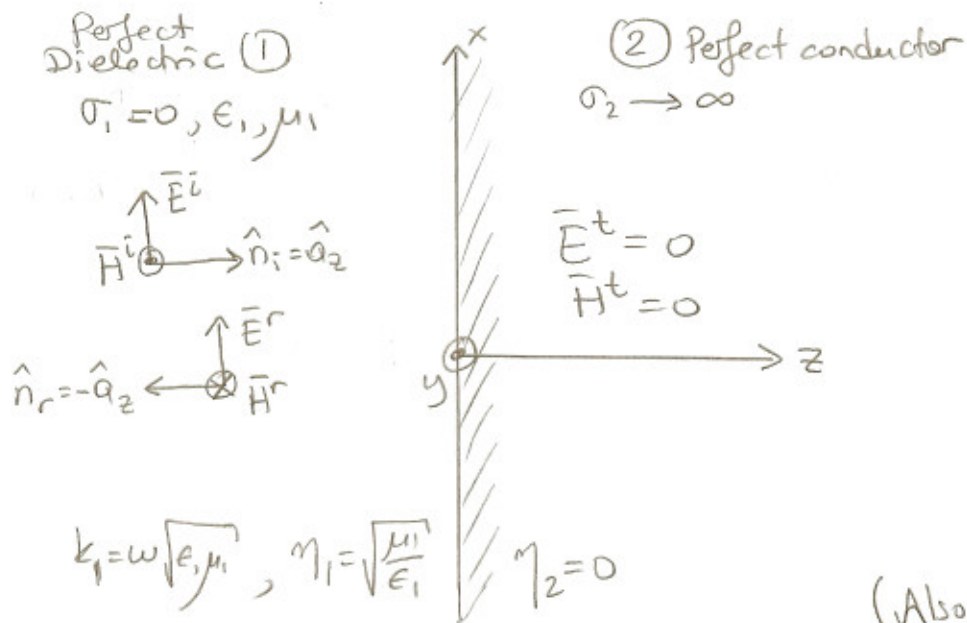


Case 3: Perfect Dielectric/Perfect Conductor Boundary



Remember, we have
 $\eta \approx \sqrt{\frac{j\omega\mu}{\sigma}}$ in good conductors

$$\lim_{\sigma \rightarrow \infty} \eta = 0 = \eta_2$$

in this medium (2)
 which is a perfect conductor with $\sigma \rightarrow \infty$.

(Also, $\lim_{\sigma \rightarrow \infty} \frac{1}{\sqrt{\pi f \mu \sigma}} = 0 \Rightarrow \delta_2 = 0$
 zero skin depth in a perfect conductor.

Then, $\Gamma = \frac{\eta_2 - \eta_1}{\eta_2 + \eta_1} \bigg|_{\eta_2=0} = -1$

and $T = \frac{2\eta_2}{\eta_2 + \eta_1} \bigg|_{\eta_2=0} = 0$

$\Gamma = -1$ and $T = 0$
 as expected
 can also be obtained by directly applying the BC's at $z=0$ as shown below:

B.C: $E_{\text{tang}1}^{\text{total}} = E_{\text{tang}2}^{\text{total}}$ at $z=0$

$$\left. \begin{aligned} \bar{E}^i &= \hat{a}_x E_1 e^{-jk_1 z} \\ \bar{E}^r &= \hat{a}_x E_2 e^{+jk_1 z} \\ \bar{E}^t &= 0 \end{aligned} \right\} \bar{E}_{\text{total}}^i = \bar{E}^i + \bar{E}^r = \hat{a}_x E_1 e^{-jk_1 z} + \hat{a}_x E_2 e^{+jk_1 z}$$

(tangential to the boundary)

$$\Rightarrow \left(E_1 e^{-jk_1 z} + E_2 e^{+jk_1 z} \right) \bigg|_{z=0} = 0$$

$$\Rightarrow E_1 + E_2 = 0$$

$$\Rightarrow E_2 = -E_1$$

$$\Rightarrow \Gamma = \frac{E_2}{E_1} = -1$$

In this case, presence of the perfect conductor causes "Standing waves" to form in the perfect dielectric side.

