Action je crois

$$S = \int_{t_i}^{t_f} \mathrm{d}t \frac{1}{2} m \mathbf{v}^2(T) \quad \text{non-relativiste}$$

$$\rightarrow -m \int_{A}^{V} d\tau = -m \int_{A}^{B} dt \sqrt{1 - \mathbf{v}^{2}} = -m \int_{A}^{B} dt \{1 - \frac{1}{2} \mathbf{v}^{+} \frac{12}{\mathbf{v}^{2}} - \dots \} = -m(t_{B} - T_{a}) + \frac{1}{2} \int_{A}^{B} d\tau m \mathbf{v}^{2} - \dots$$

$$L = -m\sqrt{1 - \mathbf{v}^2}$$

$$\mathbf{p} = \frac{\partial L}{\partial \mathbf{v}} = \frac{m\mathbf{v}}{\sqrt{1 - \mathbf{v}^2}}$$

$$H = \mathbf{p} d\mathbf{v} - L = \frac{m\mathbf{v}^2}{\sqrt{1 - v^2}} + m\sqrt{1 - v^2} = \frac{m}{\sqrt{1 - \mathbf{v}^2}} = \sqrt{\mathbf{p}^2 + m^2}$$

4-impulsion

$$p^{\mu} = (E, \mathbf{p}) = mu^{\mu}$$

$$u^{\mu} = \frac{\mathrm{d}x^{\mu}}{\mathrm{d}\tau} = \left(\frac{\mathrm{d}t}{\mathrm{d}\tau}, \frac{\mathrm{d}\mathbf{r}}{\mathrm{d}\tau}\right) = \left(\frac{1}{\sqrt{1-v^2}}, \frac{\mathbf{v}}{\sqrt{1-v^2}}\right)$$

Invarient associé au quadri-vecteur

$$p^{\mu}p_{mu} = p^2 = E^2 - \mathbf{p}^2 = m^2 \mathcal{U} = m^2$$
$$p^2 = m^2$$

$$\mathbf{v} = \frac{\mathbf{p}}{E}$$

masse nulle $\rightarrow p^2 = 0 \rightarrow T = |\mathbf{p}|$

$$p_{\pi} = (m_{\pi}, \mathbf{0})$$

$$p_\pi = p_\mu + p_\nu$$

$$p_{\nu} = p_{\pi} - p_{\mu}$$

$$p_{\nu}^2 = p_{\pi}^2 + p_{\mu}^2 - 2p_{\pi}p_{\mu}$$

$$0 = m_{\pi}^2 + m_{\mu}^2 - 2m_{\pi}E_{\mu}$$

$$E_{\mu} = \frac{m_{\pi}^2 + m_{\mu}^2}{2m_{\pi}}$$

$$E_{nu} = m_{\pi} - E_{\mu} = \frac{m_{\pi}^2 - m_{\mu}^2}{m_{\pi}} = |\mathbf{p}_{\nu}| = |\mathbf{p}_{\nu}|$$

$$|\mathbf{v}_{\mu}| = \frac{m_{\pi}^2 - m_{\mu}^2}{m_{\pi}^2 - m_{\mu}^2}$$

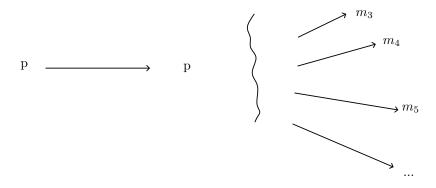


FIGURE 1 – proton incident

Énérgie de seuil? (Plus d'Énérgie cinétique à la fin)

$$E = \sum_{i=3}^{N} m_i$$

$$p_1^{\mu} + p_2^{\mu} = \sum_{i=3}^{N} p_i^{\mu}$$

$$(p_1^{\mu} + p_2^{\mu})^2 = \left[\sum_{i=3}^{N} p_i^{\mu}\right]^2 = \left(\sum_{i=3}^{N} m_i\right)^2$$

au <u>seuil</u> $p_i = (m_i, \mathbf{0})$

$$E_p = \frac{M_{\text{tot}}^2 - 2m_p^2}{2m_p}$$

L'énérige requise va comme le carré des masses.

