Computational Fluid Dynamics HW1

Almog Dobrescu

ID 214254252

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1 Inviscid Burgers Equation

The Inviscid Burgers equation, in conservation law form, is given by:

$$\frac{\partial u}{\partial t} + \frac{\partial F}{\partial x} = 0 \quad F = F_{(u)} = \frac{u^2}{2} \tag{1}$$

In non-conservation law form, is given by:

$$\frac{\partial u}{\partial t} + A \frac{\partial u}{\partial x} = 0 \quad A = \frac{\partial F}{\partial u} = u \tag{2}$$

The equation is obtained by neglecting the viscous term from the viscous Burger equation.

1.1 Boundary and Initial Conditions

$$u_{(x=0,t)} = 1.0$$

 $u_{(x=1,t)} = u_1$
 $u_{(x,t=0)} = 1 - (1 - u_1) \cdot x$ (3)

1.2 Finite Volume Formulation

$$u_i^{n+1} = u_i^n - \frac{\Delta t}{\Delta x} \left(f_{i+\frac{1}{2}}^n - f_{i-\frac{1}{2}}^n \right) \tag{4}$$

- For first-order schemes, there is no variation within a cell, and the value there is constant.
- For second-order schemes, the variation within the cell is linear.

1.3 CFL number

For the Roe method, the CFL number is defined as:

$$CFL = \frac{u\Delta t}{\Delta x} \tag{5}$$

We will want to set the maximal value of the *CFL* number. We will find the Δt at each cell and (Δt_i) and set the Δt of the current step as:

$$\Delta t = \min\left(\Delta t_i\right) \ \forall i \tag{6}$$

1.4 First Order Roe Method $(u_1 = 0.0)$

Roe scheme is based on the solution of the linear problem:

$$\frac{\partial u}{\partial t} + \bar{A}\frac{\partial u}{\partial x} = 0 \tag{7}$$

Where \bar{A} is a constant matrix that is dependent on local conditions. The matrix is constructed in a way that guarantees uniform validity across discontinuities:

1. For any u_i, u_{i+1} :

$$F_{i+1} - F_i = \bar{A} \cdot (u_{i+1} - u_i)$$

2. When $u = u_i = u_{i+1}$ then:

$$\bar{A}_{(u_i,u_{i+1})} = \bar{A}_{(u,u)} = \frac{\partial F}{\partial u} = u$$

In the case of the Burgers equation, the matrix \bar{A} is a scalar, namely, $\bar{A} = \bar{u}$. The equation becomes:

$$\frac{\partial u}{\partial t} + \bar{u}\frac{\partial u}{\partial x} = 0 \tag{8}$$

The value of \bar{u} for the cell face between i and i+1 is determined from the first conditions:

$$\bar{u} = \bar{u}_{i+\frac{1}{2}} = \frac{F_{i+1} - F_i}{u_{i+1} - u_i} = \frac{\frac{1}{2}u_{i+1}^2 - \frac{1}{2}u_i^2}{u_{i+1} - u_i} = \begin{cases} \frac{u_i + u_{i+1}}{2} & u_i \neq u_{i+1} \\ u_i & u_i = u_{i+1} \end{cases}$$
(9)

The single wave that emanates from the cell interface travels either in the positive or negative direction, depending upon the sign of $\bar{u}_{i+\frac{1}{2}}$. Define:

$$\begin{cases}
\bar{u}_{i+\frac{1}{2}}^{+} \triangleq \frac{1}{2} \left(\bar{u}_{i+\frac{1}{2}} + \left| \bar{u}_{i+\frac{1}{2}} \right| \right) \ge 0 \\
\bar{u}_{i+\frac{1}{2}} = \bar{u}_{i+\frac{1}{2}}^{+} + \bar{u}_{i+\frac{1}{2}}^{-} \\
\bar{u}_{i+\frac{1}{2}}^{-} \triangleq \frac{1}{2} \left(\bar{u}_{i+\frac{1}{2}} - \left| \bar{u}_{i+\frac{1}{2}} \right| \right) \le 0
\end{cases} (10)$$

Using the jump relation, the numerical flux at the cell interface can be evaluated by one of the following:

$$\begin{cases}
f_{i+\frac{1}{2}} - F_i = \bar{u}_{i+\frac{1}{2}}^- \cdot (u_{i+1} - u_i) \\
F_{i+1} - f_{i+\frac{1}{2}} = \bar{u}_{i+\frac{1}{2}}^+ \cdot (u_{i+1} - u_i)
\end{cases}$$
(11)

