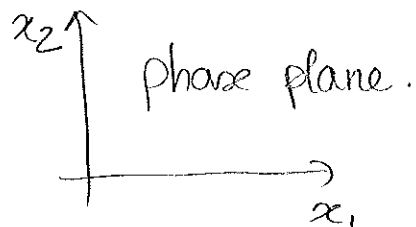


Lecture 2 Phase Plane Analysis

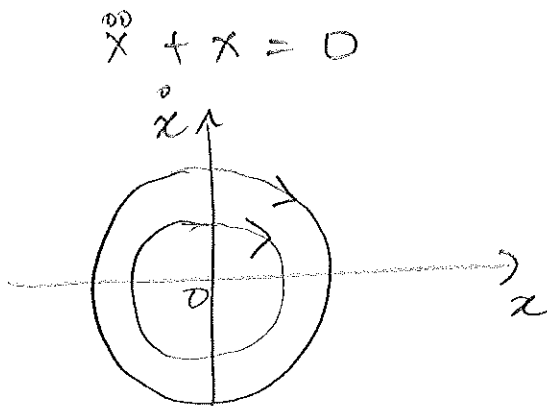
(PI)

Use : Graphically observe how 2nd order system behave for different initial conditions

states of the system $\begin{cases} \dot{x}_1 = f_1(x_1, x_2) = \text{nonlinear function} \\ \dot{x}_2 = f_2(x_1, x_2) = \text{nonlinear function} \end{cases}$



Example Spring mass system



$$\begin{aligned} x_1 &= x \\ x_2 &= \dot{x} \end{aligned}$$

$$\begin{aligned} \dot{x}_2 + x_1 &= 0 \\ \dot{x}_1 &= x_2 \\ \dot{x}_2 &= -x_1 \end{aligned}$$

Singular Points

Equilibrium pt. in the phase plane.

$$\dot{x} = 0 \Rightarrow f_1(x_1, x_2) = 0$$

$$f_2(x_1, x_2) = 0$$

Why called singular pt? slope in phase plane $\frac{dx_2}{dx_1} = \frac{\dot{x}_2}{\dot{x}_1} = \frac{f_2}{f_1}$

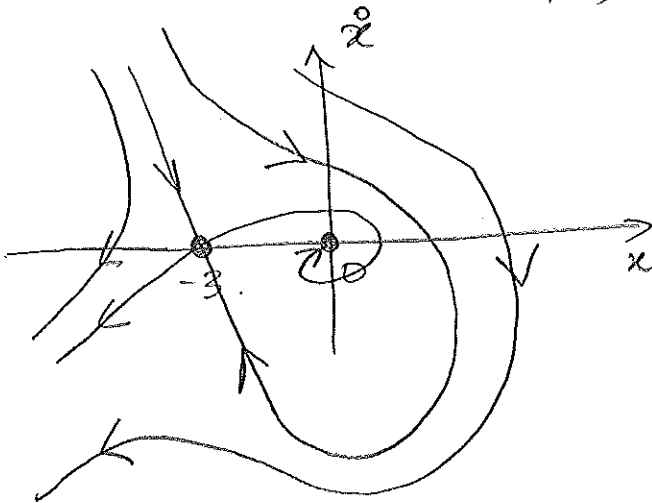
At equm pt, $\frac{dx_2}{dx_1} \Rightarrow$ indeterminate (0/0 form)

Example

(p2)

