

# Effet photoélectrique

Application de la règle d'or de Fermi !

$$\partial \mathcal{P}_{1S \rightarrow \mathbf{p}_f}(\mathbf{p}_f, t) = \frac{\pi t}{2\hbar} \int_{D_f} \rho(E_f) \|\langle \mathbf{p}_f | W | \varphi_{1S} \rangle\|^2 \partial(E_f - E_i - \hbar\omega) dE_f d\Omega$$

$$W = -\frac{q}{(c)m} e^{i\mathbf{k} \cdot \mathbf{R}} \mathbf{A}_0 \cdot \mathbf{p}$$

$kr \sim ka_0 \ll 1 \rightarrow$  Approximation dipolaire

$$\langle \mathbf{p}_f | W | \varphi_{1S} \rangle = \int d\mathbf{r}^3 \frac{e^{i\mathbf{p}_f \cdot \mathbf{r}}}{(2\pi\hbar)^{3/2}} \langle \mathbf{r} | \mathbf{D} \cdot \mathbf{E} | \varphi_{1S} \rangle$$

$$\mathbf{r} \cdot \mathbf{E} = xE_x + yE_y + zE_z = rE_0 (\cos\theta \cos\theta_0 + \sin\theta \sin\theta_0 \cos(\varphi - \varphi_0))$$

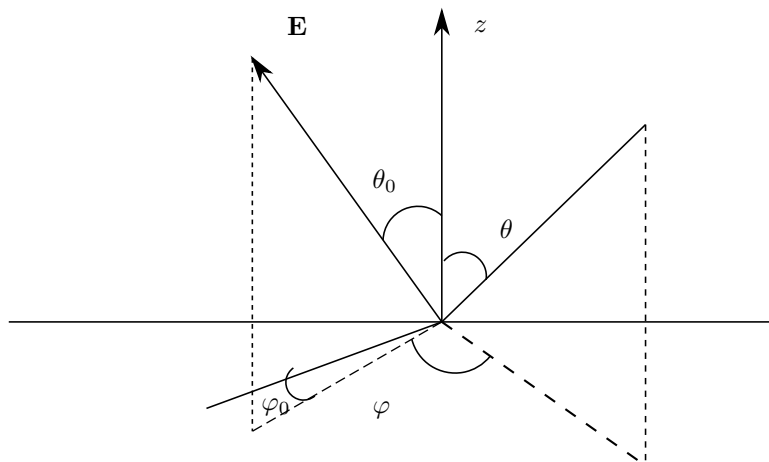


FIGURE 1 – Identification des vecteurs et angles

$$\langle \mathbf{p}_f W | \varphi_{1S} \rangle = -\frac{12qE_0}{\hbar^{2/3}\pi(2a_0^{3/2})} \times \frac{15k_f a_0^5}{[1 + k_f^2 a_0^2]^3} \cos\theta_0$$

...

$$\frac{\partial \mathcal{P}}{\partial t \partial \Omega_f} = \frac{256 a_0^3}{4 \pi \hbar} \underbrace{\frac{E_0^2 q^2}{e^2}}_{E_0^2} \left( \frac{\omega}{\omega_0} - 1 \right)^{3/2} \left( \frac{\omega}{\omega_0} \right)^6 \cos^2 \theta_0$$

