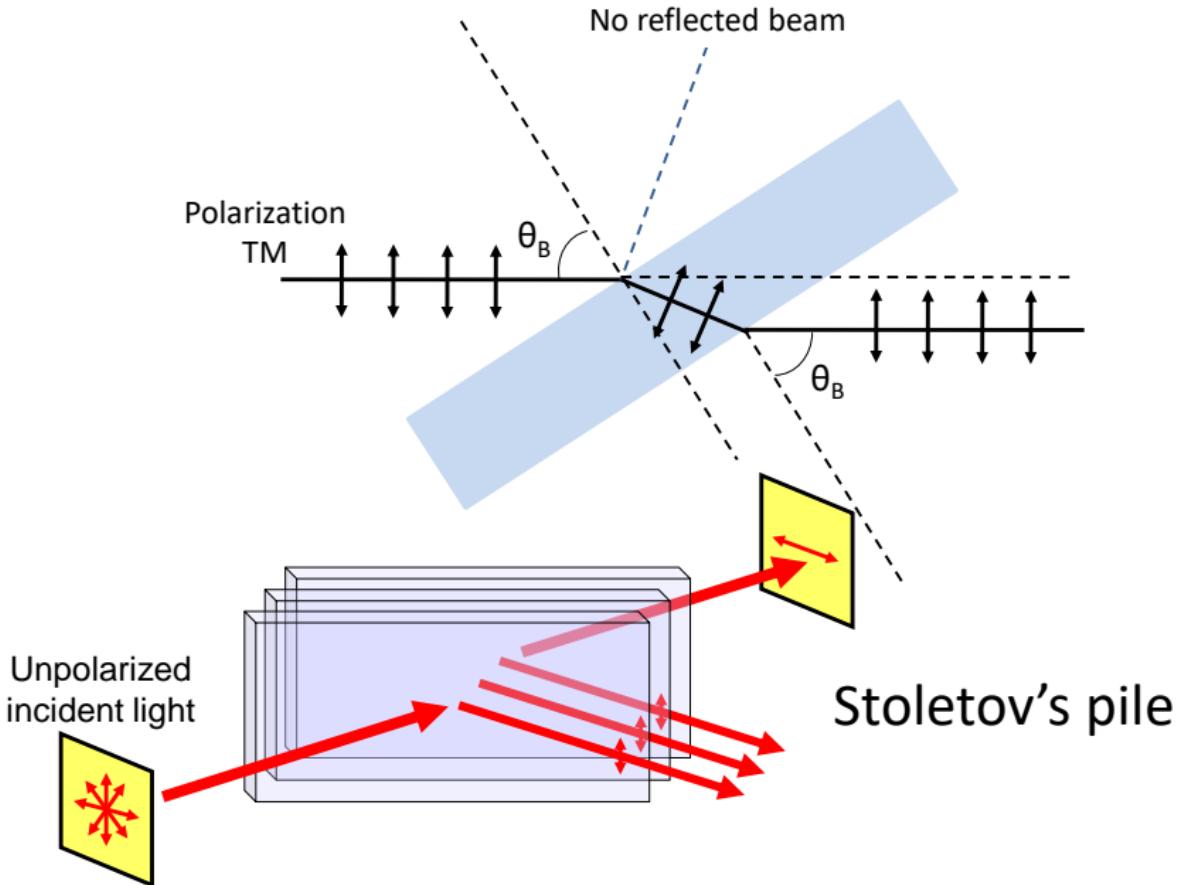
$$\theta_B = \arctg\left(\frac{n_2}{n_1}\right) = \arctg\left(\frac{1}{1.5}\right) = 33^\circ$$

The reflected and refracted rays are perpendicular!

