Computational Fluid Dynamics

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Contents

1	Mathematical Problem	2
2		2 3 3 3
3	Stability Analysis	4
4	The Computer Program	6
5	2000 0200	7 7 7 7
6	Conclusions	8

1 Mathematical Problem

The Burgers Equation:

$$\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial y} = \frac{\partial}{\partial y} \left(\mu \frac{\partial u}{\partial y} \right) \tag{1}$$

By neglection the convective term, the parabolic model equation is obtained:

$$\frac{\partial u}{\partial t} = \frac{\partial}{\partial y} \left(\mu \frac{\partial u}{\partial y} \right) \tag{2}$$

$$\frac{\partial u}{\partial t} = \mu \frac{\partial^2 u}{\partial y^2} \tag{3}$$

$$\mu = 1.0 \tag{4}$$

Boundary and initial conditions:

$$u_{(y_0,t)} = u_0$$
 $u_{(y_1,t)} = u_1$ $u_{(y,t=0)} = f_{(y)}$ (5)

2 Numerical Scheme

A general explicit-implicit scheme for constant μ is given by:

$$u_i^{n+1} = u_i^n + \frac{\mu \Delta t}{\Delta y^2} \delta_{yy} \left[(1 - \alpha) u_i^n + \alpha u_i^{n+1} \right]$$

$$\tag{6}$$

where:

$$\alpha = \begin{cases} 0 & \text{Explicit} \\ \frac{1}{2} & \text{Crank-Nicolson} \\ 1 & \text{Implicit} \end{cases}$$
 (7)

and the order is:

$$\left[\Delta x^2, \Delta\left(\frac{1}{2} - \alpha\right)\right] \tag{8}$$

In delta form:

$$\left(I - \alpha \frac{\mu \Delta t}{\Delta y^2} \delta_{yy}\right) \Delta u_i^n = \frac{\mu \Delta t}{\Delta y^2} \delta_{yy} u_i^n \tag{9}$$

Applying the operators:

$$\Delta u_i^n - \alpha \frac{\mu \Delta t}{\Delta y^2} \delta_{yy} \left(\Delta u_i^n \right) = \frac{\mu \Delta t}{\Delta y^2} \delta_{yy} \left(u_i^n \right)$$
(10)

$$\Delta u_i^n - \alpha \frac{\mu \Delta t}{\Delta y^2} \delta_y \left(\Delta u_{\left(y_i + \frac{\Delta y}{2}\right)}^n - \Delta u_{\left(y_i - \frac{\Delta y}{2}\right)}^n \right) = \frac{\mu \Delta t}{\Delta y^2} \delta_y \left(u_{\left(y_i + \frac{\Delta y}{2}\right)}^n - u_{\left(y_i - \frac{\Delta y}{2}\right)}^n \right) \tag{11}$$

$$\Delta u_i^n - \alpha \frac{\mu \Delta t}{\Delta y^2} \left(\Delta u_{(y_i + \Delta y)}^n - 2\Delta u_{(y_i)}^n + \Delta u_{(y_i - \Delta y)}^n \right) = \frac{\mu \Delta t}{\Delta y^2} \left(u_{(y_i + \Delta y)}^n - 2u_{(y_i)}^n + u_{(y_i - \Delta y)}^n \right)$$
(12)

$$\Delta u_i^n - \alpha \frac{\mu \Delta t}{\Delta y^2} \left(\Delta u_{i+1}^n - 2\Delta u_i^n + \Delta u_{i-1}^n \right) = \frac{\mu \Delta t}{\Delta y^2} \left(u_{i+1}^n - 2u_i^n + u_{i-1}^n \right)$$
 (13)

$$A_i \Delta u_{i-1}^n + B_i \Delta u_i^n + C_i \Delta u_{i+1}^n = D_i \tag{14}$$

where:

$$A_i' = -\alpha \frac{\mu \Delta t}{\Delta y^2} \tag{15}$$

$$B_i' = 1 + 2\alpha \frac{\mu \Delta t}{\Delta y^2} \tag{16}$$

$$C_i' = -\alpha \frac{\mu \Delta t}{\Delta y^2} \tag{17}$$

$$D_i' = \frac{\mu \Delta t}{\Delta y^2} \left(u_{i+1}^n - 2u_i^n + u_{i-1}^n \right)$$
 (18)

