

A Moving Mesh Finite Element Method for Modelling Defects in Liquid Crystals

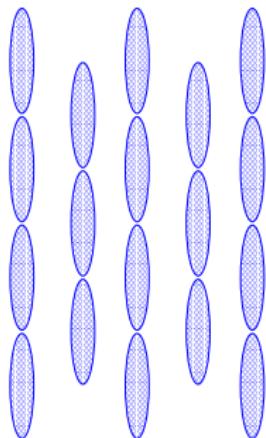
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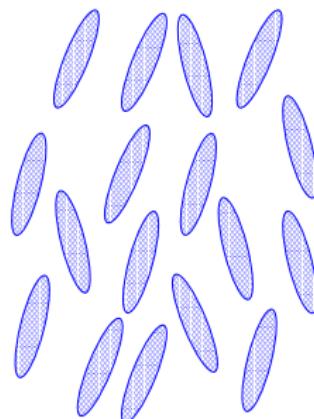


Joint work with Craig MacDonald and John Mackenzie

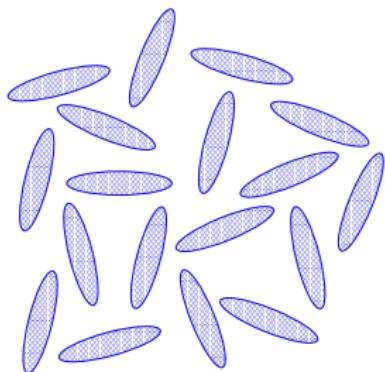
Liquid crystals



solid



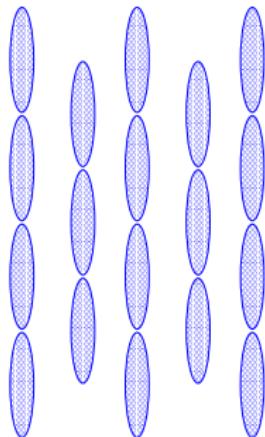
liquid crystal



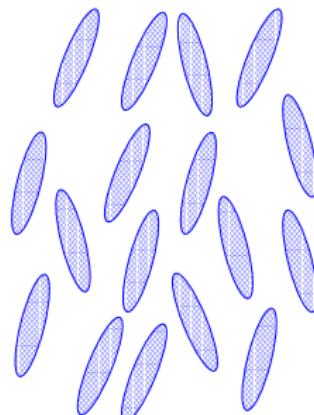
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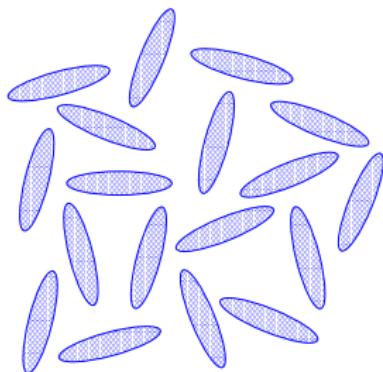
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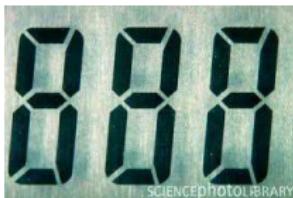
- Liquid crystals occur between solid crystal and isotropic liquid states.
- They may have different **equilibrium** configurations, but naturally prefer states with **minimum** energy.

Liquid Crystal Displays

- **IDEA:** force switching between **stable** states by altering applied voltage, magnetic field, boundary conditions, . . .

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- Used in a wide range of LCDs.



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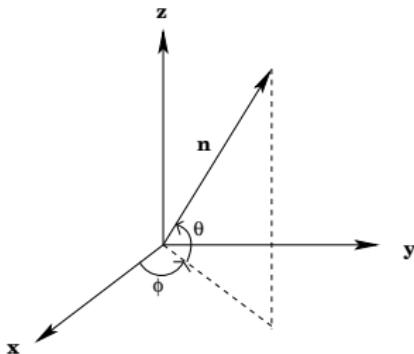
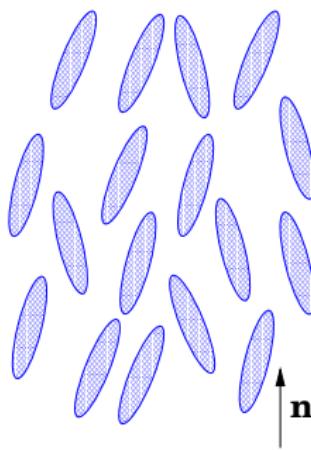
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Motivation

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- Defects typically induce distortion over very small length scales as compared to the size of the cell: this poses significant challenges for standard numerical modelling techniques.

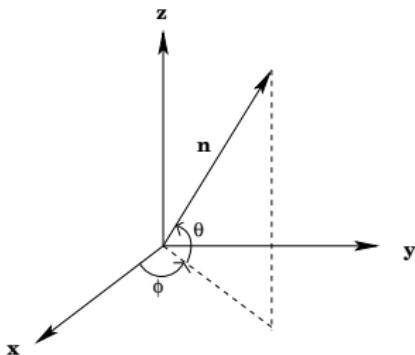
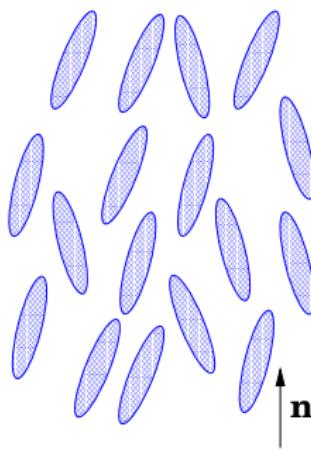
- Defects in a liquid crystal can arise due to external factors such as applied electric or magnetic fields, or the constraining geometry of the liquid crystal cell.
- Understanding the formation and dynamics of defects is important in the design and control of liquid crystal devices.
- Defects typically induce distortion over very small length scales as compared to the size of the cell: this poses significant challenges for standard numerical modelling techniques.
- In this talk we present a finite-element based adaptive moving mesh model for tracking defect movement.