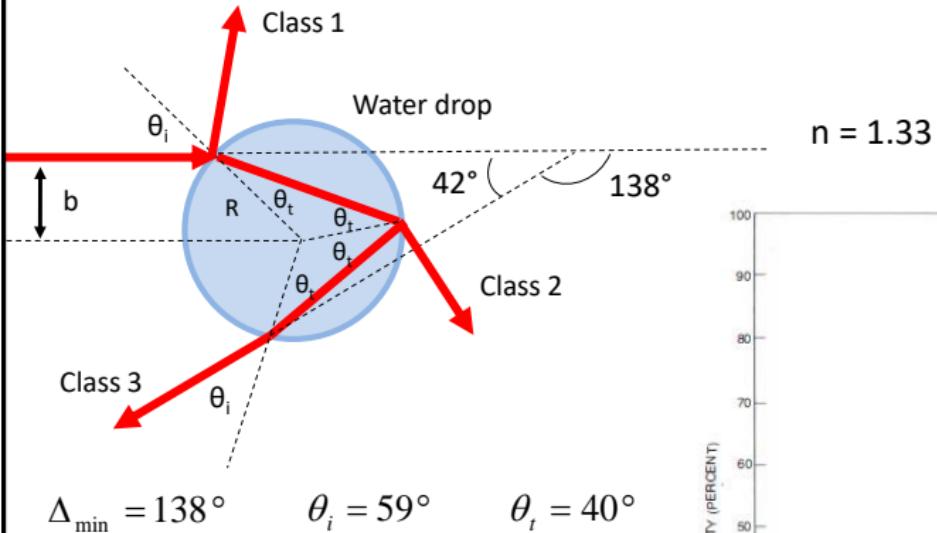
Brewster's window



Brewster's window



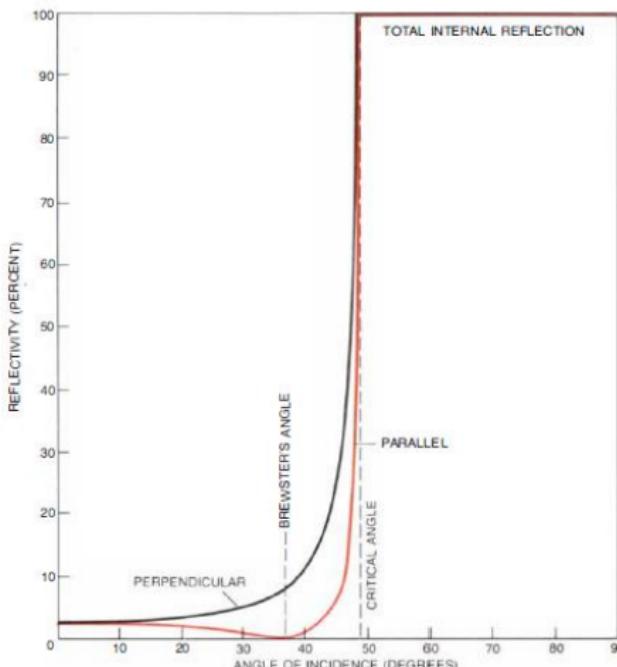
Rainbow



Brewster's angle

$$\text{drop-air} \quad \theta_B = \arctg \left(\frac{1}{1.33} \right) \cong 36.9^\circ$$

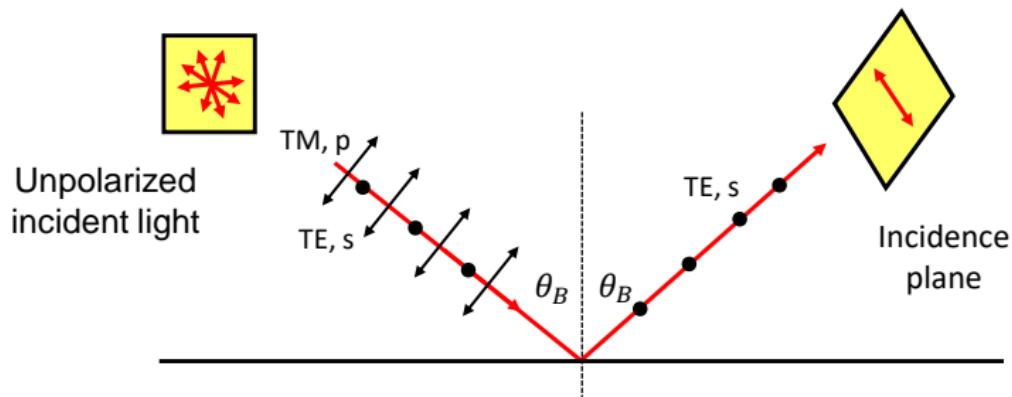
Class 3 rays are TE (s) polarized.