and advancing the solution with:

$$u_i^{n+1} = u_i^n + \Delta u_i^n \tag{19}$$

To reduce problems of big flouting point numbers, define r:

$$r \triangleq \frac{\mu \Delta t}{\Delta y^2} \tag{20}$$

After dividing by r:

$$A_i = -\alpha \tag{21}$$

$$B_i = \frac{1}{r} + 2\alpha \tag{22}$$

$$C_i = -\alpha \tag{23}$$

$$C_{i} = -\alpha$$

$$D_{i} = (u_{i+1}^{n} - 2u_{i}^{n} + u_{i-1}^{n})$$
(23)
$$(24)$$

2.1For $\alpha = 0$

When $\alpha = 0$ The sheeme simplifies to:

$$\Delta u_i^n = rD_i \tag{25}$$

2.2For $\alpha \neq 0$

When $\alpha \neq 0$ in order to calculate Δu_i^n it is needed to invert matrix as follows:

$$\begin{pmatrix}
B_{1} & C_{1} & 0 & \cdots & \cdots & 0 \\
A_{2} & B_{2} & C_{2} & 0 & \cdots & \cdots & 0 \\
0 & \ddots & \ddots & \ddots & 0 & \cdots & 0 \\
0 & 0 & A_{i} & B_{i} & C_{i} & 0 & 0 \\
0 & \cdots & 0 & \ddots & \ddots & \ddots & 0 \\
0 & \cdots & \cdots & 0 & A_{N-2} & B_{N-2} & C_{N-2} \\
0 & 0 & \cdots & \cdots & 0 & A_{N-1} & B_{N-1}
\end{pmatrix}
\begin{pmatrix}
\Delta u_{1} \\
\Delta u_{2} \\
\vdots \\
\omega \\
\omega \\
\Delta u_{N-2} \\
\Delta u_{N-1}
\end{pmatrix} =
\begin{pmatrix}
D_{1} - A_{1} \cdot u_{0} \\
D_{2} \\
\vdots \\
\omega \\
\vdots \\
\Delta u_{N-2} \\
D_{N-2} \\
D_{N-1} - C_{N-1} \cdot u_{N}
\end{pmatrix}$$
(26)

In Our Case 2.3

$$y_0 = 0$$
 $u_0 = 0$ $y_1 = 1$ $u_1 = 1$ $f_{(y)} = 1$ $N = 11$ (27)

The scheme considered converged when the size of the RHS drops by 6 order of magnitude.

3 Stability Analysis

According to the "von Neumann Stability Analysis", the error between the differential and numerical equation is given by:

$$\varepsilon = N - D \tag{28}$$

where:

$$\varepsilon(y,t) = \sum_{m} c_m(t)e^{ik_m y} \tag{29}$$

since the analysis is linear, one can choose one element from the series. Morover, the error becomes:

$$\varepsilon^n = Z^n e^{iky} \tag{30}$$

The constant" k" is the wave number theat is related to the initial error.

$$Z = e^{\lambda \Delta t} \Rightarrow Z^n = e^{\lambda n \Delta t} = e^{\lambda t}$$
 (31)

Any element of the error therefore takes the form:

$$\varepsilon^n = e^{\lambda t} e^{iky} \tag{32}$$

The ratio between consecutive errors is the amplification factor, names "G":

$$G \triangleq \frac{\varepsilon^{n+1}}{\varepsilon^n} = e^{\lambda \Delta t} \tag{33}$$

The condition for Stability:

$$|G| < 1 \tag{34}$$

The error must satisfies the finite difference equation:

$$\Delta \varepsilon_i^n - \alpha \frac{\mu \Delta t}{\Delta y^2} \left(\Delta \varepsilon_{i+1}^n - 2\Delta \varepsilon_i^n + \Delta \varepsilon_{i-1}^n \right) = \frac{\mu \Delta t}{\Delta y^2} \left(\varepsilon_{i+1}^n - 2\varepsilon_i^n + \varepsilon_{i-1}^n \right)$$
 (35)