Director-based model



- **Director:** represents average direction of molecular alignment

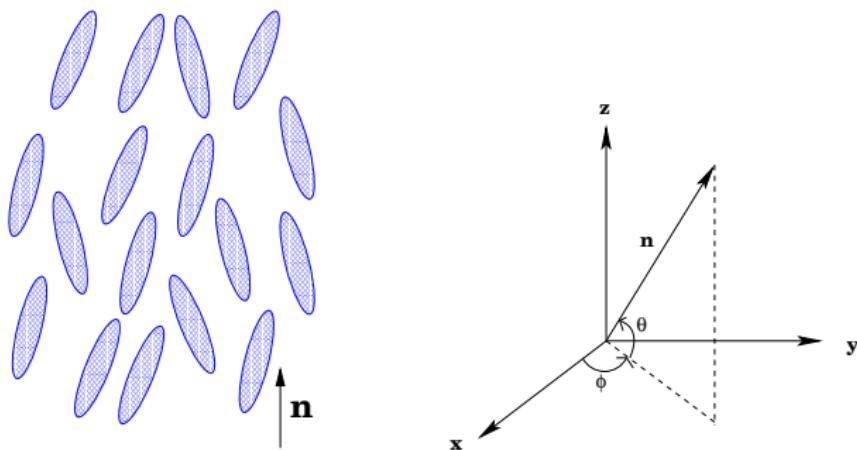
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- Represent using unit vectors with $\mathbf{n} = -\mathbf{n}$:

$$\mathbf{n} = (\cos \theta \cos \phi, \cos \theta \sin \phi, \sin \theta)$$

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- Represent using unit vectors with $\mathbf{n} = -\mathbf{n}$:
$$\mathbf{n} = (\cos \theta \cos \phi, \cos \theta \sin \phi, \sin \theta)$$
- Model with **Leslie-Ericksen** dynamic theory.

Liquid crystal model: \mathbf{Q} -tensor theory

- Describe orientation of each molecule by a single vector \mathbf{u} in direction of its main axis.

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- Represent **average** orientation by **symmetric** and **traceless** order tensor

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- With orthogonal eigenframe $\{\mathbf{l}, \mathbf{m}, \mathbf{n}\}$, write

$$\mathbf{Q} = S \left(\mathbf{n} \otimes \mathbf{n} - \frac{1}{3} \mathbf{I} \right) + T (\mathbf{m} \otimes \mathbf{m} - \mathbf{l} \otimes \mathbf{l})$$

where **S**, **T** are **uniaxial** and **biaxial** order parameters.

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- Consider **uniaxial** molecular distribution ($T = 0$) where \mathbf{n} is the liquid crystal **director**.

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- **Five** unknowns for PDE model:

$$q_1, q_2, q_3, q_4, q_5.$$

\mathbf{Q} -tensor equations

- Minimise the free energy:

$$F = \int_V F_{bulk}(\mathbf{Q}, \nabla \mathbf{Q}) dv + \int_S F_{surface}(\mathbf{Q}) dS$$

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- Derive expressions for individual energy contributions in terms of \mathbf{Q} , $\nabla \mathbf{Q}$.
- With **strong anchoring** (Dirichlet boundary conditions), there is no contribution from surface energy.
- Solutions with **least** energy are physically relevant: solve **Euler-Lagrange** equations.

Bulk energies

- **Elastic** energy: induced by distorting the **Q**-tensor in space.

$$F_{elastic} = \frac{1}{2}L_1(\operatorname{div} \mathbf{Q})^2 + \frac{1}{2}L_2|\nabla \times \mathbf{Q}|^2$$

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$$F_{thermotropic} = \frac{1}{2}A(T - T^*) \operatorname{tr} \mathbf{Q}^2 - \frac{\sqrt{6}}{3}B \operatorname{tr} \mathbf{Q}^3 + \frac{1}{4}C(\operatorname{tr} \mathbf{Q}^2)^2$$

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- **Electrostatic** energy: due to an applied electric field \mathbf{E} (electric potential U with $\mathbf{E} = -\nabla U$).

$$F_{\text{electrostatic}} = -\frac{1}{2}\epsilon_0 \mathbf{E} \cdot \boldsymbol{\epsilon} \mathbf{E} - (\bar{\epsilon} \operatorname{div} \mathbf{Q}) \cdot \mathbf{E}$$

Derivation of time-dependent PDEs

- Use a **dissipation function** with viscosity coefficient ν .

$$\mathcal{D} = \frac{\nu}{2} \text{tr} \left[\left(\frac{\partial \mathbf{Q}}{\partial t} \right)^2 \right] = \nu(\dot{q}_1 \dot{q}_4 + \dot{q}_1^2 + \dot{q}_2^2 + \dot{q}_3^2 + \dot{q}_4^2 + \dot{q}_5^2)$$

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- Obtain \mathbf{Q} -tensor PDEs (for $i = 1, \dots, 5$ and $j = 1, 2, 3$):

$$\frac{\partial \mathcal{D}}{\partial \dot{q}_i} = \nabla \cdot \hat{\mathbf{r}}_i - \hat{f}_i$$

$$(\hat{\mathbf{r}}_i)_j = \frac{\partial F_{bulk}}{\partial q_{i,j}}, \quad q_{i,j} = \frac{\partial q_i}{\partial x_j}, \quad \hat{f}_i = \frac{\partial F_{bulk}}{\partial q_i}$$

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- Combining equations and manipulating terms gives

$$\frac{\partial q_i}{\partial t} = \nabla \cdot \Gamma_i - f_i, \quad i = 1, \dots, 5.$$

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SUMMARY

- Final time-dependent physical PDEs (PPDEs) are

$$\frac{\partial q_i}{\partial t} = \nabla \cdot \boldsymbol{\Gamma}_i - f_i \quad i = 1, \dots, 5$$

$$\nabla \cdot \mathbf{D} = 0$$

- 6 PDEs in 6 unknowns $(q_1, q_2, q_3, q_4, q_5, U)$

- Three common forms of grid adaptivity in finite elements:
 - ***h*-refinement:** initially uniform mesh is locally **coarsened or refined** by inclusion or deletion of mesh points, normally based on *a posteriori* error estimates
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- Focus here on **Moving Mesh PDE model**.

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- Define mesh velocity

$$\dot{\mathbf{x}}(\mathbf{x}, t) = \frac{\partial \mathbf{x}}{\partial t} \Big|_{\xi} (\mathcal{A}_t^{-1}(\mathbf{x}))$$

and apply the Chain Rule to get

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- Additional convection-like term due to the mesh movement.