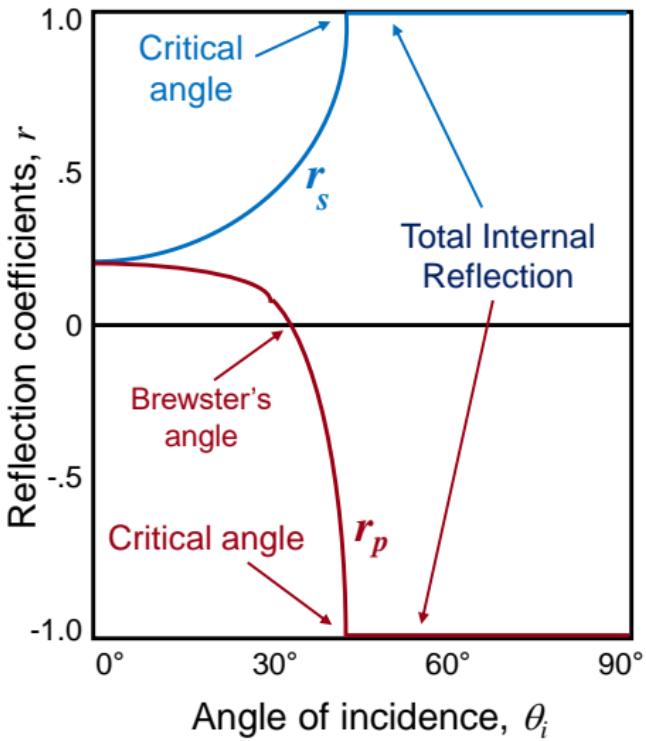


Brewster's angle

D: describe a method based on Brewster's angle to identify the transmission axis of an unknown linear polarizer.

R: at Brewster's angle $r_p = 0$, thus the reflected beam will be TE (s) polarized. If the linear polarizer is aligned to minimize the reflected beam, its transmittance axis will be orthogonal with respect to the polarization direction.



(glass-air) $n < 1$ 

$$\sin \theta_c = \frac{n_2}{n_1} = n < 1$$

Critical angle

$$\theta_c = \arcsin\left(\frac{n_2}{n_1}\right)$$

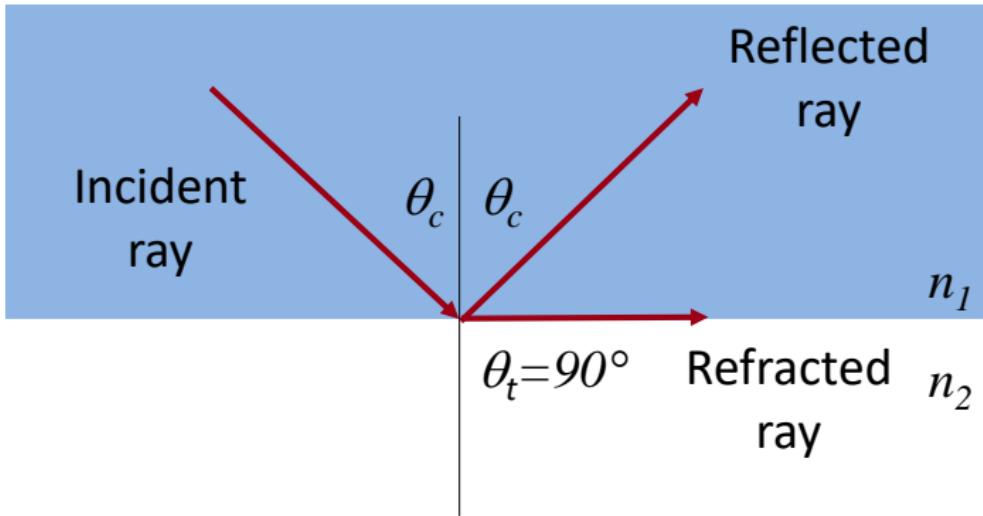
$$\theta_i \geq \theta_c \quad R = R_s = R_p = 1$$