To see this, let's write the total fields in medium ①:

$$\bar{\mathbf{E}}_{\text{①}}^{\text{total}} = \bar{\mathbf{E}}^i + \bar{\mathbf{E}}^r = \hat{\mathbf{a}}_x E_1 e^{-jk_1 z} + \hat{\mathbf{a}}_x \underbrace{E_2}_{\substack{\text{reflected} \\ E_2 = -E_1}} e^{+jk_1 z}$$

$$\bar{\mathbf{E}}_{\text{①}}^{\text{total}} = \hat{\mathbf{a}}_x E_1 \underbrace{(e^{-jk_1 z} - e^{+jk_1 z})}_{-2j \sin(k_1 z)} \Rightarrow \boxed{\bar{\mathbf{E}}_{\text{①}}^{\text{total}} = -j 2 E_1 \sin(k_1 z) \hat{\mathbf{a}}_x}$$

in phasor domain.

Assuming E_1 to be real for computational simplicity,

$$\bar{\mathcal{E}}_{\text{①}}^{\text{total}}(t) = \text{Re} \{ \bar{\mathbf{E}}_{\text{①}}^{\text{total}} e^{j\omega t} \} = \text{Re} \{ -j 2 E_1 \sin(k_1 z) (\cos \omega t + j \sin \omega t) \hat{\mathbf{a}}_x \}$$

$$\boxed{\bar{\mathcal{E}}_{\text{①}}^{\text{total}}(t) = 2 E_1 \sin(k_1 z) \sin(\omega t) \hat{\mathbf{a}}_x}$$

in time-domain

Similarly, the total H-field in medium ① can be obtained as:

$$\bar{\mathbf{H}}_{\text{①}}^{\text{total}} = \bar{\mathbf{H}}^i + \bar{\mathbf{H}}^r = \hat{\mathbf{a}}_y \frac{E_1}{\eta_1} (e^{-jk_1 z} - \underbrace{e^{+jk_1 z}}_{-1}) = \hat{\mathbf{a}}_y \frac{E_1}{\eta_1} \underbrace{(e^{-jk_1 z} + e^{+jk_1 z})}_{2 \cos(k_1 z)}$$

$$\boxed{\bar{\mathbf{H}}_{\text{①}}^{\text{total}} = \frac{2 E_1}{\eta_1} \cos(k_1 z) \hat{\mathbf{a}}_y}$$

in phasor domain

$$\boxed{\bar{\mathcal{H}}_{\text{①}}^{\text{total}} = \frac{2 E_1}{\eta_1} \cos(k_1 z) \cos \omega t \hat{\mathbf{a}}_y}$$

in time-domain

Now, let's plot amplitudes $\bar{\mathcal{E}}_x^{\text{tot}}$ and $\bar{\mathcal{H}}_y^{\text{tot}}$ of medium ① as a function of distance z for several time instances

(for $\omega t = 0, \frac{\pi}{4}, \frac{\pi}{2}, \frac{3\pi}{4}, \pi, \frac{5\pi}{4}, \frac{3\pi}{2}, 2\pi$)

