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Abstract—Some dynamic model's are greatly simplified by implementantion of Langrangian equations. The use of a Langrangian overhauls Newton style internal mechanics and focuses on

1 Introduction: Inverted Pendu-Lum Problem

An inverted pendulum is visualized in the figure below. With a swinging pendulum bob of mass m_2 , a sliding base of mass m_1 , and a pendulum of length L.

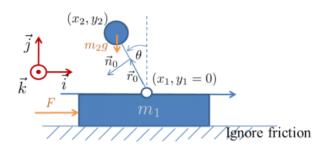


Fig. 1. Inverted Pendulum

Variables x_1 and θ will be used to describe the system. Describing the bob's coordinates in terms of those variables is the following:

$$x_2 = x_1 - L\sin(\theta)$$

$$y_2 = y_1 + L\cos(\theta) = L\cos(\theta)$$
(1)

$$\dot{x}_2 = \dot{x}_1 - L\cos(\theta)\dot{\theta}
\dot{y}_2 = \dot{y}_1 - L\sin(\theta)\dot{\theta} = -L\sin(\theta)\dot{\theta}$$
(2)

Having described the bob and mass in terms of (x_1, θ) the system needs to be solved for \ddot{x} , and $\ddot{\theta}$. Normally, these functions would be solved for from several force summation equations, but in this system, we solve using a Langrangian to simplify solutions.

2 Langrangian Solution

A Lagrangian is ideal for systems of conservative forces as it is dependent on energy. The general form, where KE is Kinetic Energy and PE is Potential Energy, is the following:

$$\mathcal{L} = KE - PE$$

$$\mathcal{L} = \frac{1}{2}m_1\dot{x_1}^2 + \frac{1}{2}m_2(\dot{x_2}^2 + \dot{y_2}^2) - m_2gy_2$$
(3)

The Euler-Lagrangian Equation is the following, and will be used to solve the system:

$$\frac{d}{dt}(\frac{\partial \mathcal{L}}{\partial \dot{x_1}}) - \frac{\partial \mathcal{L}}{\partial x_1} = F \tag{4}$$

$$\frac{d}{dt}(\frac{\partial \mathcal{L}}{\partial \dot{\theta}}) - \frac{\partial \mathcal{L}}{\partial \theta} = 0 \tag{5}$$

Use eq. (1) and (2) to substitute for $\dot{x_2}$, $\dot{y_2}$, into the Langrangian equation (3) and then differentiate i.e. solve the Euler-Lagrangian Equation:

$$\mathcal{L} = \frac{1}{2}m_1\dot{x_1}^2 + \frac{1}{2}m_2((\dot{x_1} - L\cos(\theta)\dot{\theta})^2 + (\dot{y_1} - L\sin(\theta)\dot{\theta})^2) - m_2g(-L\sin(\theta)\dot{\theta})$$
(6)

$$\mathcal{L} = \frac{1}{2}m_1\dot{x_1}^2 + \frac{1}{2}((\dot{x_1}^2 - L^2\dot{\theta}^2) - 2L\dot{x_1}\cos(\theta)\dot{\theta}) - m_2g(-L\sin(\theta)\dot{\theta})$$
(7)

$$\frac{\partial \mathcal{L}}{\partial \dot{x_1}} = m_1 \dot{x_1} + \frac{1}{2} (2\dot{x_1} + 2L\theta \cos(\theta))$$

$$= \dot{x_1} (m_1 + m_2) + \dot{\theta} m_2 L \cos(\theta)$$
(8)

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$$\frac{d}{dt}(\frac{\partial \mathcal{L}}{\partial \dot{x_1}}) = \ddot{x_1}(m_1 + m_2) + m_2 L^2(\dot{\theta}^2 \sin(\theta) - \ddot{\theta} \cos(\theta))$$
(9)

$$\frac{\partial \mathcal{L}}{\partial x_1} = 0 \tag{10}$$

Putting the two terms of the Euler-Lagrangian eq (4) together, we derive that:

$$\ddot{x}_1(m_1 + m_2) + m_2 L(\dot{\theta}^2 \sin \theta - \ddot{\theta} \cos \theta) = F \quad (11)$$

Then the second Euler-Lagrangian equation (5) is solved:

$$\frac{\partial \mathcal{L}}{\partial \dot{\theta}} = \frac{1}{2} m_2 (2L^2 \dot{\theta} - 2L \cos \theta \dot{x_1}) \tag{12}$$

$$\frac{d}{dt}\frac{\partial \mathcal{L}}{\partial \dot{\theta}} = \frac{1}{2}m_2(2L^2\ddot{\theta} - 2L\ddot{x_1}\cos\theta + 2L\sin\theta\dot{x_1}\dot{\theta})$$
(13)

$$\frac{\partial \mathcal{L}}{\partial \theta} = \frac{1}{2} m_2 (2L \sin \theta \dot{x_1} \dot{\theta} + 2gL \sin \theta)$$
 (14)

Subtracting the two terms for the second Euler-Lagrangian equation (6) gives:

$$0 = m_2 L^2 \ddot{\theta} - m_2 L \cos \theta \ddot{x} - m_2 g L \sin \theta \qquad (15)$$

The project manual describes these solutions equations in the form of:

$$[M] \begin{bmatrix} \ddot{x_1} \\ \ddot{\theta} \end{bmatrix} = B(\theta, \dot{\theta}, x_1.\dot{x_1}, F)$$

Where M is a 2x2 matrix. We take the solved Euler-Lagrangian solutions (11) and (15) and put them in a form that can be converted easily to this matrix algebra.

$$(m_1 + m_2)\ddot{x_1} - m_2L^2\cos\theta\ddot{\theta} = F - m_2L^2\dot{\theta}^2\sin\theta$$

$$-m_2L^2\cos\theta\ddot{x_1} + m_2L^2\ddot{\theta} = m_2qL\sin\theta$$
(15)

Our solutions for \ddot{x} and $\ddot{\theta}$ from a Lagrangian approached are then expressed in the matrix form:

$$\begin{bmatrix} m_1 + m_2 & -m_2 L^2 \cos \theta \\ -m_2 L^2 \cos \theta & m_2 L^2 \end{bmatrix} \begin{bmatrix} \ddot{x} \\ \ddot{\theta} \end{bmatrix} = \begin{bmatrix} F - m_2 L \dot{\theta}^2 \sin \theta \\ m_2 g L \sin \theta \end{bmatrix}$$
(16)

These accleration equations can then solve for position using Matlab integration like ODE45().



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