**Total Internal
Reflection (TIR)**

$$n_{air} \sim 1 < n_{glass} \sim 1.5$$

$$\theta_c \approx 42^\circ$$

Critical angle



Critical
angle

$$\theta_c = \arcsin\left(\frac{n_2}{n_1}\right) = \arcsin\left(\frac{1}{1.5}\right) = 42^\circ$$

$$n = \frac{n_2}{n_1} < 1 \quad \vartheta_i > \vartheta_c \quad \text{Critical angle}$$

$$\sin^2 \vartheta_t + \cos^2 \vartheta_t = 1 \quad \sin \vartheta_i = n \sin \vartheta_t$$

$$r_s = \frac{\cos \vartheta_i - n \cos \vartheta_t}{\cos \vartheta_i + n \cos \vartheta_t} = \frac{\cos \vartheta_i - i \sqrt{\sin^2 \vartheta_i - n^2}}{\cos \vartheta_i + i \sqrt{\sin^2 \vartheta_i - n^2}} = e^{i \phi_s}$$

$$r_p = \frac{-n \cos \vartheta_i + \cos \vartheta_t}{n \cos \vartheta_i + \cos \vartheta_t} = \frac{-n^2 \cos \vartheta_i + i \sqrt{\sin^2 \vartheta_i - n^2}}{n^2 \cos \vartheta_i + i \sqrt{\sin^2 \vartheta_i - n^2}} = -e^{i \phi_p}$$

$$R = R_s = R_p = |r_s|^2 = |r_p|^2 = 1 \quad \text{Total Internal Reflection}$$

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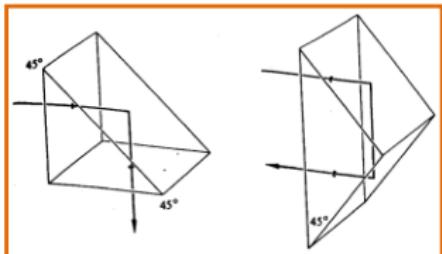
Total Internal Reflection (TIR)

Optics and
Laser Physics
T. Cesca

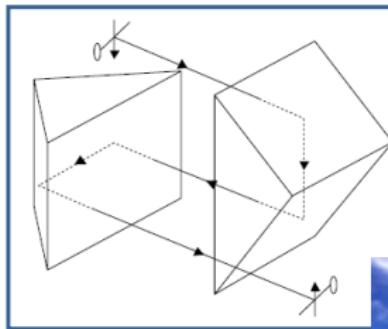


Applications of TIR

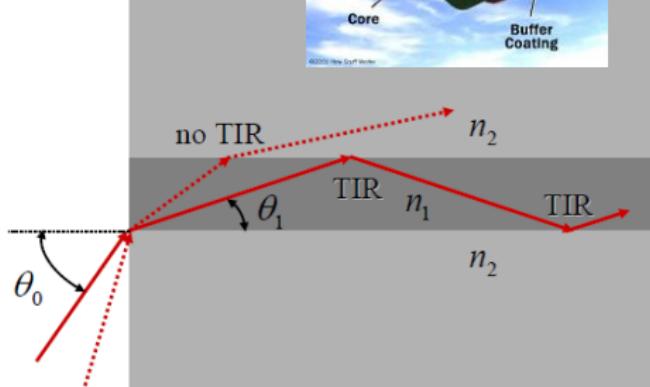
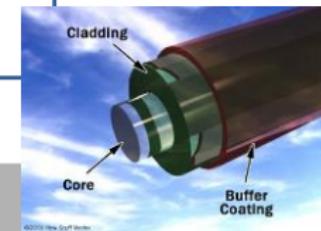
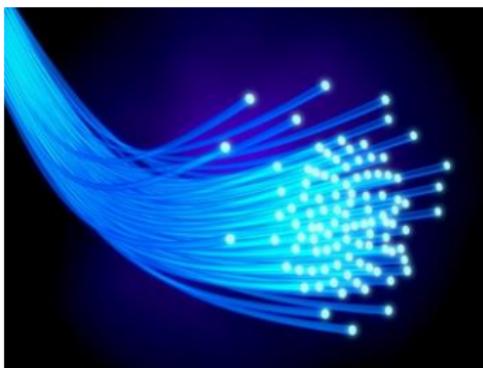
Beam steerers



