Théorie des champs quantique

Classique
$$\rightarrow$$
 Quantique
$$[A, B]_p \rightarrow \frac{1}{i\hbar} [\hat{A}, \hat{B}]$$

$$L = \int d\mathcal{L} = \int d^3r \left\{ \dot{\phi}^2 - m^2 \phi^2 - (\nabla \phi)^2 \right\}$$

$$H = \int d^3\tau \left\{ \pi(\mathbf{r}) + m^2 \phi^2 + (\nabla \phi)^2 \right\} \quad \text{où} \quad \pi(\mathbf{r}) = \frac{\partial \mathcal{L}}{\partial \phi} \dot{\phi}$$

$$[\phi(\mathbf{r}), \pi(\mathbf{r}')]_p = \delta(\mathbf{r} - \mathbf{r}') \xrightarrow{M.Q.} [\phi(\mathbf{r}), \pi(\mathbf{r}')] = i\delta(\mathbf{r} - \mathbf{r}') \qquad \hbar = 1$$

On introduit un transformé de Fourrier pour profiter de la symétrie de translation :

$$\phi(\mathbf{r}) = \frac{1}{\nu} \sum_{\mathbf{p}} \phi_{\mathbf{p}} e^{i\mathbf{p}\cdot\mathbf{r}}$$
, $\phi_{\mathbf{p}} = \int d^3r \phi(\mathbf{r}) e^{-i\mathbf{p}\cdot\mathbf{r}}$

L'opérateur, une fois passé dans la TF, n'est plus hermitien. Sa conjugason hermitien préserve quand même une expression simple :

$$\phi^{\dagger}(\mathbf{r}) = \phi(\mathbf{r}) \qquad \phi_{\mathbf{p}^{\dagger} = \phi_{-\mathbf{p}}}$$

Pareil pour π

$$\pi^{\dagger}(\mathbf{r}) = \pi(\mathbf{r})$$
 $\pi_{\mathbf{p}^{\dagger} = \pi_{-\mathbf{p}}}$

$$H = \frac{1}{2} \int d^3 r \frac{1}{\nu^2} \sum_{\mathbf{p}, \mathbf{p}'} \left\{ \pi_{\mathbf{p}} \pi_{\mathbf{p}'} e^{i(\mathbf{p} \cdot \mathbf{p}')} + \cdots \right\}$$

L'intégrale sur les exponentielles donne des $\delta(p-p')$

$$\begin{split} H &= \frac{1}{2} \frac{1}{\nu} \sum_{p} \left\{ \pi_{\mathbf{p}} \pi_{-\mathbf{p}} + \left(m^2 - \mathbf{p}^2 \right) \phi_{\mathbf{p}} \phi_{-\mathbf{p}} \right\} \\ H &= \frac{1}{2} \frac{1}{\nu} \sum_{p} \left\{ \pi_{\mathbf{p}}^{\dagger} \pi_{\mathbf{p}} + \omega_{p}^2 \phi_{\mathbf{p}}^{\dagger} \phi_{\mathbf{p}} \right\} \end{split}$$

$$[\phi_{\mathbf{p}}, \pi_{\mathbf{p}'}] = \int d^3r d^3r' e^{-i(\mathbf{p}\cdot\mathbf{r}+\mathbf{p}'\cdot\mathbf{r}')} [\phi(r), \pi(r')] = \cdots = i\nu\delta_{\mathbf{p}, -\mathbf{p}'}$$

$$[\phi_{\mathbf{p}}, \phi_{\mathbf{p}'}] = 0 \qquad [\pi_{\mathbf{p}}, \pi_{\mathbf{p}'}] = 0$$

On introduit alors

$$a_{\mathbf{p}} = \sqrt{\frac{\omega_{\mathbf{p}}}{2\nu}} \left(\phi_{\mathbf{p}} + i \frac{\pi_{\mathbf{p}}}{\omega_{\mathbf{p}}} \right) \qquad a_{\mathbf{p}}^{\dagger} = \sqrt{\frac{\omega_{\mathbf{p}}}{2\nu}} \left(\phi_{\mathbf{p}}^{\dagger} - i \frac{\pi_{\mathbf{p}}^{\dagger}}{\omega_{\mathbf{p}}} \right)$$
$$[a_{\mathbf{p}}, a_{\mathbf{q}}^{\dagger}] = \dots = \delta_{p,q}$$
$$H = \sum_{\mathbf{p}} \omega_{\mathbf{p}} \left(a_{\mathbf{p}}^{\dagger} a_{\mathbf{p}} + \frac{1}{2} \right)$$