

Dynamics and Stability of Multiphase Flows
HW4

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1 Velocity Profile $U_{(z)}$ as a Function of z

Consider a liquid jet ejected from a circular orifice of radius a (point A). The jet has flux Q of fluid with density ρ that exits with velocity U_0 . The surface tension between the fluid and the air with pressure P_0 is σ .

By assuming laminar, inviscid, incompressible, axisymmetric and steady flow, we get the Bernoulli equation, between point A and arbitrary point downstream B:

$$\frac{1}{2}\rho U_0^2 + \rho g z + P_A = \frac{1}{2}\rho U_{(z)}^2 + P_B \quad (1)$$

The Laplace pressure for a jet:

$$\Delta P = \sigma \left(\frac{1}{R_1} + \frac{1}{R_2} \right) \quad \left\{ \begin{array}{l} R_1 = R_{1(z)} = \text{inner radius (normal to the flow)} \\ R_2 = R_{2(z)} = \text{outer radius (parallel to the flow)} \end{array} \right. \quad (2)$$

$$\Downarrow$$

$$P_A = P_0 + \sigma \left(\frac{1}{a} \right) \quad P_B = P_0 + \sigma \left(\frac{1}{R_1} + \frac{1}{R_2} \right)$$

substituting into the Bernoulli equation:

$$\frac{1}{2}\rho U_0^2 + \rho g z + P_0 + \frac{\sigma}{a} = \frac{1}{2}\rho U_{(z)}^2 + P_0 + \sigma \left(\frac{1}{R_1} + \frac{1}{R_2} \right) \quad (3)$$

Isolating $U_{(z)}$ we get:

$$\frac{U_{(z)}}{U_0} = \sqrt{1 + \frac{2gz}{U_0^2} + \frac{2\sigma}{\rho U_0^2} \left(\frac{1}{a} - \frac{1}{R_1} - \frac{1}{R_2} \right)} \quad (4)$$

and in dimensionless form:

$$\frac{U_{(z)}}{U_0} = \sqrt{1 + \underbrace{\frac{2z}{Fr^2 a}}_{\text{acceleration due to gravity}} + \underbrace{\frac{2}{We} \left(1 - \frac{a}{R_1} - \frac{a}{R_2} \right)}_{\text{deceleration due to surface tension}}} \quad (5)$$

where:

- $We = \frac{\rho U_0^2 a}{\sigma}$ is the Webber #
- $Fr = \frac{U_0^2}{ga}$ is the Froude #

Let's consider the curvature of the wall of the jet parallel to the flow:

$$\underbrace{\kappa}_{\text{principal curvature}} = - \frac{R_{1zz}}{(1 + R_{1z}^2)^{\frac{3}{2}}} \quad \underbrace{=}_{\text{linearization}} - R_{1zz} \quad (6)$$

The principal curvature equals one over the principal radius so:

$$\kappa = \frac{1}{R_2} = - \frac{\partial^2 R_1}{\partial z^2} \quad (7)$$

So, the velocity profile $U_{(z)}$ as a function of z is therefore:

$$\frac{U_{(z)}}{U_0} = \sqrt{1 + \frac{2z}{Fr^2 a} + \frac{2}{We} \left(1 - \frac{a}{R_1} + a \frac{\partial^2 R_1}{\partial z^2} \right)} \quad (8)$$



2 ODE of Jet Radius $r_{(z)}$

2.1 Formulate the ODE

From conservation of flux:

$$\begin{aligned}
 Q &= 2\pi \int_0^{R_1} U_{(z)} r dr = \pi R_1^2 U_{(z)} \underbrace{=}_{z=0} \pi a^2 U_0 \\
 Q_{(z)} &= \pi U_{(z)} R_1^2 \underbrace{=}_{\substack{\text{conservation} \\ \text{of mass}}} \pi a^2 U_0 \\
 &\Downarrow \\
 \frac{R_1}{a} &= \sqrt{\frac{U_0}{U_{(z)}}}
 \end{aligned} \tag{9}$$

After substituting $U_{(z)}$:

$$R_1 = \frac{a^2 U_0}{\sqrt{U_0^2 + 2gz + \frac{2\sigma}{\rho} \left(\frac{1}{a} - \frac{1}{R_1} + \frac{\partial^2 R_1}{\partial z^2} \right)}} \tag{10}$$

2.2 Boundary Conditions

Assume BC:

1. $R_{1(z=0)} = a$
2. $\left. \frac{\partial R_1}{\partial z} \right|_{z=0} = 0$

The first conditions is given.