Unités naturelles

$$\hbar c = 197 \text{MeV} \cdot \text{fm}$$

$$\frac{\hbar c}{m_e c^2} = \frac{197 {\rm MeV fm}}{0,511 {\rm MeV}} = 400 {\rm fm} \quad {\rm longueur~d'onde~de~Compton}$$

Constante de structure fine

$$\alpha = \frac{e^2}{4\pi\hbar c} \sim \frac{1}{137}$$

Heaviside-Lorentz

$$\epsilon_0 = 1$$
 $\mu_0 = 1$

$$\frac{\alpha \hbar}{m_e c} = \frac{e^2}{4\pi m_e c^2} = \mbox{ rayon classique de l'éléctron}$$

$$\frac{\hbar}{m_e\alpha c} = \frac{4\pi\hbar c\hbar}{e^2m_ec} = \frac{4\pi\hbar^2}{m_ee^2} = \text{ rayon de bohr}$$

$$\psi(\mathbf{r}) = \frac{1}{\sqrt{\nu}} e^{i\mathbf{p}\cdot\mathbf{r}}$$

Condition au limite périodiques à l'univers (une boîte bien sûr)

$$\begin{split} e^{ip_xL_x} &= 1 \\ \mathbf{p} &= 2\pi \left(\frac{n_x}{L_x}, \frac{n_y}{L_y}, \frac{n_z}{L_z}\right) \end{split}$$

où $n \in \mathbb{Z}$

$$\Delta p_x = \frac{2\pi}{L_x} \leftrightarrow \Delta n_x = 1$$

$$\Delta p_x \Delta p_y \Delta p_z = \frac{(2\pi)^3}{L_x L_y L_z} = \frac{(2\pi)^3}{\nu}$$

$$\sum_{\mathbf{p}} f_{\mathbf{p}} = \nu \int \frac{\mathrm{d}P}{(2\pi)^3} f_{\mathbf{p}}$$

$$\langle \mathbf{p}' | \mathbf{p} \rangle = \delta_{\mathbf{p}\mathbf{p}'} \& \sum_{\mathbf{p}} |\mathbf{p} \rangle \langle \mathbf{p}| = 1$$

$$\langle \mathbf{r} | \mathbf{p} \rangle = \frac{1}{\sqrt{\nu}} e^{i\mathbf{p} \cdot \mathbf{r}}$$

$$\delta_{\mathbf{p}\mathbf{p}'} \to \frac{(2\pi)^3}{\nu} \delta(\mathbf{p} - \mathbf{p}')$$
Normalisation continue
$$\begin{cases} \langle \mathbf{p}' | \mathbf{p} \rangle = (2\pi)^3 \delta(\mathbf{p} - \mathbf{p}'') \\ \int \frac{\mathrm{d}^3 P}{(2\pi)^3} \langle \mathbf{p} | \mathbf{p} \rangle = 1 \end{cases}$$

On a le problème que d^3P n'est pas invarient de Lorentz d^3pdp^0 en revanche l'est

$$\mathrm{d}^3\gamma\mathrm{d}t = \mathrm{d}^4x = \mathrm{d}^4x'$$

Le Jacobien

$$J=1$$

$$\int \frac{\mathrm{d}^3 P}{(2\pi)^3} \to \int \frac{\mathrm{d}^4}{(2\pi)^3} \delta(p^2 - m) \theta(p^0)$$
$$\int \frac{\mathrm{d}^4 P}{(2\pi)^3} \delta\left((p^0 - E_{\mathbf{p}})(p^0 + E_{\mathbf{p}})\right) \Theta(p^0)$$

$$E_{\mathbf{p}} = \sqrt{\mathbf{p}^2 + m^2}$$

$$\delta(\beta x) = \frac{1}{|\beta|}\delta(x)$$

$$\begin{split} \int \frac{\mathrm{d}^4 P}{(2\pi^3) 2E_p} \delta(p^0 - & E_p) = \int \frac{\mathrm{d}^3 P}{(2\pi)^3 2E_p} \\ \text{Normalisation relativiste} \begin{cases} \int \frac{\mathrm{d}^3 P}{(2\pi)^3 2E_p} \left\langle \mathbf{p} | \mathbf{p} \right\rangle = \mathbb{1} \\ \left\langle \mathbf{p} | \mathbf{p}' \right\rangle = 2E_{\mathbf{p}} \delta(\mathbf{p} - \mathbf{p}'(2\pi)^2) \end{cases} \end{split}$$