The numerical flux may then be written in the following symmetric form:

$$f_{i+\frac{1}{2}} = \frac{F_i + F_{i+1}}{2} - \frac{1}{2} \left(\bar{u}_{i+\frac{1}{2}}^+ - \bar{u}_{i+\frac{1}{2}}^- \right) (u_{i+1} - u_i)$$
OR:
$$f_{i+\frac{1}{2}} = \frac{F_i + F_{i+1}}{2} - \frac{1}{2} \left| \bar{u}_{i+\frac{1}{2}} \right| (u_{i+1} - u_i)$$
(12)

1.4.1 Effect of CFL

1.5 Second Order Roe $(u_1 = 0.5)$

The first-order accurate Roe method interface flux function will be denoted like this:

$$f_{i+\frac{1}{2}}^{\text{Roe},1} = f_{(u_i,u_{i+1})}$$

The second order accurate Roe takes the form:

$$f_{i+\frac{1}{2}}^{\mathrm{Roe},2} = f_{\left(u_{i+1}^{l},u_{i+1}^{r}\right)}$$

Hence:

$$f_{i+\frac{1}{2}}^{\text{Roe},2} = \frac{1}{2} \left(F_{\left(u_{1+\frac{1}{2}}^{l}\right)} + F_{\left(u_{1+\frac{1}{2}}^{r}\right)} - \left| \bar{u}_{i+\frac{1}{2}} \right| \left(u_{1+\frac{1}{2}}^{r} - u_{1+\frac{1}{2}}^{l} \right) \right)$$

$$\bar{u}_{1+\frac{1}{2}} = \frac{F_{\left(u_{1+\frac{1}{2}}^{r}\right)} - F_{\left(u_{1+\frac{1}{2}}^{l}\right)}}{u_{1+\frac{1}{2}}^{r} - u_{1+\frac{1}{2}}^{l}} = \frac{u_{i+\frac{1}{2}}^{l} + u_{i+\frac{1}{2}}^{r}}{2}$$

$$(13)$$

1.5.1 Without Limiters

The interface values without limiters are evaluated as:

$$\begin{cases}
 u_{i+\frac{1}{2}}^{l} = u_{i} + \frac{1-k}{4} \delta u_{i-\frac{1}{2}} + \frac{1+k}{4} \delta u_{i+\frac{1}{2}} \\
 u_{i+\frac{1}{2}}^{r} = u_{i+1} - \frac{1+k}{4} \delta u_{i+\frac{1}{2}} - \frac{1-k}{4} \delta u_{i+\frac{3}{2}}
\end{cases}$$

$$\delta u_{i} \triangleq u_{i+\frac{1}{2}} - u_{i-\frac{1}{2}} \qquad (14)$$

The parameter k determines the scheme:

$$k = \begin{cases} -1 & \text{upwind} \\ 1 & \text{central} \end{cases}$$

1.5.2 With Limiters

The interface values with limiters are evaluated as:

$$\begin{cases} u_{i+\frac{1}{2}}^{l} &= u_{i} + \frac{1-k}{4}\overline{\delta^{+}}u_{i-\frac{1}{2}} + \frac{1+k}{4}\overline{\delta^{-}}u_{i+\frac{1}{2}} \\ u_{i+\frac{1}{2}}^{r} &= u_{i+1} - \frac{1+k}{4}\overline{\delta^{+}}u_{i+\frac{1}{2}} - \frac{1-k}{4}\overline{\delta^{-}}u_{i+\frac{3}{2}} \end{cases} \overline{\delta^{\pm}}u \text{ are limited slopes}$$
 (15)

 $\overline{\delta}$ is an operator such that $\overline{\delta}u_i = \psi \delta u_i$, where $\psi(r)$ is a limiter function and:

$$r^{\pm} = \begin{cases} r_{1+\frac{1}{2}}^{+} & \triangleq \frac{u_{i+2} - u_{i+1}}{u_{i+1} - u_{i}} = \frac{\Delta u_{i+1}}{\Delta u_{i}} \\ r_{1+\frac{1}{2}}^{-} & \triangleq \frac{u_{i} - u_{i-1}}{u_{i+1} - u_{i}} = \frac{\nabla u_{i}}{\nabla u_{i+1}} \end{cases}$$
(16)

There are many types of limiters. For example, the van Albada limiter:

$$\psi\left(r\right) = \frac{r+r^2}{1+r^2} \tag{17}$$

1.5.3 Effect of CFL

1.5.4 Effect of Limiter

2 Generalized Burgers Equation

The generalized Burgers equation is given by:

$$\frac{\partial u}{\partial t} + (c + bu) \frac{\partial u}{\partial x} = \mu \frac{\partial^2 u}{\partial x^2} \tag{18}$$