Finite elements for the physical PDEs

- Final set of six coupled PDEs ($i = 1, \dots, 5$):

$$\frac{\partial q_i}{\partial t} \Big|_{\xi} - \dot{\mathbf{x}} \cdot \nabla q = \nabla \cdot \Gamma_i - f_i, \quad \nabla \cdot \mathbf{D} = 0$$

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- Find $q_{ih}(t)$, U_h such that for test functions v_h

$$\frac{d}{dt} \int_{\Omega} q_{ih} v_h \, d\mathbf{x} - \int_{\Omega} (\nabla \cdot (\dot{\mathbf{x}} q_{ih})) v_h \, d\mathbf{x} = \int_{\Omega} \Gamma_{ih} \cdot \nabla v_h \, d\mathbf{x} - \int_{\Omega} f_{ih} v_h \, d\mathbf{x},$$

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- Non-linear differential algebraic system ($i = 1, \dots, 5$)

$$\frac{d}{dt} (M(t) \mathbf{q}_i(t)) = \mathbf{G}_i(t, \mathbf{q}_i(t), \mathbf{u}(t)), \quad \mathbf{C}(\mathbf{q}_i(t), \mathbf{u}(t)) = \mathbf{0}.$$

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$$I[\xi] = \frac{1}{2} \int_{\Omega_t} [(\nabla \xi)^T G^{-1} (\nabla \xi) + (\nabla \eta)^T G^{-1} (\nabla \eta)] d\mathbf{x}$$

with 2×2 symmetric positive definite monitor matrix G .

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- For robustness, evolve mesh via gradient flow equations

$$\frac{\partial \xi}{\partial t} = \frac{P}{\tau} \nabla \cdot (G^{-1} \nabla \xi), \quad \frac{\partial \eta}{\partial t} = \frac{P}{\tau} \nabla \cdot (G^{-1} \nabla \eta).$$

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- User-specified parameters:

- positive temporal smoothing parameter τ ,
- positive function spatial balancing parameter $P(\mathbf{x}, t)$.

Final form of MMPDE

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- In practice, interchange variable roles in MMPDE to obtain

$$\tau \frac{\partial \mathbf{x}}{\partial t} = P(a\mathbf{x}_{\xi\xi} + b\mathbf{x}_{\xi\eta} + c\mathbf{x}_{\eta\eta} + d\mathbf{x}_\xi + e\mathbf{x}_\eta)$$

$$a = \frac{1}{w} \frac{x_\eta^2 + y_\eta^2}{J^2}, \quad b = -\frac{2}{w} \frac{(x_\xi x_\eta + y_\xi y_\eta)}{J^2}, \quad c = \frac{1}{w} \frac{x_\xi^2 + y_\xi^2}{J^2},$$

$$d = \frac{1}{(wJ)^2} [w_\xi(x_\eta^2 + y_\eta^2) - w_\eta(x_\xi x_\eta + y_\xi y_\eta)],$$

$$e = \frac{1}{(wJ)^2} [-w_\xi(x_\xi x_\eta + y_\xi y_\eta) + w_\eta(x_\xi^2 + y_\xi^2)].$$

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- Solve resulting linear systems using iterative method BiCGSTAB with Incomplete LU preconditioner.
- Adaptive time-stepping based on computed solutions of PPDEs and MMPDE.

Overview of full algorithm

Set an initial uniform mesh Δ_N^0 . Set the initial guess \mathbf{q}_i^0 .
Select an initial Δt^0 . Set $n = 0$.

while ($t^n < t^{\max}$);

- Evaluate monitor function at time t^n .
- Integrate **MMPDE** forward in time to obtain new grid Δ_N^{n+1} .
- Integrate **PPDEs** forward using SDIRK2 to obtain \mathbf{q}_i^{n+1} , \mathbf{u}^{n+1} .
- $n := n + 1$.

end while.

Choosing the monitor function

- Consider three different forms of monitor function:
 - AL. Based on a measure of the **arc-length** of \mathcal{T} :

$$w(\mathcal{T}(\mathbf{x}, t)) = \left(1 + |\nabla \mathcal{T}(\mathbf{x}, t)|^2\right)^{\frac{1}{2}}$$

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- BM1**. Based on **first-order partial derivatives** of \mathcal{T} :

$$w(\mathcal{T}(\mathbf{x}, t)) = \alpha(\mathbf{x}, t) + |\nabla \mathcal{T}(\mathbf{x}, t)|^{\frac{1}{m}}$$

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- BM1**. Based on **first-order partial derivatives** of \mathcal{T} :

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- BM2**. Based on **second-order partial derivatives** of \mathcal{T} :

$$w(\mathcal{T}(\mathbf{x}, t)) = \alpha(\mathbf{x}, t) + \left(\sqrt{\left(\frac{\partial^2 \mathcal{T}}{\partial x^2}\right)^2 + 2 \left(\frac{\partial^2 \mathcal{T}}{\partial x \partial y}\right)^2 + \left(\frac{\partial^2 \mathcal{T}}{\partial y^2}\right)^2} \right)^{\frac{1}{m}}$$

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- Scaling parameters α and m regulate **mesh clustering**.

Choosing the input function

- Consider two different forms of input function:
 - Scalar order parameter. Based on the trace of \mathbf{Q}^2

$$\mathcal{T}(\mathbf{x}, t) = \text{tr}(\mathbf{Q}^2)$$

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- Both have extrema at the centre of a defect and vary rapidly in the immediate neighbourhood of the defect centre.

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- PPDEs **non-dimensionalised** with respect to lengths and energies.

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Numerical experiments

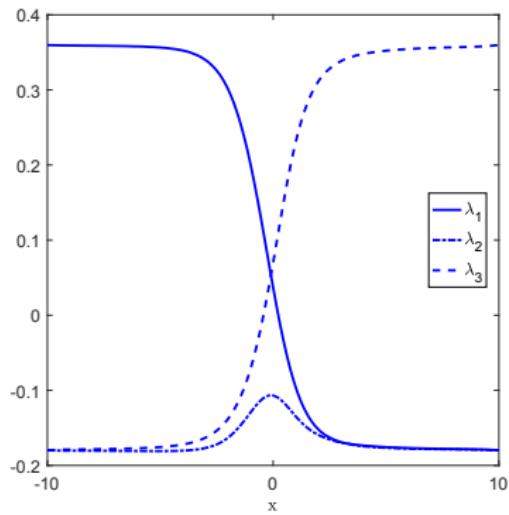
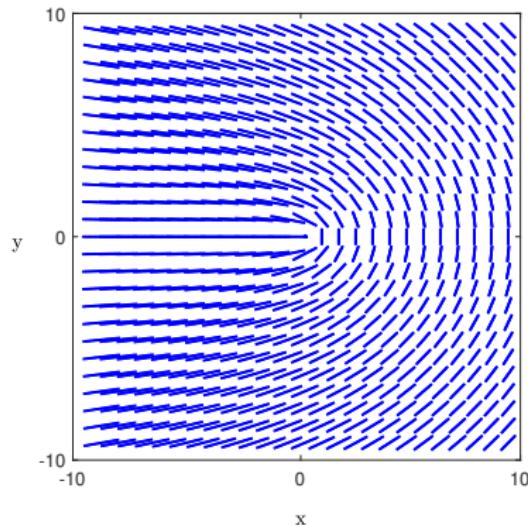
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- All experiments carried out in MATLAB.