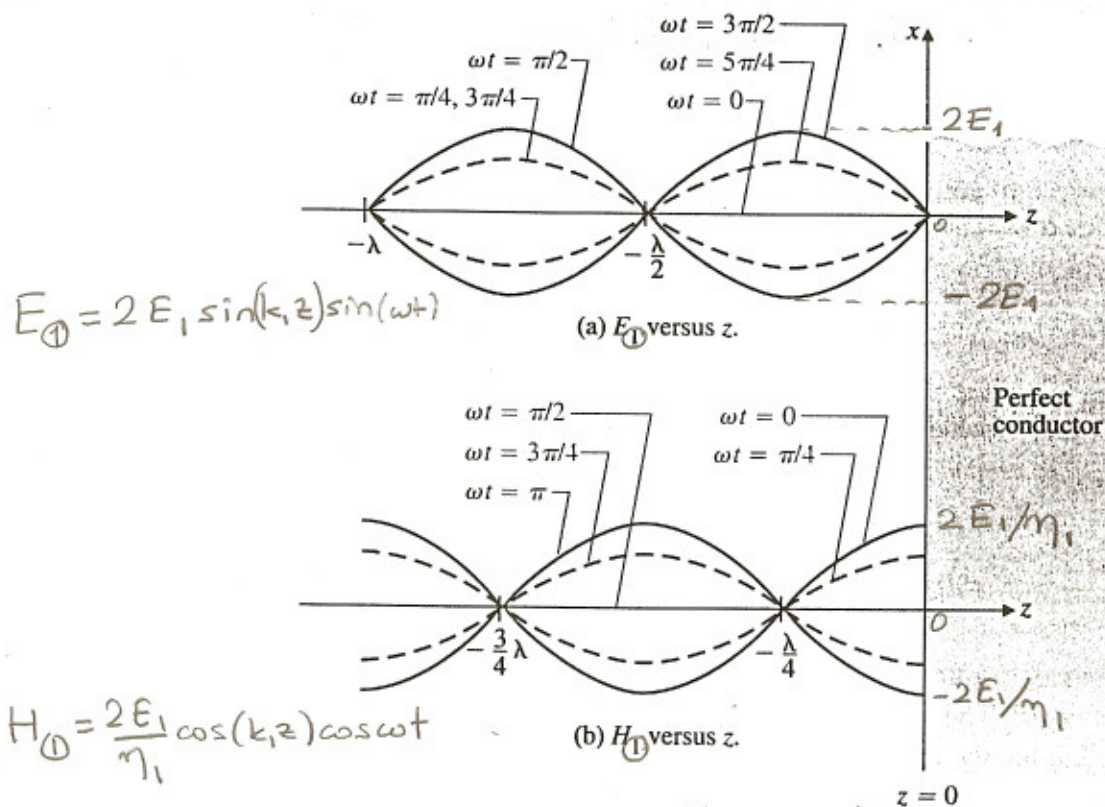


FIGURE 7-9 Standing waves of $\mathbf{E}_0 = \mathbf{a}_x E_0$ and $\mathbf{H}_0 = \mathbf{a}_y H_0$ for several values of ωt .

(*) From the superposition of oppositely traveling waves, STANDING WAVES are formed in the dielectric medium (i.e. for $z < 0$).

ωt	$\sin \omega t$	$\cos \omega t$	E_0	H_0
$0, 2\pi$	0	1	0	$\frac{2E_1}{\eta_1} \cos k_1 z$
$\pi/4$	$1/\sqrt{2}$	$1/\sqrt{2}$	$2E_1 \frac{1}{\sqrt{2}} \sin k_1 z$	$\frac{2E_1}{\eta_1} \frac{1}{\sqrt{2}} \cos k_1 z$
$\pi/2$	1	0	$2E_1 \sin k_1 z$	0
$3\pi/4$	$1/\sqrt{2}$	$-1/\sqrt{2}$	$2E_1 \frac{1}{\sqrt{2}} \sin k_1 z$	$-\frac{2E_1}{\eta_1} \frac{1}{\sqrt{2}} \cos k_1 z$
π	0	-1	0	$-\frac{2E_1}{\eta_1} \cos k_1 z$
$5\pi/4$	$-1/\sqrt{2}$	$-1/\sqrt{2}$	$-2E_1 \frac{1}{\sqrt{2}} \sin k_1 z$	$-\frac{2E_1}{\eta_1} \frac{1}{\sqrt{2}} \cos k_1 z$
$3\pi/2$	-1	0	$-2E_1 \sin k_1 z$	0

Basic observations about the STANDING WAVES

- ① Zeros (nulls) of $\bar{E}_0^{\text{tot}}(t)$ and Maxima/minima of $\bar{H}_0^{\text{tot}}(t)$ } occur when $k_1 z = -n\pi$ or using $k_1 = \frac{2\pi}{\lambda}$ when $z = -n \frac{\lambda}{2}$ ($n=0,1,2,3,\dots$)

Similarly,

- Maxima/minima of $\bar{E}_0^{\text{tot}}(t)$ and Zeros (nulls) of $\bar{H}_0^{\text{tot}}(t)$ } occur when $k_1 z = -(2n+1)\frac{\pi}{2}$ or $z = -(2n+1)\frac{\lambda}{4}$ ($n=0,1,2,3,\dots$)

- ② \bar{E}_0^{total} always becomes zero on the perfect conductor boundary, i.e., at $z=0$ due to the $\sin(k_1 z)$ term.
- ③ \bar{H}_0^{total} takes its either maximum or minimum value (depending on the value of $\cos \omega t$ term at a given time instant) at the boundary $z=0$ due to the $\cos(k_1 z)$ term.
- ④ The surface current density vector \bar{J}_s at $z=0$ boundary can be computed using the boundary condition on tangential H-field (taking into account the fact that $\bar{H}=0$ in the perfect conductor medium).

Diagram: A vertical line at $z=0$ separates the "perfect dielectric" region (left) from the "perfect conductor (PEC)" region (right). A unit vector $\hat{n} = -\hat{a}_z$ points from the PEC into the dielectric. The boundary is labeled "boundary at $z=0$ ".