Where:

$$\Delta \varepsilon_i^n = \varepsilon_i^{n+1} - \varepsilon_i^n = e^{\lambda(t+\Delta t)} e^{iky} - e^{\lambda t} e^{iky} = e^{\lambda t} e^{iky} \left(e^{\lambda \Delta t} - 1 \right)$$
(36)

$$\Delta \varepsilon_{i+1}^n = \varepsilon_{i+1}^{n+1} - \varepsilon_{i+1}^n = e^{\lambda(t+\Delta t)} e^{ik(y+\Delta y)} - e^{\lambda t} e^{ik(y+\Delta y)} = e^{\lambda t} e^{ik(y+\Delta y)} \left(e^{\lambda \Delta t} - 1 \right)$$
(37)

$$\Delta \varepsilon_{i-1}^n = \varepsilon_{i-1}^{n+1} - \varepsilon_{i-1}^n = e^{\lambda(t+\Delta t)}e^{ik(y-\Delta y)} - e^{\lambda t}e^{ik(y-\Delta y)} = e^{\lambda t}e^{ik(y-\Delta y)} \left(e^{\lambda \Delta t} - 1\right) \tag{38}$$

Substitut the values:

$$\left(e^{\lambda t}e^{iky} - \alpha \frac{\mu \Delta t}{\Delta y^2} \left(e^{\lambda t}e^{ik(y+\Delta y)} - 2e^{\lambda t}e^{iky} + e^{\lambda t}e^{ik(y-\Delta y)}\right)\right) \left(e^{\lambda \Delta t} - 1\right) =
= \frac{\mu \Delta t}{\Delta y^2} \left(e^{\lambda t}e^{ik(y+\Delta y)} - 2e^{\lambda t}e^{iky} + e^{\lambda t}e^{ik(y-\Delta y)}\right)$$
(39)

Dividing by $e^{\lambda t}e^{iky}$:

$$\left(1 - \alpha \frac{\mu \Delta t}{\Delta y^2} \left(e^{ik\Delta y} - 2 + e^{-ik\Delta y}\right)\right) \left(e^{\lambda t} - 1\right) = \frac{\mu \Delta t}{\Delta y^2} \left(e^{ik\Delta y} - 2 + e^{-ik\Delta y}\right) \tag{40}$$

Define:

$$r \triangleq \frac{\mu \Delta t}{\Delta y^2}, \quad A \triangleq e^{ik\Delta y} - 2 + e^{ik\Delta y} = 2\cos\beta - 2, \quad \beta \triangleq k\Delta y$$
 (41)

After substitution and rearranging we get:

$$e^{\lambda t} - 1 = \frac{rA}{1 - \alpha rA} \tag{42}$$

For stability:

$$|G| < 1 \Rightarrow \left| 1 + \frac{rA}{1 - \alpha rA} \right| < 1 \tag{44}$$

For $\alpha = 0$:

$$G = 1 + rA$$

$$|1 - 4r| < 1$$

$$-1 < 1 - 4r < 1$$

$$-2 < -4r < 0$$

$$0 < 4r < 2$$

$$0 < r < \frac{1}{2}$$

$$\Delta t < \frac{\Delta y^2}{2\mu} \tag{45}$$

For $\alpha = 1$:

$$G = \frac{1}{1 - rA}$$

$$\left| \frac{1}{1 - rA} \right| < 1 \quad A \le 0$$

$$\downarrow \downarrow$$

$$(46)$$

smaller for every r (Δt)

For every α and $\Delta t \to \infty$:

$$\lim_{\Delta t \to \infty} |G| = \frac{1 - \alpha}{\alpha} \tag{47}$$

Which means that stability for infinite time step:

$$\alpha \begin{cases} <\frac{1}{2} & \text{Unstable} \\ =\frac{1}{2} & \text{Neutrally stable} \\ >\frac{1}{2} & Stable \end{cases} \tag{48}$$

4 The Computer Program

5 Results

5.1 $\alpha = 0$ - Explicit Scheme

From Eq.45 we get the stability limit of Δt :

$$\Delta t < \frac{\Delta y^2}{2\mu} = \frac{0.1^2}{2 \cdot 1} = 5 \cdot 10^{-3} \tag{49}$$

5.2 $\alpha = \frac{1}{2}$ - Crank-Nicolson Scheme

5.3 $\alpha = 1$ - Implicit Scheme

6 Conclusions