Where:

$$c = \frac{1}{2}$$
 $b = -1$ $\mu = [0.001, 0.25]$

The equation can also be presented as:

$$\frac{\partial u}{\partial t} + \frac{\partial \bar{F}}{\partial x} = 0 \qquad \bar{F} = \underbrace{cu + \frac{bu^2}{2}}_{F} - \underbrace{\mu \frac{\partial u}{\partial x}}_{F_{\nu}} \tag{19}$$

In non-conservation law form, is given by:

$$\frac{\partial u}{\partial t} + A \frac{\partial u}{\partial x} = 0 \qquad A = \frac{\partial \bar{F}}{\partial u} = c + bu - \mu \frac{\partial}{\partial u} \left(\frac{\partial u}{\partial x} \right) \tag{20}$$

The generalized Burgers equation has a stationary solution:

$$u = -\frac{c}{b} \left(1 + \tanh\left(\frac{c(x - x_0)}{2\mu}\right) \right) \tag{21}$$

2.1 Domain and Computational Mesh

Using 41 grid points with $\Delta x = 1$ and computing until t = 18.0. $\Delta t = [0.5, 1.0]$.

2.2 Boundary and Initial Conditions

2.2.1 Initial Conditions

$$u_{(x,t=0)} = \frac{1}{2} \left(1 + \tanh \left(250 \left(x - 20 \right) \right) \right) \tag{22}$$

2.2.2 Boundary Conditions

Using Dirichlet boundary conditions:

$$u_{(x=0,t)} = 0$$
 $u_{(x=40,t)} = 1$ (23)

2.3 First Order Roe Method (explicit)

As written above for the inviscid Burgers equation (1.4), Roes scheme is based on the solution of the linear problem:

$$\frac{\partial u}{\partial t} + \bar{A}\frac{\partial u}{\partial x} = \mu \frac{\partial^2 u}{\partial x^2} \qquad \bar{A} = \frac{\partial F}{\partial u} \tag{24}$$

In the case of the Burgers equation, the matrix \bar{A} is a scalar.

$$\bar{A} = \bar{A}_{i+\frac{1}{2}} = \frac{F_{i+1} - F_i}{u_{i+1} - u_i} = \begin{cases} \frac{c(u_{i+1} - u_i) + \frac{b}{2}(u_{i+1}^2 - u_i^2)}{u_{i+1} - u_i} & u_i \neq u_{i+1} \\ A_i & u_i = u_{i+1} \end{cases}$$
(25)

The numerical flux at the cell interface:

$$\bar{f}_{i+\frac{1}{2}} = \frac{F_i + F_{i+1}}{2} - \frac{1}{2} \left(\bar{A}_{i+\frac{1}{2}}^+ - \bar{A}_{i+\frac{1}{2}}^- \right) (u_{i+1} - u_i) \tag{26}$$

Where:

$$\begin{cases}
\bar{A}_{i+\frac{1}{2}}^{+} \triangleq \frac{1}{2} \left(\bar{A}i + \frac{1}{2} + \left| \bar{A}_{i+\frac{1}{2}} \right| \right) \ge 0 \\
\bar{A}_{i+\frac{1}{2}} = \bar{A}_{i+\frac{1}{2}}^{+} + \bar{A}_{i+\frac{1}{2}}^{-} \\
\bar{A}_{i+\frac{1}{2}}^{-} \triangleq \frac{1}{2} \left(\bar{A}_{i+\frac{1}{2}} - \left| \bar{A}_{i+\frac{1}{2}} \right| \right) \le 0
\end{cases} (27)$$

And finally:

$$u_i^{n+1} = u_i^n - \frac{\Delta t}{\Delta x} \left(\bar{f}_{i+\frac{1}{2}}^n - \bar{f}_{i-\frac{1}{2}}^n \right) + \mu \frac{\Delta t}{(\Delta x)^2} \left(u_{i+1}^n - 2u_i^n + u_{i-1}^n \right)$$
 (28)

2.4 MacCormack Method

The original MacCormack method applied to Burgers equaiton results in:

Predictor:
$$u_i^{\overline{n+1}} = u_i^n - \Delta t \frac{\Delta F_i^n}{\Delta x} + r\delta^2 u_i^n$$
Corrector:
$$u_i^{n+1} = \frac{1}{2} \left(u_i^n + u_i^{\overline{n+1}} - \Delta t \frac{\nabla F_i^{\overline{n+1}}}{\Delta x} \right) + r\delta^2 u_i^{\overline{n+1}}$$
(29)

Where:

•
$$r = \frac{\mu \Delta t}{(\Delta x)^2}$$

•
$$\delta^2 u_i = u_{i+1} - 2u_i + u_{i-1}$$

$$\bullet \ \Delta f = f_{i+1} - f_i$$

$$\bullet \ \nabla f = f_i - f_{i-1}$$

2.4.1 Stability

The stability condition for MacCormack's method is:

$$\Delta t \le \frac{\left(\Delta x\right)^2}{|u|\,\Delta x + \mu}\tag{30}$$

2.5 Beam and Warming

Beam and Warming inroduced the Delta form and their method is more efficient:

$$\left(I + \theta \left(\Delta t \frac{D_0 A_i^n}{2\Delta x} - \frac{\Delta t \mu}{(\Delta x)^2} \delta^2\right)\right) \Delta u_i^n = \underbrace{-\Delta t \frac{F_{i+1}^n - F_{i-1}^n}{2\Delta x} + \Delta t \mu \frac{u_{i+1}^n - 2u_i^n + u_{i-1}^n}{(\Delta x)^2}}_{\text{RHS}_i^n}$$

$$\left\{\begin{array}{l} \theta = 1 & \text{first order} \\ \theta = 0.5 & \text{second order} \end{array}\right.$$
(31)

Deriving the matrix to invert:

$$\left(I + \theta \Delta t \frac{D_0 A_i^n}{2\Delta x}\right) \Delta u_i^n - \theta \frac{\Delta t \mu}{(\Delta x)^2} \delta^2 \Delta u_i^n = \text{RHS}_i^n$$

$$\left(I + \theta \Delta t \frac{D_0 A_i^n}{2\Delta x}\right) \Delta u_i^n - \theta \frac{\Delta t \mu}{(\Delta x)^2} \left(\Delta u_{i+1}^n - 2\Delta u_i^n + \Delta u_{i-1}^n\right) = \text{RHS}_i^n$$

$$\Delta u_i^n + \theta \frac{\Delta t}{2\Delta x} \left(A_{i-1}^n \Delta u_{i-1}^n - A_{i+1}^n \Delta u_{i+1}^n\right) - \theta \frac{\Delta t \mu}{(\Delta x)^2} \left(\Delta u_{i+1}^n - 2\Delta u_i^n + \Delta u_{i-1}^n\right) = \text{RHS}_i^n$$

$$A_i \Delta u_{i-1}^n + B_i \Delta u_i^n + C_i \Delta u_{i+1}^n = D_i$$

$$A_{i} = -\theta \frac{\mu \Delta t}{(\Delta x)^{2}} + \theta \frac{\Delta t}{2\Delta x} A_{i-1}^{n}$$

$$B_{i} = 1 + 2\theta \frac{\mu \Delta t}{(\Delta x)^{2}}$$

$$C_{i} = -\alpha \frac{\mu \Delta t}{(\Delta x)^{2}} - \theta \frac{\Delta t}{2\Delta x} A_{i-1}^{n}$$

$$D_{i} = RHS_{i}^{n}$$

$$(32)$$

and advancing the solution with:

$$u_i^{n+1} = u_i^n + \Delta u_i^n \tag{33}$$

In order to calculate Δu_i^n it is needed to invert matrix as follows:

$$\begin{pmatrix}
B_1 & C_1 & 0 & \cdots & \cdots & 0 \\
A_2 & B_2 & C_2 & 0 & \cdots & \cdots & 0 \\
0 & \ddots & \ddots & \ddots & 0 & \cdots & 0 \\
0 & 0 & A_i & B_i & C_i & 0 & 0 \\
0 & \cdots & 0 & \ddots & \ddots & \ddots & 0 \\
0 & \cdots & \cdots & 0 & A_{N-2} & B_{N-2} & C_{N-2} \\
0 & 0 & \cdots & \cdots & 0 & A_{N-1} & B_{N-1}
\end{pmatrix}
\begin{pmatrix}
\Delta u_1^n \\
\Delta u_2^n \\
\vdots \\
\vdots \\
\Delta u_{N-2}^n \\
\Delta u_{N-1}^n
\end{pmatrix} =
\begin{pmatrix}
D_1 - A_1 \cdot u_0 \\
D_2 \\
\vdots \\
\vdots \\
\Delta u_{N-2} \\
D_{N-2} \\
D_{N-1} - C_{N-1} \cdot u_N
\end{pmatrix}$$
(34)

The Beam and Warming method is extremely dispersive and therefore artiricial viscousity must be explicitly added. Beam and Warming used the following artificial viscosity term:

$$-\frac{w}{8}\left(u_{i+2}^n - 4u_{i+1}^n + 6u_i^n - 4u_{i-1}^n + u_{i-2}^n\right) \qquad 0 < w \le 1$$
(35)

which can be added to the RHS with no change in accuracy