B.C: $\hat{n} \times (\bar{H}_0 - \bar{H}_{\text{in}}) = \bar{J}_s \Rightarrow \bar{J}_s = -\hat{a}_z \times \bar{H}_0^{\text{tot}}|_{z=0}$

$\Rightarrow \bar{J}_s = \frac{2E_1}{\eta_1} \cos \omega t \hat{a}_x \text{ (A/m)}$

⑤ Distance between two successive nodes is $\frac{\lambda}{2}$

Distance " " " maxima is $\frac{\lambda}{2}$

Distance " " " minima is $\frac{\lambda}{2}$

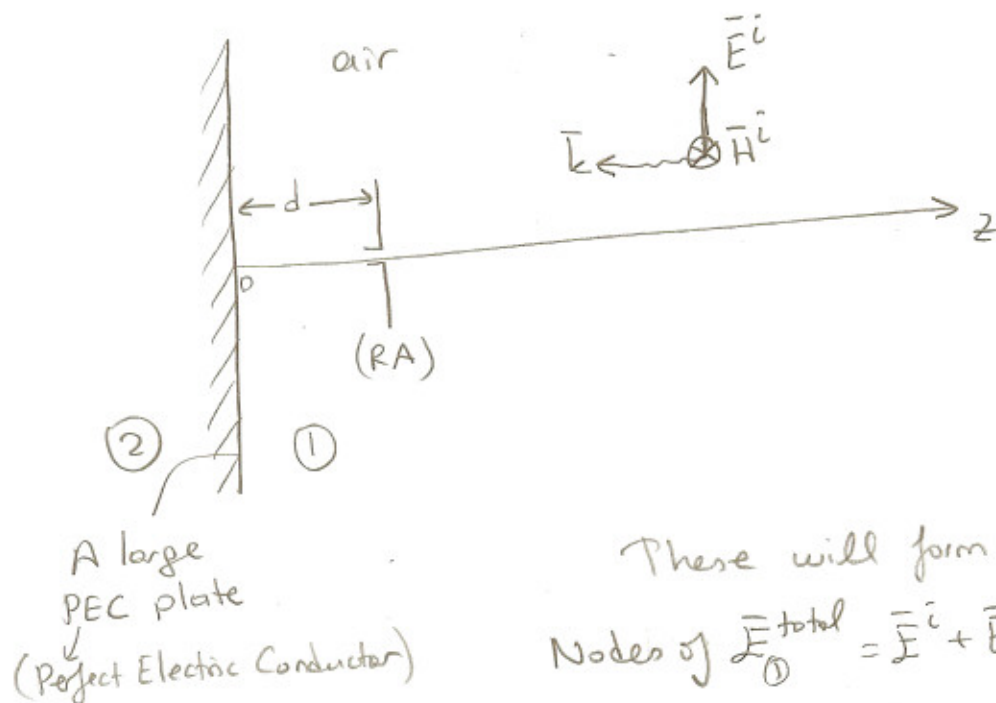
Distance between a node and the closest extremum point (minimum or maximum) is $\frac{\lambda}{4}$.

⑥ Standing waves $\bar{E}_{(1)}^{\text{tot}}$ and $\bar{H}_{(1)}^{\text{tot}}$ are in time quadrature (having $\sin \omega t$ and $\cos \omega t$ terms) and in space quadrature (i.e, shifted by a quarter wavelength ($\frac{\lambda}{4}$) with respect to each other due to the $\sin k_z z$ and $\cos k_z z$ terms).

Problem: In the figure below, a receiving antenna (RA) is placed "d" meters in front of a large PEC plate. The antenna responds to the total field $\vec{E}^{\text{total}} = \vec{E}^i + \vec{E}^r$.

If the incoming EM radiation is a u.p.w oscillating at $f = 100 \text{ MHz}$,

- For which d values, antenna picks up the max. field?
- For which d values, antenna senses no field at all?



Soln: We know that there will be both incident and reflected field in air due to the presence of PEC plate.

These will form standing waves in air region.

Nodes of $\vec{E}^{\text{total}} = \vec{E}^i + \vec{E}^r$ occur at $z = n \frac{\lambda}{2}$

Extremum points of \vec{E}^{total} occur at $z = (2n+1) \frac{\lambda}{4}$

where $\lambda = \frac{c}{f} = \frac{3 \times 10^8}{100 \times 10^6} = 3 \text{ m}$. $\lambda = 3 \text{ m}$ $\rightarrow \lambda/2 = 1.5 \text{ m}$
 $\rightarrow \lambda/4 = 0.75 \text{ m}$

So, (a) for $d = 0, 1.5 \text{ m}, 3 \text{ m}, 4.5 \text{ m}, 6 \text{ m}, \dots$ RA senses zero intensity

(b) For $d = 0.75 \text{ m}, 2.25 \text{ m}, 3.75 \text{ m}, 5.25 \text{ m}, \dots$ RA senses maximum intensity.