The second condition is derived from the fact that at the begining, the jet looks like a cylinder.

3 Numerical Solver

The selected method is the Runge-Kutta-Merson method.

Isolating the second derivative of R_1 :

$$\begin{aligned}\frac{a^4}{R_1^4} &= 1 + \frac{2z}{\text{Fr}^2 a} + \frac{2}{\text{We}} \left(1 - \frac{a}{R_1} + a \frac{\partial^2 R_1}{\partial z^2} \right) \\ \frac{\text{We} \cdot a^4}{2R_1^4} &= \frac{\text{We}}{2} + \frac{\text{We}z}{\text{Fr}^2 a} + 1 - \frac{a}{R_1} + a \frac{\partial^2 R_1}{\partial z^2} \\ \frac{\partial^2 R_1}{\partial z^2} &= \frac{\text{We} \cdot a^3}{2R_1^4} - \frac{\text{We}}{2a} - \frac{\text{We} \cdot z}{\text{Fr}^2 a^2} - \frac{1}{a} + \frac{1}{R_1}\end{aligned}\tag{11}$$

Defining the system of equations:

$$x = R_1 \quad y = \frac{\partial R_1}{\partial z}\tag{12}$$

$$\begin{cases} \frac{\partial x}{\partial z} = g(z, x, y) \\ \frac{\partial y}{\partial z} = f(z, x, y) \end{cases} \quad \text{where} \quad \begin{cases} g = y \\ f = \frac{\text{We} \cdot a^3}{2x^4} - \frac{\text{We}}{2a} - \frac{\text{We} \cdot z}{\text{Fr}^2 a^2} - \frac{1}{a} + \frac{1}{x} \end{cases}\tag{13}$$

The Runge-Kutta-Merson method is defined as:

$$\begin{aligned}\begin{cases} x_{i+1} &= x_i + \frac{1}{6} (m_1 + 4m_4 + m_5) \\ y_{i+1} &= y_i + \frac{1}{6} (k_1 + 4k_4 + k_5) \end{cases} \\ \begin{cases} m_1 &= hg(z_i, x_i, y_i) \\ k_1 &= hf(z_i, x_i, y_i) \end{cases} \\ \begin{cases} m_2 &= hg(z_i + \frac{1}{3}h, x_i + \frac{1}{3}m_1, y_i + \frac{1}{3}k_1) \\ k_2 &= hf(z_i + \frac{1}{3}h, x_i + \frac{1}{3}m_1, y_i + \frac{1}{3}k_1) \end{cases} \\ \begin{cases} m_3 &= hg(z_i + \frac{1}{3}h, x_i + \frac{1}{6}(m_1 + m_2), y_i + \frac{1}{6}(k_1 + k_2)) \\ k_3 &= hf(z_i + \frac{1}{3}h, x_i + \frac{1}{6}(m_1 + m_2), y_i + \frac{1}{6}(k_1 + k_2)) \end{cases} \\ \begin{cases} m_4 &= hg(z_i + \frac{1}{2}h, x_i + \frac{1}{8}(m_1 + 3m_3), y_i + \frac{1}{8}(k_1 + 3k_3)) \\ k_4 &= hf(z_i + \frac{1}{2}h, x_i + \frac{1}{8}(m_1 + 3m_3), y_i + \frac{1}{8}(k_1 + 3k_3)) \end{cases} \\ \begin{cases} m_5 &= hg(z_i + h, x_i + \frac{1}{2}(m_1 - 3m_3 + 4m_4), y_i + \frac{1}{2}(k_1 - 3k_3 + 4k_4)) \\ k_5 &= hf(z_i + h, x_i + \frac{1}{2}(m_1 - 3m_3 + 4m_4), y_i + \frac{1}{2}(k_1 - 3k_3 + 4k_4)) \end{cases}\end{aligned}\tag{14}$$

4 Solution

$$U_0 = 1 \left[\frac{\text{m}}{\text{sec}} \right] \quad a = 0.05 \text{ [m]} \quad \rho = 998 \left[\frac{\text{kg}}{\text{m}^3} \right] \quad \sigma = 0.0728 \left[\frac{\text{N}}{\text{m}} \right] \quad g = 9.81 \left[\frac{\text{m}}{\text{sec}^2} \right] \quad (15)$$