Porro's prisms



Optical fibers



Snell's law

$$n_0 \sin \theta_0 = n_1 \sin \theta_1 = n_1 \cos \theta_1^{\text{int}} = n_1 \sqrt{1 - (\sin \theta_1^{\text{int}})^2} = n_1 \sqrt{1 - \left(\frac{n_2 \sin \theta_2}{n_1} \right)^2} =$$

$$= n_1 \sqrt{1 - \left(\frac{n_2}{n_1} \right)^2} = \sqrt{n_1^2 - n_2^2}$$

$\theta_2 = 90^\circ \quad (\text{TIR})$

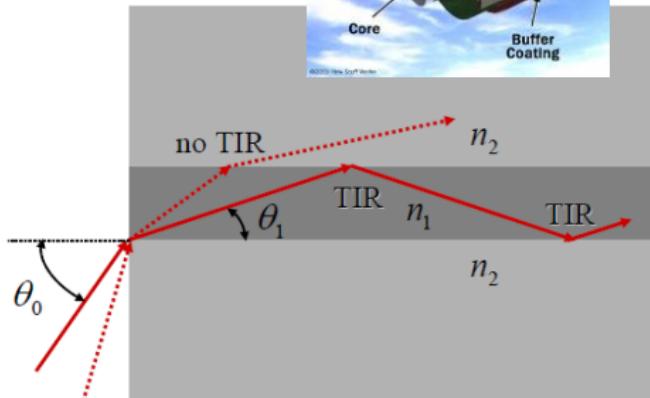
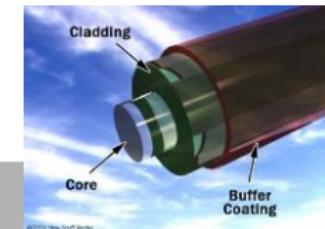
Optical fibers

$$n_1 = n_{\text{core}} > n_{\text{cladding}} = n_2$$

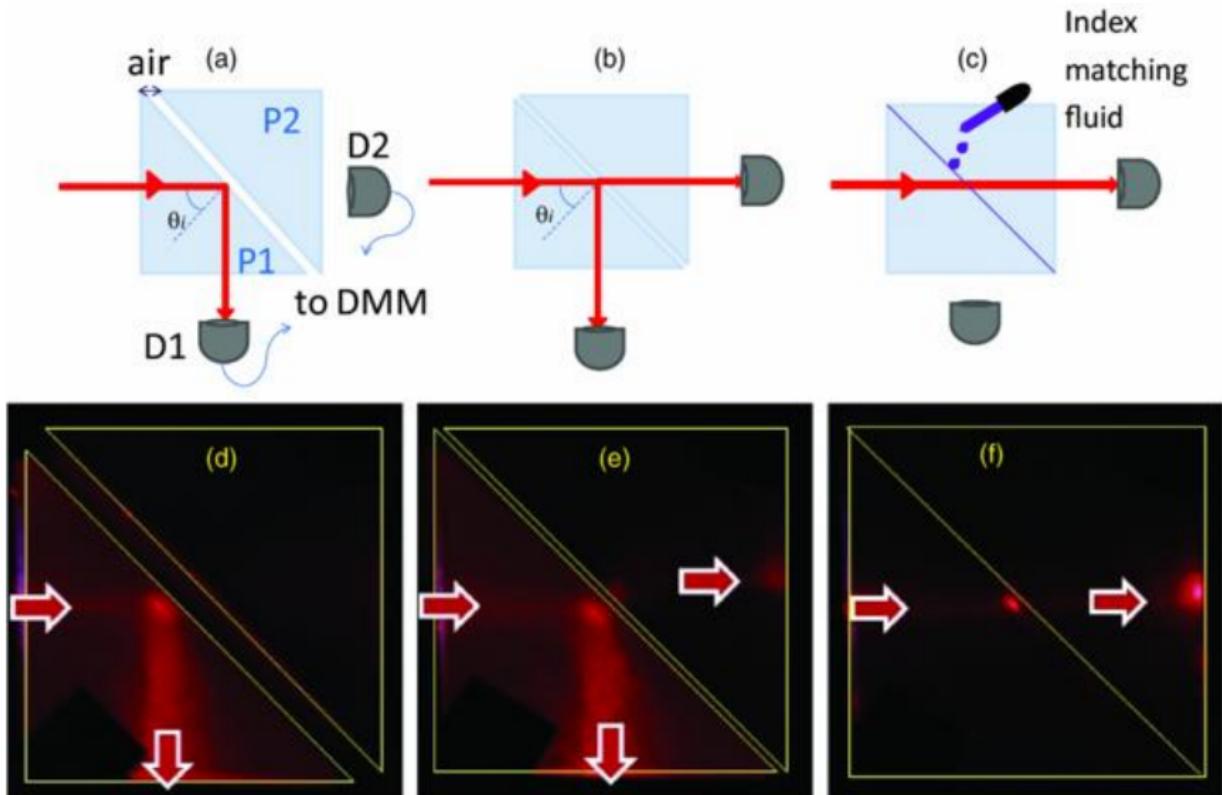
θ_0 = acceptance angle
(Numerical Aperture, NA)