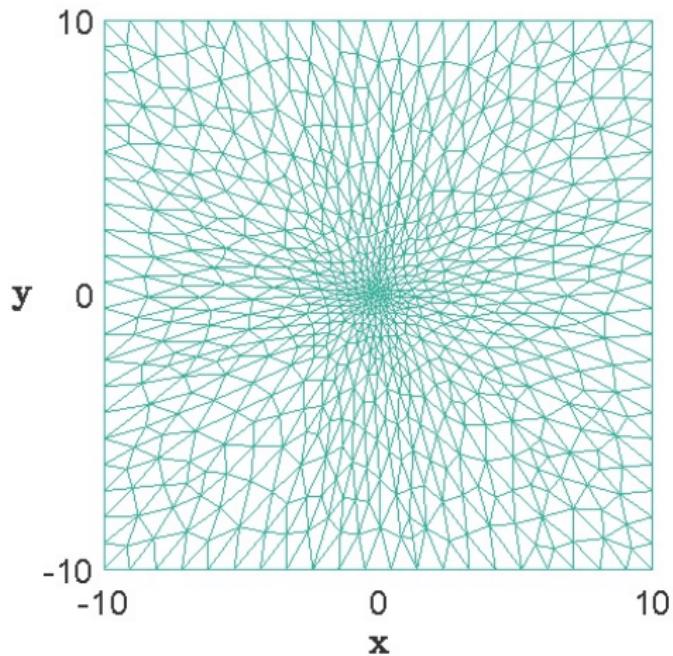
Test problem 1: stationary defect

- Director field of 1/2 defect and eigenvalue exchange along the line $y = 0$.



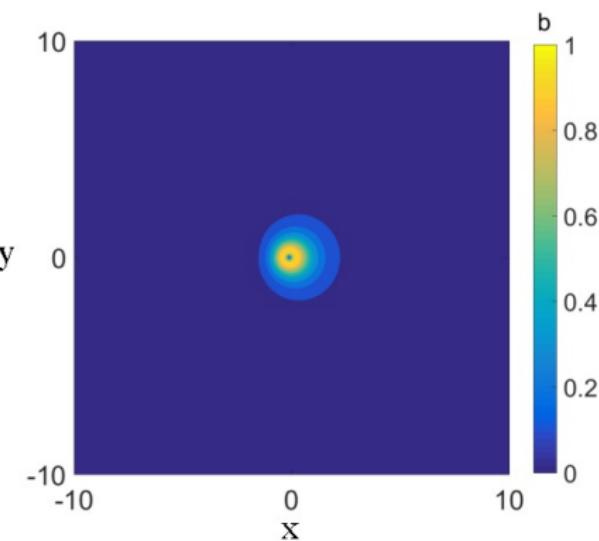
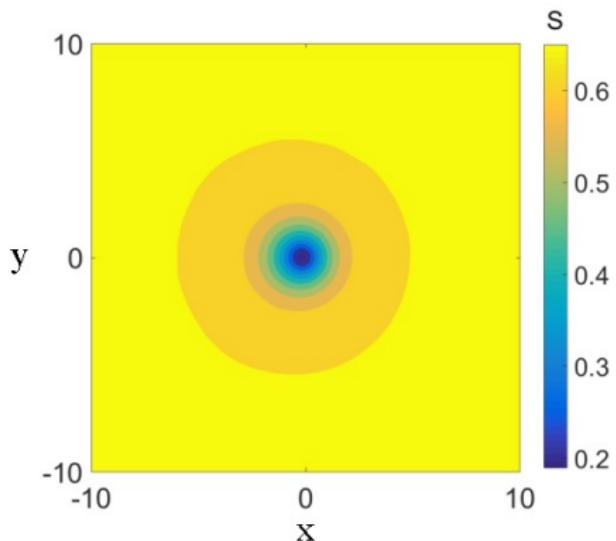
Typical adapted grid

- Sample adapted grid with 1388 quadratic elements.



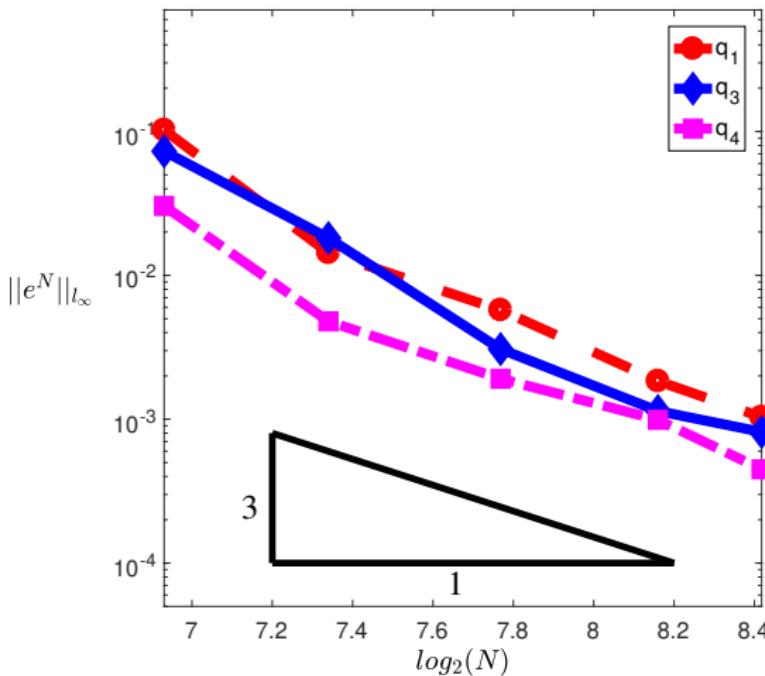
Typical solutions

- Scalar order parameter S (left) and biaxiality (b) (right).