$$h = 10^{-5} \text{ [m]}$$

$$z = [0, h, 2h, \dots, 5] \text{ [m]}$$

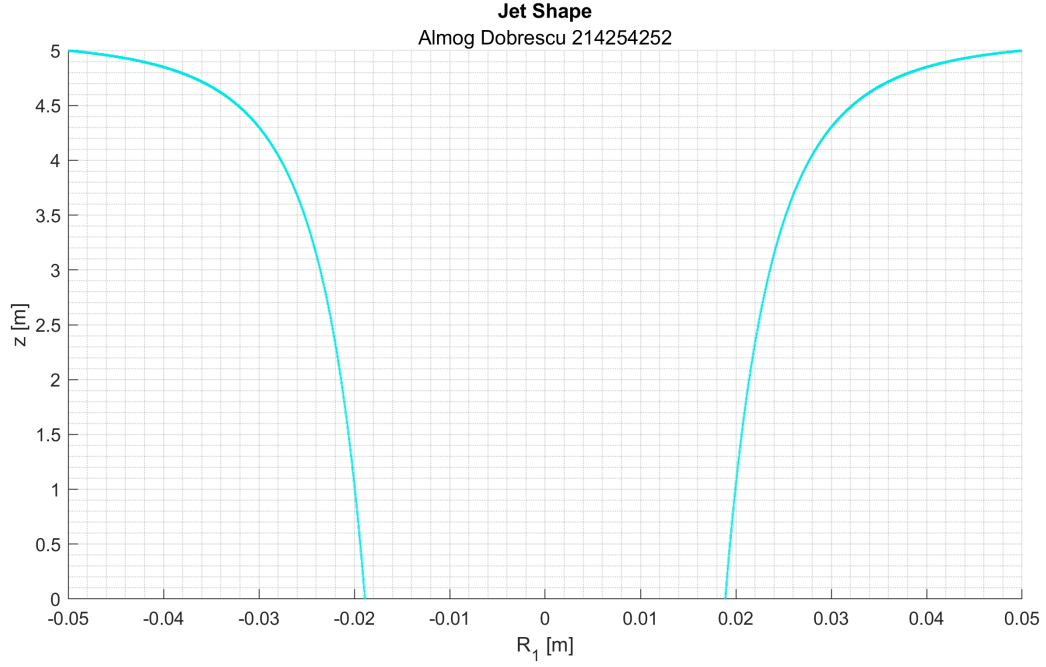


Figure 1: The jet shape

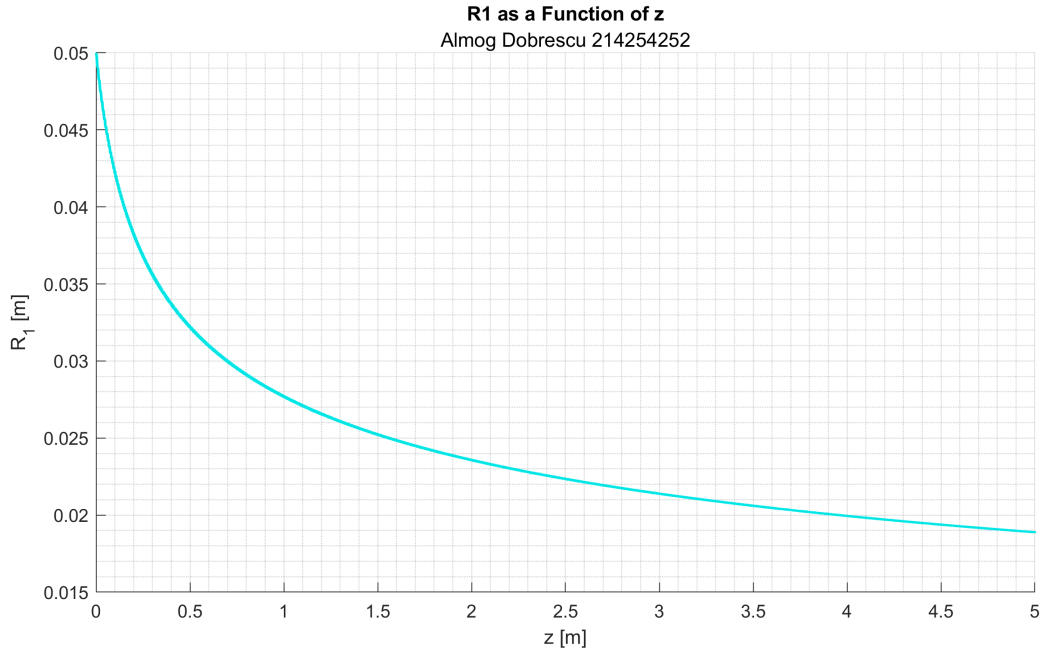


Figure 2: R_1 as a function of z

5 Compering The Results