

# Épisode 2

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Spineurs, bases et représentations

$$\text{ECOC} : X, Y, Z, S_z, (S^2) : \mathcal{E}_{\vec{r}} \otimes \mathcal{E}_s = \mathcal{E} \quad |\vec{r}, s\rangle \quad (1)$$

$$\text{ECOC} : P_x, P_y, P_z, S_z; |\vec{p}, s\rangle \quad (2)$$

$$\text{ECOC} : H_0, \mathbf{L}^2, L_z, S_z; |n, l, m, s\rangle \quad (3)$$

Relation de fermeture dans  $\mathcal{E}$  :

$$\begin{aligned} 1 &= 1_{\vec{r}} \otimes 1_S = \int d^3r |\vec{r}\rangle \langle \vec{r}| \otimes \sum_{\epsilon} |\epsilon\rangle \langle \epsilon| \\ \implies 1 &= \sum_{\epsilon} \int d^3r |\vec{r}\epsilon\rangle \langle \vec{r}, \epsilon| \end{aligned}$$

Preuve très similaire pour les autres bases.

$$|\psi\rangle = 1 |\psi\rangle = \sum_{\epsilon} \int d^3r |\vec{r}, \epsilon\rangle \underbrace{\langle \vec{r}, \epsilon | \psi \rangle}_{\Psi_{\epsilon}(\vec{r})}$$

Représentation matricielle :

$$\begin{aligned} |\psi\rangle &= \int d^3r \begin{pmatrix} \psi_+(\vec{r}) \\ \psi_-(\vec{r}) \end{pmatrix} |\vec{r}\rangle \\ \langle \vec{r} | \psi \rangle &= \begin{pmatrix} \psi_+(\vec{r}) \\ \psi_-(\vec{r}) \end{pmatrix} = [\psi] \text{ (Spineur!)} \\ \langle \psi | &= \int d^3r \begin{pmatrix} \psi_+^*(\vec{r}) & \psi_-^*(\vec{r}) \end{pmatrix} \langle \vec{r} | \\ |\psi\rangle = 1 |\psi\rangle &= \sum_{\epsilon} \sum_{n,l,m} |n, l, m, \epsilon\rangle \overbrace{\langle n, l, m, \epsilon | \psi \rangle}^{C_{n,l,m,\epsilon}} \end{aligned} \quad (4)$$

si

$$|\vec{r}\rangle \langle n, l, m| = R_{n,l}(r) Y_l^m(\theta, \phi)$$

$$\langle \vec{r} | \psi \rangle = \sum_{n,l,m} \sum_{\epsilon} \underbrace{\langle \vec{r} | n, l, m \rangle}_{R_{n,l}(r) Y_l^m(\theta, \phi)} |\epsilon\rangle C_{n,l,m,\epsilon} = \sum_{n,l,m} \begin{pmatrix} c_{n,l,m,+R_{n,l}(r) Y_l^m(\theta, \phi)} \\ c_{n,l,m,-R_{n,l}(r) Y_l^m(\theta, \phi)} \end{pmatrix}$$

Norme

$$\langle \psi | \psi \rangle = \int d^3 r [\psi^*][\psi]$$

Produit interieur

$$\langle \psi | \phi \rangle = \int d^3 r [\psi^*][\phi]$$

Élément de matrice

$$\langle \Psi | \mathbb{K} A \mathbb{K} | \Phi \rangle = \sum_{\epsilon, \epsilon'} \int d^3 r d^3 r' \underbrace{\langle \psi | \vec{r}', \epsilon' \rangle}_{\psi_{\epsilon}^*(\vec{r}')} \underbrace{\langle \vec{r}', \epsilon' | A | \vec{r}, \epsilon \rangle}_{A_{\epsilon' \epsilon}(\vec{r}', \vec{r})} \underbrace{\langle \vec{r}, \epsilon | \psi \rangle}_{\psi_{\epsilon}(\vec{r})} = \int d^3 r d^3 r' [\psi^*][\mathbb{K} A][\phi]$$

$$L_z \rightarrow_{|\vec{r}\rangle} \frac{\hbar}{i} \frac{\partial}{\partial \varphi} \rightarrow_{\mathcal{E}_{\vec{r}} \otimes \mathcal{E}_{\epsilon}} = \frac{\hbar}{i} \begin{pmatrix} \frac{\partial}{\partial \phi} & 0 \\ 0 & \frac{\partial}{\partial \phi} \end{pmatrix}$$

Mesure

La quatrième postula reste valable :

$|\psi\rangle$  : vecteur d'état

$$\mathcal{P}(\underbrace{a_n}_{\text{val dicrete d'un obs}}) = |\langle \varphi_n | \psi \rangle|^2$$

$$d\mathcal{P}(\underbrace{\alpha}_{\text{val continue d'un obs}}) = |\langle \omega_{\alpha} | \psi \rangle|^2 d\alpha$$

$$\mathcal{P}(a_n) = \sum_{i=1}^{g_n} |\langle \varphi_n^i | \psi \rangle|^2$$

Dans notre cas, qui est une combinaisons de discret et continue, on a :

$$d\mathcal{P}(\vec{r}, \pm) = |\langle \vec{r}, \pm | \psi \rangle|^2 d^3 r$$

$$\mathcal{P}_{\pm} = \int d\mathcal{P} = \int d^3 r |\psi(\vec{r})|^2$$

$$\text{Si } [\psi] = \begin{pmatrix} \psi_+(r, \theta, \varphi) \\ \psi_-(r, \theta, \varphi) \end{pmatrix}$$

$$\mathcal{P}_{\tilde{L}^2} = \left| \int \sum_{l', m'} Y_l^{m*} a_{l', m', +}(r) Y_{l'}^{m'} d\Omega \right|^2 + \left| \int \sum_{l', m'} Y_l^{m*} a_{l', m', -}(r) Y_{l'}^{m'} d\Omega \right|^2$$