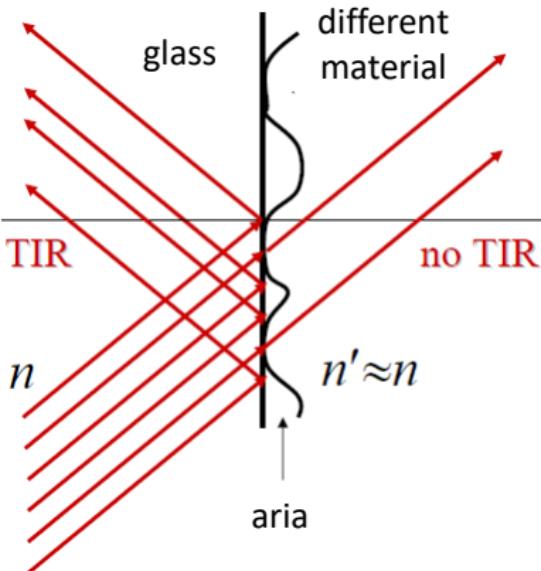
$$\sin \theta_0 = \sqrt{n_1^2 - n_2^2}$$

High index contrast (n_1/n_2): high NA



Frustrated Total Internal Reflection (FTIR)

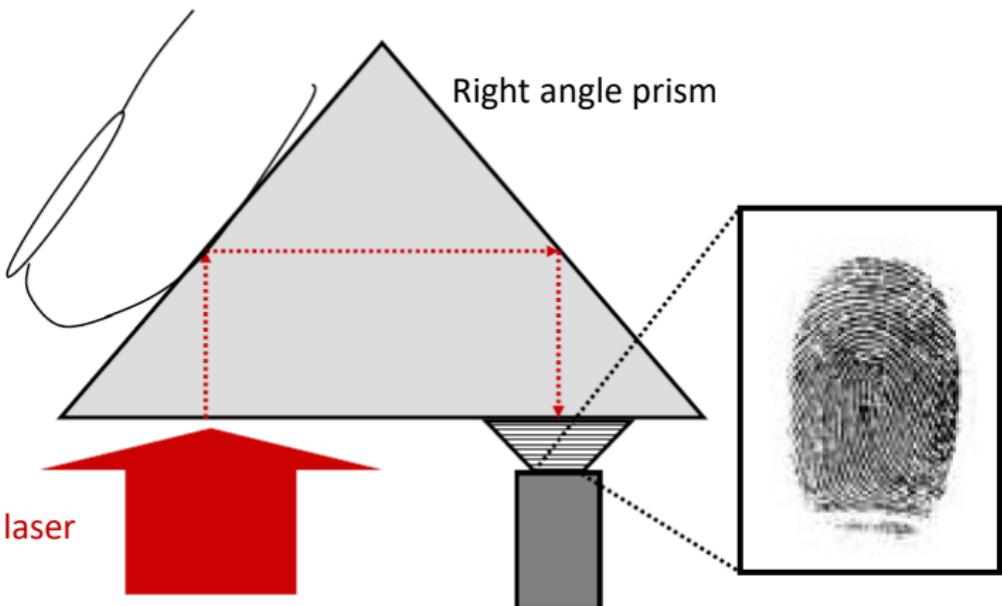




Angle of incidence >
critical angle for TIR

Where the two surfaces are in contact the
total internal reflection is frustrated.

