

## Action je crois

$$S = \int_{t_i}^{t_f} dt \frac{1}{2} m \mathbf{v}^2(T) \quad \text{non-relativiste}$$

$$\rightarrow -m \int_A^V d\tau = -m \int_A^B dt \sqrt{1 - \mathbf{v}^2} = -m \int_A^B dt \left\{ 1 - \frac{1}{2} \mathbf{v}^2 + \frac{12}{\mathbf{v}^2} - \dots \right\} = -m(t_B - T_a) + \frac{1}{2} \int_A^B d\tau m \mathbf{v}^2 - \dots$$

$$L = -m \sqrt{1 - \mathbf{v}^2}$$

$$\mathbf{p} = \frac{\partial L}{\partial \mathbf{v}} = \frac{m \mathbf{v}}{\sqrt{1 - \mathbf{v}^2}}$$

$$H = \mathbf{p} d\mathbf{v} - L = \frac{m \mathbf{v}^2}{\sqrt{1 - v^2}} + m \sqrt{1 - v^2} = \frac{m}{\sqrt{1 - \mathbf{v}^2}} = \sqrt{\mathbf{p}^2 + m^2}$$

### 4-impulsion

$$p^\mu = (E, \mathbf{p}) = m u^\mu$$

$$u^\mu = \frac{dx^\mu}{d\tau} = \left( \frac{dt}{d\tau}, \frac{d\mathbf{r}}{d\tau} \right) = \left( \frac{1}{\sqrt{1 - v^2}}, \frac{\mathbf{v}}{\sqrt{1 - v^2}} \right)$$

Invariant associé au quadri-vecteur

$$p^\mu p_{\mu} = p^2 = E^2 - \mathbf{p}^2 = m^2 \overset{1}{\cancel{v^2}} = m^2$$

$$p^2 = m^2$$

$$\mathbf{v} = \frac{\mathbf{p}}{E}$$

$$\text{masse nulle} \rightarrow p^2 = 0 \rightarrow T = |\mathbf{p}|$$

$$p_\pi = (m_\pi, \mathbf{0})$$

$$p_\pi = p_\mu + p_\nu$$

$$p_\nu = p_\pi - p_\mu$$

$$p_\nu^2 = p_\pi^2 + p_\mu^2 - 2p_\pi p_\mu$$

$$0 = m_\pi^2 + m_\mu^2 - 2m_\pi E_\mu$$

$$E_\mu = \frac{m_\pi^2 + m_\mu^2}{2m_\pi}$$

$$E_{nu} = m_\pi - E_\mu = \frac{m_\pi^2 - m_\mu^2}{m_\pi} = |\mathbf{p}_\nu| = |\mathbf{p}_\nu|$$

$$|\mathbf{v}_\mu| = \frac{m_\pi^2 - m_\mu^2}{m_\pi^2 - m_\mu^2}$$

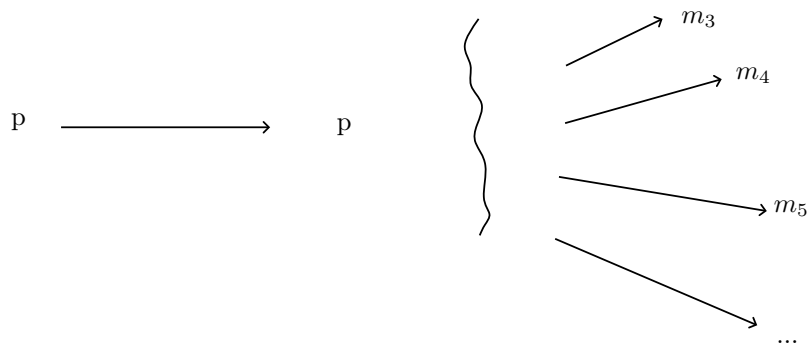


FIGURE 1 – proton incident

Énergie de seuil ? (Plus d'Énergie cinétique à la fin)

$$E = \sum_{i=3}^N m_i$$

$$p_1^\mu + p_2^\mu = \sum_{i=3}^N p_i^\mu$$

$$(p_1^\mu + p_2^\mu)^2 = \left[ \sum_{i=3}^N p_i^\mu \right]^2 = \left( \sum_{i=3}^N m_i \right)^2$$

au seuil  $p_i = (m_i, \mathbf{0})$

$$E_p = \frac{M_{\text{tot}}^2 - 2m_p^2}{2m_p}$$

L'énergie requise va comme le carré des masses.