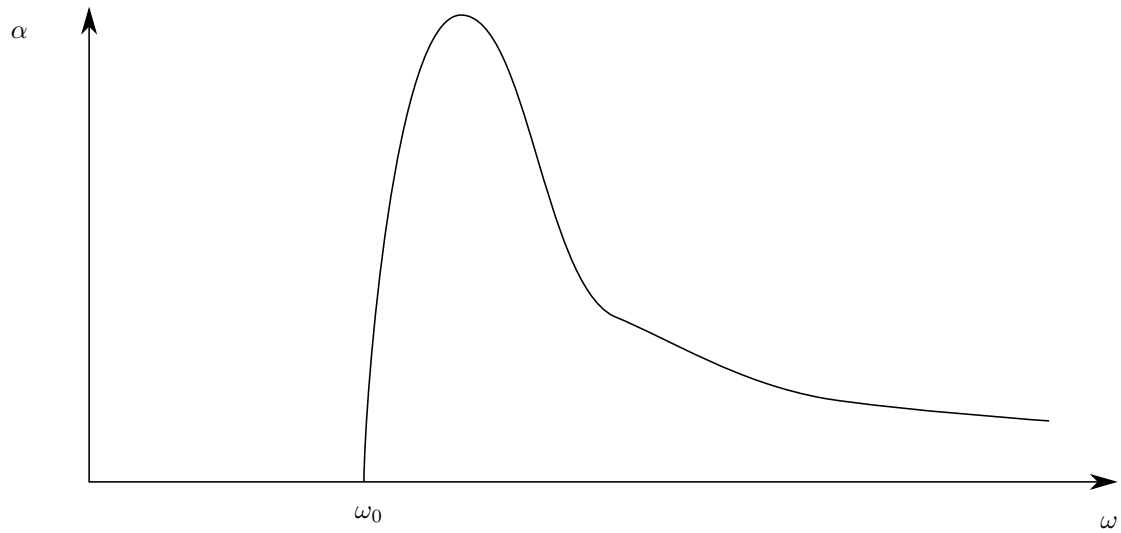


FIGURE 2 – Transition de l’effet photoélectrique

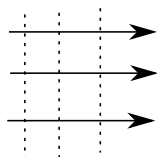
On remarque que l’expression est imaginaire pour  $\omega < \omega_0$

## Théorie de la diffusion (élastique)

$$dn \propto F_i d\Omega$$

$$dn = (\theta, \varphi) F_i d\Omega$$

$\sigma(\theta, \varphi)$  : Section efficace de diffusion



Flux :  $F_i$  (part/surface temps)

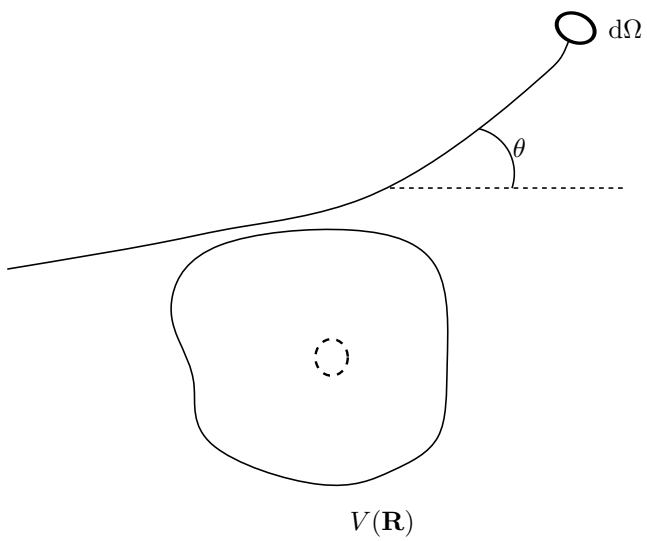


FIGURE 3 – Diffusion

$\square$  = surface (barr =  $10^{-24}\text{cm}^2$ )

Considérations physiques

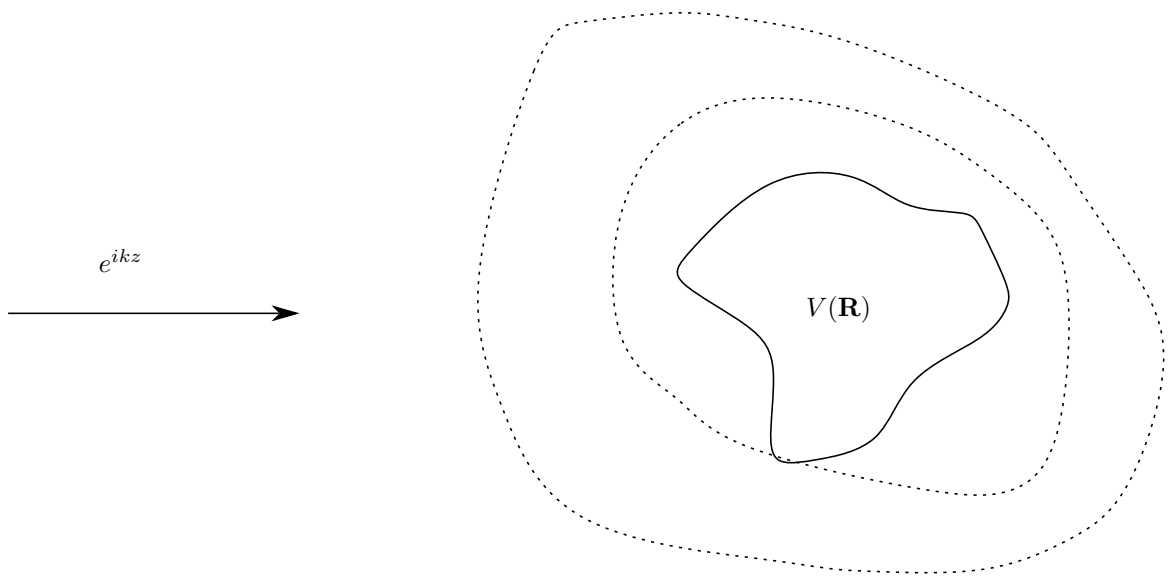


FIGURE 4 – Onde plane incidente