Sensors of fingerprints



See TIR from a fingerprint valley and FTIR from a ridge.

This works because the ridges are higher than the evanescent wave penetration.

TM (p): horizontal
TE (s): vertical

Incidence plane: horizontal plane

$\begin{pmatrix} A \\ B \end{pmatrix}$ Jones vector

Reflection matrix

$$M_R = \begin{pmatrix} -r_p & 0 \\ 0 & r_s \end{pmatrix}$$

Transmission matrix

$$M_T = \begin{pmatrix} t_p & 0 \\ 0 & t_s \end{pmatrix}$$

$$\begin{pmatrix} A' \\ B' \end{pmatrix} = \begin{pmatrix} -r_p & 0 \\ 0 & r_s \end{pmatrix} \begin{pmatrix} A \\ B \end{pmatrix}$$

$$\begin{pmatrix} A' \\ B' \end{pmatrix} = \begin{pmatrix} t_p & 0 \\ 0 & t_s \end{pmatrix} \begin{pmatrix} A \\ B \end{pmatrix}$$

At **grazing** incidence: $M_R = \pm \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$

- (+): internal reflection, $n < 1$
- (-): external reflection, $n > 1$

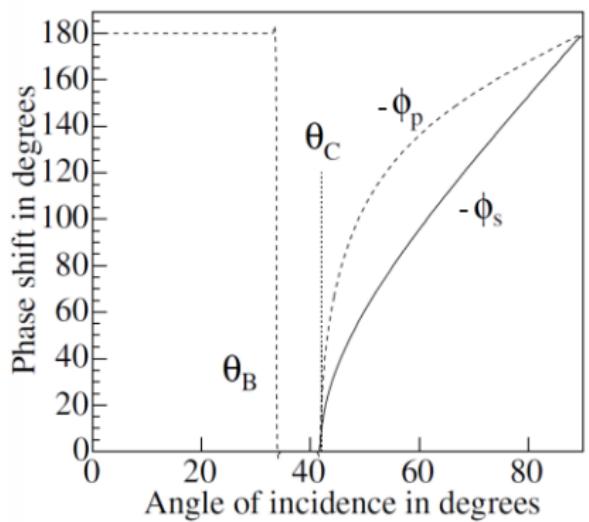
At **TIR**: $\begin{pmatrix} A' \\ B' \end{pmatrix} = \begin{pmatrix} e^{i\phi_p} & 0 \\ 0 & e^{i\phi_s} \end{pmatrix} \begin{pmatrix} A \\ B \end{pmatrix} = e^{i\phi_p} \begin{pmatrix} A \\ Be^{i\Delta} \end{pmatrix}$

$\boxed{\Delta = \phi_s - \phi_p}$

dephasing

At TIR the reflected wave is typically elliptically polarized!

Phase shift at TIR



$$r_s = \frac{\cos \theta_i - i\sqrt{\sin^2 \theta_i - n^2}}{\cos \theta_i + i\sqrt{\sin^2 \theta_i - n^2}} = e^{i\phi_s}$$

$$r_p = \frac{-n^2 \cos \theta_i + i\sqrt{\sin^2 \theta_i - n^2}}{n^2 \cos \theta_i + i\sqrt{\sin^2 \theta_i - n^2}} = -e^{i\phi_p}$$

$$\Delta = \phi_s - \phi_p$$

$$\operatorname{tg} \frac{\Delta}{2} = \frac{\cos \theta_i \sqrt{\sin^2 \theta_i - n^2}}{\sin^2 \theta_i}$$

Fresnel's romb

$$\tan \frac{\Delta}{2} = \frac{\cos \theta_i \sqrt{\sin^2 \theta_i - n^2}}{\sin^2 \theta_i}$$