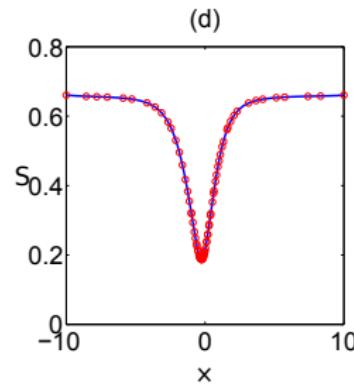
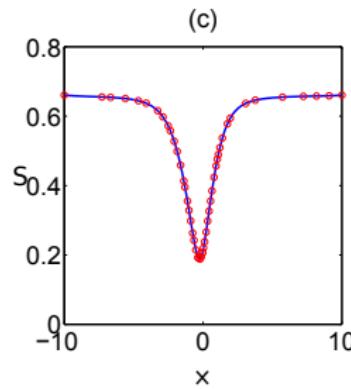
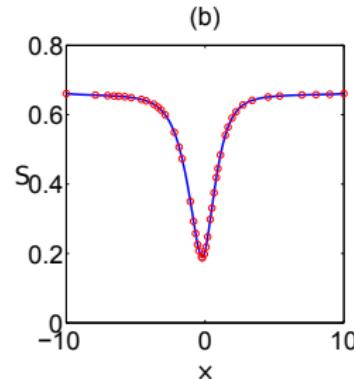
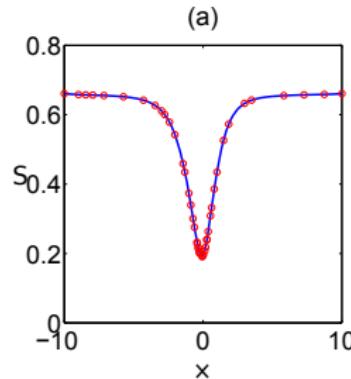
Estimated rate of spatial convergence

- ℓ_∞ error compared with reference solution is $O(N^{-3})$.



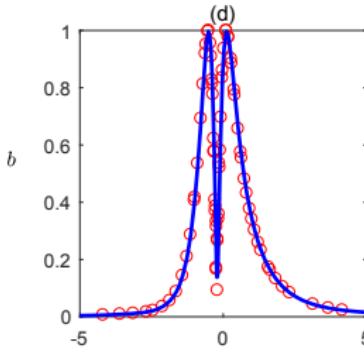
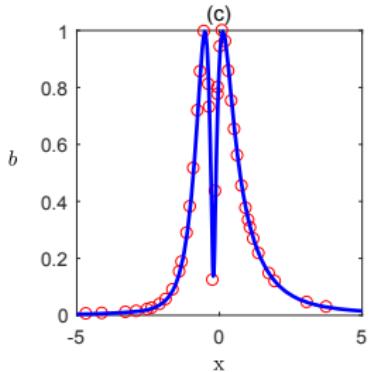
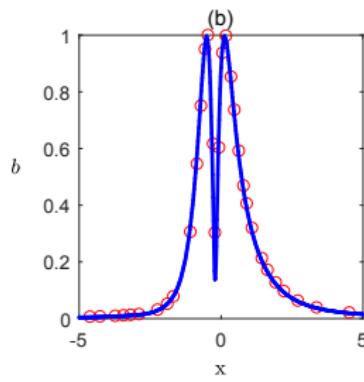
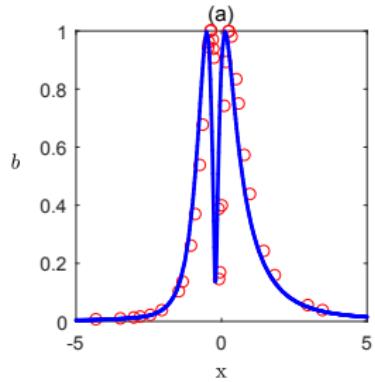
Scalar order parameter along line $y = 0$

- (a) AL; (b) BM1a; (c) BM1b; (d) BM2b



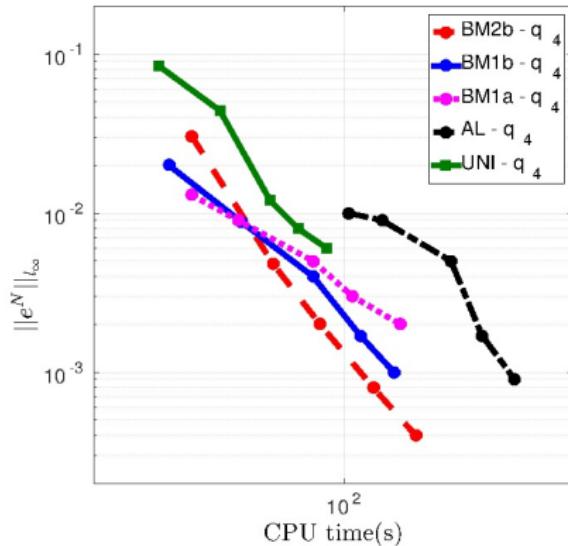
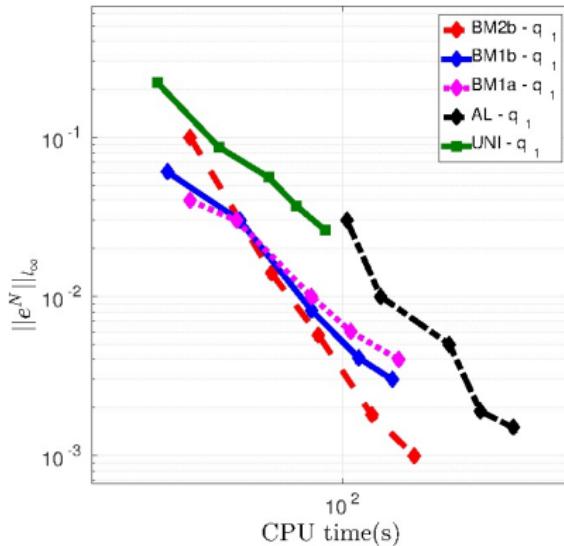
Biaxiality along line $y = 0$