#### Unités naturelles

$$\hbar c = 197 \text{MeV} \cdot \text{fm}$$

$$\frac{\hbar c}{m_e c^2} = \frac{197 \text{MeV} \cdot \text{fm}}{0,511 \text{MeV}} = 400 \text{fm} \quad \text{longueur d'onde de Compton}$$

Constante de structure fine

$$\alpha = \frac{e^2}{4\pi\hbar c} \sim \frac{1}{137}$$

#### Heaviside-Lorentz

$$\epsilon_0 = 1 \quad \mu_0 = 1$$

$$\frac{\alpha \hbar}{m_e c} = \frac{e^2}{4\pi m_e c^2} = \text{rayon classique de l'électron}$$

$$\frac{\hbar}{m_e \alpha c} = \frac{4\pi \hbar c \hbar}{e^2 m_e c} = \frac{4\pi \hbar^2}{m_e e^2} = \text{rayon de bohr}$$

onde plane

$$\psi(\mathbf{r}) = \frac{1}{\sqrt{\nu}} e^{i\mathbf{p} \cdot \mathbf{r}}$$

Condition au limite périodiques à l'univers (une boîte bien sûr)

$$e^{ip_x L_x} = 1$$

$$\mathbf{p} = 2\pi \left( \frac{n_x}{L_x}, \frac{n_y}{L_y}, \frac{n_z}{L_z} \right)$$

où  $n \in \mathbb{Z}$

$$\Delta p_x = \frac{2\pi}{L_x} \leftrightarrow \Delta n_x = 1$$

$$\Delta p_x \Delta p_y \Delta p_z = \frac{(2\pi)^3}{L_x L_y L_z} = \frac{(2\pi)^3}{\nu}$$

$$\sum_{\mathbf{p}} f_{\mathbf{p}} = \nu \int \frac{dP}{(2\pi)^3} f_{\mathbf{p}}$$

$$\langle \mathbf{p}' | \mathbf{p} \rangle = \delta_{\mathbf{p}\mathbf{p}'} \& \sum_{\mathbf{p}} |\mathbf{p}\rangle \langle \mathbf{p}| = \mathbb{1}$$

$$\langle \mathbf{r} | \mathbf{p} \rangle = \frac{1}{\sqrt{\nu}} e^{i\mathbf{p} \cdot \mathbf{r}}$$

$$\delta_{\mathbf{p}\mathbf{p}'} \rightarrow \frac{(2\pi)^3}{\nu} \delta(\mathbf{p} - \mathbf{p}')$$

$$\text{Normalisation continue} \begin{cases} \langle \mathbf{p}' | \mathbf{p} \rangle = (2\pi)^3 \delta(\mathbf{p} - \mathbf{p}') \\ \int \frac{d^3 P}{(2\pi)^3} \langle \mathbf{p} | \mathbf{p} \rangle = \mathbb{1} \end{cases}$$

On a le problème que  $d^3 P$  n'est pas invariant de Lorentz

$d^3 p dp^0$  en revanche l'est

$$d^3 \gamma dt = d^4 x = d^4 x'$$

Le Jacobien

$$J = 1$$

$$\int \frac{d^3 P}{(2\pi)^3} \rightarrow \int \frac{d^4}{(2\pi)^3} \delta(p^2 - m^2) \theta(p^0)$$

$$\int \frac{d^4 P}{(2\pi)^3} \delta((p^0 - E_{\mathbf{p}})(p^0 + E_{\mathbf{p}})) \Theta(p^0)$$

$$E_{\mathbf{p}} = \sqrt{\mathbf{p}^2 + m^2}$$

$$\delta(\beta x) = \frac{1}{|\beta|} \delta(x)$$

$$\int \frac{d^4 P}{(2\pi)^3 2E_p} \delta(p^0 - E_p) = \int \frac{d^3 P}{(2\pi)^3 2E_p}$$

$$\text{Normalisation relativiste} \begin{cases} \int \frac{d^3 P}{(2\pi)^3 2E_p} \langle \mathbf{p} | \mathbf{p} \rangle = \mathbb{1} \\ \langle \mathbf{p} | \mathbf{p}' \rangle = 2E_{\mathbf{p}} \delta(\mathbf{p} - \mathbf{p}') (2\pi)^2 \end{cases}$$