$$\ddot{x} + 0.6 \dot{x} + 3x + x^2 = 0$$

Two singular pts: $(0,0)$
 $(-3,0)$



Symmetry in Phase-Plane Portraits

$$\ddot{x} + f(x, \dot{x}) = 0 \quad \text{2nd order dynamics}$$

$$\frac{dx_2}{dx_1} = \frac{-f(x_1, x_2)}{\dot{x}}$$

$$x_1 = x$$

$$x_2 = \dot{x}$$

$$\dot{x}_2 = f(x_1, x_2)$$

Symmetry about x_1 axis

$$f(x_1, x_2) = f(x_1, -x_2)$$

Symmetry about x_2 axis

$$f(x_1, x_2) = -f(-x_1, x_2)$$

Symmetry about origin

$$f(x_1, x_2) = -f(-x_1, -x_2)$$

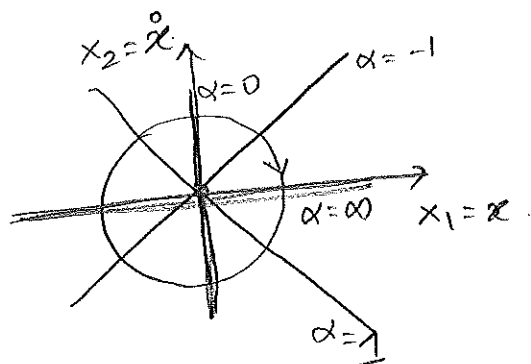
Method of isoclines to construct phase portrait

$$\frac{dx_2}{dx_1} = \frac{f_2(x_1, x_2)}{f_1(x_1, x_2)} = \alpha \Rightarrow \text{locus of pts w/ slope } \alpha$$

$$f_2(x_1, x_2) = \alpha f_1(x_1, x_2)$$

For mass-spring system

$$\frac{dx_2}{dx_1} = -\frac{x_1}{x_2}$$

Phase plane analysis for linear system

second order

$$\begin{cases} \ddot{x}_1 = ax_1 + bx_2 \\ \ddot{x}_2 = cx_1 + dx_2 \end{cases}$$

can be
rewritten as

$$\ddot{x}_1 \pm (a+d)\dot{x}_1 + (cb-ad)x_1 = 0$$

Solution $x(t) = k_1 e^{\lambda_1 t} + k_2 e^{\lambda_2 t}, \lambda_1 \neq \lambda_2$ generalize

$$x(t) = (k_1 + k_2 t) e^{\lambda_1 t}, \lambda_1 = \lambda_2$$

where λ_1 & λ_2 are solutions of characteristic eqn

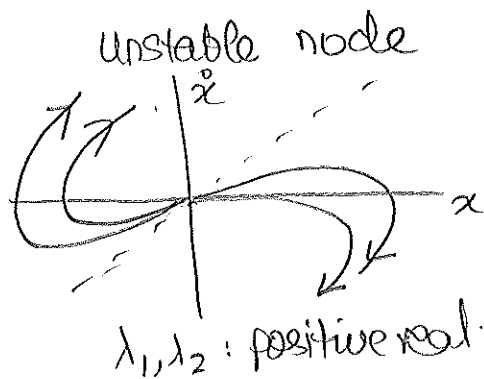
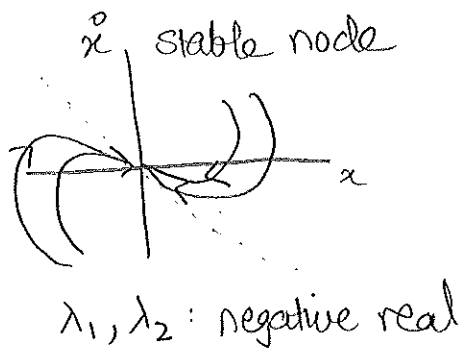
$$s^2 + as + b = 0$$

Roots are

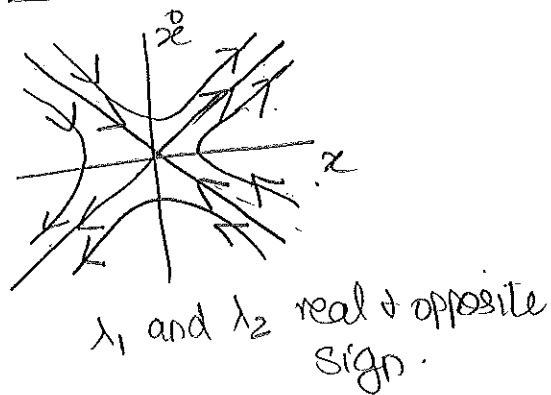
$$\lambda_1 = \frac{-a + \sqrt{a^2 - 4b}}{2} \quad \lambda_2 = \frac{-a - \sqrt{a^2 - 4b}}{2}$$