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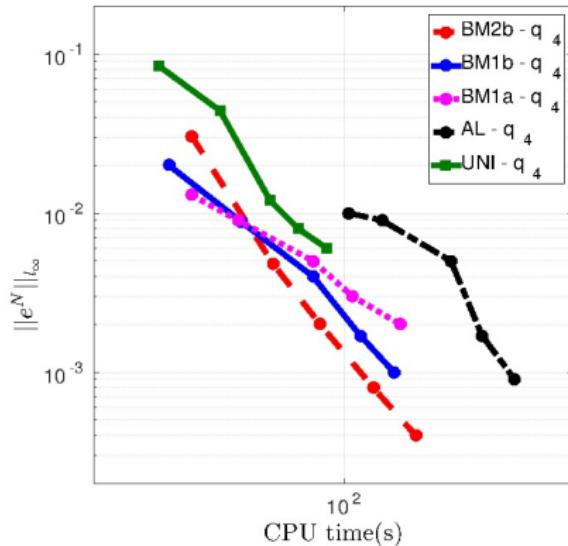
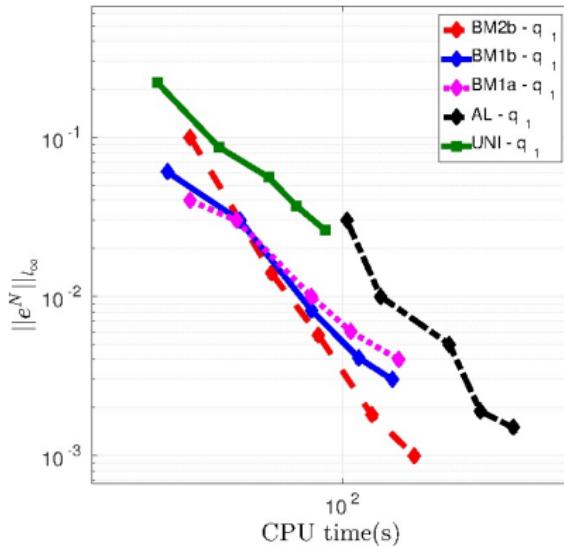
Comparing computational costs

- CPU time versus ℓ_∞ error for different grid sizes.



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- BM2b established as combination of choice.

Test problem 2: 2D Pi-cell

- Two-dimensional **Pi-cell** geometry.

Zhang, Chung, Wang and Bos, *Liquid Crystals* 34(2), 2007

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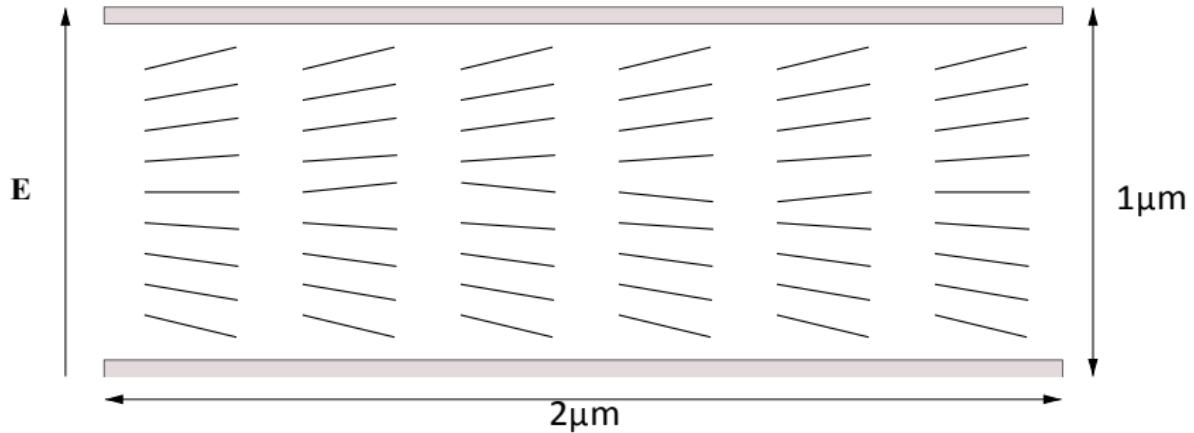
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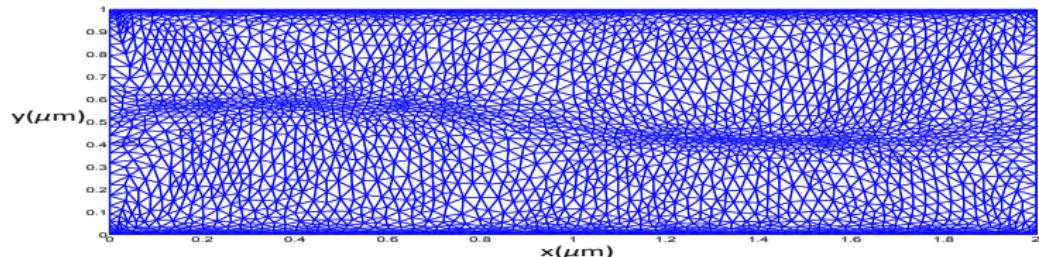
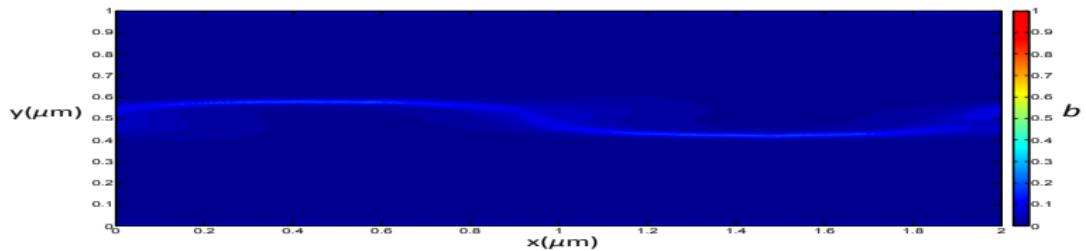
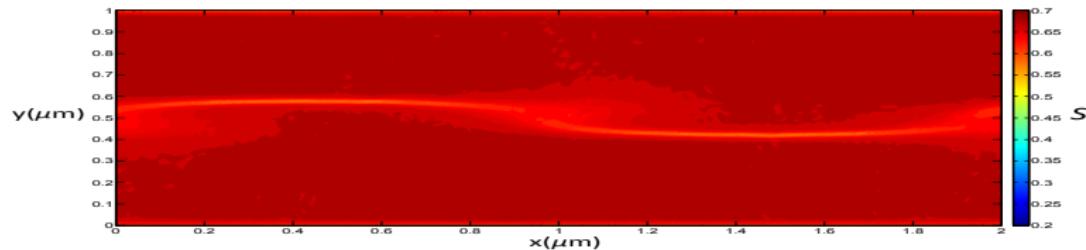
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- Perturbation fixed only at $t = 0$ for one time step, but introduces **solution gradients** in two dimensions.

Pi-cell geometry

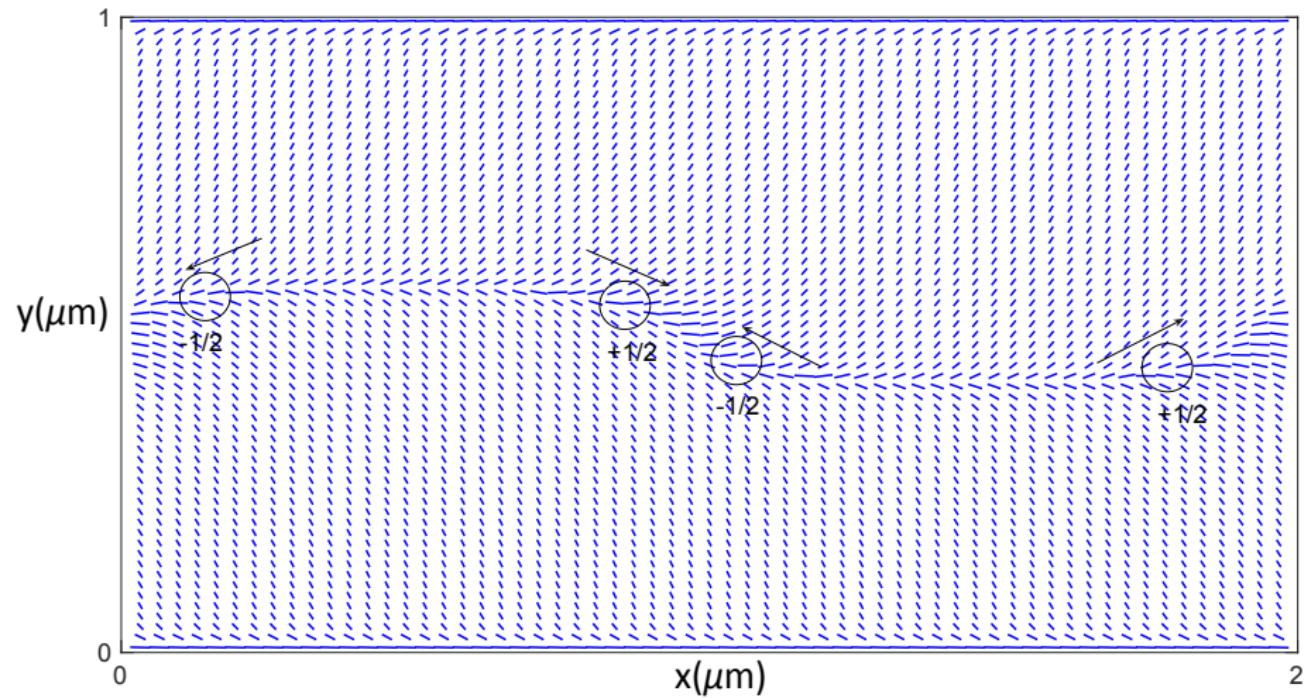
- Pre-tilt angle $\theta = \pm 6^\circ$ at boundaries.
- Electric field strength $18V\mu\text{m}^{-1}$.



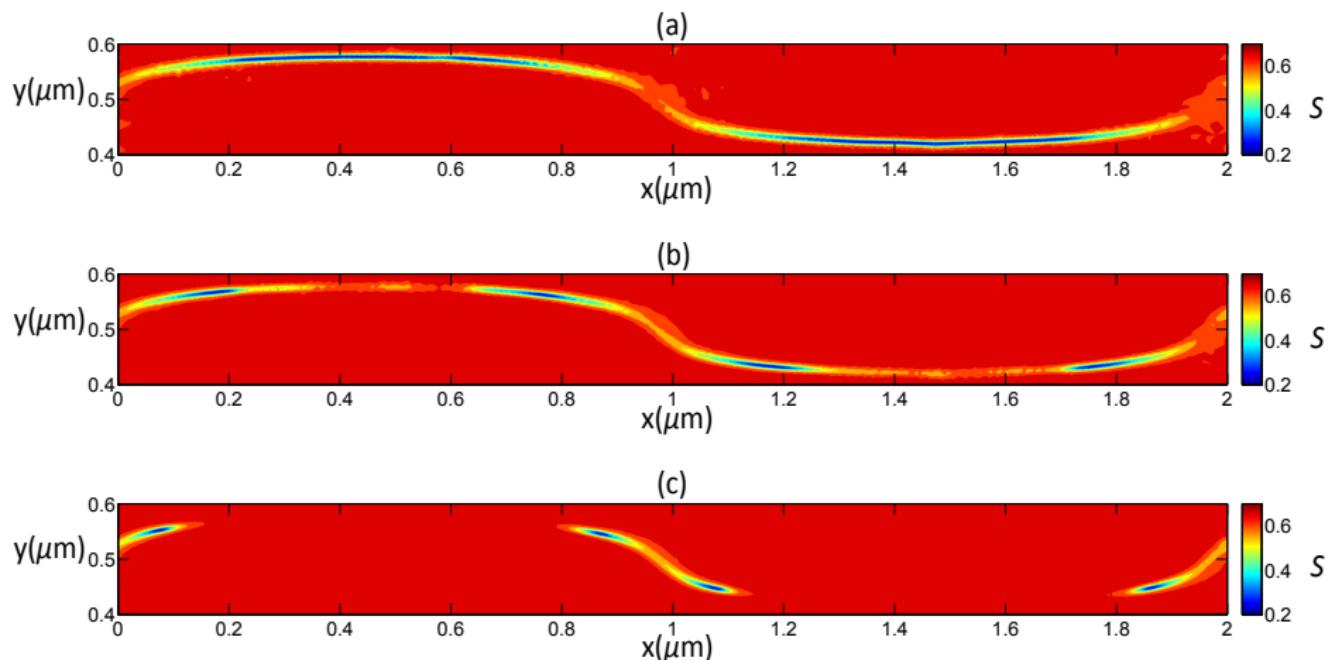
S , biaxiality and mesh after $12\mu\text{s}$