There is only one equilibrium pt. ($x=0$)

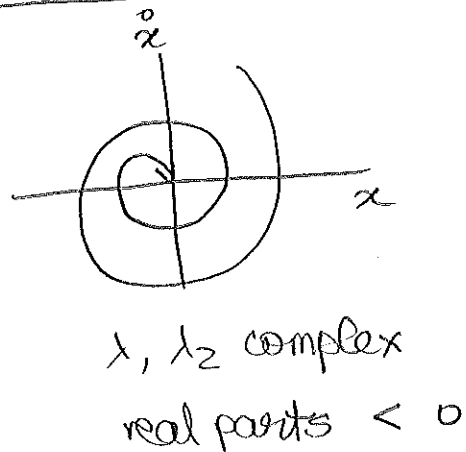
However, system behavior is different depending on λ_1, λ_2 .



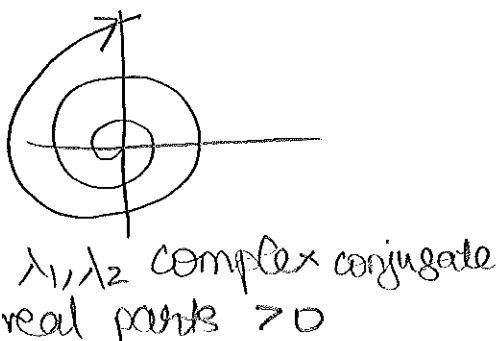
saddle pt.



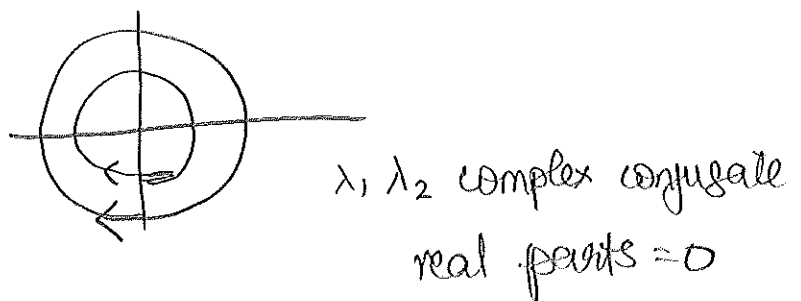
stable focus



unstable focus



center point



Phase Plane Analysis of Nonlinear Systems

(p5)

- Local behavior can be approximated by linear system
- can exhibit multiple eqm pts + limit cycles

Limit cycle \rightarrow Isolated closed curve.

\swarrow
nearby curves
converging or diverging
from it

\downarrow
periodic \rightarrow motion starting on the
curve stays on it forever.

Stable limit cycle \rightarrow All trajectories nearby converge to it as $t \rightarrow \infty$

Unstable " " \rightarrow " " " " diverge from it as $t \rightarrow \infty$

Semistable " " \rightarrow " " " either converge or
diverge as $t \rightarrow \infty$

Three theorems on limit cycles

Poincaré : If a limit cycle exists in the 2nd order

autonomous system, then $N = S + 1$

\downarrow
(no notion of time) $N = \#$ of nodes, center, foci enclosed by limit cycle

$\dot{x} = f(x)$

$S = \#$ of enclosed saddle points

Utility for us: Limit cycle must enclose at least one eqm pt.

Poincare-Bendixson : If a trajectory of a 2nd order autonomous system remains in a finite region J_2 , then one of the following is true: (p6)

(a) the trajectory goes to an equm pt.

(b) " " " tends to an asymptotically stable limit cycle

(c) trajectory is itself a limit cycle.

Bendixson : For the nonlinear system $\dot{x}_1 = f_1(x_1, x_2)$
 $\dot{x}_2 = f_2(x_1, x_2)$

no limit cycle can exist in a region J_2 in which

$\frac{\partial f_1}{\partial x_1} + \frac{\partial f_2}{\partial x_2}$ does not vanish and does not change

sign.