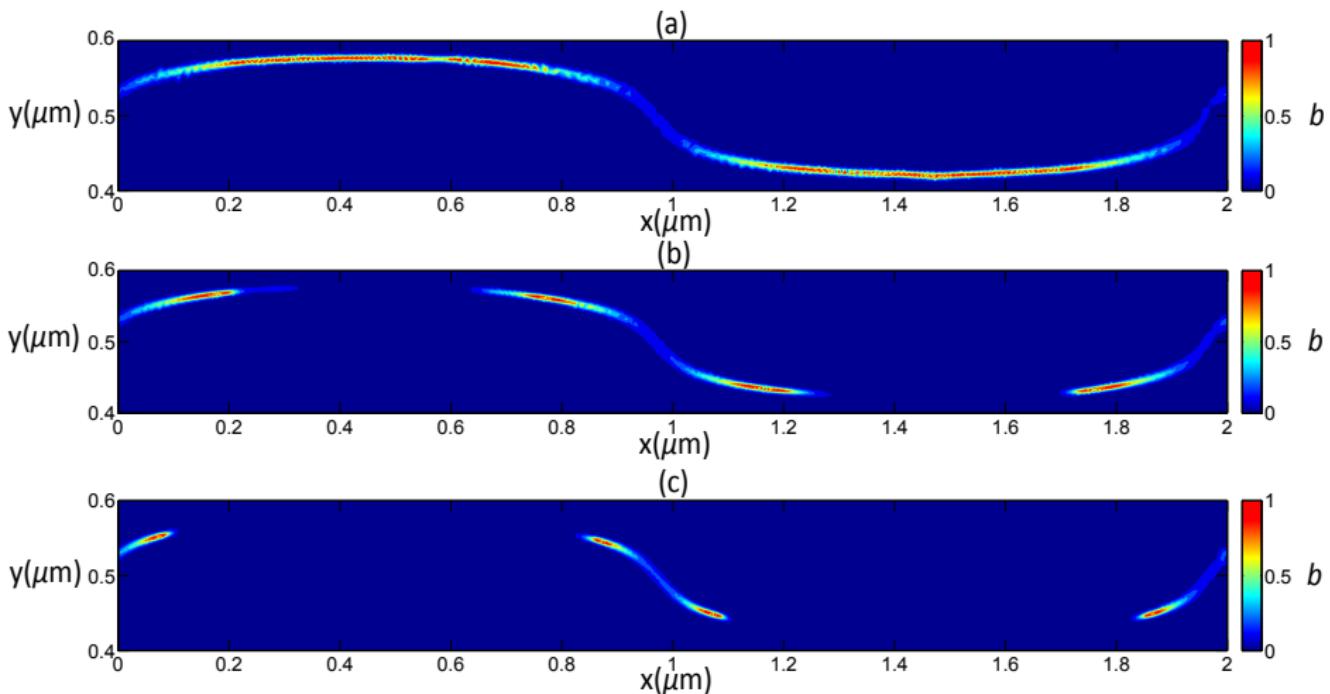
Director field after $15.5\mu\text{s}$



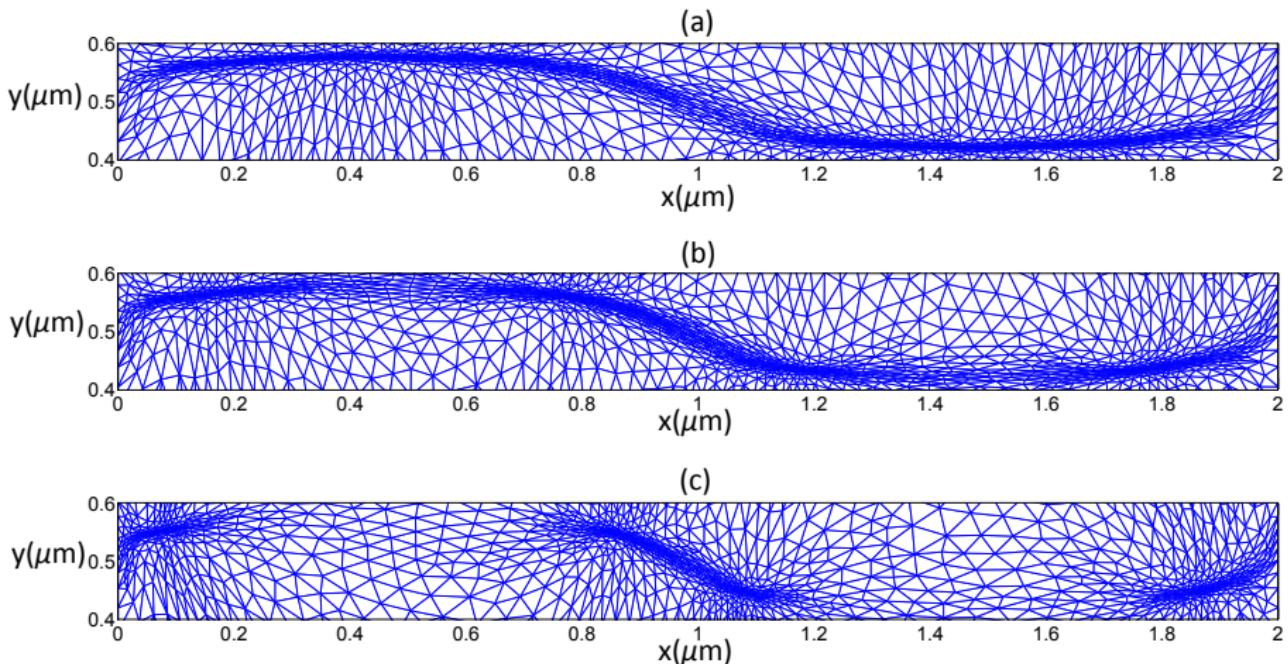
Order parameter S after (a) $15.5\mu\text{s}$ (b) $16\mu\text{s}$ and (c) $17\mu\text{s}$



Biaxiality after (a) $15.5\mu\text{s}$ (b) $16\mu\text{s}$ and (c) $17\mu\text{s}$



Adaptive mesh after (a) $15.5\mu\text{s}$ (b) $16\mu\text{s}$ and (c) $17\mu\text{s}$



Summary and future work

- We have developed a new efficient moving mesh method for Q-tensor models of liquid crystal cells.
- We have shown that biaxiality is a good choice for the monitor input function.
- We demonstrated optimal spatial convergence for a model of a static +1/2 defect.
- We resolved the movement and core details of defects in a time-dependent Pi-cell problem.
- Modelling the creation and annihilation of moving singularities on very small length and time scales is a real challenge for numerical methods.
- Future challenges involve the extension to three dimensions and more irregular geometries (e.g. the ZBD).